In 2012 the last particle predicted by the SM was detected at CERN, the Higgs Boson, yet again proving the success of the SM. But this also means no more free parameters in the SM for new particles. All interactions found obey the local $SU(3) \times SU(2) \times U(1)$ gauge symmetries and more recent data only strengthens the SM prediction. While no data was found suggesting inconsistencies with the electroweak symmetry breaking $SU(2) \times U(1) \rightarrow U(1)$ and no new exotic particles are found between 0-2 TeV 0.2. This means physicist have to search for new interactions considering different interaction ranges and strengths (Figure 0.1).

- (a) Considerably weaker than gravity with infinite range
- (b) Shorter range than the weak interaction of any strength
- (c) Range between weak interaction and nuclear force and considerably weaker than the weak interaction

There are various methods for finding BSM particles and as many theoretical models for new interactions, but these can generally be broken down in three categories.

The Model specific search where one takes a well-defined model often describing a complete UV theory and tries to find the predicted particles in measurements. The most popular example for this is the supersymmetry (SUSY) which is often associated with the search for dark matter and adds a wide range of different particles. In the search for Dark Matter one would for example look for missing transverse momentum in top quark interactions or look for resonance with the so-called two Higgs doublet models.

Another way of looking for new physics is by using simplified Models. Here one takes a well-defined model in order to describe some aspects or specific phenomenon of the UV Theory. Again a good example comes from dark matter physics the beautifully named neutralino in minimal supersymmetric standard model (MSSM) with MSSM being a low energy model of the SUSY. Here the neutralinos are electrically neutral fermions with a mass over 300 GeV and conserves the hypothetical R-parity in MSSM. Compared to the model specific search the result would give evidence only for the R-parity and the neutralino but not for other aspects of SUSY.

These first two methods depend entirely on a theoretical model. Granted these models are often well motivated by known physics or mathematical structures. But history has shown multiple times that experiments can produce unexpected results. As a recent example being the accelerating expansion of the universe by some so-called dark energy. Unexpected results are often not described by existing models, therefore a model independent search has to be done and EFT is the primary tool for this in particle physics. Taking a look at the Fermi theory (section ??) again one will realize that Fermi's theory was able to approximate the weak interaction but isn't able to describe the structure of weak interactions. Only using Fermi's

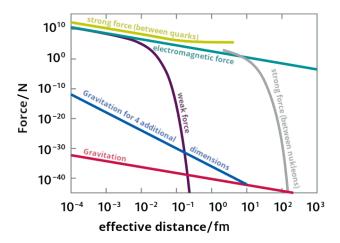


Figure 0.1: Interaction strength for different ranges of the Fundamental interactions. Although the Nuclear force isn't a fundamental it is shown in order to visualize point (c). [Teilchenwelt]

theory one would not be able to describe CP-Violation. Generally describing an exotic interaction using EFT one can only predict low order processes, but the low energy observer can't make predicting as to what he is actually describing. The results can then be used to construct an appropriate model or match an existing one and use the first to methods to validate this model.

| Model | S | ignatur | e ∫. | £ dt [fb | Mass limit | | | | Reference |
|---|--------------------------|---------------------------|--|--------------|---|-----------|-----------|--|--|
| $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_{1}^{0}$ | 0 e,μ | 2-6 jets | E _T miss | 139 | ÿ [1×, 8× Degen.] | 1.0 | 1.85 | | 2010.14293 |
| $\bar{g}\bar{g}, \bar{g} \rightarrow g\bar{g}\bar{\chi}^0_1$ | mono-jet $0 e, \mu$ | 1-3 jets 2-6 jets | E_T^{miss} E_T^{miss} E_T^{miss} | 36.1 139 | (8x Degen.) | 0.9 | | $m(\tilde{q}) \cdot m(\tilde{\chi}_1^0) = 5 \text{ GeV}$ 2.3 $m(\tilde{\chi}_1^0) = 0 \text{ GeV}$ | 2102.10874 2010.14293 |
| | | | / | | · | Forbidder | 1.15-1.9 | $m(\tilde{\chi}_1^0) = 1000 \text{ GeV}$ | 2010.14293 |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}W\tilde{\chi}_{1}^{0}$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_{1}^{0}$ | 1 e, μ ee, μμ | 2-6 jets 2 jets | E_{τ}^{miss} | 139 36.1 | š | | 1.2 | 2.2 m(x̄ ⁰)<600 GeV m(x̄)-m(x̄ ⁰)=50 GeV | 2101.01629 1805.11381 |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}^0$ | 0 e, μ SS e, μ | 7-11 jets 6 jets | E_T^{miss} | 139 | | | 1.15 | | 2008.06032 1909.08457 |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_{1}^{0}$ | 0-1 e, µ | 3 b | E_T^{miss} | 79.8 | s f | | | 2.25 m(x ⁰) < 200 GeV | ATLAS-CONF-2018-041 |
| | SS e, µ | 6 jets | | 139 | ? | | 1.25 | m(g)-m(ξ ⁰ ₁)=300 GeV | 1909.08457 |
| $\bar{b}_1\bar{b}_1$ | 0 e, μ | 2 b | E_T^{miss} | 139 | δ ₁ δ ₁ | 0.68 | 1.255 | $m(\tilde{X}_{1}^{0})<400 \text{ GeV}$ 10 GeV< $\Delta m(\tilde{b}_{1}, \tilde{X}_{1}^{0})<20 \text{ GeV}$ | 2101.12527 2101.12527 |
| $\tilde{b}_1\tilde{b}_1, \tilde{b}_1{\rightarrow}b\tilde{\chi}^0_2 \rightarrow bh\tilde{\chi}^0_1$ | 0 e, μ 2 τ | 6 b 2 b | E ^{miss} E ^{miss} | 139 139 | Forbidden | 0.13-0.85 | 0.23-1.35 | $\Delta m(\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0})=130 \text{ GeV}, m(\tilde{\chi}_{1}^{0})=100 \text{ GeV}$ $\Delta m(\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0})=130 \text{ GeV}, m(\tilde{\chi}_{1}^{0})=0 \text{ GeV}$ | 1908.03122 ATLAS-CONF-2020-031 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$ | 0-1 e, μ | ≥ 1 jet | E ^{miss} | 139 | i, | | 1.25 | $m(\tilde{\mathcal{X}}_1^0)=1 \text{ GeV}$ | 2004.14060,2012.03799 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wh\tilde{\chi}_1^0$ | 1 e, μ | 3 jets/1 b | E_T^{miss} | 139 | Forbidd | | | m(₹ ⁰ ₁)=500 GeV | 2012.03799 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$ | 1-2 т | 2 jets/1 b | E ^{miss} E ^{miss} | 139 | h. | Forbidden | 1.4 | m(₹₁)=800 GeV | ATLAS-CONF-2021-008 1805.01649 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$ | $0 e, \mu$ $0 e, \mu$ | 2 c mono-jet | E ^{Iniss} | 36.1 139 | i i | 0.85 | | $m(\tilde{x}_1^0)=0 \text{ GeV}$ $m(\tilde{x}_1)=m(\tilde{x}_1^0)=5 \text{ GeV}$ | 1805.01649 2102.10874 |
| $\tilde{t}_1\tilde{t}_1$, $\tilde{t}_1 \rightarrow t\tilde{\chi}_2^0$, $\tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$ $\tilde{t}_2\tilde{t}_3$, $\tilde{t}_3 \rightarrow \tilde{t}_1 + Z$ | 1-2 e, μ 3 e, μ | 1-4 b | E_T^{miss} E_T^{miss} | 139 139 | 5.000 | | 7-1.18 | m(ξ ⁰ ₂)=500 GeV | 2006.05880 2006.05880 |
| | | | E _T | | Forbidde | | | $m(\tilde{\chi}_1^0)$ =360 GeV, $m(\tilde{r}_1)$ - $m(\tilde{\chi}_1^0)$ = 40 GeV | |
| $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via WZ | Multiple ℓ/jet ee, μμ | s ≥ ljet | $E_T^{\rm fniss}$ | 139 139 | $\tilde{\chi}_{1}^{a}/\tilde{\chi}_{2}^{a}$ $\tilde{\chi}_{1}^{a}/\tilde{\chi}_{2}^{a}$ 0.205 | 0.96 | | $m(\tilde{\chi}_1^0)=0$, wino-bino $m(\tilde{\chi}_1^0)=f$ GeV, wino-bino | 2106.01676, ATLAS-CONF-2021-0 1911.12606 |
| $\tilde{X}_{1}^{\pm}\tilde{X}_{1}^{\mp}$ via WW | $2e, \mu$ | | E_T^{miss} | 139 | 0.42 | | | $m(\hat{t}_1^0)=0$, wine-bine | 1908.08215 |
| $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via Wh | Multiple ℓ/jet | S | E _T miss | 139 | Forbidden | 1. | | m(t1)=70 GeV, wino-bino | 2004.10894, ATLAS-CONF-2021-0 |
| $\tilde{X}_{1}^{\pm}\tilde{X}_{1}^{\mp}$ via $\tilde{\ell}_{L}/\tilde{\nu}$ $\tilde{\tau}\tilde{\tau}$, $\tilde{\tau} \rightarrow \tau \tilde{X}_{1}^{0}$ | 2 e, μ 2 τ | | Emiss Emiss | 139 139 | († († _L , † _{R,L}) 0.16-0.3 0.12-0.39 | 1.0 | , | $m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\ell}_1^n)+m(\tilde{\ell}_1^n))$ $m(\tilde{\ell}_1^n)=0$ | 1908.08215 1911.06660 |
| $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_1^0$ | 2 e, μ ee, μμ | 0 jets ≥ 1 jet | Emiss Emiss | 139 | 0.256 | 0.7 | | $m(\tilde{\chi}_1^0)=0$ $m(\tilde{\chi}_1^0)=10$ GeV | 1908.08215 1911.12606 |
| $\tilde{H}\tilde{H}, \tilde{H}\rightarrow h\tilde{G}/Z\tilde{G}$ | 0 ε.μ | ≥ 3 b | Emiss Emiss | 36.1 | ŭ 0.13-0.23 | 0.29-0.88 | | $BR(\tilde{\mathcal{C}}_{i}^{0} \rightarrow h\tilde{G})=1$ | 1806.04030 |
| | 4 e, μ 0 e, μ | 0 jets ≥ 2 large jet | Effiss S Effiss | 139 139 | Ĥ Ĥ | 0.55 | | $BR(\tilde{X}_{1}^{0} \rightarrow Z\tilde{G})=1$ $BR(\tilde{X}_{1}^{0} \rightarrow Z\tilde{G})=1$ | 2103.11684 ATLAS-CONF-2021-022 |
| Direct $\bar{X}_{1}^{+}\bar{X}_{1}^{-}$ prod., long-lived \bar{X}_{1}^{\pm} | Disapp. trk | 1 jet | E ^{miss} | 139 | ₹* 1 0.21 | 0.66 | | Pure Wino | ATLAS-CONF-2021-015 |
| 0.11.401.1. | | Multiple | | | 1 | | | Pure higgsino | ATLAS-CONF-2021-015 |
| Stable § R-hadron Metastable § R-hadron, §→aavviii | | Multiple | | 36.1 36.1 | ē ē [π(ē) =10 ns. 0.2 ns] | | | 2.0 m \tilde{\chi}_1^0 =100 GeV | 1902.01636,1808.04095 1710.04901.1808.04095 |
| $\tilde{\ell}\tilde{\ell}$, $\tilde{\ell} \rightarrow \ell \tilde{G}$ | Displ. lep | шопоріс | E_T^{miss} | 139 | i, p | 0.7 | | $\tau(\hat{\ell}) = 0.1 \text{ ns}$ | 2011.07812 |
| | | | , | | 0.34 | | | $\tau(\tilde{\ell}) = 0.1 \text{ ns}$ | 2011.07812 |
| $\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{1}^{\mp}/\tilde{\chi}_{1}^{0}$, $\tilde{\chi}_{1}^{\pm}\rightarrow Z\ell\rightarrow\ell\ell\ell$ | 3 e, µ | | | 139 | $\tilde{\mathcal{K}}_{1}^{*}/\tilde{\mathcal{K}}_{1}^{0} = [BR(Z\tau)-1, BR(Ze)-1]$ | 0.625 | | Pure Wino | 2011.10543 |
| $\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{1}^{\mp}/\tilde{\chi}_{2}^{0} \rightarrow WW/Z\ell\ell\ell\ell\nu\nu$ | $4 e, \mu$ | 0 jets | E_T^{miss} | 139 | $\tilde{X}_{1}^{2}/\tilde{X}_{2}^{0} = [\lambda_{03} \neq 0, \lambda_{122} \neq 0]$ | 0.95 | 1.55 | m(t ⁰ 1)=200 GeV | 2103.11684 |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow qqq$ $\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow tbs$ | | 4-5 large jet Multiple | s | 36.1 36.1 | ğ [m(ξ ⁰ ₁)=200 GeV, 1100 GeV] f [λ'' _{1,1} =2e-4, 1e-2] | 0.55 1.0 | 1.3 1. | Earge λ ₁₁₂ m(√ ₁)=200 GeV, bino-like | 1804.03568 ATLAS-CONF-2018-003 |
| $ii, i \rightarrow ix_1, x_1 \rightarrow ibs$ $ii, i \rightarrow b\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow bbs$ | | ≥ 4b | | 139 | Forbldo | | 05 | m(r ₁)=200 GeV, bino-like m(r̄ ₁)=500 GeV | 2010.01015 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | | 2 jets + 2 b | | 36.7 | $i_1 = [qq, bs]$ 0.42 | 0.61 | | | 1710.07171 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$ | 2 e, μ 1 μ | 2 b DV | | 36.1 136 | [1] [1e-10< X <1e-8, 3e-10< X <3e-9] | 1.0 | 0.4-1.45 | $BR(\vec{r}_1 \rightarrow be/b\mu) > 20\%$ $BR(\vec{r}_1 \rightarrow a\mu) = 100\%$, $cos\theta_i = 1$ | 1710.05544 2003.11956 |
| $\tilde{\chi}_{1}^{\pm}/\tilde{\chi}_{2}^{0}/\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1,2}^{0} \rightarrow tbs, \tilde{\chi}_{1}^{+} \rightarrow bbs$ | 1-2 e, μ | ≥6 jets | | 139 | 0.2-0.32 | 1.0 | 1.0 | Pure higgsino | ATLAS-CONF-2021-007 |
| | | | | | | | | | |
| | | | | | | | Ĭ. | | I |

Figure 0.2: Example for SUSY as BSM Model. All particles shown here have been rejected as they should have been found using described methods in BSM search. [.07.06.2021]