

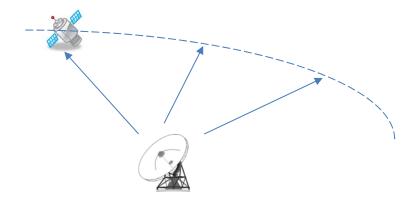
Orbit Determination I





What is Orbit Determination?

- Now that we understand the process of sensing space objects as part of a passive tracking campaign, we'll focus on orbit determination
- In short, orbit determination (OD) involves using observation data of a space object obtained from one or more sensors to accurately calculate & predict the object's orbit
- OD draws on a branch of mathematics known as state estimation



- State estimation is a technique whereby the current and past behavior of a dynamic system is discerned, based on observations of measurable quantities related to the system, in order that its future behavior may be predicted
- The dynamics of a system (i.e. its behavior over time) is often represented by a vector of "states"
 - Example would be a particle in 3-D motion \rightarrow states consist of its position components ([x y z]^T) & velocity components ([v_x v_y v_z]^T)
 - Together these 6 qty's form the state vector
- While a particle is an idealization, many practical system are modeled in this way, e.g. aircraft, ships, missiles, & space objects

 Change in the state vector with time is often described by a set of ordinary differential equations (ODEs) known as "state equations":

$$\dot{\overline{X}}(t) = \overline{F}(\overline{X}(t))$$

- Both vectors in the above equation are of dimension $n \times 1$, where n is the number of states
- In the case where the states consist of an object's position & velocity, $n = 6 \rightarrow \overline{F}$ is the vector of "state rates":
- The system's motion can be discerned by propagating (i.e. numerically integrating forward in time) the ODEs from some initial conditions \bar{X}_0
- So how can the initial conditions be determined in a given scenario?

- Unfortunately, in the case of most systems, their states can't be measured directly
- Qty's that can be measured, which are functions of the states, are often called "measurements," "outputs," or "observations"
- The functional relationship between measurements & states are usually governed by static (non-ODE) equations known as "measurement equations":

$$\overline{\mathbf{Y}}(\mathbf{t}) = \overline{\mathbf{G}}(\overline{\mathbf{X}}(\mathbf{t}))$$

 \overline{Y} and \overline{G} are $p \times 1$ vectors, where p is the total # of measurements

- A good example of a measurement-state relationship is a Pitot-static tube on an aircraft
- Total pressure in the atmosphere is the sum of static & dynamic pressure:

$$p_t = p + \frac{1}{2}\rho v^2$$

- Aircraft velocity can't be measured directly, but a Pitotstatic tube measures static & total pressure
- If air density is also known, velocity can be solved for:

$$v = \sqrt{\frac{2}{\rho}(p_t - p)}$$

- Let's extend these concepts of states & measurements to orbital mechanics of space objects
- A practical state vector for an Earth-orbiting space object is:
 - The object's position vector from Earth's center, expressed in ECI coordinates (denote this $[x \ y \ z]^T$)
 - The object's inertial velocity vector, also expressed in ECI coordinates (denote this $[v_x \ v_v \ v_z]^T$)
- For such objects, the 2-body ODEs of motion are well known:

$$\ddot{r} = -\frac{\mu \bar{r}}{r^3}$$

• where μ is Earth's gravitational constant & r is the object's instantaneous orbit radius (i.e. magnitude of its position vector from Earth's center): $r = \sqrt{x^2 + y^2 + z^2}$

 Assuming 2-body motion, the state equations for an Earth-orbiting space object can then be written as

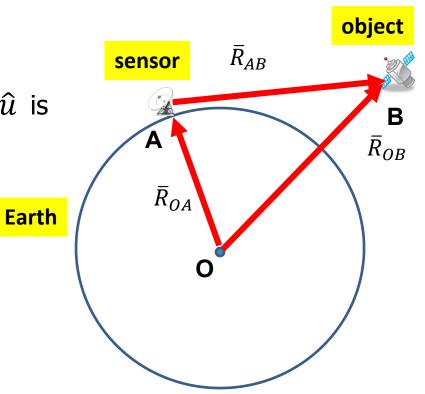
$$\frac{d}{dt} \begin{bmatrix} x \\ y \\ z \\ v_x \\ v_y \\ v_z \end{bmatrix} = \begin{bmatrix} v_x \\ v_y \\ -\frac{\mu x}{r^3} \\ -\frac{\mu y}{r^3} \\ -\frac{\mu z}{r^3} \end{bmatrix}$$

$$\dot{\overline{X}}(t) \quad \overline{F}(\overline{X}(t))$$

- Let us now derive the measurement equations
- Recall last week we examined sensor-to-object range & sensor-toobject line-of-sight (see sketch below)
- Let us write the vector $ar{R}_{AB}$ as $ho \hat{u}$

where ρ is sensor-to-object range & \hat{u} is sensor-to-object LOS:

$$\begin{bmatrix} u_{xTH} \\ u_{yTH} \\ u_{zTH} \end{bmatrix} = \begin{bmatrix} cosElcosAz \\ cosElsinAz \\ sinEl \end{bmatrix}$$



• We can then rewrite the vector equation $\bar{R}_{OB} = \bar{R}_{OA} + \bar{R}_{AB}$ as

$$\bar{r} = \bar{R}_{OA} + \rho \hat{u} \text{ OR } \bar{r} - \bar{R}_{OA} = \rho \hat{u}$$

- Where \bar{r} is the object's position vector from Earth's center (recall this comprises the first 3 states of the object's state vector)
- Typically \bar{r} is expressed in ECI coordinates
- Note that, if the sensor's lat/long location & sidereal time are known, \bar{R}_{OA} and \hat{u} can also be expressed in ECI coordinates (using coordinate transformations we've learned); write their ECI expressions as

$$\bar{R}_{OA} = \begin{bmatrix} R_{OAx} \\ R_{OAy} \\ R_{OAz} \end{bmatrix} \qquad \hat{u} = \begin{bmatrix} u_x \\ u_y \\ u_z \end{bmatrix}$$

• Therefore, ho is the magnitude of $ar{R}_{AB}$:

$$\rho = |\bar{r} - \bar{R}_{OA}| = \sqrt{(x - R_{OAx})^2 + (y - R_{OAy})^2 + (z - R_{OAz})^2}$$

• And \hat{u} is the unit vector in the direction of \bar{R}_{AB} :

$$\hat{u} = \frac{\bar{r} - \bar{R}_{OA}}{|\bar{r} - \bar{R}_{OA}|} \longrightarrow u_x = \frac{x - R_{OAx}}{\sqrt{(x - R_{OAx})^2 + (y - R_{OAy})^2 + (z - R_{OAz})^2}}$$

$$u_y = \frac{y - R_{OAy}}{\sqrt{(x - R_{OAx})^2 + (y - R_{OAy})^2 + (z - R_{OAz})^2}}$$

- Note that, because LOS is a unit vector, if any 2 of the 3 components (u_x, u_y, u_z) are known, the 3rd component is automatically known
 - Thus, we only obtain 2 (not 3) independent equations from a LOS measurement
 - Any 2 of the 3 component equations will suffice; here we arbitrarily choose the u_x & u_y equations

 Thus, we obtain one measurement equation for each radar measurement taken:

$$\rho = \sqrt{(x - R_{OAx})^2 + (y - R_{OAy})^2 + (z - R_{OAz})^2}$$

$$\uparrow \qquad \qquad \uparrow \qquad \qquad \downarrow \qquad$$

 & two measurement equations for each optical sensor measurement taken:

$$\begin{bmatrix} u_{x} \\ u_{y} \end{bmatrix} = \begin{bmatrix} \frac{x - R_{OAx}}{\sqrt{(x - R_{OAx})^{2} + (y - R_{OAy})^{2} + (z - R_{OAz})^{2}}} \\ \frac{y - R_{OAy}}{\sqrt{(x - R_{OAx})^{2} + (y - R_{OAy})^{2} + (z - R_{OAz})^{2}}} \end{bmatrix}$$

$$\overline{Y}(t)$$

$$\overline{G}(\overline{X}(t))$$

- Note that when we display the one range measurement equation, y is a scalar qty & g is a scalar function
- Whereas, when we display the two LOS measurement equations, Y & G
 are vectors
- What is the actual dimension of Y& G, & will they contain range equations, LOS equations, or both?
 - The answer depends on the total # and types of measurements to be processed
 - Some OD scenarios involve data from a single sensor, or multiple sensors of the same type (all RF or all optical) → these are "range-only" or "anglesonly" scenarios
 - Other scenarios involve data from multiple sensors of both types (RF & optical) → "range & angle" scenarios
 - This will become important later on

- We now have the fundamental building blocks for orbit determination: a model of the space object's dynamics & the measurement-state relationships
- OD is typically a 2-step process
 - Initial orbit determination (IOD) involves obtaining a decent (& fairly quick) solution from a minimal # of measurements
 - Precise orbit determination (POD) involves processing (usually) a large # of measurements, using an IOD solution as an initial guess (or "starter" solution)
- How do these 2 steps differ?
 - IOD normally considered a "deterministic" process, assuming no measurement or modeling error & simply fitting the measurements to the assumed dynamics
 - Whereas POD has the ability to accommodate error sources (e.g. sensor & modeling error) to find a statistical "best guess" of the object's orbit

This week we will focus on IOD

Initial Orbit Determination

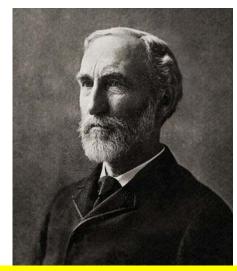
- IOD has a centuries-long history
- Early mathematicians & astronomers devised IOD methods to determine orbits of the Moon & deep space objects (e.g. planets, comets, asteroids)
- Based primarily on naked-eye observations (or with rudimentary telescopes)



Pierre-Simon Laplace



Carl Friedrich Gauss



Josiah Willard Gibbs

Initial Orbit Determination

- Since the dawn of the Space Age (Sputnik, 1957), IOD has become increasingly important for SSA
- Basically, IOD is necessary whenever a new object is detected
 - New launch of a satellite
 - Debris created from a breakup/collision event
 - Occasionally an object may be "lost" from the catalog, then "return"
- Some IOD methods based only on range measurements ("range-only"), while others based only on LOS measurements ("angles-only")
 - A few methods require both range & angles

Initial Orbit Determination

- Range-only vs angles-only IOD? Here are some considerations:
 - Radar sensors typically have sufficient signal power to sense objects in LEO, but very few can "see" beyond LEO
 - Telescopes can see celestial bodies in deep space (of course those bodies are massively large), but also mid- to large-size Earth-orbiting objects out to GEO
 - Telescopes require slewing (rotating) to track any non-GEO object as it moves across the sky; to track an object whose pass is particularly fast (e.g. LEO) may require unreasonably large slew rates
 - Result, in terms of SSA, is that low-altitude objects (e.g. LEO) tend to be sensed via radar, & high-altitude objects (e.g. GEO) tend to be sensed optically
- In this course, we will focus on angles-only OD
- We will now explore Laplace's angles-only IOD method, which he developed in 1780

- Requires 3 LOS measurements
 - Makes sense because each LOS measurement yields 2 equations → 3
 measurements yield 6 equations to solve 6 unknowns (object's position & velocity states)
- Begin by taking the 1st & 2nd time derivatives of $\bar{r} = \bar{R}_{OA} + \rho \hat{u}$ to obtain:

$$\dot{\bar{r}} = \dot{\bar{R}}_{OA} + \dot{\rho}\hat{u} + \rho\dot{\hat{u}}$$

$$\ddot{r} = \ddot{R}_{OA} + \ddot{\rho}\hat{u} + 2\dot{\rho}\dot{\hat{u}} + \rho \ddot{\hat{u}}$$

• Assuming 2-body motion, we can substitute $\ddot{r} = -\frac{\mu \bar{r}}{r^3}$ into the left-hand side:

$$-\frac{\mu \bar{r}}{r^3} = \ddot{\bar{R}}_{OA} + \ddot{\rho}\hat{u} + 2\dot{\rho}\dot{\hat{u}} + \rho \ddot{\hat{u}}$$

• Now substituting $\bar{r} = \bar{R}_{OA} + \rho \hat{u}$ into the left-hand side:

$$-\frac{\mu(\bar{R}_{OA} + \rho \hat{u})}{r^3} = \ddot{\bar{R}}_{OA} + \ddot{\rho}\hat{u} + 2\dot{\rho}\dot{\hat{u}} + \rho \ddot{\hat{u}}$$

• Collecting the terms involving ρ & its derivatives on one side:

$$\ddot{\rho}\hat{u} + 2\dot{\rho}\dot{\hat{u}} + \rho(\ddot{\hat{u}} + \frac{\mu}{r^3}\hat{u}) = -\ddot{\bar{R}}_{OA} - \frac{\mu R_{OA}}{r^3}$$

- These derived relationships are true at any time; but at which specific time(s) are we interested in evaluating the relationships?
- We will decide this later; for now, we will simply refer to this specific time as "t*"

- Let us assume for now that ρ & its derivatives are the only unknowns at t^* , & all other qty's are known
- We can then write the previous relationship in matrix form as:

$$\begin{bmatrix} | & | & | \\ \widehat{u} & 2\dot{\widehat{u}} & \ddot{\widehat{u}} + \frac{\mu}{r^3} \widehat{u} \end{bmatrix} \begin{bmatrix} \ddot{\rho} \\ \dot{\rho} \\ \rho \end{bmatrix} = \begin{bmatrix} -\ddot{R}_{OA} - \frac{\mu \bar{R}_{OA}}{r^3} \\ | & | & | & | \end{bmatrix}$$
3x3 matrix

3x1 vector

3x1 vector

(Note that \hat{u} , $2\dot{\hat{u}}$, & $\ddot{\hat{u}}$ + $\frac{\mu}{r^3}\hat{u}$ are each 3x1 vectors, so

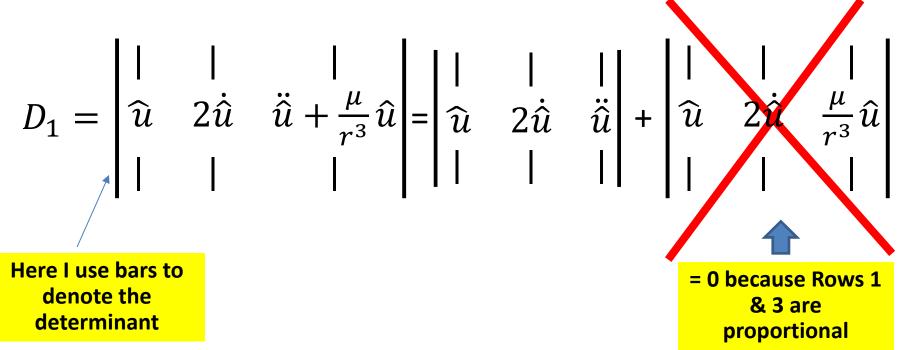
concatenating them together produces a 3x3 matrix)

We can use Cramer's rule to solve ρ :

where
$$D_1$$
 is the determinant of
$$\begin{bmatrix} | & | & | \\ \widehat{u} & 2\dot{\widehat{u}} & \ddot{\widehat{u}} + \frac{\mu}{r^3}\widehat{u} \end{bmatrix}$$

And
$$D_2$$
 is the determinant of
$$\begin{bmatrix} | & | & | \\ \widehat{u} & 2\dot{\widehat{u}} & -\ddot{\overline{R}}_{OA} - \frac{\mu \overline{R}_{OA}}{r^3} \\ | & | & | \end{bmatrix}$$

Note that D₁ can be simplified to



$$= \begin{vmatrix} 1 & 1 & 1 \\ \widehat{u} & 2\hat{u} & \ddot{u} \end{vmatrix} = 2 \begin{vmatrix} 1 & 1 & 1 \\ \widehat{u} & \dot{u} & \ddot{u} \end{vmatrix}$$

• And D_2 can be rewritten as

If we define

$$D_2^* = \begin{vmatrix} | & | & | & | \\ \widehat{u} & \dot{\widehat{u}} & \ddot{R}_{OA} \\ | & | & | \end{vmatrix} \text{ and } D_2^{**} = \begin{vmatrix} | & | & | \\ \widehat{u} & \dot{\widehat{u}} & \bar{R}_{OA} \\ | & | & | \end{vmatrix}$$

- Then D_2 becomes $D_2 = -2D_2^* 2\frac{\mu}{r^3}D_2^{**}$
- Thus, ρ can be expressed as $\rho = \frac{D_2}{D_1} = -2\frac{D_2^*}{D_1} 2\frac{\mu}{r^3}\frac{D_2^{**}}{D_1}$
- ho is then a function of \hat{u} , $\dot{\hat{u}}$, $\ddot{\hat{u}}$, r, $ar{R}_{OA}$, and $ar{R}_{OA}$
- We will now derive expressions for some of these qty's

- Recall that this method requires LOS measurements at 3 times; call these LOS vectors \hat{u}_1 , \hat{u}_2 , \hat{u}_3 at times t_1 , t_2 , t_3
- Thus, if we wish to evaluate our expression for ρ at $t^* = t_1$, t_2 , or t_3 , \hat{u} at these times is known; what about \hat{u} and \hat{u} ?
- We can use Lagrange interpolation to fit the LOS history over the span of measurements (between $t_1 \& t_3$) to a 2nd-order polynomial:

$$\hat{u}(t) = \frac{(t - t_2)(t - t_3)}{(t_1 - t_2)(t_1 - t_3)} \hat{u}_1 + \frac{(t - t_1)(t - t_3)}{(t_2 - t_1)(t_2 - t_3)} \hat{u}_2 + \frac{(t - t_1)(t - t_2)}{(t_3 - t_1)(t_3 - t_2)} \hat{u}_3$$







Note that evaluating this formula at t_1 yields \hat{u}_1 , at t_2 yields \hat{u}_2 , & at t_3 yields \hat{u}_3 (as expected)

Taking the 1st & 2nd time derivatives of this expression yields:

$$\dot{\hat{u}}(t) = \frac{2t - t_2 - t_3}{(t_1 - t_2)(t_1 - t_3)} \hat{u}_1 + \frac{2t - t_1 - t_3}{(t_2 - t_1)(t_2 - t_3)} \hat{u}_2 + \frac{2t - t_1 - t_2}{(t_3 - t_1)(t_3 - t_2)} \hat{u}_3$$

$$\ddot{\hat{u}}(t) = \frac{2}{(t_1 - t_2)(t_1 - t_3)} \hat{u}_1 + \frac{2}{(t_2 - t_1)(t_2 - t_3)} \hat{u}_2 + \frac{2}{(t_3 - t_1)(t_3 - t_2)} \hat{u}_3$$

- The expectation is that these formulas should provide a decent approximation for \hat{u} and \hat{u} in the time range from t_1 to t_3
- If we insert the values for our measurements times t_1 , t_2 , & t_3 into these formulas, we see that $\hat{\vec{u}}$ is constant & $\hat{\vec{u}}$ is linear in time
- Thus, \hat{u} & \hat{u} can be approximated at any time from t_1 to t_3 by the above formulas

- Now what about \bar{R}_{OA} and $\ddot{\bar{R}}_{OA}$?
- \bar{R}_{OA} is the position vector from Earth's center to the sensor
 - Generally, a sensor's location on Earth is precisely known (often expressed in latitude & longitude)
 - From latitude & longitude, we know how to calculate the sensor's ECEF position vector, then if sidereal time at a particular instant is known, we can calculate the sensor's ECI position vector as well
- \bar{R}_{OA} is the acceleration vector of this point on the Earth
 - Assuming Earth is a rotating sphere, kinematics tells us that this acceleration is $\overline{\omega}_E \ x \ (\overline{\omega}_E \ x \ \overline{R}_{OA})$, where $\overline{\omega}_E$ is the rotational velocity vector of the Earth about its axis

- All vectors in these calculations should be expressed in ECI coordinates; so how do we express \ddot{R}_{OA} in ECI?
- ECI expression of $\overline{\omega}_E$ is $\begin{bmatrix} 0 \\ 0 \\ \omega_E \end{bmatrix}$, where $\omega_E = 360^\circ$ per sidereal day
- ECI expression of \bar{R}_{OA} is $\begin{bmatrix} R_{OAx} \\ R_{OAy} \\ R_{OAz} \end{bmatrix}$ (discussed on Slide 10 how to calculate this)
- Inserting into $\overline{\omega}_E x (\overline{\omega}_E x \overline{R}_{OA})$ will yield a vector in ECI

- Recall that ρ at a particular time has been shown to be a function of \hat{u} , $\dot{\hat{u}}$, \ddot{u} , r, \bar{R}_{OA} , and \ddot{R}_{OA}
- If we choose to cacluate ρ at one of the measurement times $(t_1, t_2, \text{ or } t_3)$, all of these qty's are known (i.e. we know how to cacluate them) except r
- Let us now take the formula $\bar{r} = \bar{R}_{OA} + \rho \hat{u}$ & dot it with itself:

$$\bar{r} \cdot \bar{r} = (\bar{R}_{OA} + \rho \hat{u}) \cdot (\bar{R}_{OA} + \rho \hat{u}) = \rho^2 \hat{u} \cdot \hat{u} + 2\rho \bar{R}_{OA} \cdot \hat{u} + \bar{R}_{OA} \cdot \bar{R}_{OA}$$

$$\rightarrow r^2 = \rho^2 + 2\rho \bar{R}_{OA} \cdot \hat{u} + R_{OA}^2$$

- Note that, for a spherical Earth, $R_{OA} = R_E$ (Earth's radius)
- Also, because \bar{R}_{OA} & \hat{u} are known at t_{1} , t_{2} , & t_{3} , $\bar{R}_{OA} \cdot \hat{u}$ is known \rightarrow refer to this qty as "C":

$$r^2 = \rho^2 + 2C\rho + R_E^2$$

• Inserting $\rho = -2\frac{D_2^*}{D_1} - 2\frac{\mu}{r^3}\frac{D_2^{**}}{D_1}$ into this formula yields:

$$r^{2} = 4\left(\frac{D_{2}^{*}}{D_{1}} + \frac{\mu}{r^{3}} \frac{D_{2}^{**}}{D_{1}}\right)^{2} - 4C\left(\frac{D_{2}^{*}}{D_{1}} + \frac{\mu}{r^{3}} \frac{D_{2}^{**}}{D_{1}}\right) + R_{E}^{2}$$

Expanding, multiplying through to eliminate r from the denominator,
 & grouping powers of r yields:

$$r^{8} + \left(4C\frac{D_{2}^{*}}{D_{1}} - 4\frac{D_{2}^{*2}}{D_{1}^{2}} - R_{E}^{2}\right)r^{6} + \mu\left(4C\frac{D_{2}^{**}}{D_{1}} - 8\frac{D_{2}^{*}D_{2}^{**}}{D_{1}^{2}}\right)r^{3} - 4\mu^{2}\frac{D_{2}^{**2}}{D_{1}^{2}} = 0$$

- This is an 8th-order polynomial in *r* (object's orbit radius, i.e. magnitude of its position vector)
- Once we decide the time at which we wish to evaluate the polynomial, we insert all the known qty's at that time to calculate the coefficients → we can then solve for the 8 roots
- Depending on the accuracy of the method, one of these roots should be at least fairly close to the true value of r at that time, & it is this root that we choose to be the "solution" value of r

So how does this value of r yield the orbit solution?

- Insert r into $\rho = -2\frac{D_2^*}{D_1} 2\frac{\mu}{r^3}\frac{D_2^{**}}{D_1}$ to yield ρ , then insert ρ into $\bar{r} = \bar{R}_{OA} + \rho\hat{u}$ to yield \bar{r} , whose coordinates are (x, y, z)
- We've now solved 3 of the object's 6 orbital states; what about the other 3 states v_x , v_y , v_z ?

First we solve p from

$$\begin{bmatrix} | & | & | \\ \hat{u} & 2\dot{\hat{u}} & \ddot{\hat{u}} + \frac{\mu}{r^3}\hat{u} \end{bmatrix} \begin{bmatrix} \ddot{\rho} \\ \dot{\rho} \\ \rho \end{bmatrix} = \begin{bmatrix} -\ddot{R}_{OA} - \frac{\mu\bar{R}_{OA}}{r^3} \\ | & | \end{bmatrix}$$

Use Cramer's rule as before:

$$D_{1} = \begin{vmatrix} | & | & | & | \\ \widehat{u} & 2\dot{\hat{u}} & \ddot{\hat{u}} + \frac{\mu}{r^{3}} \widehat{u} \\ | & | & | & | \end{vmatrix} = 2 \begin{vmatrix} | & | & | & | \\ \widehat{u} & \dot{\hat{u}} & \ddot{\hat{u}} \\ | & | & | & | \end{vmatrix}, D_{2} = \begin{vmatrix} | & | & | & | \\ \widehat{u} & -\ddot{R}_{OA} - \frac{\mu \bar{R}_{OA}}{r^{3}} & \ddot{\hat{u}} + \frac{\mu}{r^{3}} \widehat{u} \end{vmatrix} \rightarrow \dot{\rho} = \frac{D_{2}}{D_{1}}$$



This determinant may be simplified if desired (as in the case where we solved for ρ)

- Next we solve $\dot{\bar{R}}_{OA}$, the velocity vector of the point on the Earth where the sensor lies
 - Assuming Earth is a rotating sphere, kinematics tells us that this velocity is $\overline{\omega}_E \propto \overline{R}_{OA}$ (be sure to calculate in ECI coordinates, as discussed previously)
- Insert $\dot{\rho}$ & $\dot{\bar{R}}_{OA}$ into $\dot{\bar{r}}=\dot{\bar{R}}_{OA}+\dot{\rho}\hat{u}+\rho\hat{u}$ to yield $\dot{\bar{r}}$, whose coordinates are (v_x , v_v , v_z)
- Thus ends the Laplace method!

A FEW ISSUES...

- 1) Which of the 8 roots is the "correct" value of r to choose?
- Begin by process of elimination:
 - Eliminate complex roots (come in conjugate pairs)
 - Eliminate values too small ($< R_E$) or large (beyond the radius of any reasonable Earth orbit)
- If more than 1 candidate value remain, these values will likely be significantly different from one another
 - If we have at least a vague a priori idea of the object's orbit radius (e.g. LEO, GPS, GEO), we should be able to select the correct value of r

2) What accuracy do we expect from Laplace's method?

- There are 2 approximations employed by the method:
 - 2-body dynamics assumed for the object's motion
 - Lagrange interpolation to fit the LOS history (& its derivatives) between t₁ & t₃
- Therefore, the more accurately these approximations represent the object's true orbital motion, the more accurate the Laplace orbit solution will be
- Example scenarios where we expect more accuracy are:
 - Shorter time spans → more accurate interpolation, & less time for propagation error to build up
 - Orbits with less perturbation force (GEO vs LEO)

- 3) At what time(s) in the object's orbit is it appropriate to evaluate the polynomial (& therefore solve the object's orbit)?
- Technically, we can evaluate the polynomial at any time
 - Coefficients are functions of \hat{u} , $\dot{\hat{u}}$, $\ddot{\bar{u}}$, \bar{R}_{OA} , and $\ddot{\bar{R}}_{OA}$
 - Lagrange interpolation provides formulas for \hat{u} , $\dot{\hat{u}}$, and $\dot{\hat{u}}$, which can be used to evaluate these qty's at any time
 - \bar{R}_{OA} and $\ddot{\bar{R}}_{OA}$ can also be easily determined at any time
- Practically, it is best to evaluate the polynomial at the time(s) where the Lagrange interpolation provides the best approximation to LOS
 - This approximation is generally much better between $t_1 \& t_3$ than outside of this range
 - Most IOD implementations choose t_2 to evaluate the polynomial because it is between the endpoints of the approximation (t_1 & t_3)

These 3 issues have been the subject of extensive research \rightarrow the answer in each case tends to be very scenario-dependent

IOD Wrap-up

- Though Laplace's method is centuries old, it is still used today as a practical means of optical IOD (for both deepspace & Earth-orbiting objects)
- Other angles-only IOD methods include:
 - Gauss' method (1800)
 - Double-R method (Escobal, 1965)
 - Taff's method (1983)
 - Gooding's method (1993)