# Ground state energy of quantum dots using the coupled cluster method

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Something about coupled-cluster... Preferably doubles.

### CONTENTS

	CONTENTS		They are the eigenfunctions to the one-body part of the
			Hamiltonian and are thus harmonic oscillator functions.
I.	Introduction	1	Note that $ \Phi_0\rangle$ consists of the N first occupied single particle functions. We will denote the occupied indices
II.	Theory	1	with $i, j, k, l, \dots \in \{1, \dots, N\}$ , the virtual states with
	A. Second quantization	1	$a,b,c,d,\dots \in \{N+1,\dots,L\}$ and general indices with
	B. The coupled cluster approximation	1	$p,q,r,s,\dots \in \{1,\dots,L\}.$
	1. Energy of the coupled cluster approxima	tion 2	

2

### INTRODUCTION

References

In this project we will study the ground state energy of quantum dots.

#### II. THEORY

In this project we will study a system of N interacting electrons. We will be looking at a Hamiltonian consisting of a one-body and a two-body part. The one-body part is given by

$$h(\mathbf{r}_i) = -\frac{1}{2}\nabla_i^2 + \frac{1}{2}\omega^2 \mathbf{r}_i^2,\tag{1}$$

where we use natural units  $\hbar = c = e = 1$  and set the mass to unity. The two-body part is the Coulomb interaction potential.

$$v(\mathbf{r}_i, \mathbf{r}_j) = \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}.$$
 (2)

We thus get the total Hamiltonian

$$H = \sum_{i=1}^{N} h(\mathbf{r}_i) + \sum_{i < j}^{N} v(\mathbf{r}_i, \mathbf{r}_j).$$
 (3)

Next we define the reference Slater determinant as

$$|\Phi_0\rangle \equiv |1, 2, \dots, N\rangle,\tag{4}$$

where  $|p\rangle$  are the single particle functions that satisfies the equation

$$h|p\rangle = \varepsilon_p|p\rangle. \tag{5}$$

### Second quantization

We now move over to a use second quantization. Employing the creation operators,  $a_p^{\dagger}$ , and the destruction operators,  $a_p$ , we can write the Hamiltonian as

$$H = \sum_{pq} h_q^p a_p^{\dagger} a_q + \sum_{pqrs} w_{rs}^{pq} a_p^{\dagger} a_q^{\dagger} a_s a_r, \tag{6}$$

where the sums are general indices over all L basis states and the matrix elements are defined as

$$h_q^p \equiv \langle p|h|q\rangle,\tag{7}$$

$$w_{rs}^{pq} \equiv \langle pq|v|rs\rangle. \tag{8}$$

In the case of the eigenstates we have that the one-body operator reduces to a diagonal matrix.

$$h_a^p = \langle p|h|q\rangle = \varepsilon_q \langle p|q\rangle = \varepsilon_q \delta_{pq}, \tag{9}$$

where  $\delta_{pq}$  is the Kronecker-Delta.

### The coupled cluster approximation

We approximate the true wavefunction,  $|\Psi\rangle$ , of the system by the coupled cluster wavefunction,  $|\Psi_{\rm CC}\rangle$ , defined

$$|\Psi_{\rm CC}\rangle \equiv e^T |\Phi_0\rangle = \left(\sum_{i=0}^n \frac{1}{n!} T^n\right) |\Phi_0\rangle,$$
 (10)

where the *cluster operator*, T, is given by a sum of pexcitation operators

$$T = T_1 + T_2 + \dots \tag{11}$$

$$= \sum_{ia} t_a^a a_a^{\dagger} a_i + \left(\frac{1}{2!}\right)^2 \sum_{ijab} t_{ij}^{ab} a_a^{\dagger} a_b^{\dagger} a_i a_j + \dots$$
 (12)

We will be looking at the doubles approximation, that is

$$T = T_{\text{CCD}} \equiv T_2. \tag{13}$$

<sup>\*</sup> Project code: https://github.com/Schoyen/FYS4411

## 1. Energy of the coupled cluster approximation

When we're going to compute the energy of a system using the coupled cluster approximation we would ideally want to find the expectation value of the energy using the coupled cluster wavefunction.

$$E = \langle \Psi_{\rm CC} | H | \Psi_{\rm CC} \rangle. \tag{14}$$

As it turns out, this is an uncomfortable way of finding the energy as  $T \neq T^{\dagger}$ . Instead we will define what we call the *similarity transformed Hamiltonian*. We plug the coupled cluster wavefunction into the Schrödinger equation.

$$H|\Psi_{\rm CC}\rangle = E|\Psi_{\rm CC}\rangle.$$
 (15)

Next, we left multiply with the inverse of the cluster expansion, i.e.,

$$e^{-T}H|\Psi_{\rm CC}\rangle = e^{-T}E|\Psi_{\rm CC}\rangle = E|\Phi_0\rangle.$$
 (16)

Projecting this equation on the reference state we get

$$E = \langle \Phi_0 | e^{-T} H | \Psi_{\rm CC} \rangle = \langle \Phi_0 | e^{-T} H e^T | \Phi_0 \rangle, \tag{17}$$

where in the latter inner-product we have located the similarity transformed Hamiltonian defined by

$$\bar{H} \equiv e^{-T} H e^{T}. \tag{18}$$