

# Ground state energy of quantum dots using the coupled cluster method

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Something about coupled-cluster... Preferably doubles.

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Next we define the reference Slater determinant as

$$|\Phi_0\rangle \equiv |1, 2, \dots, N\rangle, \quad (4)$$

where  $|p\rangle$  are the single particle functions that satisfies the equation

$$h|p\rangle = \varepsilon_p|p\rangle. \quad (5)$$

## I. INTRODUCTION

In this project we will study the ground state energy of quantum dots.

## II. THEORY

In this project we will study a system of  $N$  interacting electrons. We will be looking at a Hamiltonian consisting of a one-body and a two-body part. The one-body part is given by

$$h(\mathbf{r}_i) = -\frac{1}{2}\nabla_i^2 + \frac{1}{2}\omega^2\mathbf{r}_i^2, \quad (1)$$

where we use natural units  $\hbar = c = e = 1$  and set the mass to unity. The two-body part is the Coulomb interaction potential.

$$v(\mathbf{r}_i, \mathbf{r}_j) = \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}. \quad (2)$$

We thus get the total Hamiltonian

$$H = \sum_{i=1}^N h(\mathbf{r}_i) + \sum_{i<j}^N v(\mathbf{r}_i, \mathbf{r}_j). \quad (3)$$

\* Project code: <https://github.com/Schoyen/FYS4411>

They are the eigenfunctions to the one-body part of the Hamiltonian and are thus harmonic oscillator functions. Note that  $|\Phi_0\rangle$  consists of the  $N$  first occupied single particle functions. We will denote the occupied indices with  $i, j, k, l, \dots \in \{1, \dots, N\}$ , the virtual states with  $a, b, c, d, \dots \in \{N+1, \dots, L\}$  and general indices with  $p, q, r, s, \dots \in \{1, \dots, L\}$ .

### A. Second quantization

We now move over to a use second quantization. Employing the creation operators,  $a_p^\dagger$ , and the destruction operators,  $a_p$ , we can write the Hamiltonian as

$$H = \sum_{pq} h_q^p a_p^\dagger a_q + \sum_{pqrs} w_{rs}^{pq} a_p^\dagger a_q^\dagger a_s a_r, \quad (6)$$

where the sums are general indices over all  $L$  basis states and the matrix elements are defined as

$$h_q^p \equiv \langle p|h|q\rangle, \quad (7)$$

$$w_{rs}^{pq} \equiv \langle pq|v|rs\rangle. \quad (8)$$

In the case of the eigenstates we have that the one-body operator reduces to a diagonal matrix.

$$h_q^p = \langle p|h|q\rangle = \varepsilon_q \langle p|q\rangle = \varepsilon_q \delta_{pq}, \quad (9)$$

where  $\delta_{pq}$  is the Kronecker-Delta.