

THE POTATO AND ITS DYNAMICS

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Abstract

This is an abstract text.

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Chapter 1

Introduction

Start your chapter by writing something smart. Then go get coffee.

Part I

Theory

Chapter 2

Hartree-Fock theory

Because its all about that energy, 'bout that energy, 'bout that energy!

Part II

Appendices

Appendix A

Hartree-Fock

Hartree-Fock appendix.

Appendix B

Reformulating the amplitude equations as matrix products

We will in this appendix show how to formulate the tensor contractions occurring in the coupled cluster equations as matrix products. The reason we wish to do this is to be able to perform these contractions as dot products (or matrix products) as there exists highly optimized code performing these operations, e.g., BLAS¹.

To be able to treat tensors of rank > 2 as matrices we have to create *compound indices* by stacking the dimensions after one another. For instance, by looking at the tensor $g \in \mathbb{C}^{I \times J \times K \times L}$, where we denote a single element by g_{ijkl} . Here g is a tensor of rank 4. By creating compound indices $\tilde{I} = IJ$ and $\tilde{K} = KL$ we can create a new tensor $\tilde{g} = \mathbb{C}^{\tilde{I} \times \tilde{K}}$ of rank 2 (represented as a matrix). Using the indices $\tilde{i} = iJ + j$ and $\tilde{k} = kL + l$ we now construct \tilde{g} in such a way that $\tilde{g}_{\tilde{i}\tilde{k}} = g_{ijkl}$.

It is also possible to create compound indices of more than two indices. For instance; choosing $\tilde{J} = JKL$ and setting $\tilde{j} = jKL + kL + l$ we can construct $\bar{g} = \mathbb{C}^{I \times \tilde{J}}$ where $\bar{g}_{i\tilde{j}} = g_{ijkl}$.

For the sake of brevity and clarity we will in the following avoid renaming the compound indices and their sizes, but we will instead indicate with a comma where we construct new indices.

B.1 Reformulating the CCD equations

B.2 Reformulating the CCSD equations

We use the expressions for the CCSD equations derived by Gauss et al.[1]. We start with the *effective double excitation amplitudes* found at the bottom of table

¹BLAS can be found here: <http://www.netlib.org/blas/>

3 in their article. Note that we rename $\tilde{\tau} \rightarrow \xi$ thus reserving the twiddle for intermediate calculations.

$$\tau_{ij}^{ab} = t_{ij}^{ab} + \frac{1}{2}P(ij)P(ab)t_i^a t_j^b \quad (\text{B.1})$$

$$\implies \tau_{ab,ij} = t_{ab,ij} + \frac{1}{2}P(ij)P(ab)(t_{a,i}t_{b,j})_{ab,ij}, \quad (\text{B.2})$$

$$\xi_{ij}^{ab} = t_{ij}^{ab} + \frac{1}{4}P(ij)P(ab)t_i^a t_j^b \quad (\text{B.3})$$

$$\implies \xi_{ab,ij} = t_{ab,ij} + \frac{1}{4}P(ij)P(ab)(t_{a,i}t_{b,j})_{ab,ij}. \quad (\text{B.4})$$

Next we look at the one-body intermediates found at the top of table 3 in the article by Gauss et al.[1]. We use the notation

$$u_{ef}^{am} \equiv \langle am || ef \rangle, \quad (\text{B.5})$$

that is, we treat the matrix elements u as the antisymmetric matrix elements of the two-body operator.

$$F_{ae} = f_{ae} - \frac{1}{2}f_{me}t_m^a + t_m^f u_{ef}^{am} - \frac{1}{2}\xi_{mn}^{af} u_{ef}^{mn} \quad (\text{B.6})$$

$$\implies F_{a,e} = f_{a,e} - \frac{1}{2}t_{a,m}f_{m,e} + (t_{fm}\tilde{u}_{fm,ae})_{a,e} - \frac{1}{2}\xi_{a,fmn}\tilde{u}_{fmn,e}, \quad (\text{B.7})$$

$$F_{mi} = f_{mi} + \frac{1}{2}f_{me}t_i^e + t_n^e u_{ie}^{mn} + \frac{1}{2}\xi_{in}^{ef} u_{ef}^{mn} \quad (\text{B.8})$$

$$\implies F_{m,i} = f_{m,i} + \frac{1}{2}f_{m,e}t_{e,i} + (t_{en}\tilde{u}_{en,mi})_{m,i} + \frac{1}{2}\tilde{u}_{m,nef}\tilde{\xi}_{nef,i}, \quad (\text{B.9})$$

$$F_{me} = f_{me} + t_n^f u_{ef}^{mn} \quad (\text{B.10})$$

$$\implies F_{m,e} = f_{m,e} + (t_{fn}\tilde{u}_{fn,me})_{m,e}. \quad (\text{B.11})$$

Bibliography

- [1] Jürgen Gauss and John F Stanton. “Coupled-cluster calculations of nuclear magnetic resonance chemical shifts”. In: *The Journal of chemical physics* 103.9 (1995), pp. 3561–3577.