### COMPUTATIONAL EFFICIENCY OF INFERENCE METHODS FOR 1 2 STOCHASTIC DIFFERENTIAL EQUATION DYNAMICS

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1. Introduction. Stochastic differential equations with respect to a Wiener 4 process (SDEs) are useful implicit descriptions of certain Markov processes. They 5have been applied to modelling financial markets [5, 1], projecting the paths of space-6 craft [1], and describing biological phenomenon [8]. They also can be used to build 7stochasticity into state space models [9]. 8

9 Inference of SDEs has remained a challenge - in particular, the task of computationally efficient calculation of gradients of a loss function. Early methods included 10Monte Carlo forward differentiation of parameters of the drift and noise functions, 11 however, these scale poorly with the number of parameters, which can be very high 12when these functions are high dimensional or if a deeper or wider neural network is 13 necessary for example. A more computationally efficient method in this case is Monte 14 Carlo backpropogation through the SDE solver given a realization of the Wiener 15 process [?]. However, this method requires one to save the forward path of the SDE which can be memory intensive for long time series. Adjoint methods instead define a backwards SDE that one may solve to calculate gradients. Early adjoint methods 18 19 considered each step of an Euler-Maruyama SDE solver a constrained optimization problem and were able to define the backwards SDEs by taking derivatives: one ef-20 fort defined the backwards SDE for use in calculating monte carlo estimates of the 21 gradient [6], while another used quadrature to approximately integrate over all paths 22[2]. Recently, the adjoint method has been extended for the use of any Stratonovich 23 24SDE solver [9].

Here I explore these ideas by implementing a high performance version of both 25 this later adjoint method, and backpropagation. To achieve low level control of the 26 SDE solver, I implement my own solver using the Euler-Heun method. I compare 27the accuracy of the gradients of this adjoint method, and backpropogation as well as 28 benchmark their computational cost and ability to fit data. I will demonstrate the 29convergence of my solver and gradient estimates to the correct values. Then I will 30 show that my implementation of gradient estimation using backpropogation is more 31 computationally efficient and results in higher quality inference than my implementa-32 tion of the adjoint method for gradient calculation. This investigation represents my 33 edification in the solving and inference of SDEs, as well as an investigation into the 34 35 factors that effect the utility of these two gradient estimation methods.

#### 2. Results. 36

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#### 37 2.1. Euler-Heun method solver.

**2.1.1.** Setup. I consider here the SDE, from time 0 to T,

$$dX_t = f(X_t, t)dt + g(X_t, t) \circ dW_t$$

where  $X_t$  has d dimensions and the Wiener process W has m dimensions. To inves-38

39 tigate how the architecture of an SDE solver may affect the performance of gradient

estimation techniques, and to easily implement backpropogation, I implement my own 40

SDE solver. 41

> The naive implementation of a solver for a Stratonovich integral is known as the Euler-Heun. In probability, as the mesh size of the grating  $0 = t_0 < \cdots < t_n = T$ 1

goes to 0,

$$\sum_{i=0}^{n-1} \frac{g(X_{t_{i+1}}, t_{i+1}) + g(X_{t_i}, t_i)}{2} (W_{t_{i+1}} - W_{t_i}) \to \int_0^T g(X_t, t) \circ dW_t.$$

- 42 The Euler-Heun (EH) method SDE solver approximates the integral for a time step
- 43  $\delta t$  via a second order Runga-Kutta (RK) method to approximate the average of the
- 44 drift and diffusion evaluated at the start and stop time of each step:

45 
$$H_{i+1} = X_i + f(X_i, t_i)\delta t + g(X_i, t_i)\Delta W_i$$

46 
$$X_{i+1} = X_i + \frac{\delta t}{2} \left( f(X_i, t_i) + f(H_{i+1}, t_{i+1}) \right)$$

47 
$$+\frac{1}{2}\left(g(X_{i},t_{i})+g(H_{i+1},t_{i+1})\right)\Delta W_{i}$$

$$\Delta W_i \sim \sqrt{\delta t N(0,1)}.$$

The EH method is strong order 0.5 - if  $X_T^{\delta t}$  is approximated using the EH method with step size  $\delta t$ , then

52 (2.1) 
$$||X_T^{\delta t} - X_T||_{L^2} \lesssim \delta t^{0.5}.$$

**2.1.2.** Results. Throughout, I will use one layer neural networks with tanh ac-53 tivation as the drift and diffusion functions f and q. As well, I will consider a two dimensional SDE with 2 dimensions of diffusion, i.e. d = m = 2. g will not be assumed diagonal or commutative. To allow for efficient evaluation of the neural networks, I 56 used the @einsum macro, which avoids boundschecking by checking compatibility of 57the dimensions of tensors before the multiplication. I used function barriers when 58unpacking parameters to allow the function to specialize on the input type (functions f\_NN! and g\_NN! in the code). Finally, to avoid repeated memory allocation for eval-60 61 uations of intermediate values in the evaluation of f or q, I preallocate the memory to store these values; I then pass the variables associated to this memory as a parameter 62 to the functions f\_NN! and g\_NN! and mutate it to these intermediate values (variables 63 t1 and t2 in the code). 64

I implement a high performance EH solver for SDEs that takes as input, a starting point, a parametrized drift function, a parametrized diffusion function, and parameters (function my\_SDEsolve! in code) (Fig 1A). To minimize memory allocation per step, I preallocated all the memory necessary to store intermediate values for each step and mutate these variables (function EH! in code).

One step of the solver, given the step size of the Wiener process, was benchmarked 70 71 as having 0 memory allocations and taking 1.790  $\mu$ s on my machine (section "Time Euler-Heun" in the code). Much of this computational complexity came from the 72need to apply the costly tanh non-linearity: the time was almost halved to 1.080  $\mu$ s when using a ReLu activation. Below however, I only consider the use of the 74 tanh non-linearity unless stated otherwise. To evaluate the convergence of my solver, 75I used the Euler-Heun solver implemented in DifferentialEquations.jl with step size 76  $2^{-16}$  to approximate  $X_T$  (in this case I used T = 1). Since the solver had a fixed step 77 size, it was easy to save its noise and use it in my solver to calculate  $X_T^{\delta t}$  with step sizes  $\delta t$  ranging from  $2^{-7}$  to  $2^{-12}$  (section "EH is order 0.5" in the code). Empirical 78 79 estimation of the  $L^2$  distance reveal that my method converges to the correct solution 80 and a linear regression reveals that it is of approximate order 0.522, similar to the 81 82 theoretical 0.5 described in equation 2.1.

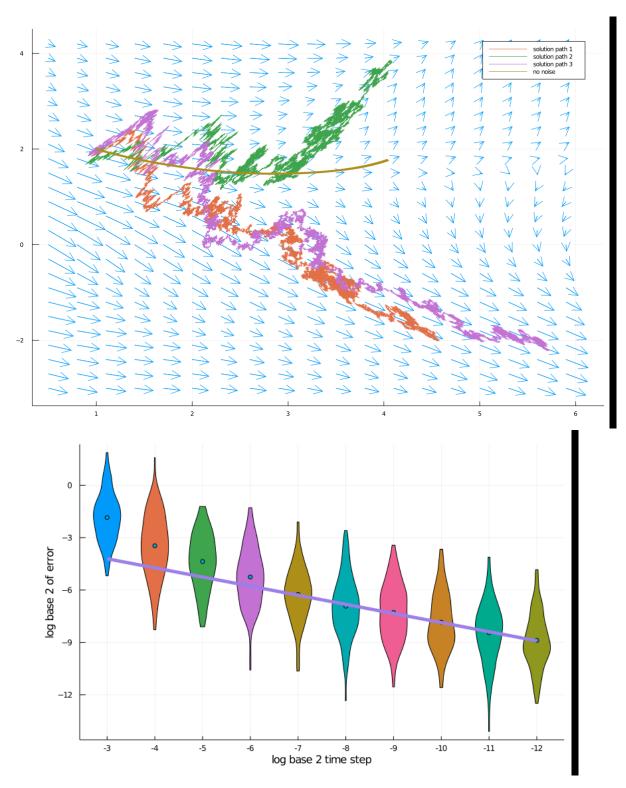


FIG. 1. A: Example paths starting at the point (1, 2) given by my Euler-Heun solver given a instantiation of the drift (blue arrows) and diffusion functions. The solution to the ordinary differential equation without the drift is also shown. In this example, the stochasticity causes some paths to cross a separatrix and depart significantly from the ODE solution. B: Converge of the  $L^2$ error defined in equation 2.1. A linear regression of the last 6 points gives a slope of -0.522 and an  $R^2$  of 0.995, close to the theoretical -0.5. The violin plots contain 100 point each.

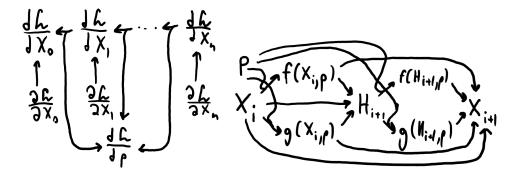


FIG. 2. A: A schematic of the calculation of the gradient of the loss with respect to the parameters and starting point by backpropogation through he solver. B: The computation graph of a single Euler-Heun step. I reverse it to calculate the pullback.

# 83 2.2. Backpropogation through the solver.

2.2.1. Setup. Here I consider the problem of approximating gradients of some loss function  $\mathcal{L}$  by backpropogation through the solver. Say  $(X_i)_{i=0}^n$  are the steps of the SDE iterator, i.e.  $X_{i+1} = EH(X_i, p)$  where EH is one step of the solver. Call  $\mathcal{B}_{EH(X,p)}^p$  the *p* component of the backpropogation of the Euler-Heun step and  $\mathcal{B}_{EH(X,p)}^X$  the *X* component. Then

89 (2.2) 
$$\frac{d\mathcal{L}}{dX_i} = \frac{\partial \mathcal{L}}{\partial X_i} + \mathcal{B}_{EH(X_i,p)}^X \left(\frac{d\mathcal{L}}{dX_{i+1}}\right)$$

90 (2.3) 
$$\frac{d\mathcal{L}}{dp} = \sum_{i=1}^{n} \mathcal{B}_{EH(X_{i-1},p)}^{p} \left(\frac{d\mathcal{L}}{dX_{i}}\right)$$

92 (Fig 2A). To calculate  $\mathcal{B}_{EH(X,p)}^X$  and  $\mathcal{B}_{EH(X,p)}^p$  I consider the computation graph in 93 fig 2B.

**2.2.2. Results.** To generate data to define a loss function, I initialize the parameters of f and g to random values and sample values from a path of the solution (section "Gen data" in the code). Bellow I consider the  $\ell^2$  loss with respect to this data.

When writing the pullback of the neural networks, I again used the @einsum macro (functions pullback\_f! and pullback\_g! in the code). When these functions were called, I again preallocate memory to store intermediate calculations and pass these variables as parameters. I also again use function barriers when unpacking parameters.

I implement a high performance pullback of the SDE solver that takes as input, the forward path, a parametrized drift function, a parametrized diffusion function, pullbacks of the drift and diffusion that add to the derivatives with respect to the parameters instead of overwrite them, and parameters (function my\_SDEbackprop! in code).

108 One iteration of the Euler-Heun pullback, given the step size of the Wiener 109 process, was benchmarked as having 0 memory allocations and taking 18.799  $\mu$ s on my 110 machine. The derivatives calculated by the backpropogation algorithm are validated 111 by those calculated by finite differencing (section "Backprop" in code) (Fig 3). Small 112 differences between the gradient calculated by backpropogation and that calculated

by finite differencing are likely due to the finite differencing algorithm only returning 113

114values up to a certain precision or precision loss when calculating gradients through

a tanh function; in particular, the step size doesn't affect the agreement between the 115

two methods. 116

### 117 2.3. Adjoint method.

**2.3.1.** Setup. I now implement the adjoint method of gradient estimation described in [9]. Call  $\mathcal{B}_{f(X,p)}$  the pullback of the drift function and  $\mathcal{B}_{g_i(X,p)}$  the pullback of the *i*-th column diffusion function. Consider the backwards SDE defined in [9]

$$d(A_t, B_t) = \mathcal{B}_{f(X,p)}(A_t) + \sum_{i=1}^m \mathcal{B}_{g_i(X,p)}(A_t) \circ d\widetilde{W}_{i,t}$$

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With  $A_T = \frac{\partial \mathcal{L}}{\partial X_T}$ ,  $B_T = \vec{0}$  and A jumping at time t' by a magnitude of  $\frac{\partial \mathcal{L}}{\partial X_{t'}}$ . Then, by the results in [9],  $A_0 = \frac{d\mathcal{L}}{dX_0}$  and  $B_0 = \frac{d\mathcal{L}}{dp}$ . It is also shown that the forward path X satisfies the backwards SDE 119

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121 (2.4) 
$$dX_t = -f(X_t, p)dt - g(X_t, p) \circ d\tilde{W}_t$$

Crucially, the stochastic integrals in these results are of the Stranovich interpretation. 122

Thus, the EH method may be used to evaluate them. 123

Solving such a backwards SDE is identical to solving the forwards SDE

$$dY_t = -f(Y_t, p)dt - g(Y_t, p) \circ dV_t$$

where  $Y_t = X_{T-t}$  and  $V_t = -W_{T-t}$ . The major difficulty in implementing this model 124is that the forwards and backwards noise must be the same. One solution for fixed 125step-size solvers is to save the path of the noise and reuse it when going backwards. 126 [9] avoids this by implementing using what they call a Brownian tree. Here, a single 127 random seed determines the value of the Wiener process path at all points  $(i2^{-m})_{i=0}^{2^m}$ 128 for a given m. The computation cost of querying the Brownian tree scales linearly 129with m. 130

**2.3.2.** Results. I use the implementation of the Brownian tree at http://github. 131com/SciML/DiffEqNoiseProcess.jl/pull/65. 132

133 To implement a high performance version of the adjoint method, I simply plugged 134in the pullback function for the neural networks f and q into my SDE solver (section "Backwards SDE" in the code). I show that the forward and reverse solutions to the 135 SDE are the same as described by equation 2.4 (Fig 4). Furthermore, I show that 136the gradients of the adjoint method agree with those calculated by finite differencing, 137 with the agreement becoming better as the step size used in the solver decreases, and 138 thus the solution becomes more accurate (Fig 3). 139

At a step size of  $2^{-12}$ , the Brownian tree algorithm takes 5.750  $\mu$ s and makes 140 100 allocations of 10 kB. This is almost an order of magnitude more resources than a 141 step of the solver. Since the Euler-Heun method is a fixed step size solver and saving 142the path does not incur a significant memory cost, hereafter, I opt instead to simply 143144 save the noise from the forward pass for use in the backwards pass. As well, instead of using the forward solution as calculated in the backwards pass, I simply save the 145forwards pass, incurring minimal memory cost, as pass this path to the backwards 146path; I expect this to increase the stability of the solver. 147

#### 148 2.4. Comparing the adjoint method and backpropogation.

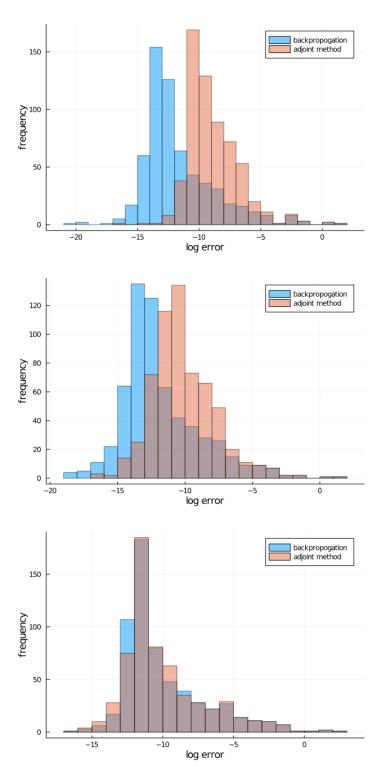


FIG. 3. Fractional error of parameter gradients calculated using back propagation or the adjoint method in comparison to those calculated by finite differencing for step sizes A:  $2^{-10}$ , B:  $2^{-12}$ , C:  $2^{-14}$ . These plots are each for a single realization of the Wiener process and plot all parameters of f, g and the starting point. Notice the x axis scaling is not the same between plots (sorry!).

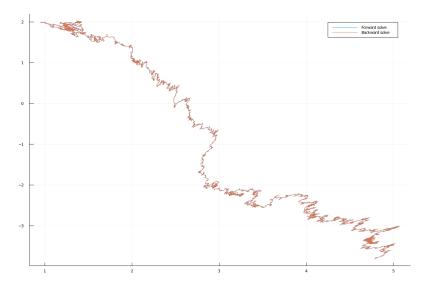


FIG. 4. Shown here is the forwards and backwards solution to an SDE for a given realization of the Wiener process. The two paths are superimposed and thus visually identical.

**2.4.1. Performance.** Here I will compare the computational complexity and
 accuracy of the gradient calculation by the adjoint method and by backpropogation
 through the solver.

For a step size of  $2^{-12}$  I was able to optimize the backpropogation method to take 86.573 ms while the adjoint method took 554.085 ms with an order of magnitude more allocations. The most time consuming step in the later case is the calculation of the pullback of the diffusion function. In particular, my implementation updates the diffusion matrix column by column requiring one to take slices of the diffusion matrix. These views are allocated to the heap. This represents an opportunity for optimization by predefining the views and reusing them step to step.

As was shown above, both methods agree with parameter gradient estimated 159by finite differencing with the agreement increasing for the adjoint method as the 160 step size decreased. The result is similar when comparing the gradients of the two 161 methods directly (Fig 5A). In principle, the gradients calculated by finite differencing 162 and backpropogation are not exactly the gradients of the loss function as the forward 163solve is inexact. However, there is no reason a priori to believe the adjoint method 164should have more accurate gradient estimates; in fact, one would expect that the 165 error from the backwards solve add to that of the forwards solve to give an even less 166 accurate estimate. Thus, those values in Fig 5A may be interpreted as lower bound for 167 the error of the gradients. Clearly, due to the increased computational efficiency and 168 more accurate gradient estimates, my implementation of backpropogation through 169the solver is superior to my implementation of the adjoint method. 170

As a curiosity, it is worth mentioning that path of  $(A_t)_t$  as calculated by the adjoint method is very similar to  $(\frac{d\mathcal{L}}{dX_{T-t}})_t$  (Fig 5B), although it is expected that these two paths have similar start and end points  $(\frac{d\mathcal{L}}{dX_T})_t$  and  $(\frac{d\mathcal{L}}{dX_0})_t$  respectively.

174 **2.4.2.** Inference. Here I will investigate how these two method perform in an 175 inference task. I will attempt to minimize the  $\ell^2$  error to the data generated above 176 via (stochastic) gradient descent (section "Training loop" in the code). I had better

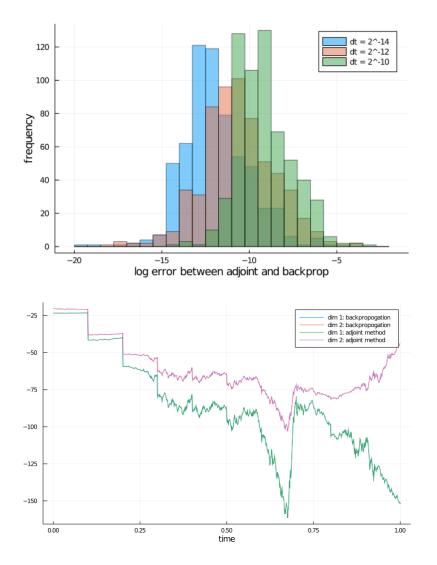


FIG. 5. A: Fractional error of parameter gradients calculated using back propagation in comparison to those calculated by adjoint method for step sizes  $2^{-10}$ ,  $2^{-12}$ ,  $2^{-14}$ . B: Two dimensions of the paths of the adjoints  $A_t$  and  $(\frac{d\mathcal{L}}{dX_{T-t}})_t$  for a particular realization of the Wiener process. The two paths are superimposed and visually identical.

results using a ReLu activation in the neural networks than tanh so below I will consider fitting f and g with ReLu activations. As well, inference of the starting point made the problem much more difficult, so in this setting, I considered the starting point known.

Because of the stochasticity of the gradients, inference was very difficult. Parameters would explode or the magnitude of g would stabilize at a large value if the learning rate was too high. To avoid divergence, it was necessary to average the gradients of the parameters over large batches of realizations of the Wiener process. Here I used a batch size of 100 and a learning rate of 0.0005. I also observed better results when initializing with f and g near 0 but with the values of the first hidden layer positive,

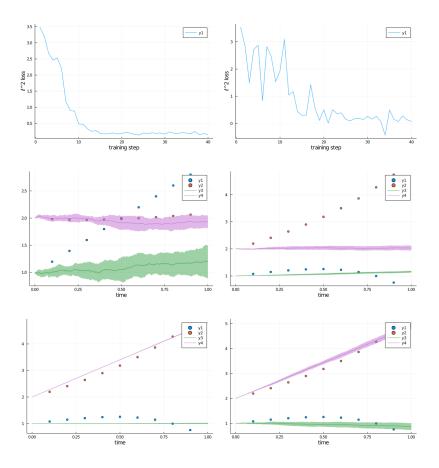


FIG. 6. The first column shows inference using backpropagation through the solver to estimate gradients while the second shows the use of the adjoint method to estimate gradients. For each set of data, both models were trained for 40 steps, the first column at a learning rate of 0.0005 and the second at 0.0003. The first row shows the loss decreasing stochastically with increased training. The middle and last row show the data (points) and the average and standard deviation of 10 paths of the SDE; the middle row shows paths from the SDE with untrained f and g, and the last row uses the trained f and g.

187 so as not to be set to 0 by the activation function.

I thorough comparison of the two methods would require tuning of the inference 188procedure which is outside the scope of this report. However, I noticed that when 189 using gradients calculated by backpropogation, inference with the above parameters 190 would usually converge while when using the adjoint method the sometimes diverged. 191 Decreasing the learning rate slightly to 0.0003 usually resulted in convergence for 192the adjoint method. However, both methods were eventually able to converge to 193solutions (Fig 6). The solution arrived at when using the adjoint method had a 194195higher magnitude g when stochasticity is not necessary to explain the data and had a 196 much more stochastic loss curve; this is possibly a by-product of less accurate gradient estimates. Neither solution fit the data satisfactorily despite the loss plateauing, likely 197 because of many of the nodes of the neural network being turned off by the ReLu 198 activation. Further tuning of the architectures of f and g would likely lead to higher 199200 quality fits.

**3. Discussion.** Here I have implemented the EH SDE solver, backpropogation through the solver, and implemented the adjoint method of calculating integrals described in [9]; I was able to use these techniques to fit an SDE to data. My preliminary results show that the use of the Brownian tree represents a significant computational cost in the evaluation of gradients using the adjoint method, and that the more accurate gradient estimates of the backpropogation method resulted in higher quality inference.

While backpropogation through an EH step should in principle cost more function evaluations (one extra pullback of f) the particulars of my implementation meant that the backpropogation method was significantly more computationally efficient than the adjoint method for gradient calculation. To evaluate accurately the computational complexities of these models, a more careful implementation of the adjoint method is necessary.

Multiple ideas important to the comparison of the adjoint method and backpropogation were not explored in this report.

1) Most important is the selection of the SDE solver.

a) The EH method has fixed timesteps. This allowed me to save the forward path and noise at the time-points visited in the forwards pass for use in the backwards path, significantly decreasing the computational complexity of the adjoint method while likely increasing its stability. Many solvers use adaptive step sizes however to increase solution stability. This negates the benefit in stability and speed I saw in my analysis.

223 b) When efficiency is measured as number of calculations required to reach a certain error rate, the stochastic Runga-Kutta (SRK) method described in [12] leads 224 to very efficient solvers; in particular, methods that are more efficient than the EH 225method. By running these SRK methods with two tableau, one may calculate esti-226mated error to create adaptive SDE solvers as well [11]. These family of methods are 227 more efficient given more stringent requirements on the form of the diffusion term g: 228 229 the SRA methods are the most efficient, assuming additive noise (i.e. q(x,t) = q(t)); the SRID methods are less efficient, only assuming the diffusion function is diagonal, 230i.e. the noise in each dimension is independent; the SRIC methods are less efficient as-231 suming a commutativity property; and finally, the general SRI methods, which make 232no assumption about the form of the noise, are least efficient as they must calculate 233iterated stochastic integrals [13]. The form of the noise for the backwards SDE may 234235differ from that of the forwards, requiring a less efficient solver. In particular, when the forward SDE has diagonal noise, the backwards SDE is only guaranteed to have 236 commutative noise. However, when the forwards SDE has additive noise, the pull 237back of the noise does as well. Thus it is possible that the adjoint method may only 238239be optimal for certain SDEs.

c) Here I only considered simple neural SDEs without stiffness. Stiff SDEs require 240 the application of solvers with large stability regions. To achieve this, implicit methods 241 have often been used. Backpropogation through implicit methods is extremely costly. 242 This fact has previously been described as an advantage of the adjoint method for 243 gradient estimation [3]. However, recent work has also developed stability optimised 244SRK methods SOSRI and SOSRI2 [10]; these methods are stable enough to solve stiff 245246 SDEs, are significantly more computationally efficient than implicit methods, and are much easier to backpropogate through. As well, these methods are restricted to SDEs 247 with diagonal noise however and thus cannot be used for the backwards pass of the 248 adjoint method. This potentially puts the adjoint method at a disadvantage when 249

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251d) Here I considered an SDE that did not need to be run with small time steps 252for the solver to get accurate estimates of the solution. SDEs where smaller step sizes are necessary would increase the memory burden, making the adjoint method 253more appealing. However, adaptive methods that decrease the step size of the solver 254commensurately with the estimated error of the solver may decrease the number 255of steps in the path and thus incur a smaller burden on memory. As well, more 256computationally efficient methods, with fewer calls to the Wiener process and function 257evaluations would incur a smaller memory cost to the solver. Finally, not every 258intermediate calculation need be saved to implement backpropogation: some may be 259 recalculated, such as  $H_{i+1}$  in my implementation of backpropogation. 260

201 2) One potential solution to the instability of inference of SDEs is regulariza-202 tion by penalizing the magnitude of higher order derivatives of the solutions or the 203 derivatives of the dynamics as suggested in [4, 7]. While the former solution is not 204 easily translatable to the SDE case where solutions are a.s. nowhere differentiable, 205 the latter may lead to more stable inference. This could make the instability of the 206 adjoint method less of a detriment during inference.

267 3) In [9] it was recommended that latent SDEs be trained using the evidence 268 lower bound (ELBO) as an objective. This objective avoids the collapse of the noise during training as seen in Fig 6. To make sense of the ELBO, one must be able to 269 compute the KL divergence to a prior. Using Girsanov's theorem, one may find the 270KL divergence between a prior SDE with the same diffusion, but possibly different 271drift, and the learned SDE (when the diffusions are different, it is possible that the 272273supports of the laws of the solutions are non-overlapping and thus have divergence 274infinity - for example solutions for SDEs with different additive diffusions a.s. have paths with different quadratic variation). The Bayesian setup would normally require 275that one fix the prior, this would result in fixing the diffusion. Optimizing the ELBO with respect to the diffusion is, however, still mathematically possible (9) section 2779.7). How the treatment of the diffusion as a learnable function or hyper-parameter 278279affects the quality of the latent representation is unexplored; it is thus possible that the adjoint method may be superior in the situation where one must only learn the 280drift. 281

Both methods of gradient calculation explored in this report have their theoretical advantages. However, an enumeration of use case and class of SDE is complex and results are subject to change as new SDE solvers are implemented. However, such challenges must be overcome to progress towards scalable inference of SDEs.

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