

INITIAL CONDITIONS CONTINUITY OF A NUMERICAL APPROXIMATION FOR KOLMOGOROV EQUATIONS *

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Abstract. We provide theory to characterizes the stability respect to initial conditions of a weak numerical scheme to approximate the solution of a particular family of SPDEs. Our approach consists in solving the associated Kolmogorov equation of the underlying SPDE whit a spectral method. We illustrate our results with numerical experiments.

Key words. Stability, spectral methods, Kolmogorov equation, stochastic parabolic equations.

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1. Introduction. Stochastic Partial Differential Equations (SPDEs) are important tools in modeling complex phenomena, they arise in many fields of knowledge like Physics, Biology, Economy, Finance, etc. Develop efficient numerical methods for simulating SPDEs is very important but also very difficult and challenging.

The Fokker-Planck-Kolmogorov (FPK) equation is a partial differential equation that describes the time evolution of the probability density function of the velocity of a particle under the influence of drag forces and random forces, it is a kind of continuity equation for densities. Citing [8] “parabolic equations on Hilbert spaces appear in mathematical physics to model systems with infinitely many degrees of freedom. Typical examples are provided by spin configurations in statistical mechanics and by crystals in solid state theory. Infinite-dimensional parabolic equations provide an analytic description of infinite dimensional diffusion processes in such branches of applied mathematics as population biology, fluid dynamics, and mathematical finance.” This kind of equations have been deeply studied in the last years, see for instance [2, 4, 6] and the references therein.

Try to finding analytical solutions of FPK associated with SPDEs results impractical. Thus, work with efficient and accurate numerical schemes is crucial. In this way, the spectral methods play an essential role to obtain better schemes—under certain conditions; this sort of methods are more accurate than finite differences of finite elements and need fewer grid points. Here the adjective “better” would be under accuracy, consistency, stability, and other targets properties. In this work, we explore the ability of the method reported in [9] to preserve the continuity respect to initial conditions. That is, if a given problem satisfies certain regularity conditions, then two of its solution remain closed if its initial function conditions are close. So, we desire that a numerical method reproduce this behavior and if it is the case, we say that an underlying method is stable in this context.

Our main contribution is the characterization of mild conditions to assure the

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continuity respect to initial function conditions to a family of SPDEs and the stability of a regarding weak spectral approximation. To the best of our knowledge, this paper is the first in report numerical stability theory for Kolmogorov equations in infinite dimensions.

The stability theory for spectral methods is still under construction and is an active research area. We mention the seminal works of L.N. Trefethen and M.R. Trummer [17], D. Gottlieb et. al. [12] as reference for the deterministic case, and N. Li, J. Fiordilino, and X. Feng, [15] A. Lang, A. Petersson, and A. Thalhammer, [14] for the stochastic version.

This paper is organized as follows. In Section 2 we review the Fokker-Plank-Kolmogorov equation associated with SPDEs in a separable Hilbert space. Section 3 provides conditions to assure stability respect initial conditions and in Section 4 we illustrate our results with numerical experiments.

2. Kolmogorov equations for SPDEs in Hilbert spaces. Let \mathcal{H} be a separable infinite-dimensional Hilbert space with inner product $(\cdot, \cdot)_{\mathcal{H}}$ and norm $\|\cdot\|_{\mathcal{H}}$. We define a Gaussian measure μ with mean zero and nuclear covariance operator Λ with $Tr(\Lambda) < +\infty$.

We focus on the following Kolmogorov equation

$$(2.1) \quad \frac{\partial u}{\partial t} = \frac{1}{2}Tr(QD^2u) + \langle Ax, Du \rangle_{\mathcal{H}} + \langle B(x), Du \rangle_{\mathcal{H}}, \quad x \in D(A).$$

Several authors have proved results on existence and uniqueness of the solution of the Kolmogorov equations, see for instance Da Prato [4] for a survey, Da Prato-Debussche [5] for the Burgers equation, Barbu-Da Prato [1] for the 2D Navier-Stokes stochastic flow in a channel.

2.1. On the Ornstein-Uhlenbeck semigroup. Following [3], in \mathcal{H} we define a Gaussian measure μ with mean zero and nuclear covariance operator Λ with $Tr(\Lambda) < +\infty$ and since $\Lambda : \mathcal{H} \mapsto \mathcal{H}$ is a positive definite, self-adjoint operator then its square-root operator $\Lambda^{1/2}$ is a positive definite, self-adjoint Hilbert-Schmidt operator on \mathcal{H} .

Define the inner product $(g, h)_0 := (\Lambda^{-1/2}g, \Lambda^{-1/2}h)_{\mathcal{H}}$, for $g, h \in \Lambda^{1/2}\mathcal{H}$. Let \mathcal{H}_0 denote the Hilbert subspace of \mathcal{H} , which is the completion of $\Lambda^{1/2}\mathcal{H}$ with respect to the norm $\|g\|_0 := (g, g)_0^{1/2}$. Then \mathcal{H}_0 is dense in \mathcal{H} and the inclusion map $i : \mathcal{H}_0 \hookrightarrow \mathcal{H}$ is compact. The triple $(i, \mathcal{H}_0, \mathcal{H})$ forms an abstract Wiener space.

Let $\mathbb{H} = L^2(\mathcal{H}, \mu)$ denote the Hilbert space of Borel measurable functionals on the probability space with inner product

$$[\Phi, \Psi]_{\mathbb{H}} := \int_{\mathcal{H}} \Phi(v)\Psi(v)\mu(dv), \quad \text{for } \Phi, \Psi \in \mathbb{H},$$

and norm $\|\Phi\|_{\mathbb{H}} := [\Phi, \Phi]_{\mathbb{H}}^{1/2}$. We choose a basis system $\{\varphi_k\}$ for \mathcal{H} .

A functional $\Phi : \mathcal{H} \mapsto \mathbb{R}$, is said to be a smooth simple functional (or a cylinder functional) if there exists a C^∞ -function ϕ on \mathbb{R}^n and n -continuous linear functional l_1, \dots, l_n on \mathcal{H} such that for $h \in \mathcal{H}$

$$(2.2) \quad \Phi(h) = \phi(h_1, \dots, h_n) \quad \text{where} \quad h_i = l_i(h), \quad i = 1, \dots, n.$$

The set of all such functionals will be denoted by $\mathcal{S}(\mathbb{H})$. Denote by $P_k(x)$ the Hermite polynomial of degree k taking values in \mathbb{R} . Then, $P_k(x)$ is given by the following

80 formula

$$81 \quad P_k(x) = \frac{(-1)^k}{(k!)^{1/2}} e^{\frac{x^2}{2}} \frac{d^k}{dx^k} e^{-\frac{x^2}{2}}$$

82 with $P_0 = 1$. It is well-known that $\{P_k(\cdot)\}_{k \in \mathbb{N}}$ is a complete orthonormal system for
 83 $L^2(\mathbb{R}, \mu_1(dx))$ with $\mu_1(dx) = \frac{1}{\sqrt{2\pi}} e^{-\frac{x^2}{2}} dx$. Define the set of infinite multi-index as

$$84 \quad \mathcal{J} = \left\{ \alpha = (\alpha_i, i \geq 1) \mid \alpha_i \in \mathbb{N} \cup \{0\}, \quad |\alpha| := \sum_{i=1}^{\infty} \alpha_i < +\infty \right\}.$$

85 For $\mathbf{n} \in \mathcal{J}$ define the *Hermite polynomial functionals* on \mathcal{H} by

$$86 \quad (2.3) \quad H_{\mathbf{n}}(h) = \prod_{i=1}^{\infty} P_{n_i}(l_i(h)), \quad h \in \mathcal{H}_0, \quad \mathbf{n} \in \mathcal{J},$$

88 and where $l_i(h) = \langle h, \Lambda^{-1/2} \varphi_i \rangle_{\mathcal{H}}$, $i = 1, 2, \dots$ where $P_n(\xi)$ is the usual Hermite
 89 polynomial for $\xi \in \mathbb{R}$ and $n \in \mathbb{N}$.

90 *Remark 2.1.* Notice that $l_i(h)$ is defined only for $h \in \mathcal{H}_0$. However, regarding h
 91 as a μ -random variable in \mathcal{H} , we have $\mathbb{E}(l_i(h)) = \|\varphi_i\|^2 = 1$ and then $l_k(h)$ can be
 92 defined μ -a.e. $h \in \mathcal{H}$, similar to defining a stochastic integral.

93 It is possible to identify the Hermite polynomial functionals defined in (2.3), for
 94 $h \in \mathcal{H}_0$, as a deterministic version of the Wick polynomials defined on the canonical
 95 Wiener space (for further details see [13] for instance).

96 We have the following result (See Theorems 9.1.5 and 9.1.7 in Da Prato-Zabczyk
 97 [8] or Lemma 3.1 in chapter 9 from Chow [3]).

98 **LEMMA 2.2.** *For $h \in \mathcal{H}$ let $l_i(h) = \langle h, \Lambda^{-1/2} \varphi_i \rangle_{\mathcal{H}}$, $i = 1, 2, \dots$. The set $\{H_{\mathbf{n}}\}$ of*
 99 *all Hermite polynomials on \mathcal{H} forms a complete orthonormal system for \mathbb{H} . Hence the*
 100 *set of all functionals are dense in \mathbb{H} . Moreover, we have the direct sum decomposition:*
 101 $\mathbb{H} = \bigoplus_{j=0}^{\infty} K_j$, *where K_j is the subspace of \mathbb{H} spanned by $\{H_{\mathbf{n}} : |\mathbf{n}| = j\}$.*

102 Let Φ be a smooth simple functional given by (2.2). Then the Fréchet derivatives,
 103 $D\Phi = \Phi'$ and $D_2\Phi = \Phi''$ in \mathcal{H} can be computed as follows:

$$104 \quad \begin{aligned} (D\Phi(h), v) &= \sum_{k=1}^n [\partial_k \phi(h_1, \dots, h_n)] l_k(v) \\ (D^2\Phi(h), v) &= \sum_{j,k=1}^n [\partial_j \partial_k \phi(h_1, \dots, h_n)] l_j(v) l_k(v), \end{aligned}$$

105 for any $u, v \in \mathcal{H}$, where $\partial_k \phi = \frac{\partial}{\partial h_k} \phi$. Similarly, for $m > 2$, $D^m \Phi(h)$ is a m -linear form
 106 on \mathcal{H}^m with inner product $(\cdot, \cdot)_m$. We have $[D^m \Phi(h)](v_1, \dots, v_m) = (D^m \Phi(h), v_1 \otimes$
 107 $\dots \otimes v_m)_m$, for $h, v_1, \dots, v_m \in \mathcal{H}$. Consider the following linear stochastic equation

$$108 \quad (2.4) \quad du_t = Au_t dt + dW_t, \quad u_0 = h \in \mathcal{H}.$$

109 Where $A : \mathcal{D}(A) \subset \mathcal{H} \rightarrow \mathcal{H}$ is the infinitesimal generator of a strongly continuous
 110 semigroup e^{tA} in \mathcal{H} . W_t is a Q -Wiener process in \mathcal{H} . Chow in [3, Lemma 9.4.1] has
 111 shown the following result.

112 **LEMMA 2.3.** *Suppose that A and Q satisfy the following:*

1. $A : \mathcal{D}(A) \subset \mathcal{H} \rightarrow \mathcal{H}$ is self-adjoint and there is $\beta > 0$ such that

$$\langle Av, v \rangle_{\mathcal{H}} \leq -\beta \|v\|_{\mathcal{H}}^2 \quad \forall v \in \mathcal{H}.$$

2. A commutes with Q in $\mathcal{D}(A) \subset \mathcal{H}$.

Then (2.4) has a unique invariant measure μ which is a Gaussian measure on \mathcal{H} with zero mean and covariance operator $\Lambda = \frac{1}{2}Q(-A)^{-1} = \frac{1}{2}(-A)^{-1}Q$.

Suppose that A and Q have the same eigenfunctions e_k with eigenvalues λ_k and ρ_k respectively.

It is well-known (See for instance Da Prato and Zabczyk [8]) that the solution of (2.4) is a time-homogeneous Markov process with transition operator P_t defined for $\Phi \in \mathbb{H}$ given by

$$(P_t \Phi)(h) = \int_{\mathcal{H}} \Phi(v) \mu_t^h(dv) = \mathbb{E}[\Phi(u_t^h)].$$

Let $\Phi \in \mathcal{S}(\mathbb{H})$ be a smooth simple functional. By setting $\varphi_k = e_k$ in (2.2), it takes the form $\Phi(h) = \phi(l_1(h), \dots, l_n(h))$, where $l_k(h) = (h, \Lambda^{-1/2} e_k)$. Define a differential operator \mathcal{A}_0 on $\mathcal{S}(\mathbb{H})$ by

$$\mathcal{A}_0 \Phi(v) = \frac{1}{2} \text{Tr}[RD^2 \Phi(v)] + \langle Av, D\Phi(v) \rangle, \quad v \in H,$$

which is well defined, since $D\Phi \in D(A)$ and $\langle Av, D\Phi(v) \rangle = (v, AD\Phi(v))_{\mathcal{H}}$.

The following results have been proved in [3].

LEMMA 2.4. Let P_t be the transition operator as defined by (2.4). Then the following properties hold:

1. $P_t : \mathcal{S}(\mathbb{H}) \rightarrow \mathcal{S}(\mathbb{H})$ for $t \geq 0$.
2. $\{P_t, t \geq 0\}$ is a strongly continuous semigroup on $\mathcal{S}(\mathbb{H})$ so that, for any $\Phi \in \mathcal{S}(\mathbb{H})$, we have $P_0 = I$, $P_{t+s}\Phi = P_t P_s \Phi$, for all $t, s \geq 0$, and $\lim_{t \downarrow 0} P_t \Phi = \Phi$.
3. \mathcal{A}_0 is the infinitesimal generator of P_t so that, for each $\Phi \in \mathcal{S}(\mathbb{H})$,

$$\lim_{t \downarrow 0} \frac{1}{t} (P_t - I) \Phi = \mathcal{A}_0 \Phi.$$

□

LEMMA 2.5. Let $H_n(h)$ be a Hermite polynomial functional given by (2.3). Then the following hold:

$$\mathcal{A}_0 H_n(h) = -\lambda_n H_n(h),$$

$$P_t H_n(h) = \exp\{-\lambda_n t\} H_n(h),$$

for any $\mathbf{n} \in \mathcal{J}$ and $h \in H$, where $\lambda_n = \sum_{i=1}^{\infty} n_i \lambda_i$.

The following Theorem is a Green formula that we will need forward. Its proof can be seen, for instance, in [3, Thm. 3.3, Ch. 9].

THEOREM 2.6. Let $\Phi \in \mathcal{S}(\mathbb{H})$ be a smooth simple functional and let $\mu \sim N(0, \Lambda)$ be a Gaussian measure in \mathcal{H} . Then, for any $g, h \in \mathcal{H}$ the following formula holds

$$\int_{\mathcal{H}} (\Lambda h, D\Phi(v))_{\mathcal{H}} \mu(dv) = \int_{\mathcal{H}} (v, h)_{\mathcal{H}} \Phi(v) \mu(dv).$$

LEMMA 2.7. Assume the conditions for Lemma 2.5 hold. Then, for any $\Phi, \Psi \in \mathcal{S}(\mathbb{H})$, the following Green's formula holds:

$$(2.10) \quad \int_{\mathcal{H}} (\mathcal{A}_0 \Phi) \Psi d\mu = \int_{\mathcal{H}} \Phi (\mathcal{A}_0 \Psi) d\mu = -\frac{1}{2} \int_{\mathcal{H}} (QD\Phi, D\Psi) d\mu .$$

By Lemma 2.2, for $\Phi \in \mathbb{H}$, it can be represented as

$$(2.11) \quad \Phi(v) = \sum_{n=0}^{\infty} \phi_{\mathbf{n}} H_{\mathbf{n}}(v),$$

where $n = |\mathbf{n}|$ and $\mathbf{n} \in \mathcal{J}$. Notice that we can think in \mathbf{n} as a vector of r dimension, i.e. $\mathbf{n} = (n_1, \dots, n_r)$. Let $\alpha_{\mathbf{n}} = \alpha_{n_1} \cdots \alpha_{n_r}$ be a sequence of positive numbers with $\alpha_{\mathbf{n}} > 0$, such that $\alpha_{\mathbf{n}} \rightarrow \infty$ as $n \rightarrow \infty$. Define

$$|||\Phi|||_{k,\alpha} = \left[\sum_{\mathbf{n}} (1 + \alpha_{\mathbf{n}})^k |\phi_{\mathbf{n}}|^2 \right]^{1/2},$$

$$|||\Phi|||_{0,\alpha} = |||\Phi||| = \left[\sum_{\mathbf{n}} |\phi_{\mathbf{n}}|^2 \right]^{1/2},$$

which is $L^2(\mu)$ -norm of Φ . For the given sequence $\alpha = \{\alpha_n\}$, let $\mathbb{H}_{k,\alpha}$ denote the completion of $\mathcal{S}(\mathbb{H})$ with respect to the norm $|||\cdot|||_{k,\alpha}$. Then $\mathbb{H}_{k,\alpha}$ is called a Gauss-Sobolev space of order k with parameter α . The dual space of $\mathbb{H}_{k,\alpha}$ is $\mathbb{H}_{-k,\alpha}$. From now on, we will fix the sequence $\alpha_{\mathbf{n}} = \lambda_{\mathbf{n}}$, where $\lambda_{\mathbf{n}}$ is given in Lemma 2.5. We shall simply denote $\mathbb{H}_{k,\alpha}$ by \mathbb{H}_k and $|||\Phi|||_{k,\alpha}$ by $|||\Phi|||_k$.

The following results ensure the existence of an extension for the operator \mathcal{A}_0 to a domain containing \mathbb{H}_2 . Their proofs can be found in [3] for instance.

THEOREM 2.8. Let the conditions on A and Q in Lemma 2.3 hold. Then $P_t : \mathbb{H} \rightarrow \mathbb{H}$, for $t \geq 0$, is a contraction semigroup with the infinitesimal generator \tilde{A} . The domain of \tilde{A} contains \mathbb{H}_2 and we have $\tilde{A} = \mathcal{A}_0$ in $\mathcal{S}(\mathbb{H})$.

THEOREM 2.9. Let the conditions for Theorem 2.8 hold true. The differential operator \mathcal{A}_0 defined by (2.6) in $\mathcal{S}(\mathbb{H})$ can be extended to be a self-adjoint linear operator A in \mathbb{H} with domain \mathbb{H}_2 .

Since both \tilde{A} and A are extensions of \mathcal{A}_0 to a domain containing \mathbb{H}_2 , they must coincide there.

Given the Gauss-Sobolev space \mathbb{H}_k with norm $|||\cdot|||_k$ we denote its dual space by \mathbb{H}_{-k} with norm $|||\cdot|||_{-k}$. Thus, we have the inclusions, $\mathbb{H}_k \subset \mathbb{H} \subset \mathbb{H}_k$. We denote the duality between \mathbb{H}_k and \mathbb{H}_{-k} by $\langle\langle \Psi, \Phi \rangle\rangle_k$, $\Phi \in \mathbb{H}_k$, $\Psi \in \mathbb{H}_{-k}$. We also set $\mathbb{H}_0 = \mathbb{H}$, with $|||\cdot|||_0 = |||\cdot|||$ and $\langle\langle \cdot, \cdot \rangle\rangle_1 = \langle\langle \cdot, \cdot \rangle\rangle$, $\langle\langle \cdot, \cdot \rangle\rangle_0 = [\cdot, \cdot]$.

2.2. A non linear Kolmogorov equation. Consider the following Kolmogorov equation,

$$\frac{\partial}{\partial t} \Psi(v, t) = \mathcal{A} \Psi(v, t) + \langle B(v), D\Psi(v, t) \rangle_{\mathcal{H}}, \quad \text{a.e. } v \in \mathbb{H}_2,$$

$$\Psi(v, 0) = \phi(v),$$

where, as defined in Theorem 2.8, $\mathcal{A} : \mathbb{H}_2 \rightarrow \mathbb{H}$ is given by

$$(2.12) \quad \mathcal{A} \Phi = \frac{1}{2} \text{Tr}[RD^2 \Phi(v)] + \langle Av, D\Phi(v) \rangle .$$

Hypothesis on B will be specified latter. For now, we will consider that it is a locally Lipschitz function. The additional term $\langle B(v), D\Psi(v, t) \rangle_{\mathcal{H}}$ is defined μ -a.e. $v \in \mathbb{H}_2$. We will allow the initial datum ϕ will be in \mathbb{H} .

We will study a mild solution of the equation (2.12). Let $\lambda > 0$ be a parameter. By changing Ψ to $e^{\lambda t}\Psi$ in (2.12) we get the following equation:

$$\begin{aligned} \frac{\partial}{\partial t} \Psi(v, t) &= \mathcal{A}_\lambda \Psi(v, t) + \langle B(v), D\Psi(v, t) \rangle_{\mathcal{H}}, \quad \text{a.e. } v \in \mathbb{H}_2, \\ \Psi(v, 0) &= \phi(v), \end{aligned}$$

where $\mathcal{A}_\lambda = \mathcal{A} - \lambda I$, with I the identity operator in \mathbb{H} . Clearly, the problems (2.12) and (2.2) are equivalent, as far for the existence and uniqueness questions are concerned. We will work on the problem (2.2).

Denote by P_t the semigroup with infinitesimal generator \mathcal{A}_λ . The existence of P_t is ensured by the Theorem 2.8. Then, we can rewrite the equation (2.2) in an integral form by using the semigroup P_t

$$(2.13) \quad \Psi(v, t) = e^{-\lambda t}(P_t \phi)(v) + \int_0^t e^{-\lambda(t-s)} [P_{t-s}(B, D\Psi_s)](v) ds,$$

where we denote $\phi = \phi(\cdot)$ and $\Psi_s = \Psi(\cdot, s)$. Chow [3] had proved the following lemma.

LEMMA 2.10. *Let $\Psi \in L^2((0, T); \mathbb{H})$ for some $T > 0$. Then, for any $\lambda > 0$ there exists $C_\lambda > 0$ such that*

$$(2.14) \quad \left\| \int_0^t e^{-\lambda(t-s)} P_{t-s} \Psi_s ds \right\|^2 \leq C_\lambda \int_0^t \|\Psi_s\|_{-1}^2 ds, \quad 0 < t \leq T.$$

We now prove the following theorem on existence and uniqueness of a mild solution to (2.2).

THEOREM 2.11. *Suppose that $B : \mathcal{H} \rightarrow \mathcal{H}_0$ satisfies $(B, D\Phi) \in L^2((0, T); \mathbb{H})$ for any $\Phi \in \mathbb{H}$ and*

$$\sup_{v \in \mathcal{H}} \|\Lambda^{-1/2} B(v)\|_{\mathcal{H}} < +\infty.$$

Then, B satisfies

$$(2.15) \quad \left\| (B(v), D\Phi(v)) \right\|_{-1}^2 \leq C \|\Phi(v)\|^2 \quad \text{for any } \Phi \in \mathbb{H}, \quad v \in \mathbb{H}_2,$$

for some $C > 0$. Moreover, for $\Phi \in \mathbb{H}$, the initial-value problem (2.2) has a unique mild solution $\Psi \in C((0, T); \mathbb{H})$.

For the part of the existence and uniqueness of the solution we will adapt the proof of the Theorem 5.2 in Chapter 9 from [3].

Proof. First we will prove (2.15). We have

$$\left\| (B(v), D\Phi(v)) \right\|_{-1}^2 = \sum_{\mathbf{n}} (1 + \lambda_{\mathbf{n}})^{-1} |\phi_{\mathbf{n}}|^2,$$

with

$$(2.16) \quad \phi_{\mathbf{n}} = \left((B(v), D\Phi(V))_{\mathcal{H}}, H_{\mathbf{n}}(v) \right)_{\mathbb{H}} = \int_{\mathcal{H}} (B(v), D\Phi(v))_{\mathcal{H}} H_{\mathbf{n}}(v) \mu(dv).$$

By the [Theorem 2.6](#), in particular [\(2.9\)](#), we have

$$\int_{\mathcal{H}} (\Lambda h, D\Phi(v))_{\mathcal{H}} \mu(dv) = \int_{\mathcal{H}} (v, h)_{\mathcal{H}} \Phi(v) \mu(dv),$$

for all $\Phi \in \mathcal{S}(\mathbb{H})$, $g, h \in \mathcal{H}$ and $\mu \sim N(0, \Lambda)$. Then, in particular, in each direction $H_{\mathbf{n}}$ this formula is still true, so we have

$$\int_{\mathcal{H}} (\Lambda h, D\Phi(v))_{\mathcal{H}} H_{\mathbf{n}}(v) \mu(dv) = \int_{\mathcal{H}} (v, h)_{\mathcal{H}} \Phi(v) H_{\mathbf{n}}(v) \mu(dv) .$$

218 Then, applying this last equality to [\(2.16\)](#) we get

$$\begin{aligned} 219 \quad \phi_n &= \int_{\mathcal{H}} (\Lambda[\Lambda^{-1}B(v)], D\Phi(v))_{\mathcal{H}} H_{\mathbf{n}}(v) \mu(dv) \\ 220 \quad &= \int_{\mathcal{H}} (\Lambda^{-1}B(v), v)_{\mathcal{H}} \Phi(v) H_{\mathbf{n}}(v) \mu(dv) \\ 221 \quad &= \int_{\mathcal{H}} (\Lambda^{-1/2}B(v), \Lambda^{1/2}v)_{\mathcal{H}} \Phi(v) H_{\mathbf{n}}(v) \mu(dv) . \\ 222 \end{aligned}$$

223 Thus,

$$\begin{aligned} 224 \quad (2.17) \quad |\phi_n|^2 &= \left| \int_{\mathcal{H}} (\Lambda^{-1/2}B(v), \Lambda^{1/2}v)_{\mathcal{H}} \Phi(v) H_{\mathbf{n}}(v) \mu(dv) \right|^2 \\ &\leq \int_{\mathcal{H}} |(\Lambda^{-1/2}B(v), \Lambda^{1/2}v)_{\mathcal{H}}|^2 |H_{\mathbf{n}}(v)|^2 \mu(dv) \int_{\mathcal{H}} |\Phi(v)|^2 \mu(dv) . \end{aligned}$$

225 We now focus on the first integral. Let I_1 be the first integral of [\(2.17\)](#). Then,

$$\begin{aligned} 226 \quad I_1 &\leq \int_{\mathcal{H}} \|\Lambda^{-1/2}B(v)\|_{\mathcal{H}}^2 \|\Lambda^{1/2}v\|_{\mathcal{H}}^2 |H_{\mathbf{n}}(v)|^2 \mu(dv) \\ 227 \quad &\leq \sup_{v \in \mathcal{H}} \|\Lambda^{-1/2}B(v)\|_{\mathcal{H}}^2 \int_{\mathcal{H}} \|\Lambda^{1/2}v\|_{\mathcal{H}}^2 |H_{\mathbf{n}}(v)|^2 \mu(dv) \\ 228 \quad &\leq C \int_{\mathcal{H}} \|v\|_{\mathcal{H}}^2 |H_{\mathbf{n}}(v)|^2 \mu(dv) \\ 229 \quad &\leq C . \end{aligned}$$

231 The last inequality follows by using proposition 9.2.10 in page 198 from [\[7\]](#). Then, by
232 using this bound on [\(2.17\)](#) we have.

$$\begin{aligned} 233 \quad |\phi_n|^2 &\leq C \int_{\mathcal{H}} |\Phi(v)|^2 \mu(dv) \\ 234 \quad &\leq C \|\Phi(v)\|^2 . \end{aligned}$$

Thus,

$$\| (B(v), D\Phi(v)) \|_{-1}^2 \leq C \| \Phi(v) \|^2 \sum_{\mathbf{n}} (1 + \lambda_{\mathbf{n}})^{-1} \leq C \| \Phi(v) \|^2 ,$$

236 which proves [\(2.15\)](#).

We now prove the existence and uniqueness of a solution to the initial-value problem (2.2). Let \mathbb{X}_T denote the Banach space $\mathcal{C}([0, T]; \mathbb{H})$ with the sup-norm

$$\|\Psi\|_T := \sup_{0 \leq t \leq T} \|\Psi\|.$$

In \mathbb{X}_T define the linear operator \mathbb{Q} as

$$\mathbb{Q}\Psi = e^{-\lambda t} P_t \Phi + \int_0^t e^{-\lambda(t-s)} P_{t-s}(B, D\Psi_s) ds, \quad \text{for any } \Psi \in \mathbb{X}_T.$$

By Theorem 2.8 P_t is a contraction semigroup, then using this fact and Lemma 2.10 we have

$$\begin{aligned} \|\mathbb{Q}\Psi\|^2 &\leq 2 \left[\|e^{-\lambda t} P_t \Phi\|^2 + \left\| \int_0^t e^{-\lambda(t-s)} P_{t-s}(B, D\Psi_s) ds \right\|^2 \right] \\ &\leq 2 \left[\|\Phi\|^2 + C_\lambda \int_0^t \|(B, D\Psi_s)\|_{-1}^2 ds \right] \\ &\leq 2\|\Phi\|^2 + C_1 \int_0^t \|\Psi_s\|^2 ds, \end{aligned}$$

for some $C_1 > 0$. Hence, $\|\mathbb{Q}\Psi\|_T \leq C(1 + \|\Psi\|_T)$, with $C = C(\Phi, \lambda, T)$. Then, the map $\mathbb{Q} : \mathbb{X}_T \rightarrow \mathbb{X}_T$ is well defined. We now show that is a contraction for a small t . Let $\Psi, \Psi' \in \mathbb{X}_T$. Then

$$\begin{aligned} \|\mathbb{Q}\Psi - \mathbb{Q}\Psi'\|^2 &= \left\| \int_0^t e^{-\lambda(t-s)} P_{t-s}[(B, D\Psi_s) - (B, D\Psi'_s)] ds \right\|^2 \\ &\leq C_\lambda \int_0^t \|(B, D\Psi_s - D\Psi'_s)\|_{-1}^2 ds \\ &\leq C_2 \int_0^t \|\Psi_s - \Psi'_s\|^2 ds, \end{aligned}$$

for some $C_2 > 0$.

It follows that $\|\mathbb{Q}\Psi - \mathbb{Q}\Psi'\|_T \leq \sqrt{C_2 T} \|\Psi - \Psi'\|_T$. Then, for small T , \mathbb{Q} is a contraction on \mathbb{X}_T . Hence the Cauchy problem (2.2) has a unique mild solution. \square

We now prove a theorem on the dependence on initial conditions for the mild solution of (2.2).

THEOREM 2.12. *Suppose that $B : \mathcal{H} \rightarrow \mathcal{H}_0$ satisfies $(B, D\Phi) \in L^2((0, T); \mathbb{H})$ for any $\Phi \in \mathbb{H}$ and*

$$(2.18) \quad \sup_{v \in \mathcal{H}} \|\Lambda^{-1/2} B(v)\|_{\mathcal{H}} < +\infty.$$

Then, the unique mild solution $\Psi \in C((0, T); \mathbb{H})$ for (2.2) depends continuously on the initial conditions.

Proof. We know, with the assumption (2.18), that the existence of a unique mild solution for (2.2) is guaranteed by Theorem 2.11. We will denote by Ψ_t^φ its mild solution at time t with initial condition φ :

$$\Psi_t^\varphi = e^{-\lambda t} P_t \varphi + \int_0^t e^{-\lambda(t-s)} P_{t-s}(B, D\Psi_s^\varphi) ds.$$

Then,

$$\begin{aligned}\Psi_t^\varphi - \Phi_t^\psi &= e^{-\lambda t} P_t \varphi - e^{-\lambda t} P_t \psi + \int_0^t e^{-\lambda(t-s)} P_{t-s} (B, D\Psi_s^\varphi - D\Phi_s^\psi) ds \\ &= e^{-\lambda t} P_t (\varphi - \psi) + \int_0^t e^{-\lambda(t-s)} P_{t-s} (B, D\Psi_s^\varphi - D\Phi_s^\psi) ds.\end{aligned}$$

From this expression we get

$$\begin{aligned}\|\Psi_t^\varphi - \Phi_t^\psi\|^2 &\leq \|e^{-\lambda t} P_t (\varphi - \psi)\|^2 + \left\| \int_0^t e^{-\lambda(t-s)} P_{t-s} (B, D\Psi_s^\varphi - D\Phi_s^\psi) ds \right\|^2 \\ &\leq \|\varphi - \psi\|^2 + C_\lambda \int_0^t \|(B, D\Psi_s^\varphi - D\Phi_s^\psi)\|_{-1}^2 ds \\ &\leq \|\varphi - \psi\|^2 + C_2 \int_0^t \|\Psi_s^\varphi - \Phi_s^\psi\|^2 ds.\end{aligned}$$

Thus, by Gronwall's inequality we obtain

$$(2.19) \quad \|\Psi_t^\varphi - \Phi_t^\psi\|^2 \leq \exp(C_2 t) \|\varphi - \psi\|^2,$$

which implies, $\|\Psi_t^\varphi - \Phi_t^\psi\| \leq \exp(Ct) \|\varphi - \psi\|$. This completes the proof. \square

3. Numerical stability respect to initial conditions. In this section, we prove the continuity with respect to the initial conditions for a numerical approximation of the Kolmogorov equation associated with an SPDE. Here we understand that a numerical scheme is stable respect to initial conditions if this method reproduces the same behavior when the continuous problem satisfies continuity respect initial conditions.

Consider the stochastic differential equation in \mathcal{H}

$$(3.1) \quad dX_t = AX_t dt + B(X_t) dt + \sqrt{Q} dW_t,$$

where the operator $A : \mathcal{D}(A) \subset \mathcal{H} \rightarrow \mathcal{H}$ is the infinitesimal generator of a strongly continuous semigroup e^{tA} in \mathcal{H} , Q is a bounded operator from another Hilbert space \mathcal{U} to \mathcal{H} and $B : \mathcal{D}(B) \subset \mathcal{H} \rightarrow \mathcal{H}$ is a nonlinear mapping.

The equation (3.1) can be associated to a Kolmogorov equation in the next way, we define

$$(3.2) \quad u(t, x) = \mathbb{E}[\varphi(X_t^x)],$$

where $\varphi : \mathcal{H} \rightarrow \mathbb{R}$ and X_t^x is the solution to (3.1) with initial conditions $X_0 = x$ where $x \in \mathcal{H}$. Then u satisfies the Kolmogorov equation (2.1).

We use Lemma 2.2 to write the solution Ψ_t^φ as in a Fourier-Hermite decomposition:

$$(3.3) \quad \Psi_t^\varphi = \sum_{\mathbf{n} \in \mathcal{J}} u_{\mathbf{n}}(t) H_{\mathbf{n}}(x), \quad x \in \mathcal{H}, \quad t \in [0, T].$$

Note that the time-dependent coefficients $u_{\mathbf{n}}(t)$ depend on the functional and on the initial condition but it is not a function of the initial condition. First we prove an auxiliary result.

LEMMA 3.1. Set $\{P_k(\xi)\}_{k \in \mathbb{N}}$ the family of normalized Hermite polynomials in \mathbb{R} . For every $k \in \mathbb{N}$ and $\xi, \eta \in \mathbb{R}$ such that $\eta < \xi$ we have that

$$(3.4) \quad P_k(\xi) - P_k(\eta) = C(k)Pe_{k+1}(\gamma) \cdot (\xi - \eta),$$

where $\gamma \in (\eta, \xi)$ and $C(k) = \frac{(-1)^k}{(k+1)(k!)^{1/2}}$. Moreover, $Pe_k(x)$ is the unnormalized Hermite polynomial of k degree.

Proof. We know that $P_k(\xi) = \frac{(-1)^k}{(k!)^{1/2}} e^{\xi^2/2} \frac{d}{d\xi^k} e^{-\xi^2/2}$. Set $c(k) = (-1)^k (k!)^{-1/2}$, then

$$\begin{aligned} P_k(\xi) - P_k(\eta) &= c(k) \left[e^{\xi^2/2} \frac{d}{d\xi^k} e^{-\xi^2/2} - e^{\eta^2/2} \frac{d}{d\eta^k} e^{-\eta^2/2} \right] \\ &= c(k) \left[e^{x^2/2} \frac{d}{dx^k} e^{-x^2/2} \Big|_{x=\eta}^\xi \right] \\ &= c(k) \int_\eta^\xi F_k(x) dx, \end{aligned}$$

where F_k is a continuous function such that $F'_k(x) = e^{x^2/2} \frac{d}{dx^k} e^{-x^2/2}$. In fact, denoting by $Pe_k(x)$ the unnormalized Hermite polynomial of k degree, results

$$F'_k(x) = e^{x^2/2} \frac{d}{dx^k} e^{-x^2/2} = Pe_k(x),$$

and since the Hermite polynomials constitute an Appell sequence we have that

$$F'_k(x) = Pe_k(x) = \frac{1}{k+1} Pe'_{k+1}(x),$$

which implies that $F_k(x) = \frac{1}{k+1} Pe_{k+1}(x)$. Now, since $F_k(x)$ is a continuous function, then there exists $\gamma \in (\eta, \xi)$ such that

$$\int_\eta^\xi F_k(x) dx = F_k(\gamma) \cdot (\xi - \eta).$$

All these implies that $P_k(\xi) - P_k(\eta) = c(k)F_k(\gamma) \cdot (\xi - \eta)$. From this expression the lemma follows immediately. \square

We will use some technical results on the SPDE to prove the following result—the main result of this section.

THEOREM 3.2. Assume that the eigenvalues of Λ , satisfies that for every $k \in \mathbb{N}$, $\lambda_k < \lambda_{k+1} \rightarrow \infty$. Assume that the functional φ is Lipschitz. Then, the numeric approximation Ψ_t^φ (given by (3.3)) to the solution of the Kolmogorov equation $\Psi \in C((0, T); \mathbb{H})$ depends continuously on the initial conditions.

Proof. Let $x, y \in H$ be two different initial values. We want to estimate $\Psi_t^x - \Psi_t^y$. By definition,

$$(3.5) \quad \Psi_t^x = \sum_{\bar{n} \in \mathcal{I}} u_{\bar{n}}^x(t) H_{\bar{n}}(x) .$$

328 Thus,

$$\begin{aligned}
 \Psi_t^x - \Psi_t^y &= \sum_{\bar{n} \in \mathcal{J}} u_{\bar{n}}^x(t) H_{\bar{n}}(x) - \sum_{\bar{n} \in \mathcal{J}} u_{\bar{n}}^y(t) H_{\bar{n}}(y) \\
 (3.6) \quad &= \sum_{\bar{n} \in \mathcal{J}} \left[u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t) \right] H_{\bar{n}}(x) + \sum_{\bar{n} \in \mathcal{J}} u_{\bar{n}}^y(t) \left[H_{\bar{n}}(x) - H_{\bar{n}}(y) \right].
 \end{aligned}$$

330 We focus on the first term in (3.6). From the definition of the initial condition
 331 we obtain the following expression for the time-dependent coefficient

$$\begin{aligned}
 332 \quad u_{\bar{n}}^x(t) &= \int_{\mathcal{H}} H_{\bar{n}}(x) \mathbb{E}[\varphi(X_t^x)] \mu(dx) . \\
 333
 \end{aligned}$$

334 From this we get

$$\begin{aligned}
 335 \quad u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t) &= \int_{\mathcal{H}} H_{\bar{n}}(x) \mathbb{E}[\varphi(X_t^x)] \mu(dx) - \int_{\mathcal{H}} H_{\bar{n}}(y) \mathbb{E}[\varphi(X_t^y)] \mu(dy) \\
 336 \quad &= \int_{\mathcal{H} \times \mathcal{H}} H_{\bar{n}}(x) \mathbb{E}[\varphi(X_t^x)] \mu(dx) \mu(dy) \\
 337 \quad &\quad - \int_{\mathcal{H} \times \mathcal{H}} H_{\bar{n}}(y) \mathbb{E}[\varphi(X_t^y)] \mu(dx) \mu(dy) \\
 338 \quad &= \int_{\mathcal{H} \times \mathcal{H}} H_{\bar{n}}(x) \left(\mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)] \right) \mu(dx) \mu(dy) \\
 339 \quad &\quad + \int_{\mathcal{H} \times \mathcal{H}} \left(H_{\bar{n}}(x) - H_{\bar{n}}(y) \right) \mathbb{E}[\varphi(X_t^y)] \mu(dx) \mu(dy) . \\
 340
 \end{aligned}$$

341 Then, by the Cauchy-Schwartz inequality, we obtain

$$\begin{aligned}
 |u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t)|^2 &\leq \left| \int_{\mathcal{H} \times \mathcal{H}} H_{\bar{n}}(x) (\mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)]) \mu(dx) \mu(dy) \right|^2 \\
 &\quad + \left| \int_{\mathcal{H} \times \mathcal{H}} (H_{\bar{n}}(x) - H_{\bar{n}}(y)) \mathbb{E}[\varphi(X_t^y)] \mu(dx) \mu(dy) \right|^2 \\
 &\leq \int_{\mathcal{H} \times \mathcal{H}} H_{\bar{n}}^2(x) \mu(dx) \mu(dy) \\
 &\quad \times \int_{\mathcal{H} \times \mathcal{H}} |\mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)]|^2 \mu(dx) \mu(dy) \\
 342 \quad (3.7) \quad &\quad + \int_{\mathcal{H} \times \mathcal{H}} \mathbb{E}^2[\varphi(X_t^y)] \mu(dx) \mu(dy) \\
 &\quad \times \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
 &= \int_{\mathcal{H} \times \mathcal{H}} |\mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)]|^2 \mu(dx) \mu(dy) \\
 &\quad + \int_{\mathcal{H} \times \mathcal{H}} \mathbb{E}^2[\varphi(X_t^y)] \mu(dx) \mu(dy) \\
 &\quad \times \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) .
 \end{aligned}$$

343 We now estimate the norm of the expression (3.6) with the help of (3.7).

$$\begin{aligned}
& \|\Psi_t^x - \Psi_t^y\|_{(L^2(\mathcal{H}, \mu))^2}^2 = \int_{\mathcal{H} \times \mathcal{H}} |\Psi_t^x - \Psi_t^y|^2 \mu(dx) \mu(dy) \\
& \leq \int_{\mathcal{H} \times \mathcal{H}} \left| \sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t)] H_{\bar{n}}(x) \right|^2 \mu(dx) \mu(dy) \\
& \quad + \int_{\mathcal{H} \times \mathcal{H}} \left| \sum_{\bar{n} \in \mathcal{J}} u_{\bar{n}}^y(t) [H_{\bar{n}}(x) - H_{\bar{n}}(y)] \right|^2 \mu(dx) \mu(dy) \\
& \leq \int_{\mathcal{H} \times \mathcal{H}} \sum_{\bar{n} \in \mathcal{J}} |u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t)|^2 H_{\bar{n}}^2(x) \mu(dx) \mu(dy) \\
& \quad + \int_{\mathcal{H} \times \mathcal{H}} \sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^y(t)]^2 \sum_{\bar{n} \in \mathcal{J}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
& = \sum_{\bar{n} \in \mathcal{J}} |u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t)|^2 \int_{\mathcal{H} \times \mathcal{H}} H_{\bar{n}}^2(x) \mu(dx) \mu(dy) \\
& \quad + \sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^y(t)]^2 \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
& = \sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^y(t)]^2 \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
& \quad + \sum_{\bar{n} \in \mathcal{J}} |u_{\bar{n}}^x(t) - u_{\bar{n}}^y(t)|^2 \\
& = \int_{\mathcal{H} \times \mathcal{H}} \left| \mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)] \right|^2 \mu(dx) \mu(dy) \\
& \quad + \int_{\mathcal{H} \times \mathcal{H}} \mathbb{E}^2[\varphi(X_t^y)] \mu(dx) \mu(dy) \\
& \quad \times \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
& \quad + \sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^y(t)]^2 \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) .
\end{aligned}$$

344 (3.8)

345 Note that $\mathbb{E}^2[\varphi(X_t^y)] = u^2(t, x) \in L^2(\mathcal{H}, \mu)$, therefore the first integral in the
 346 second term is a continuous bounded function of t . Moreover, $\sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^y(t)]^2$ is
 347 the $L^2(\mathcal{H}, \mu)$ -norm of the function $u(t, x)$, then the series converges and it is also a
 348 continuous bounded function of t . Thus, from (3.8) we get

$$\begin{aligned}
& \|\Psi_t^x - \Psi_t^y\|_{(L^2(\mathcal{H}, \mu))^2}^2 \leq \int_{\mathcal{H} \times \mathcal{H}} \left| \mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)] \right|^2 \mu(dx) \mu(dy) \\
& \quad + f(t) \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) ,
\end{aligned}$$

349 (3.9)

350 where $f(t) = \sum_{\bar{n} \in \mathcal{J}} [u_{\bar{n}}^y(t)]^2 + \int_{\mathcal{H}} \mathbb{E}^2[\varphi(X_t^y)] \mu(dy)$.

351 From the proof of Theorem 2.12 (see (2.19)) we know that

$$\begin{aligned}
 352 \quad (3.10) \quad & \|\Psi_t^\varphi - \Phi_t^\psi\|^2 = \int_{\mathcal{H} \times \mathcal{H}} \left| \mathbb{E}[\varphi(X_t^x)] - \mathbb{E}[\varphi(X_t^y)] \right|^2 \mu(dx) \mu(dy) \\
 & \leq \exp(Ct) \int_{\mathcal{H} \times \mathcal{H}} \|x - y\|_{\mathcal{H}}^2 \mu(dx) \mu(dy) \\
 & = \exp(Ct) \|x - y\|^2.
 \end{aligned}$$

353 Therefore the first term in the right side of (3.9) is bounded by (3.10).

354 We now focus on the second term in the last inequality. Notice that for every
 355 $\bar{n} \in \mathcal{J}$ we have

$$\begin{aligned}
 356 \quad (3.11) \quad & H_{\bar{n}}(x) - H_{\bar{n}}(y) = \prod_{i=1}^{\infty} [P_{n_i}(\xi_i) - P_{n_i}(\eta_i)], \\
 357 &
 \end{aligned}$$

358 where $\xi_i = \langle x, \Lambda^{-1/2} e_i \rangle_{\mathcal{H}}$ and $\eta_i = \langle y, \Lambda^{-1/2} e_i \rangle_{\mathcal{H}}$ (see (2.3) and lines after that for
 359 the definition). Hence, applying Lemma 3.1 to equation (3.11) we have that

$$\begin{aligned}
 360 \quad (3.12) \quad & H_{\bar{n}}(x) - H_{\bar{n}}(y) = \prod_{i=1}^{\infty} C(i) P e_{i+1}(\gamma_i) \cdot (\xi_i - \eta_i) \\
 & = \prod_{i=1}^{\infty} C(i) P e_{i+1}(\gamma_i) \langle x - y, \Lambda^{-1/2} e_i \rangle_{\mathcal{H}},
 \end{aligned}$$

361 here $\gamma_i \in (\xi_i \wedge \eta_i, \xi_i \vee \eta_i)$ for every $i \in \mathbb{N}$. Then

$$\begin{aligned}
 & \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
 & = \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} \left| \prod_{i=1}^{\infty} C(i) P e_{i+1}(\gamma_i) \langle x - y, \Lambda^{-1/2} e_i \rangle_{\mathcal{H}} \right|^2 \mu(dx) \mu(dy) \\
 & = \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} \prod_{i=1}^{\infty} [C(i) P e_{i+1}(\gamma_i)]^2 \left| \langle x - y, \Lambda^{-1/2} e_i \rangle_{\mathcal{H}} \right|^2 \mu(dx) \mu(dy) \\
 362 \quad (3.13) \quad & \leq \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} \prod_{i=1}^{\infty} [C(i) P e_{i+1}(\gamma_i)]^2 \|x - y\|_{\mathcal{H}}^2 \|\Lambda^{-1/2} e_i\|_{\mathcal{H}}^2 \mu(dx) \mu(dy) \\
 & = \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} \prod_{i=1}^{\infty} [C(i) P e_{i+1}(\gamma_i)]^2 \|x - y\|_{\mathcal{H}}^2 \lambda_i^{-1} \|e_i\|_{\mathcal{H}}^2 \mu(dx) \mu(dy) \\
 & = \|x - y\|_{\mathcal{H}}^2 \sum_{\bar{n} \in \mathcal{J}} \prod_{i=1}^{\infty} [C(i)]^2 \lambda_i^{-1} \int_{\mathcal{H} \times \mathcal{H}} [P e_{i+1}(\gamma_i)]^2 \mu(dx) \mu(dy).
 \end{aligned}$$

363 Recall that for every $i \in \mathbb{N}$ we have that $\gamma_i \in (\xi_i \wedge \eta_i, \xi_i \vee \eta_i)$, set $\hat{\gamma}_i \in (\xi_i \wedge \eta_i, \xi_i \vee \eta_i)$
 364 such that $P e_i^2(\gamma_i) \leq P e_{i+1}^2(\hat{\gamma}_i)$ for every $\gamma_i \in (\xi_i \wedge \eta_i, \xi_i \vee \eta_i)$, notice that the existence

365 of $\hat{\gamma}_i$ is guaranteed since $Pe_{i+1}^2(\cdot)$ is a continuous function. Then, from (3.13) we get

$$\begin{aligned}
 & \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \\
 366 \quad (3.14) \quad & \leq \|x - y\|_{\mathcal{H}}^2 \sum_{\bar{n} \in \mathcal{J}} \prod_{i=1}^{\infty} [C(i)]^2 \lambda_i^{-1} [Pe_{i+1}(\hat{\gamma}_i)]^2 \int_{\mathcal{H}} \int_{\mathcal{H}} \mu(dx) \mu(dy) \\
 & = \|x - y\|_{\mathcal{H}}^2 \sum_{\bar{n} \in \mathcal{J}} \prod_{i=1}^{\infty} [C(i)]^2 \lambda_i^{-1} [Pe_{i+1}(\hat{\gamma}_i)]^2.
 \end{aligned}$$

367 Here, we recall that $C(i) = \frac{(-1)^i}{(i+1)(i!)^{1/2}}$ then $\frac{(-1)^i}{[(i+1)!]^{1/2}} Pe_{i+1}(\hat{\gamma}_i)$ is the normalized
 368 Hermite polynomial of $i+1$ degree evaluated on $\hat{\gamma}_i$ which is bounded by a constant
 369 C for every $i \in \mathbb{N}$. Moreover, since $\lambda_k < \lambda_{k+1} \rightarrow \infty$ then this implies that

$$370 \quad (3.15) \quad \sum_{\bar{n} \in \mathcal{J}} \prod_{i=1}^{\infty} [C(i)]^2 \lambda_i^{-1} [Pe_{i+1}(\hat{\gamma}_i)]^2 \leq C \sum_{\bar{n} \in \mathcal{J}} \prod_{i=1}^{\infty} \lambda_i^{-1} (i+1)^{-1} \leq C,$$

372 where C is a finite constant. Putting together (3.13) and (3.15) we get that

$$373 \quad (3.16) \quad \sum_{\bar{n} \in \mathcal{J}} \int_{\mathcal{H} \times \mathcal{H}} |H_{\bar{n}}(x) - H_{\bar{n}}(y)|^2 \mu(dx) \mu(dy) \leq C \|x - y\|_{\mathcal{H}}.$$

375 Putting together inequalities (3.9), (3.10) and (3.16) we obtain

$$376 \quad (3.17) \quad \|\Psi_t^x - \Psi_t^y\|_{(L^2(\mathcal{H}, \mu))^2}^2 \leq \exp(Ct) \int_{\mathcal{H} \times \mathcal{H}} \|x - y\|_{\mathcal{H}}^2 \mu(dx) \mu(dy) + f(t) \|x - y\|_{\mathcal{H}}.$$

378 Now, if $\|x - y\|_{\mathcal{H}} \leq \delta$, then from (3.17) we get $\|\Psi_t^x - \Psi_t^y\|_{(L^2(\mathcal{H}, \mu))^2} \leq G(t)\delta$. \square

379 *Remark 3.3.* If we consider in addition the supremum norm on t , then from (3.17)
 380 we get

$$\begin{aligned}
 381 \quad (3.18) \quad & \sup_{0 \leq t \leq T} \|\Psi_t^x - \Psi_t^y\|_{(L^2(\mathcal{H}, \mu))^2}^2 \leq C \|x - y\|_{\mathcal{H}}^2 \sup_{0 \leq t \leq T} f(t) \\
 & + \exp(CT) \int_{\mathcal{H} \times \mathcal{H}} \|x - y\|_{\mathcal{H}}^2 \mu(dx) \mu(dy).
 \end{aligned}$$

382 Notice that $f(t)$ is differentiable and continuous, then $\sup_{0 \leq t \leq T} f(t) \leq C$, then from
 383 (3.18) we obtain

$$\begin{aligned}
 384 \quad (3.19) \quad & \sup_{0 \leq t \leq T} \|\Psi_t^x - \Psi_t^y\|_{(L^2(\mathcal{H}, \mu))^2} \leq C \|x - y\|_{\mathcal{H}} \\
 & + \exp(CT) \int_{\mathcal{H} \times \mathcal{H}} \|x - y\|_{\mathcal{H}}^2 \mu(dx) \mu(dy).
 \end{aligned}$$

385 From this inequality it is possible to show the continuously dependence on the initial
 386 conditions for this norm.

387 **4. Numerical experiments.** In this section we run numerical experiments to
 388 illustrate that our scheme preserves the underlying initial condition continuity. To
 389 this end, we solve a stochastic version of the Fisher and Burgers PDEs with two near
 390 initial function conditions $x(\xi)$, $\hat{x}(\xi)$. In [16] we provide a GitHub repository with a
 391 Python implementation to reproduce the following figures. We also provide in [10, 11],
 392 the 3D on-line plotly versions of Figures 2 and 5.

393 **Stochastic Fisher-KPP equation in an interval.** Let $\mathcal{H} = L^2(0, 1)$. We
 394 consider the stochastic Fisher-KPP equation

$$\begin{aligned} dX(t, \xi) &= [\nu \partial_\xi^2 X(t, \xi) + X(t, \xi)(1 - X(t, \xi))] dt + dW(t, \xi), \\ (4.1) \quad X(t, 0) &= X(t, 1) = 0, \quad t > 0, \\ X(0, \xi) &\in \mathcal{H}, \quad \xi \in [0, 1], \end{aligned}$$

396 in the interval $[0, 1]$ and with initial function conditions $x(\xi)$ and $\hat{x}(\xi)$. In order to
 397 fix this initial function conditions close, we use for our experiments

$$(4.2) \quad x(\xi) := \text{sech}^2(5(\xi - 0.5)), \quad \hat{x}(\xi) := \sum_{k=0}^N T_k(x(\xi)),$$

399 where $T_k(\cdot)$ denotes the Chebyshev polynomial of the first kind. That is, $\hat{x}(\cdot)$ is the
 400 Chebyshev truncated expansion of $x(\cdot)$.

401 **Figure 1** displays the plots of this initial conditions. In **Figure 2** we observe how
 402 the mentioned approximations remains close—blue color scale denotes the solution of
 403 equation 4.1 with initial function condition $x(\xi)$, while yellow color corresponds to the
 404 approximation with initial condition \hat{x} . Since we employ transparency to obtain this
 405 3D plot, the purple scale results from the closeness of the solutions. Further, **Figure 3**
 406 suggest the conclusion of **Theorem 3.2**, that is, the solutions of equation (4.1) are
 407 continuous respect to initial conditions and satisfies the estimation (3.10).

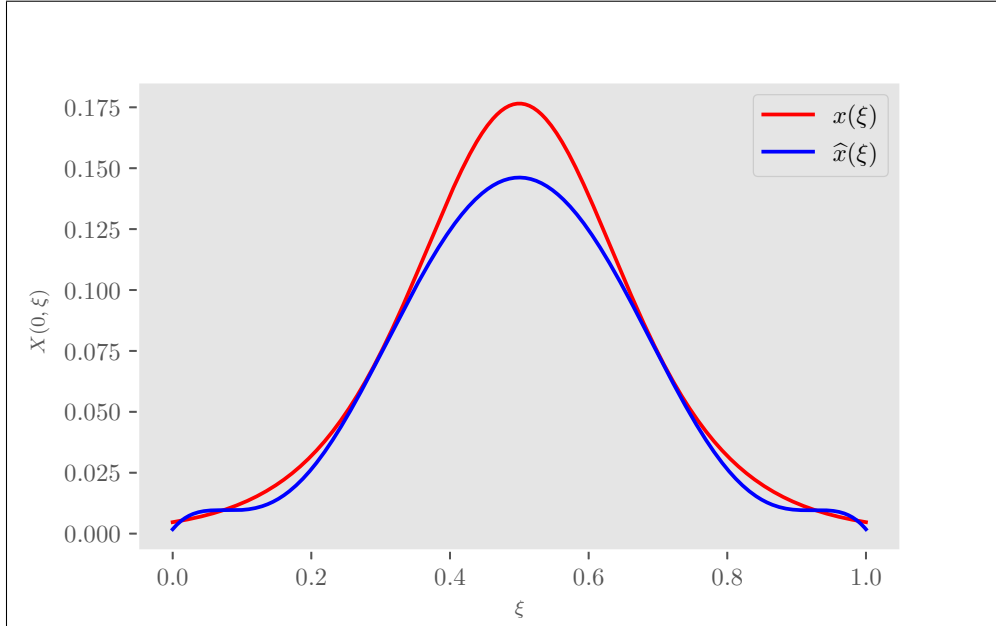


Fig. 1: Numerical Solution of the Fisher-KPP (4.1) with initial conditions x, \hat{x} at time $t = 0$.

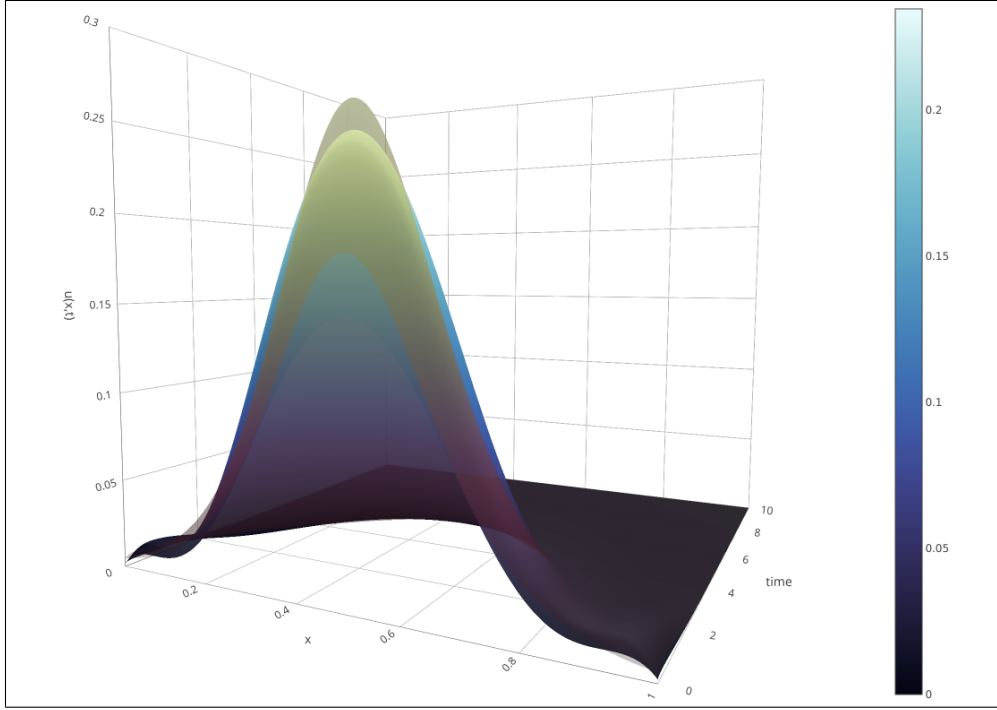


Fig. 2: Likening between two solution with closed initial conditions x, \hat{x} of the stochastic Fisher-KPP (4.1). See [11] to obtain other camera perspectives.

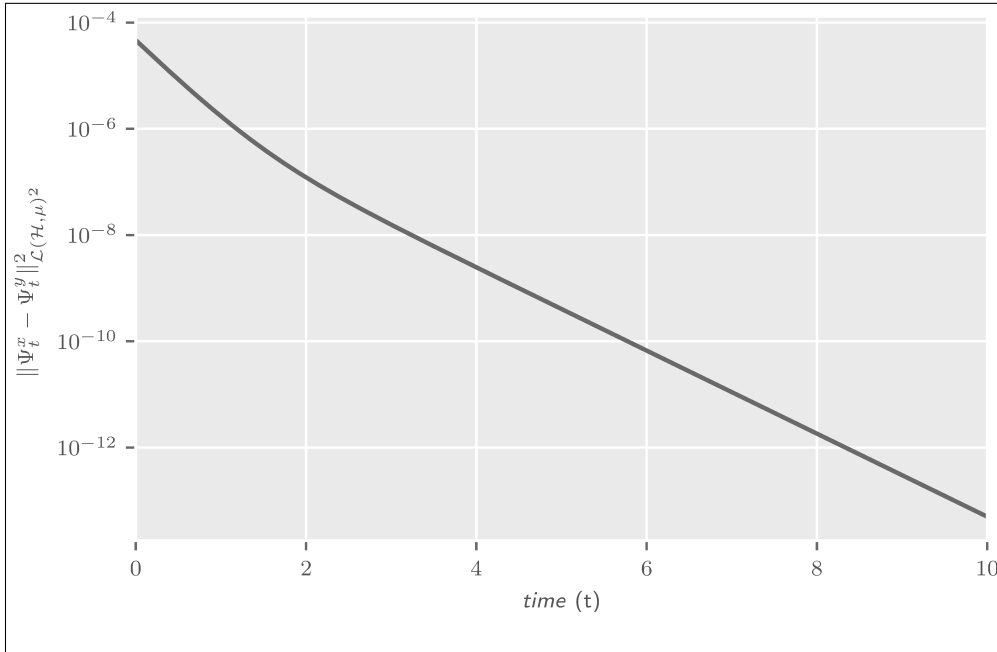


Fig. 3: $\mathcal{L}^2(\mathcal{H}, \mu)$ distance between two solutions of the stochastic Fisher PDE with initial conditions $x = x(\xi)$, and $y = \hat{x}(\xi)$.

408 **Figure 2** illustrates the distance between initial conditions. The yellow pallet with
 409 transparency and a blue scale highlight the zones where the two solutions are close.
 410 Thus, the zones where the color is purple denotes, where the two solutions of SPDEs
 411 are close. According to the $\mathcal{L}(\mathcal{H}, \mu)$ -distance between the two underlying solutions,
 412 **Figure 3** confirms the above argument.

413 **Stochastic Burgers equation.** Let $\mathcal{H} = L^2(0, 1)$, consider the stochastic Burg-
 414 ers equation in the interval $[0, 1]$

$$\begin{aligned} dX(t, \xi) &= \left[\nu \partial_\xi^2 X(t, \xi) + \frac{1}{2} \partial_\xi X^2(t, \xi) \right] dt + dW(t, \xi), \\ (4.3) \quad X(t, 0) &= X(t, 1) = 0, \quad t > 0, \\ X(0, \xi) &= x(\xi), \quad x \in \mathcal{H}. \end{aligned}$$

416 As in the above experiment, we use the initial conditions $x(\xi)$ and its truncated
 417 Chebyshev expansion

$$(4.4) \quad x(\xi) := \sin(\pi\xi), \quad \hat{x}(\xi) := \sum_{k=0}^N T_k x(\xi).$$

419 **Figures 4 to 6** illustrate a similar argument presented in the above experiment.

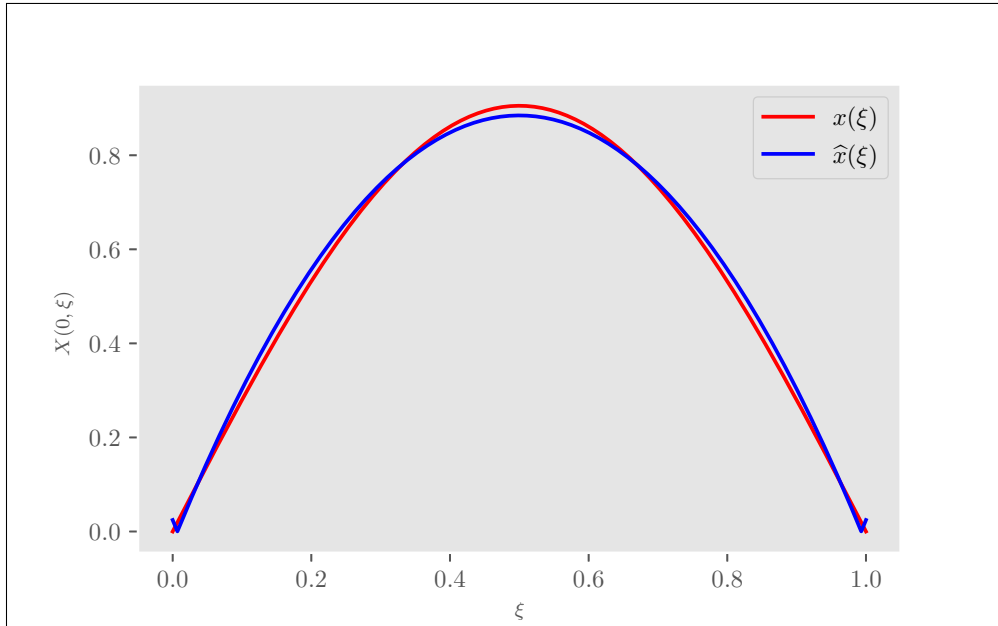


Fig. 4: Numerical Solution of the Burgers (4.3) with initial conditions $x(\xi)$, $\hat{x}(\xi)$.

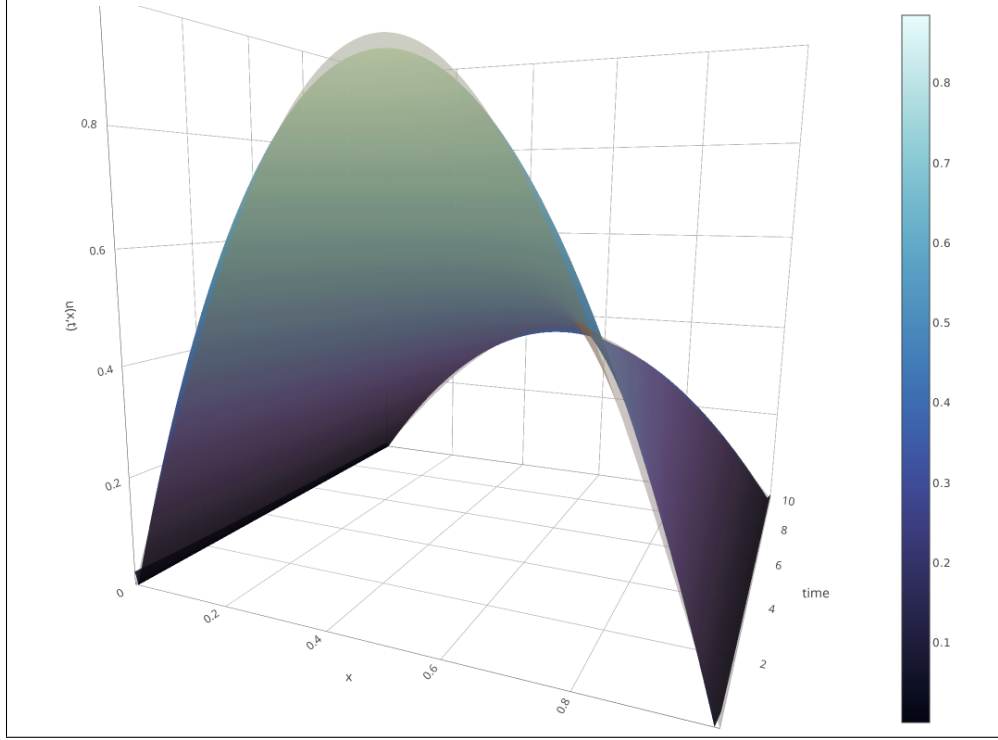


Fig. 5: Likening between two solution with closed initial conditions $x(\xi)$, and $\hat{x}(\xi)$ of the stochastic Burgers (4.3). See [11] to obtain other camera perspectives.

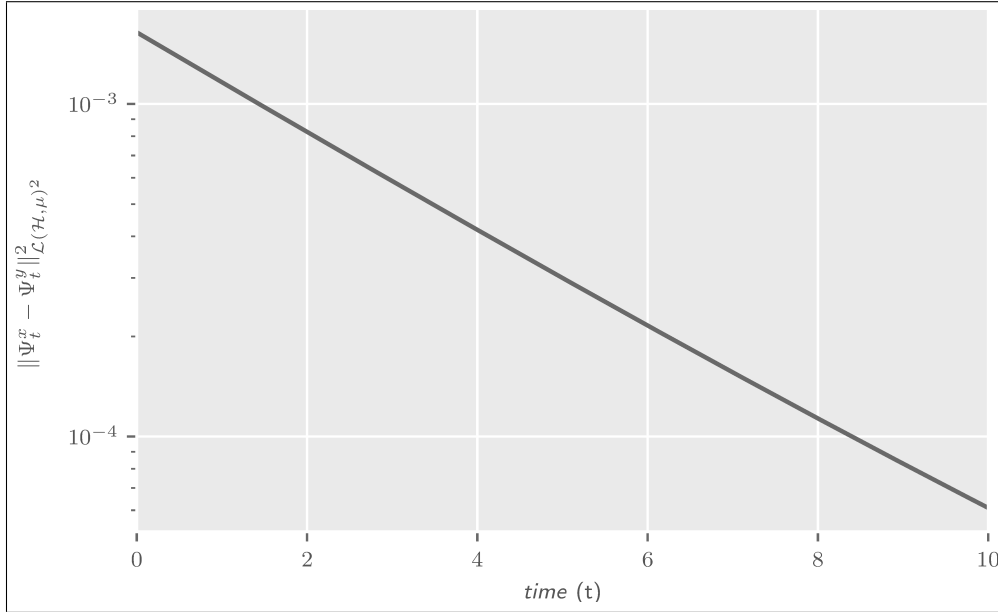


Fig. 6: Distance between two solutions of the stochastic Burgers (4.3) with initial conditions $x = x(\xi)$, and $y = \hat{x}(\xi)$.

5. Conclusions. To the best of our knowledge, our results represent the first contribution on the numeric stability respect to initial conditions of weak approximations of Kolmogorov equations in infinite dimensions. This kind of stability, combining with the weak approximation approach, would save computation time. That is, since our scheme asks specific conditions to obtain a weak numerical solution of an underlying SPDE, we convert the stochastic problem into a deterministic ODE for the first moment. This procedure overcome Montecarlo type simulations to approximate moments or distributions—simulate many realization of the numerical stochastic process to approximate distributions or moments. Further, under our setting, the regarding spectral approximation assure high precision and order of convergence. Thus we guess that our method would improve the time and save resources of computation. We are preparing another article to confirm this conjectures.

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