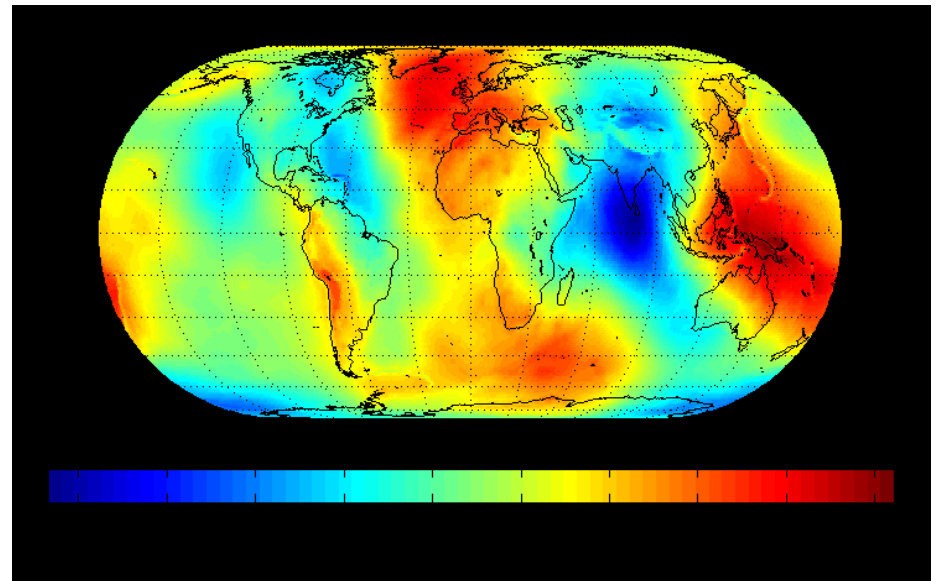
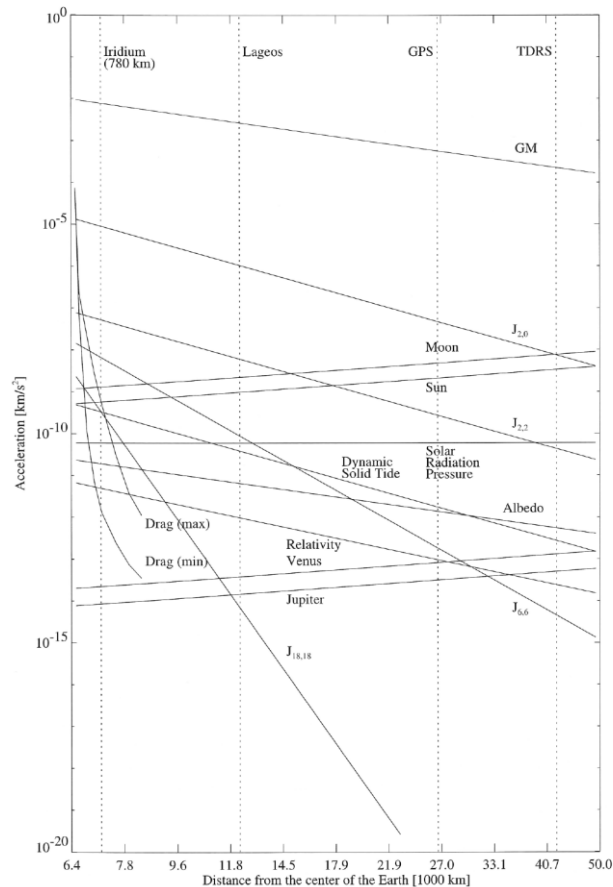


Perturbation Force Modeling



Dynamic Model Fidelity

- Consider the 2-body/Keplerian motion model for the free motion of an Earth-orbiting space object:

$$\ddot{\vec{r}} = -\frac{\mu\vec{r}}{r^3}$$

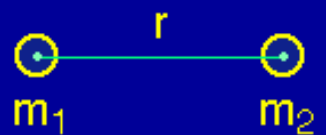
- While this is the most widely used model, it is also the lowest fidelity (i.e. most simplistic)
- While the conic section paradigm of 2-body motion approximates actual space object motion generally well over moderate time spans, various temporal & spatial errors exist
- 2-body motion based on Newton's Law of Gravitation

Earth's Gravity

- Recall Newton's Law of Gravitation from the orbital mechanics review in Module 1
- This law assumes the 2 bodies attracting each other are point masses (or equivalently, that the distance between them is much greater than either object's radius)
- 2-body motion assumes Earth's gravity is the only force acting on the object **AND** is acting according to this law

Law of Universal Gravitation

Every object in the Universe attracts every other object with a force directed along the line of centers for the two objects that is proportional to the product of their masses and inversely proportional to the square of the separation between the two objects.

$$F_g = G \frac{m_1 m_2}{r^2}$$


F_g is the gravitational force

m_1 & m_2 are the masses of the two objects

r is the separation between the objects

G is the universal gravitational constant

Earth's Gravity

- In reality, Earth-orbiting objects' motion is affected in other ways by Earth's gravity due mainly to 2 factors:
 - Earth (like all bodies) is a distributed mass
 - Earth is not a perfect sphere
- For this reason, the magnitude & direction of force due to “higher order gravity effects” varies as a function of position → the gravitational force an object experiences varies slightly throughout its orbit
- Many intricate models of Earth's gravity field have been developed, usually expressed in Cartesian or spherical coordinates → $\vec{F}(x, y, z)$ or $\vec{F}(r, \theta, \phi)$ where x, y, z or r, θ, ϕ are the coordinates of the position (i.e. point) at which we wish to calculate the force (or acceleration) due to Earth's gravity

Earth's Gravity

- One widely accepted model characterizes the potential U due to Earth's gravity
 - U often expressed as a function of ECI position: $U(x, y, z)$
 - Gradient of U is acceleration: $\frac{\partial U}{\partial x} = a_x$ $\frac{\partial U}{\partial y} = a_y$ $\frac{\partial U}{\partial z} = a_z$
- U is then modeled as a series expansion where each term U_k represents the acceleration due to a differential portion of the Earth \rightarrow terms of the form:

$$U_k = \frac{\mu}{r} \left(\frac{R_E}{r} \right)^k P_k(\sin\phi) (C_m \cos\lambda + S_m \sin\lambda)$$

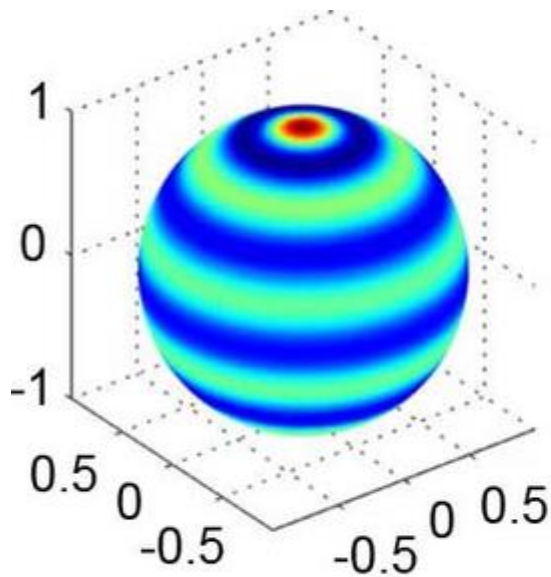
where r is distance of point from Earth's center, R_E is Earth radius, μ is Earth's gravitational constant, ϕ is latitude, λ is longitude

Earth's Gravity

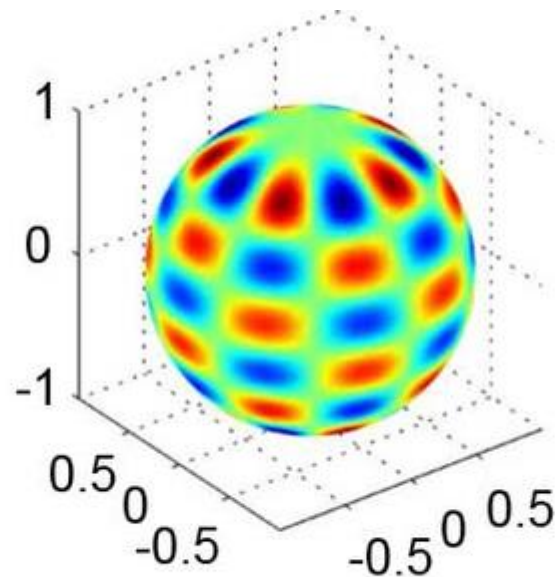
- So what are P_k , C_m , & S_m ?
- P_k represents the series of **Legendre polynomials** (or Legendre functions)
 - P_k can be derived for any k ($= 1, 2, 3, \dots$) depending on how many terms one wishes to retain in the expansion of U
 - Some math references contain tables of the first several (5, 10, etc) Legendre polynomials
 - P_k are the **latitude-dependent** terms in U
- C_m & S_m are constant coefficients representing the **longitude-dependent** terms in U
- Because we're modeling elements of a sphere (or near-sphere) & the terms involve trig functions, these are often called **spherical harmonics**

Earth's Gravity

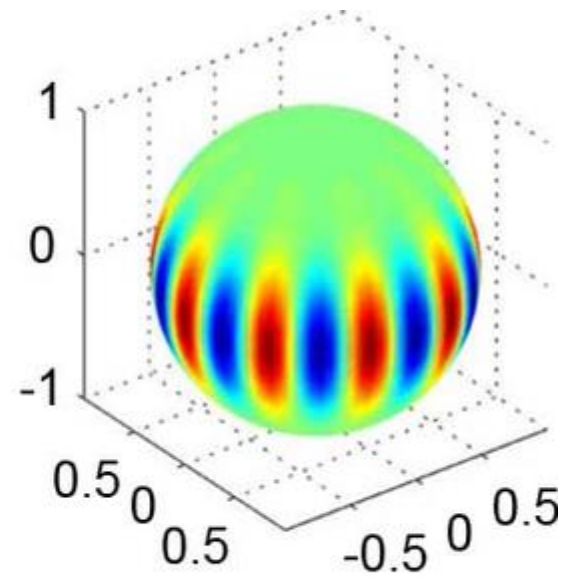
- Harmonics are divided into the following categories:
 - **Zonal**: only latitude-dependent; symmetric about Earth's polar axis
 - **Tesseral**: divides sphere into a checkerboard array
 - **Sectorial**: only longitude-dependent; symmetric about Earth's equator



Zonal



Tesseral



Sectorial

Earth's Gravity

- The dominant (& most commonly known) harmonic term is J_0 , the 2-body term
- $J_1 = 0$ (as are several of the harmonic terms)
- 2nd most dominant harmonic term is $J_2 \rightarrow$ due to Earth's **oblateness** (fatter at the equator than at the poles)
- Acceleration due to J_2 expressed as a function of ECI position by

$$a_{J2,x} = \frac{3\mu J_2 R_E^2}{2r^5} \left(\frac{5z^2}{r^2} - 1 \right) x$$
$$a_{J2,y} = \frac{3\mu J_2 R_E^2}{2r^5} \left(\frac{5z^2}{r^2} - 1 \right) y$$
$$a_{J2,z} = \frac{3\mu J_2 R_E^2}{2r^5} \left(\frac{5z^2}{r^2} - 3 \right) z$$

- **J_2 : harmonic coefficient (= 0.00108 for Earth)**

- **R_E : Earth radius**

- **$r = \sqrt{x^2 + y^2 + z^2}$**

Earth's Gravity

- So if we wish to propagate an object's motion due to 2-body & J2 forces, our ODEs are

$$\ddot{x} = -\frac{\mu x}{r^3} + a_{J2,x} = -\frac{\mu x}{r^3} + \frac{3\mu J_2 R_E^2}{2r^5} \left(\frac{5z^2}{r^2} - 1 \right) x$$

$$\ddot{y} = -\frac{\mu y}{r^3} + a_{J2,y} = -\frac{\mu y}{r^3} + \frac{3\mu J_2 R_E^2}{2r^5} \left(\frac{5z^2}{r^2} - 1 \right) y$$

$$\ddot{z} = -\frac{\mu z}{r^3} + a_{J2,z} = -\frac{\mu z}{r^3} + \frac{3\mu J_2 R_E^2}{2r^5} \left(\frac{5z^2}{r^2} - 3 \right) z$$

Effect of Earth's Gravity

- The primary effect of Earth's gravity perturbations on Earth-orbiting objects, from an SSA perspective, is the way J_2 force perturbs an object's motion
- Actual effect is fairly complex, but a simplified version is as follows:
 - Periodic oscillation (or "osculation") of all 6 orbit elements ($a, e, i, \Omega, \omega, \nu$)
 - Secular changes (i.e. changes that grow with time) to Ω, ω, ν
 - Secular changes to Ω are called **nodal precession** → gives rise to **Sun-synchronous orbits**
 - Secular changes to ω are called **perigee precession** → gives rise to **Molniya orbits**

**These effects are
detailed in the ME
595 course**

Drag

- Drag equation for Earth-orbiting objects is the same as for aircraft:

$$\bar{a}_{drag} = -\frac{1}{2} \frac{C_D A}{m} \rho v_{rel} \bar{v}_{rel} \text{ where } \bar{v}_{rel} = \bar{v} - \bar{\omega}_E \times \bar{r}$$

- **C_D : drag coefficient**
- **A : area normal to velocity vector**
- **m : mass**
- **ρ : density**
- **\bar{r} : position vector**
- **\bar{v} : velocity vector**
- **$\bar{\omega}_E$: Earth's rotation rate (= 360° /day)**

- 360° per sidereal day is approx. 7.292E-5 rad/sec, so Earth's rotational velocity vector may be expressed in the J2000 frame as $\bar{\omega}_E = 7.292\text{E-}5 \hat{k}$ rad/sec
- Drag is a **non-conservative** force → dissipates energy
- We talked last week about how drag decelerates a space object, hastening de-orbit → but how long does this take?

Drag

- So if we wish to propagate an object's motion due to 2-body & drag forces, our ODEs are

$$\ddot{x} = -\frac{\mu x}{r^3} + a_{drag,x}$$

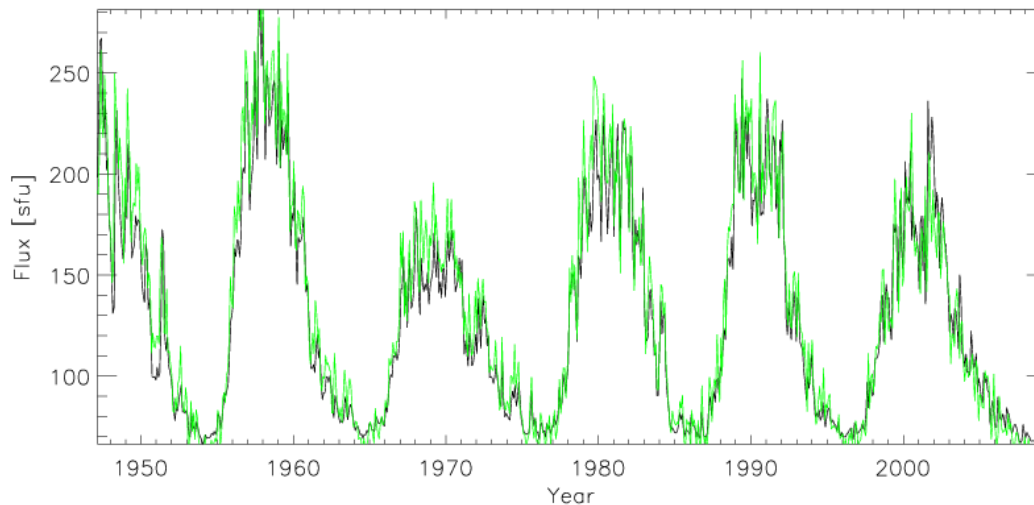
$$\ddot{y} = -\frac{\mu y}{r^3} + a_{drag,y}$$

$$\ddot{z} = -\frac{\mu z}{r^3} + a_{drag,z}$$

- Where the x , y , & z components of a_{drag} depend on the x , y , & z components of $\bar{v} - \bar{\omega}_E \times \bar{r}$ at each integration time step

Drag

- All qty's in the drag equation are uncertain to some degree
- The most uncertain (by far) is **density**, ρ
- Driven by multiple time-varying phenomena, including:
 - Geomagnetic planetary index (a_p)
 - Geomagnetic planetary index (k_p)
 - Solar flux ($F_{10.7}$) \rightarrow driven by multiple temporal effects, e.g. 11-year solar cycle, 27-day solar rotation cycle



Solar flux plot shows peaks approx. every 11 years

- Scientists can model these phenomena to some extent, but much of this behavior is not well understood, i.e., considered **random**

Drag

- Numerous density models have been derived over several decades
- Numerical or **empirical** models make extensive use of lookup tables reflecting historical data
 - Examples are Jacchia, Harris-Priester, MSIS
- Analytical or **physics** models attempt to model every phenomenon mathematically
 - Examples are GITM, TIE-GCM
- Many of these models are extremely detailed & computationally intensive
- As with most scientific models, the choice of which to use boils down to a tradeoff between accuracy & computation time

Drag

- The simplest density model is the exponential “standard” atmosphere: $\rho = \rho_0 \exp(-\frac{h-h_0}{H})$

Altitude h_{ellp} (km)	Base Altitude h_o (km)	Nominal Density ρ_o (kg/m ³)	Scale Height H (km)	Altitude h_{ellp} (km)	Base Altitude h_o (km)	Nominal Density ρ_o (kg/m ³)	Scale Height H (km)
0–25	0	1.225	7.249	150–180	150	2.070×10^{-9}	22.523
25–30	25	3.899×10^{-2}	6.349	180–200	180	5.464×10^{-10}	29.740
30–40	30	1.774×10^{-2}	6.682	200–250	200	2.789×10^{-10}	37.105
40–50	40	3.972×10^{-3}	7.554	250–300	250	7.248×10^{-11}	45.546
50–60	50	1.057×10^{-3}	8.382	300–350	300	2.418×10^{-11}	53.628
60–70	60	3.206×10^{-4}	7.714	350–400	350	9.518×10^{-12}	53.298
70–80	70	8.770×10^{-5}	6.549	400–450	400	3.725×10^{-12}	58.515
80–90	80	1.905×10^{-5}	5.799	450–500	450	1.585×10^{-12}	60.828
90–100	90	3.396×10^{-6}	5.382	500–600	500	6.967×10^{-13}	63.822
100–110	100	5.297×10^{-7}	5.877	600–700	600	1.454×10^{-13}	71.835
110–120	110	9.661×10^{-8}	7.263	700–800	700	3.614×10^{-14}	88.667
120–130	120	2.438×10^{-8}	9.473	800–900	800	1.170×10^{-14}	124.64
130–140	130	8.484×10^{-9}	12.636	900–1000	900	5.245×10^{-15}	181.05
140–150	140	3.845×10^{-9}	16.149	1000–	1000	3.019×10^{-15}	268.00

- ρ_o : reference density
- h : altitude
- h_o : reference altitude
- H : scale altitude

- Yields a rough approximation of density up to a few hundred km (actual drag does generally exhibit exponential behavior)

Effect of Drag

- The primary effect of drag, from an SSA perspective, is to de-orbit LEO objects → if there were no drag, objects would **never** naturally de-orbit
 - Magnitude of drag varies inversely with altitude
 - Effect on circular orbits is to follow a slow inward spiral toward Earth, until re-entry
 - For elliptical orbits, drag tends to slowly circularize the orbit as it spirals (b/c greatest effect of drag is at perigee of an orbit, which tends to lower its altitude at apogee)
- Due to the complexity (& “randomness”) in so many phenomena contributing to drag, any sort of orbital lifetime prediction for an object is fraught with uncertainty
 - Prediction accuracy increases as object gets closer to re-entry
- Final note about drag: as with aerodynamic drag, side force components (e.g. “lift”) exist in addition to the negative velocity force...but these components are often too small to warrant modeling

3rd-Body Gravity

- Newton's Law of Gravitation tells us that every object in the universe exerts gravitational force on every other object
- For a given satellite, out of all the objects in the universe, most of these objects' mass is too **small** &/or their distance from the satellite too **large** for their gravity to have any effect on the satellite's motion
- Therefore, even in high-fidelity orbit propagators, only a few objects' gravity (other than Earth's) is included in such models
- The two most common "3rd bodies" whose gravity is included in such models are the **Sun** & the **Moon**

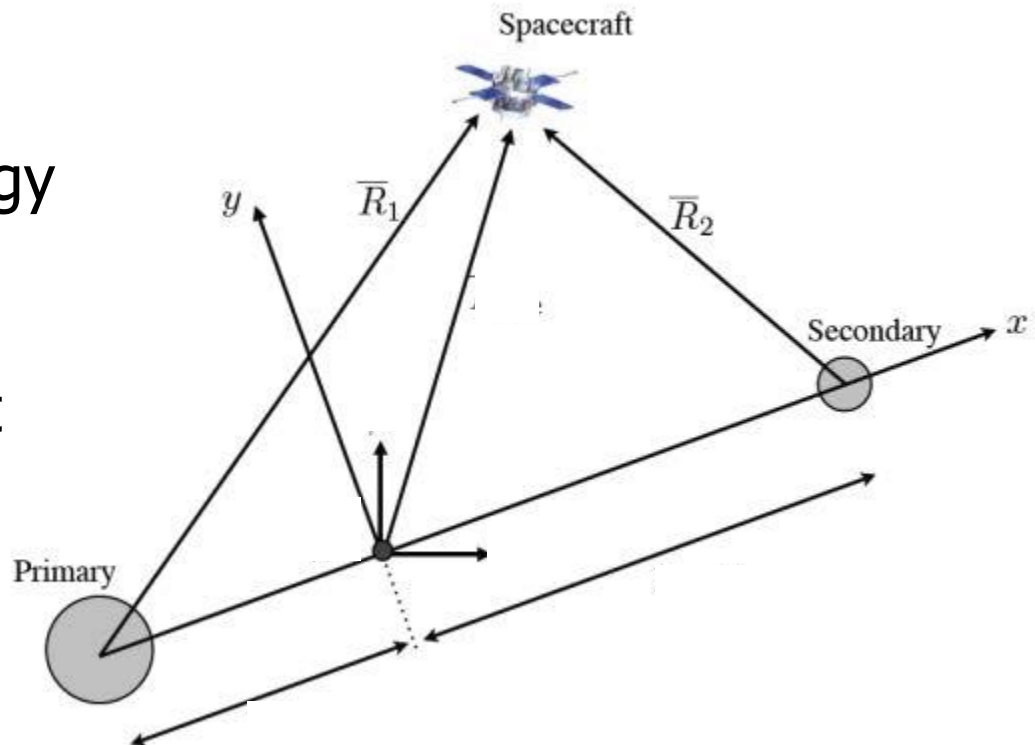
3rd-Body Gravity

- Acceleration due to a 3rd body is given by

$$a_{3rd} = -\mu_{3rd} \left[\frac{\bar{R}_2}{R_2^3} + \frac{\bar{R}_1 - \bar{R}_2}{(R_1 - R_2)^3} \right]$$

- μ_{3rd} : gravitational constant of 3rd body
- \bar{R}_1 : vector from Earth's center to object
- \bar{R}_2 : vector from center of 3rd body to object

- 3rd body forces are **conservative** (no energy dissipation), therefore they will not raise or lower an object's orbit

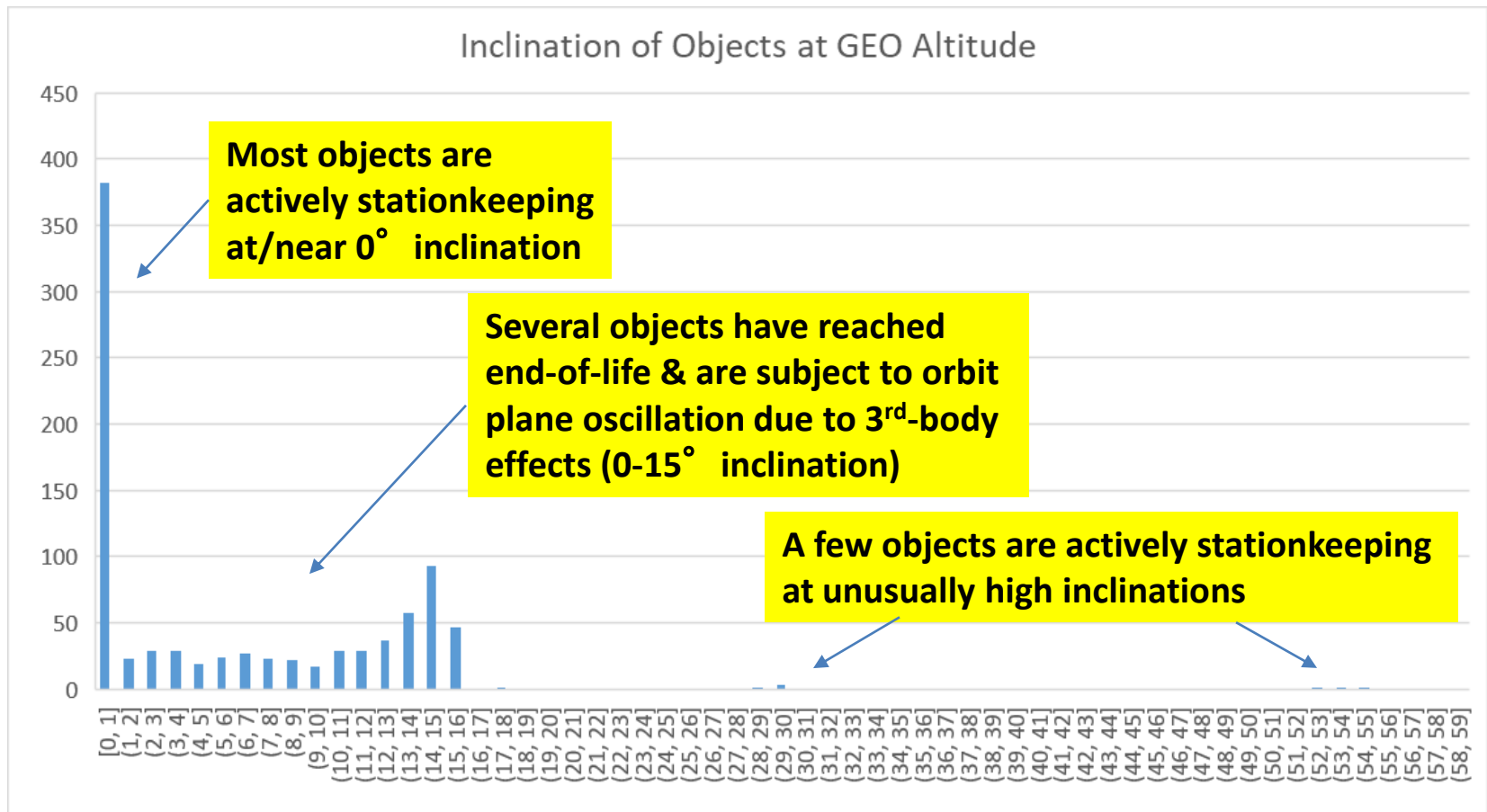


Effect of 3rd-Body Gravity

- The general effect of 3rd-body gravity is similar higher order gravity effects from Earth, i.e. periodic & secular changes to orbit elements
- One interesting 3rd-body effect, from an SSA standpoint, pertains to objects in/around GEO
 - **Orbit plane oscillation** of amplitude $\sim 15^\circ$
 - Only occurs for objects in free motion (not maneuvering or stationkeeping)
 - Most GEO objects are actively stationkeeping \rightarrow effect is only seen in objects that have reached end-of-life
 - Period of oscillation ~ 50 yrs

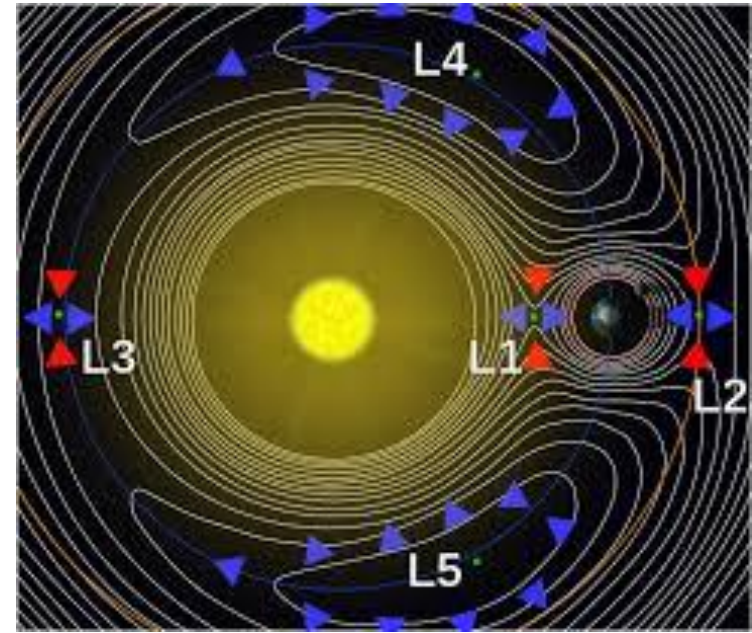
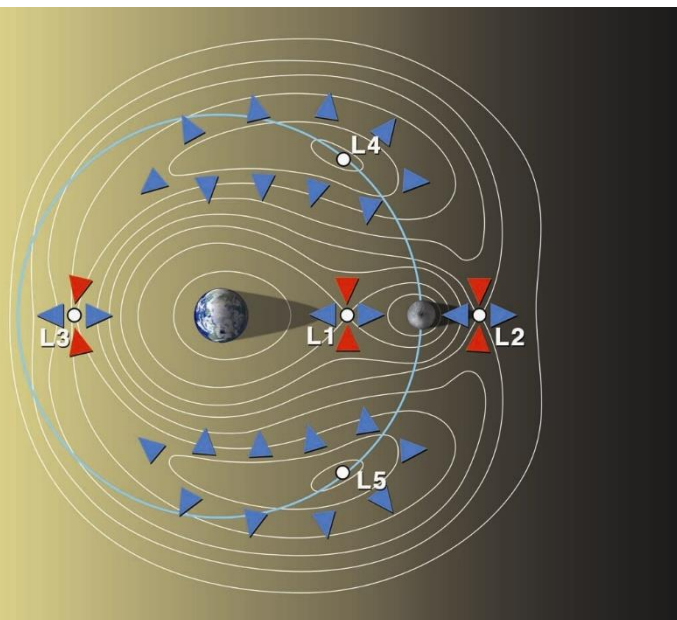
Effect of 3rd-Body Gravity

- Below is a histogram of inclination of objects in/around the GEO altitude (Source: space catalog downloaded from space-track.org)



Effect of 3rd-Body Gravity

- Another effect of 3rd-body gravity is the existence of Lagrange points
- These are points near 2 large attracting bodies where a smaller object will maintain its position relative to the large bodies
- Earth-Moon & Sun-Earth Lagrange points are shown below (5 pts each, labeled $L_1 - L_5$)

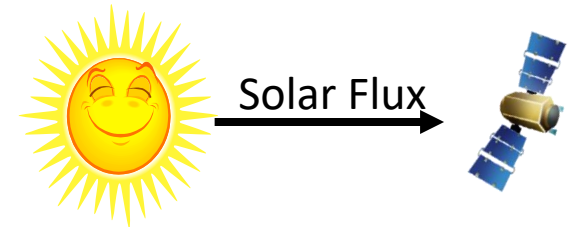


Effect of 3rd-Body Gravity

- Some interesting applications of Lagrange point missions include:
 - **Earth-Moon L_1** : Communication link at between Moon missions & ground stations on Earth
 - **Earth-Moon L_2** : Mission to explore the far side of the Moon
 - **Sun-Earth L_1** : Sun surveillance mission that is never eclipsed by Earth
 - **Sun-Earth L_2** : Mission whose objectives require eclipse conditions (would be in constant eclipse)
- For various reasons, objects in Lagrange point orbits aren't routinely tracked as part of SSA activities
 - Long distance from Earth would require large, expensive sensors
 - Very few objects currently in Lagrange point orbits → no threat of collisions, no "threatening" missions to be tracked
- However, 2 potential SSA implications of Lagrange point orbits are:
 - High-ground sensing: wide surveillance of the catalog from a strategic vantage point
 - If 1 or more Lagrange points eventually does get heavily populated, space-based sensors could be deployed in this vicinity

Solar Radiation Pressure

- Due to impingement of solar flux (i.e. photons) on an object's surface
- Many similarities between SRP & drag force
 - Both non-conservative
 - Both driven by solar phenomena
 - Both **highly** unpredictable
- Acceleration due to SRP is given by



$$\bar{a}_{SRP} = -\frac{\Phi}{c} \frac{A}{m} (\hat{u}_n \cdot \hat{u}_s) [(1 - \epsilon)\hat{u}_s + 2\epsilon(\hat{u}_n \cdot \hat{u}_s)\hat{u}_n]$$

- **Φ : solar flux**
- **c : speed of light**
- **A : object area (facing Sun)**
- **m : object mass**
- **ϵ : reflectivity coefficient**
- **\hat{u}_s : unit vector from object to Sun**
- **\hat{u}_n : unit vector normal to object surface (facing Sun)**

Effect of Solar Radiation Pressure

- Unlike drag, SRP doesn't always act against object's motion → not even purely along-track (forward or backward orbital direction)
 - To some degree, effect of SRP on object during half of its orbital path canceled out during the other half → but still has potential to add or subtract energy overall (i.e. raise or lower orbit)
- SRP basically independent of orbital altitude (since the Sun is 93 million miles from Earth)
 - Dwarfed by other perturbations (e.g. J2, drag) at lower altitudes
 - Significant (by comparison) at higher altitudes → GEO & above
- High area-to-mass ratio (HAMR) objects in GEO very sensitive to SRP
 - Since SRP is area-dependent, its effect on HAMR objects can be significant & unpredictable
 - Tracking/OD of these objects can be quite difficult!
- Concept of solar sailing has garnered interest: using high area-to-mass ratio surfaces with precise attitude control to utilize SRP for propulsion
 - Found to be feasible from GEO out to deep space/interplanetary orbits
 - Solar sail missions would provide significant tracking/OD challenges

Relative Magnitudes of Perturbations

- Below is a chart developed by Montenbruck & Gill (2012) depicting the relative magnitudes of various perturbations as a function of altitude
- Some things to note:
 - All Earth gravity forces **decrease** with altitude (consistent with Newton's Law of Gravitation)
 - SRP & 3rd body forces **increase** slightly with altitude (as object gets closer to Moon or Sun)
 - Around 600-800km altitude, SRP & 3rd body forces become **greater** than drag

