Mathematics of Physical Quantities and Units:

Dimensional Analysis and The Π Theorem

Sean Ericson

UO

Theory meeting, October 27, 2023



Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Note that we consider $\mathcal S$ as a 1-dimensional subspace of $\mathcal P$.

Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Note that we consider $\mathcal S$ as a 1-dimensional subspace of $\mathcal P$.

We can describe their relationship via the short exact sequence

Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Note that we consider ${\mathcal S}$ as a 1-dimensional subspace of ${\mathcal P}.$

We can describe their relationship via the short exact sequence

$$\mathcal{E} = \left(\mathcal{S} \xrightarrow{i} \mathcal{P} \xrightarrow{d} \mathcal{D} \right),$$

where

Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Note that we consider ${\mathcal S}$ as a 1-dimensional subspace of ${\mathcal P}.$

We can describe their relationship via the short exact sequence

$$\mathcal{E} = \left(\mathcal{S} \xrightarrow{i} \mathcal{P} \xrightarrow{d} \mathcal{D} \right),$$

where

i is the inclusion map of S into P,

Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Note that we consider ${\mathcal S}$ as a 1-dimensional subspace of ${\mathcal P}.$

We can describe their relationship via the short exact sequence

$$\mathcal{E} = \left(\mathcal{S} \xrightarrow{i} \mathcal{P} \xrightarrow{d} \mathcal{D} \right),$$

where

i is the inclusion map of S into P, d is a surjection, and

Physical quantities (\mathcal{P}) , their units (\mathcal{D}) , and the set of scalars (\mathcal{S}) all form \mathbb{Q} -vector spaces.

Note that we consider ${\mathcal S}$ as a 1-dimensional subspace of ${\mathcal P}.$

We can describe their relationship via the short exact sequence

$$\mathcal{E} = \left(\mathcal{S} \xrightarrow{i} \mathcal{P} \xrightarrow{d} \mathcal{D} \right),$$

where

i is the inclusion map of S into P, d is a surjection, and ker[d] = im[i].



Denote preimages of d via $\langle D \rangle \coloneqq d^{-1}(D)$ for $D \in \mathcal{D}$.

Denote preimages of d via $\langle D \rangle := d^{-1}(D)$ for $D \in \mathcal{D}$. i.e. "the set of quantities with units D"

Denote preimages of d via $\langle D \rangle \coloneqq d^{-1}(D)$ for $D \in \mathcal{D}$. i.e. "the set of quantities with units D" Note that $\langle 0 \rangle = \mathcal{S}$.

Denote preimages of d via $\langle D \rangle \coloneqq d^{-1}(D)$ for $D \in \mathcal{D}$. i.e. "the set of quantities with units D" Note that $\langle 0 \rangle = \mathcal{S}$.

For all $q_1, q_2 \in \langle D \rangle$, we have that $q_1 - q_2 \in \mathcal{S}$

Denote preimages of d via $\langle D \rangle \coloneqq d^{-1}(D)$ for $D \in \mathcal{D}$. i.e. "the set of quantities with units D" Note that $\langle 0 \rangle = \mathcal{S}$.

For all
$$q_1,q_2\in\langle D
angle$$
, we have that $q_1-q_2\in\mathcal{S}$ $d(q_1-q_2)=d(q_1)-d(q_2)=D-D=0$

Denote preimages of d via $\langle D \rangle \coloneqq d^{-1}(D)$ for $D \in \mathcal{D}$. i.e. "the set of quantities with units D" Note that $\langle 0 \rangle = \mathcal{S}$.

For all
$$q_1,q_2\in\langle D
angle$$
, we have that $q_1-q_2\in\mathcal{S}$ $d(q_1-q_2)=d(q_1)-d(q_2)=D-D=0$

For any $\lambda \in \mathcal{S}$ and $q \in \langle D \rangle$, we have that $\lambda + q \in \langle D \rangle$

Denote preimages of d via $\langle D \rangle \coloneqq d^{-1}(D)$ for $D \in \mathcal{D}$. i.e. "the set of quantities with units D" Note that $\langle 0 \rangle = \mathcal{S}$.

For all
$$q_1,q_2\in\langle D
angle$$
, we have that $q_1-q_2\in\mathcal{S}$ $d(q_1-q_2)=d(q_1)-d(q_2)=D-D=0$

For any
$$\lambda \in \mathcal{S}$$
 and $q \in \langle D \rangle$, we have that $\lambda + q \in \langle D \rangle$
$$d(\lambda + q) = d(\lambda) + d(q) = D + 0 = D$$

A linear map $U: \mathcal{D} \to \mathcal{P}$ such that $d(U(D)) = D \quad \forall D \in \mathcal{D}$ is called a *linear splitting*.

A linear map $U: \mathcal{D} \to \mathcal{P}$ such that $d(U(D)) = D \quad \forall D \in \mathcal{D}$ is called a *linear splitting*.

Any such splitting induces a (linear) map $V:\mathcal{P} o\mathcal{S}$ via V(q)=q-U(d(q))

A linear map $U: \mathcal{D} \to \mathcal{P}$ such that $d(U(D)) = D \quad \forall D \in \mathcal{D}$ is called a *linear splitting*.

Any such splitting induces a (linear) map $V:\mathcal{P} o\mathcal{S}$ via V(q)=q-U(d(q))

We can define the map $\chi: \mathcal{S} \times \mathcal{D} \to \mathcal{P}$ by

$$\chi(\lambda, D) = \lambda + U(D),$$

A linear map $U: \mathcal{D} \to \mathcal{P}$ such that $d(U(D)) = D \quad \forall D \in \mathcal{D}$ is called a *linear splitting*.

Any such splitting induces a (linear) map $V:\mathcal{P} o\mathcal{S}$ via V(q)=q-U(d(q))

We can define the map $\chi: \mathcal{S} \times \mathcal{D} \to \mathcal{P}$ by

$$\chi(\lambda,D)=\lambda+U(D),$$

and its inverse

$$\chi^{-1}(q) = (V(q), d(q))$$



A linear map $U: \mathcal{D} \to \mathcal{P}$ such that $d(U(D)) = D \quad \forall D \in \mathcal{D}$ is called a *linear splitting*.

Any such splitting induces a (linear) map $V:\mathcal{P} o\mathcal{S}$ via V(q)=q-U(d(q))

We can define the map $\chi: \mathcal{S} \times \mathcal{D} \to \mathcal{P}$ by

$$\chi(\lambda,D)=\lambda+U(D),$$

and its inverse

$$\chi^{-1}(q) = (V(q), d(q))$$

 χ as given is a bijection, hence $\mathcal{S} \times \mathcal{D} \simeq \mathcal{P}$.

We denote by $G(\mathcal{E})$ the group of linear maps from \mathcal{D} to \mathcal{S} , $\text{Lin}[\mathcal{D},\mathcal{S}]$.

We denote by $G(\mathcal{E})$ the group of linear maps from \mathcal{D} to \mathcal{S} , $\text{Lin}[\mathcal{D},\mathcal{S}]$.

An n-ary relation $F\subseteq \mathcal{P}^n$ is said to be $G(\mathcal{E})$ -invariant if

We denote by $G(\mathcal{E})$ the group of linear maps from \mathcal{D} to \mathcal{S} , $\text{Lin}[\mathcal{D},\mathcal{S}]$.

An n-ary relation $F \subseteq \mathcal{P}^n$ is said to be $G(\mathcal{E})$ -invariant if

$$(q_1, \cdots, q_n) \in F \implies (q_1 + \theta(d(q_1)), \cdots, q_n + \theta(d(q_n))) \in F$$

We denote by $G(\mathcal{E})$ the group of linear maps from \mathcal{D} to \mathcal{S} , $\text{Lin}[\mathcal{D},\mathcal{S}]$.

An n-ary relation $F \subseteq \mathcal{P}^n$ is said to be $G(\mathcal{E})$ -invariant if

$$(q_1, \cdots, q_n) \in F \implies (q_1 + \theta(d(q_1)), \cdots, q_n + \theta(d(q_n))) \in F$$

Example: lwh - V = 0

Statement

Statement

Consider a $G(\mathcal{E})$ -invariant $F \subseteq \langle E_1 \rangle, \cdots, \langle E_m \rangle, \langle D_1 \rangle, \cdots, \langle D_n \rangle$ where the $\{E_i\}$ are linearly independent and $D_j = \sum k_{ji} E_i$.

Statement

Consider a $G(\mathcal{E})$ -invariant

 $F \subseteq \langle E_1 \rangle, \dots, \langle E_m \rangle, \langle D_1 \rangle, \dots, \langle D_n \rangle$ where the $\{E_i\}$ are linearly independent and $D_j = \sum k_{ji} E_i$.

Then, there exists $\Phi \subseteq \mathcal{S}^n$ such that

$$(p_1,\cdots,p_m,q_1,\cdots,q_n)\in F$$

$$\implies (q_1 - \sum k_{1i}p_i, \cdots, q_n - \sum k_{ni}p_i) \in \Phi$$

Proof

Proof

Let U_0 be an arbitrary splitting of d.

Proof

Let U_0 be an arbitrary splitting of d.

Define
$$\Phi = \{(\lambda_1, \dots, \lambda_n) \in \mathcal{S}^n | (U_0(E_1), \dots, U_0(E_m), \lambda_1 + U_0(D_1), \dots, U_0(D_n)) \in F\}$$

Proof

Let U_0 be an arbitrary splitting of d.

Define
$$\Phi = \{(\lambda_1, \dots, \lambda_n) \in S^n | (U_0(E_1), \dots, U_0(E_m), \lambda_1 + U_0(D_1), \dots, U_0(D_n)) \in F\}$$

The set $\{E_i\}$ can be extended to a basis for \mathcal{D} . We can then define a new splitting U such that $U(E_i) = p_i$.

Proof

Let U_0 be an arbitrary splitting of d.

Define
$$\Phi = \{(\lambda_1, \dots, \lambda_n) \in S^n | (U_0(E_1), \dots, U_0(E_m), \lambda_1 + U_0(D_1), \dots, U_0(D_n)) \in F\}$$

The set $\{E_i\}$ can be extended to a basis for \mathcal{D} . We can then define a new splitting U such that $U(E_i) = p_i$.

Note
$$(\vec{p}, \vec{q}) \in F \implies (U(\vec{p}), \vec{q}) \in F$$

Proof

Let U_0 be an arbitrary splitting of d.

Define
$$\Phi = \{(\lambda_1, \dots, \lambda_n) \in S^n | (U_0(E_1), \dots, U_0(E_m), \lambda_1 + U_0(D_1), \dots, U_0(D_n)) \in F\}$$

The set $\{E_i\}$ can be extended to a basis for \mathcal{D} . We can then define a new splitting U such that $U(E_i) = p_i$.

Note
$$(\vec{p}, \vec{q}) \in F \implies (U(\vec{p}), \vec{q}) \in F$$

Consider $\theta: \mathcal{D} \to \mathcal{P}$ defined by $\theta(D) = U_0(D) - U(D)$.

Proof

Let U_0 be an arbitrary splitting of d.

Define
$$\Phi = \{(\lambda_1, \dots, \lambda_n) \in S^n | (U_0(E_1), \dots, U_0(E_m), \lambda_1 + U_0(D_1), \dots, U_0(D_n)) \in F\}$$

The set $\{E_i\}$ can be extended to a basis for \mathcal{D} . We can then define a new splitting U such that $U(E_i) = p_i$.

Note
$$(\vec{p}, \vec{q}) \in F \implies (U(\vec{p}), \vec{q}) \in F$$

Consider
$$\theta: \mathcal{D} \to \mathcal{P}$$
 defined by $\theta(D) = U_0(D) - U(D)$.

Note
$$\operatorname{im}[\theta] = \mathcal{S}$$

Proof

Let U_0 be an arbitrary splitting of d.

Define
$$\Phi = \{(\lambda_1, \dots, \lambda_n) \in S^n | (U_0(E_1), \dots, U_0(E_m), \lambda_1 + U_0(D_1), \dots, U_0(D_n)) \in F \}$$

The set $\{E_i\}$ can be extended to a basis for \mathcal{D} . We can then define a new splitting U such that $U(E_i) = p_i$.

Note
$$(\vec{p}, \vec{q}) \in F \implies (U(\vec{p}), \vec{q}) \in F$$

Consider
$$\theta: \mathcal{D} \to \mathcal{P}$$
 defined by $\theta(D) = U_0(D) - U(D)$.

Note
$$\mathsf{im}[\theta] = \mathcal{S}$$

Then, applying $G(\mathcal{E})$ -invariance of F implies

$$(U_0(E_1), \dots, U_0(E_m), q_1+U_0(D_1)-U(D_1), \dots, q_n+U_0(D_n)-U(D_n))$$

$$\in F$$



Proof cont.

The Π Theorem for Relations

Proof cont.

Comparing to the definition of Φ , we see that

$$(q_1-U(D_1),\cdots,q_n-U(D_n))\in\Phi,$$

and

The Π Theorem for Relations

Proof cont.

Comparing to the definition of Φ , we see that

$$(q_1-U(D_1),\cdots,q_n-U(D_n))\in\Phi,$$

and

$$U(D_i) = U(\sum k_{ij}E_j) = \sum k_{ji}U(E_j) = \sum k_{ij}p_j\Box$$

We say a function $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ is $G(\mathcal{E})$ -invariant if, for all $\theta \in G(\mathcal{E})$,

$$f(p_1, \dots, q_n) = c$$

 $\implies f(p_1 + \theta(E_1), \dots, q_n + \theta(D_n)) = c + \theta(C)$

We say a function $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ is $G(\mathcal{E})$ -invariant if, for all $\theta \in G(\mathcal{E})$,

$$f(p_1,\cdots,q_n)=c$$

$$\implies f(p_1 + \theta(E_1), \cdots, q_n + \theta(D_n)) = c + \theta(C)$$

Prop: If f is $G(\mathcal{E})$ -invariant, then $C \in \text{Span}\{E_1, \dots, E_m, D_1, \dots, D_n\}$

We say a function $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ is $G(\mathcal{E})$ -invariant if, for all $\theta \in G(\mathcal{E})$,

$$f(p_1,\cdots,q_n)=c$$

$$\implies f(p_1 + \theta(E_1), \cdots, q_n + \theta(D_n)) = c + \theta(C)$$

Prop: If f is $G(\mathcal{E})$ -invariant, then $C \in \text{Span}\{E_1, \dots, E_m, D_1, \dots, D_n\}$

Proof: Assume not, then consider θ such that

$$\theta(E_1) = \cdots = \theta(D_n) = 0$$
, but $\theta(C) \neq 0$.

We say a function $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ is $G(\mathcal{E})$ -invariant if, for all $\theta \in G(\mathcal{E})$,

$$f(p_1,\cdots,q_n)=c$$

$$\implies f(p_1 + \theta(E_1), \cdots, q_n + \theta(D_n)) = c + \theta(C)$$

Prop: If f is $G(\mathcal{E})$ -invariant, then $C \in \text{Span}\{E_1, \dots, E_m, D_1, \dots, D_n\}$

Proof: Assume not, then consider θ such that

$$\theta(E_1) = \cdots = \theta(D_n) = 0$$
, but $\theta(C) \neq 0$.

f is clearly not invariant under this map, a contradiction \square .



Statement

Statement

Let $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ be a $G(\mathcal{E})$ -invariant function.

Statement

Let $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ be a $G(\mathcal{E})$ -invariant function.

Then,

Statement

Let $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ be a $G(\mathcal{E})$ -invariant function.

Then,

1) There exists k_{ij} such that $C = \sum k_{n+1,i} E_i$

Statement

Let $f: \langle E_1 \rangle \times \cdots \times \langle E_m \rangle \times \langle D_1 \rangle \cdots \langle D_n \rangle \rightarrow \langle C \rangle$ be a $G(\mathcal{E})$ -invariant function.

Then,

- 1) There exists k_{ij} such that $C = \sum k_{n+1,i} E_i$
- 2) There exists $\phi: \mathcal{S}^n \to \mathcal{S}$ such that for all $(p_1, \dots, q_n) \in \langle E_1 \rangle \times \dots \times \langle D_n \rangle$,

$$f(\vec{p},\vec{q}) = \phi(q_1 - \sum k_{1j}p_j, \cdots, q_n - \sum k_{nj}p_j) + \sum k_{n+1,j}p_j$$

Proof

Proof

The proof of 1) follows immediately from the Obvious Proposition and the linear dependence of the D_i on the E_i .

Proof

The proof of 1) follows immediately from the Obvious Proposition and the linear dependence of the D_i on the E_i .

The proof of 2) follows from the proof the the Π Theorem for Relations:

Proof

The proof of 1) follows immediately from the Obvious Proposition and the linear dependence of the D_i on the E_i .

The proof of 2) follows from the proof the the Π Theorem for Relations:

We consider the graph of f as an (n+1)-ary relation.

Proof

The proof of 1) follows immediately from the Obvious Proposition and the linear dependence of the D_i on the E_i .

The proof of 2) follows from the proof the the Π Theorem for Relations:

We consider the graph of f as an (n+1)-ary relation. Construct $\Phi \subseteq \mathcal{S}^{n+1}$ as in the proof of the relation version

Proof

The proof of 1) follows immediately from the Obvious Proposition and the linear dependence of the D_i on the E_i .

The proof of 2) follows from the proof the the Π Theorem for Relations:

We consider the graph of f as an (n+1)-ary relation. Construct $\Phi \subseteq \mathcal{S}^{n+1}$ as in the proof of the relation version

 ϕ is then defined by

$$\Phi = (\lambda_1, \cdots, \lambda_n, \phi(\lambda_1, \cdots, \lambda_n))$$

Consider the relation

$$F(m, \hbar, \omega; x, p, H) = H - \frac{p^2}{2m} - \frac{1}{2}m\omega^2 x^2 = 0$$

Consider the relation

$$F(m, \hbar, \omega; x, p, H) = H - \frac{p^2}{2m} - \frac{1}{2}m\omega^2 x^2 = 0$$

In the basis $\mathcal{B} = \{\mathsf{Mass}, \mathsf{Length}, \mathsf{Time}\}$ of \mathcal{D} we have

Consider the relation

$$F(m, \hbar, \omega; x, p, H) = H - \frac{p^2}{2m} - \frac{1}{2}m\omega^2 x^2 = 0$$

In the basis $\mathcal{B} = \{ \mathsf{Mass}, \mathsf{Length}, \mathsf{Time} \}$ of \mathcal{D} we have

$$[m] = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \ [\hbar] = \begin{pmatrix} 1 \\ 2 \\ -1 \end{pmatrix}, \ [\omega] = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}.$$

Consider the relation

$$F(m, \hbar, \omega; x, p, H) = H - \frac{p^2}{2m} - \frac{1}{2}m\omega^2 x^2 = 0$$

In the basis $\mathcal{B} = \{\mathsf{Mass}, \mathsf{Length}, \mathsf{Time}\}\ \mathsf{of}\ \mathcal{D}\ \mathsf{we}\ \mathsf{have}$

$$[m] = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \ [\hbar] = \begin{pmatrix} 1 \\ 2 \\ -1 \end{pmatrix}, \ [\omega] = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}.$$

These vectors span the space:

$$|\mathcal{M}| = \begin{vmatrix} 1 & 1 & 0 \\ 0 & 2 & 0 \\ 0 & -1 & -1 \end{vmatrix} = -2 \neq 0$$



Example: The Harmonic Oscillator (cont.)

We can then solve for the dimensions of the variables in terms of the dimensions of the parameters:

Example: The Harmonic Oscillator (cont.)

We can then solve for the dimensions of the variables in terms of the dimensions of the parameters:

$$[x] = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}; \quad \mathcal{M}^{-1}[x] = \begin{pmatrix} -1/2 \\ 1/2 \\ -1/2 \end{pmatrix} \implies [x] = \left[\sqrt{\frac{\hbar}{m\omega}} \right]$$
$$[p] = \begin{pmatrix} 1 \\ 1 \\ -1 \end{pmatrix}; \quad \mathcal{M}^{-1}[p] = \begin{pmatrix} 1/2 \\ 1/2 \\ 1/2 \end{pmatrix} \implies [p] = \left[\sqrt{m\hbar\omega} \right]$$
$$[H] = \begin{pmatrix} 1 \\ 2 \\ -2 \end{pmatrix}; \quad \mathcal{M}^{-1}[H] = \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix} \implies [H] = [\hbar\omega]$$

Ugh, this has all only been for positive, non-vanishing quantities!

Ugh, this has all only been for positive, non-vanishing quantities!

But maybe that's ok?

Ugh, this has all only been for positive, non-vanishing quantities!

But maybe that's ok?

To what extent can we get away with only using positive quantities? (Mass is ok, but what about position/velocity/momentum etc.?)

Ugh, this has all only been for positive, non-vanishing quantities!

But maybe that's ok?

To what extent can we get away with only using positive quantities? (Mass is ok, but what about position/velocity/momentum etc.?)

Can this structure be easily extended to accommodate negative and vanishing quantities?

Ugh, this has all only been for positive, non-vanishing quantities!

But maybe that's ok?

To what extent can we get away with only using positive quantities? (Mass is ok, but what about position/velocity/momentum etc.?)

Can this structure be easily extended to accommodate negative and vanishing quantities?

The vector space description encapsulates the multiplication of quantities, but it seems we still need to put in the addition rules "by hand".

Ugh, this has all only been for positive, non-vanishing quantities!

But maybe that's ok?

To what extent can we get away with only using positive quantities? (Mass is ok, but what about position/velocity/momentum etc.?)

Can this structure be easily extended to accommodate negative and vanishing quantities?

The vector space description encapsulates the multiplication of quantities, but it seems we still need to put in the addition rules "by hand".

The Obvious Proposition seems somewhat related to this.

Ugh, this has all only been for positive, non-vanishing quantities!

But maybe that's ok?

To what extent can we get away with only using positive quantities? (Mass is ok, but what about position/velocity/momentum etc.?)

Can this structure be easily extended to accommodate negative and vanishing quantities?

The vector space description encapsulates the multiplication of quantities, but it seems we still need to put in the addition rules "by hand".

The Obvious Proposition seems somewhat related to this.

Another issue: many quantities are best described not by elements of \mathbb{R} , but rather by elements of \mathbb{R} -torsors.



"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

Technical definition: a torsor (or *principle homogenous space*) for a group G with identity element e is a set X on which G acts such that

"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

Technical definition: a torsor (or *principle homogenous space*) for a group G with identity element e is a set X on which G acts such that

1)
$$\forall x \in X$$
, $ex = x$.

"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

Technical definition: a torsor (or principle homogenous space) for a group G with identity element e is a set X on which G acts such that

- 1) $\forall x \in X$, ex = x.
- 2) $\forall x, y \in X, \exists ! \ g \in G \text{ such that } gx = y.$

"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

Technical definition: a torsor (or principle homogenous space) for a group G with identity element e is a set X on which G acts such that

- 1) $\forall x \in X$, ex = x.
- 2) $\forall x, y \in X, \exists ! g \in G \text{ such that } gx = y.$

If X is a topological group, we say that X is a topological space and that the action is continuous.

"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

Technical definition: a torsor (or principle homogenous space) for a group G with identity element e is a set X on which G acts such that

- 1) $\forall x \in X$, ex = x.
- 2) $\forall x, y \in X, \exists ! \ g \in G \text{ such that } gx = y.$

If X is a topological group, we say that X is a topological space and that the action is continuous.

If X is a Lie group, we say that X is a smooth manifold and that the action is smooth.

Bonus: Torsors

"An affine space is like a vector space that forgot its origin; a torsor is like a group that forgot its identity".

Technical definition: a torsor (or principle homogenous space) for a group G with identity element e is a set X on which G acts such that

- 1) $\forall x \in X$, ex = x.
- 2) $\forall x, y \in X, \exists ! \ g \in G \text{ such that } gx = y.$

If X is a topological group, we say that X is a topological space and that the action is continuous.

If X is a Lie group, we say that X is a smooth manifold and that the action is smooth.

If X is a algebraic group, we say that X is a algebraic variety and that the action is regular.



Examples

Examples

Energies (\mathbb{R} -torsor)

Examples

Energies (\mathbb{R} -torsor)

Phases (U(1)-torsor)

Examples

Energies (\mathbb{R} -torsor) Phases (U(1)-torsor) Dates (\mathbb{Z} -torsor)

What's up with "Natural units" $(c = \hbar = 1)$??

What's up with "Natural units" $(c = \hbar = 1)$??

The units of c and \hbar don't span the space of mass, length, time.

What's up with "Natural units" $(c = \hbar = 1)$??

The units of c and \hbar don't span the space of mass, length, time.

Implicitly use base units in which c=1 [length/time], $\hbar=1$ [mass length 2 / time], and 1eV has the same value as in SI

What's up with "Natural units" $(c = \hbar = 1)$??

The units of c and \hbar don't span the space of mass, length, time.

Implicityly use base units in which c=1 [length/time], $\hbar=1$ [mass length 2 / time], and 1eV has the same value as in SI We can then treat dimensions not as \mathbb{Q}^n , but just "mass-energy dimension" $d\in\mathbb{Q}$.

What's up with "Natural units" $(c = \hbar = 1)$??

The units of c and \hbar don't span the space of mass, length, time.

Implicityly use base units in which c=1 [length/time], $\hbar=1$ [mass length² / time], and 1eV has the same value as in SI We can then treat dimensions not as \mathbb{Q}^n , but just "mass-energy dimension" $d\in\mathbb{Q}$.

How to express this with the vector space notation above?

The quantities $\Delta\nu_{Cs}$, c, \hbar , e, and k_B span the dimensional space {mass, length, time, temperature, charge}.

The quantities $\Delta \nu_{Cs}$, c, \hbar , e, and k_B span the dimensional space {mass, length, time, temperature, charge}.

Therefore, these quantities form a basis of this space.

The quantities $\Delta \nu_{Cs}$, c, \hbar , e, and k_B span the dimensional space {mass, length, time, temperature, charge}.

Therefore, these quantities form a basis of this space.

This is the foundation of the 2019 SI redefinition.

The quantities $\Delta\nu_{Cs}$, c, \hbar , e, and k_B span the dimensional space {mass, length, time, temperature, charge}.

Therefore, these quantities form a basis of this space.

This is the foundation of the 2019 SI redefinition.

All SI base units are defined as a fixed linear transformation of these quantities.