

# All About Spinors

...on *flat* spacetime

Sean Ericson

UO

Theory meeting, June 27, 2024



# Special and General Covariance

## Some Philosophical Motivation

# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$

# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$ 
  - ▶ Expresses invariance of physical laws under isometries

# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$ 
  - ▶ Expresses invariance of physical laws under isometries
  - ▶ Implies an action of the group of isometries on the space of physical states  $\tilde{\phi}_g : \mathcal{S} \rightarrow \mathcal{S}$

# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$ 
  - ▶ Expresses invariance of physical laws under isometries
  - ▶ Implies an action of the group of isometries on the space of physical states  $\tilde{\phi}_g : \mathcal{S} \rightarrow \mathcal{S}$
- *General covariance*: For a physical entity described by a tensor field  $T^{a\cdots}_{b\cdots}$ , the equations governing the field should be of the form  $\tilde{f}(T, \partial T, \cdots, g_{ab}, \partial g_{ab}, \cdots)$

# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$ 
  - ▶ Expresses invariance of physical laws under isometries
  - ▶ Implies an action of the group of isometries on the space of physical states  $\tilde{\phi}_g : \mathcal{S} \rightarrow \mathcal{S}$
- *General covariance*: For a physical entity described by a tensor field  $T^{a\cdots}_{b\cdots}$ , the equations governing the field should be of the form  $\tilde{f}(T, \partial T, \cdots, g_{ab}, \partial g_{ab}, \cdots)$ 
  - ▶ Expresses the invariance of physical laws under diffeomorphisms

# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$ 
  - ▶ Expresses invariance of physical laws under isometries
  - ▶ Implies an action of the group of isometries on the space of physical states  $\tilde{\phi}_g : \mathcal{S} \rightarrow \mathcal{S}$
- *General covariance*: For a physical entity described by a tensor field  $T^{a\cdots}_{b\cdots}$ , the equations governing the field should be of the form  $\tilde{f}(T, \partial T, \cdots, g_{ab}, \partial g_{ab}, \cdots)$ 
  - ▶ Expresses the invariance of physical laws under diffeomorphisms
  - ▶ G.C.  $\implies$  S.C.



# Special and General Covariance

## Some Philosophical Motivation

- *Special covariance*: For any two families of inertial observers  $O$  and  $O'$  related by an isometry, any set of physical measurements observable by  $O$  is observable by  $O'$ 
  - ▶ Expresses invariance of physical laws under isometries
  - ▶ Implies an action of the group of isometries on the space of physical states  $\tilde{\phi}_g : \mathcal{S} \rightarrow \mathcal{S}$
- *General covariance*: For a physical entity described by a tensor field  $T^{a\cdots}_{b\cdots}$ , the equations governing the field should be of the form  $\tilde{f}(T, \partial T, \cdots, g_{ab}, \partial g_{ab}, \cdots)$ 
  - ▶ Expresses the invariance of physical laws under diffeomorphisms
  - ▶ G.C.  $\implies$  S.C.
- Special/General covariance  $\rightarrow$  Special/General relativity

# Minkowski Space and Quantum Theory

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$

# Minkowski Space and Quantum Theory

■ Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$

►  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$
- Associate with  $\tilde{\phi}_g$  a map  $U_g : \mathcal{H} \rightarrow \mathcal{H}$  which preserves  $\sim$

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$
- Associate with  $\tilde{\phi}_g$  a map  $U_g : \mathcal{H} \rightarrow \mathcal{H}$  which preserves  $\sim$ 
  - ▶  $U_g$  can be rephased to be (anti-)unitary

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$
- Associate with  $\tilde{\phi}_g$  a map  $U_g : \mathcal{H} \rightarrow \mathcal{H}$  which preserves  $\sim$ 
  - ▶  $U_g$  can be rephased to be (anti-)unitary
- Composition:  $\tilde{\phi}_{g_1} \circ \tilde{\phi}_{g_2} = \tilde{\phi}_{g_1 g_2} \implies U_{g_1} U_{g_2} = e^{i\theta} U_{g_1 g_2}$



# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$
- Associate with  $\tilde{\phi}_g$  a map  $U_g : \mathcal{H} \rightarrow \mathcal{H}$  which preserves  $\sim$ 
  - ▶  $U_g$  can be rephased to be (anti-)unitary
- Composition:  $\tilde{\phi}_{g_1} \circ \tilde{\phi}_{g_2} = \tilde{\phi}_{g_1 g_2} \implies U_{g_1} U_{g_2} = e^{i\theta} U_{g_1 g_2}$ 
  - ▶ Wigner ('39): can set phases so  $\theta = n\pi$ , i.e.  $U_{g_1} U_{g_2} = \pm U_{g_1 g_2}$

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$
- Associate with  $\tilde{\phi}_g$  a map  $U_g : \mathcal{H} \rightarrow \mathcal{H}$  which preserves  $\sim$ 
  - ▶  $U_g$  can be rephased to be (anti-)unitary
- Composition:  $\tilde{\phi}_{g_1} \circ \tilde{\phi}_{g_2} = \tilde{\phi}_{g_1 g_2} \implies U_{g_1} U_{g_2} = e^{i\theta} U_{g_1 g_2}$ 
  - ▶ Wigner ('39): can set phases so  $\theta = n\pi$ , i.e.  $U_{g_1} U_{g_2} = \pm U_{g_1 g_2}$
- $\mathcal{H}$  a rep. space for a unitary rep. (up to sign) of  $ISO(3, 1)^+$ !

# Minkowski Space and Quantum Theory

- Consider  $(\mathbb{R}^4, \eta_{ab})$ ,  $\mathcal{S} = \{\psi \in \mathcal{H} : |\psi|^2 = 1\} / \sim$ 
  - ▶  $\psi \sim \psi' \iff \psi = e^{i\alpha} \psi'$
- Isometry group:  $G = ISO(3, 1)^+$
- Associate with  $\tilde{\phi}_g$  a map  $U_g : \mathcal{H} \rightarrow \mathcal{H}$  which preserves  $\sim$ 
  - ▶  $U_g$  can be rephased to be (anti-)unitary
- Composition:  $\tilde{\phi}_{g_1} \circ \tilde{\phi}_{g_2} = \tilde{\phi}_{g_1 g_2} \implies U_{g_1} U_{g_2} = e^{i\theta} U_{g_1 g_2}$ 
  - ▶ Wigner ('39): can set phases so  $\theta = n\pi$ , i.e.  $U_{g_1} U_{g_2} = \pm U_{g_1 g_2}$
- $\mathcal{H}$  a rep. space for a unitary rep. (up to sign) of  $ISO(3, 1)^+!$
- Bargmann ('54): reps up to sign are *exactly* the true reps of the universal cover

# Universal Covering Spaces

# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$

# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$
- Construction:

# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$
- Construction:
  - ▶ Cut  $M$  such that it becomes simply connected with boundary

# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$
- Construction:
  - ▶ Cut  $M$  such that it becomes simply connected with boundary
  - ▶ Glue together copies to eliminate the boundaries



# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$
- Construction:
  - ▶ Cut  $M$  such that it becomes simply connected with boundary
  - ▶ Glue together copies to eliminate the boundaries
- Lie group structure of  $M$  is naturally lifted to  $\mathcal{U}(M)$

# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$
- Construction:
  - ▶ Cut  $M$  such that it becomes simply connected with boundary
  - ▶ Glue together copies to eliminate the boundaries
- Lie group structure of  $M$  is naturally lifted to  $\mathcal{U}(M)$
- Fundamental group of  $ISO(3, 1)^+$  is  $\mathbb{Z}_2 \rightarrow$  double cover

# Universal Covering Spaces

- The universal cover  $\mathcal{U}(M)$  of a topological space  $M$  is a simply connected space which covers  $M$
- Construction:
  - ▶ Cut  $M$  such that it becomes simply connected with boundary
  - ▶ Glue together copies to eliminate the boundaries
- Lie group structure of  $M$  is naturally lifted to  $\mathcal{U}(M)$
- Fundamental group of  $ISO(3, 1)^+$  is  $\mathbb{Z}_2 \rightarrow$  double cover
- In fact,  $\mathcal{U}(ISO(3, 1)^+) \cong ISL(2, \mathbb{C})$

# Spinors, Spinorial Tensors, and Spinor Space

# Spinors, Spinorial Tensors, and Spinor Space

- Let  $W \cong \mathbb{C}^2$

# Spinors, Spinorial Tensors, and Spinor Space

- Let  $W \cong \mathbb{C}^2$ 
  - ▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$

# Spinors, Spinorial Tensors, and Spinor Space

- Let  $W \cong \mathbb{C}^2$ 
  - ▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$
  - ▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$

# Spinors, Spinorial Tensors, and Spinor Space

- Let  $W \cong \mathbb{C}^2$ 
  - ▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$
  - ▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$
  - ▶ *Conjugate space*  $\overline{W}$ : the dual space of  $\overline{W}^*$



# Spinors, Spinorial Tensors, and Spinor Space

■ Let  $W \cong \mathbb{C}^2$

▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$

▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$

▶ *Conjugate space*  $\overline{W}$ : the dual space of  $\overline{W}^*$

■ Tensor of type  $(k, l; k', l')$ :

$$T_{B_1 \dots B_l B'_1 \dots B'_{l'}}^{A_1 \dots A_k A'_1 \dots A'_{k'}} : (W^*)^k \times (W)^l \times (\overline{W}^*)^{k'} \times (\overline{W})^{l'} \rightarrow \mathbb{C}$$

# Spinors, Spinorial Tensors, and Spinor Space

■ Let  $W \cong \mathbb{C}^2$

▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$

▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$

▶ *Conjugate space*  $\overline{W}$ : the dual space of  $\overline{W}^*$

■ Tensor of type  $(k, l; k', l')$ :

$$T_{B_1 \dots B_l B'_1 \dots B'_{l'}}^{A_1 \dots A_k A'_1 \dots A'_{k'}} : (W^*)^k \times (W)^l \times (\overline{W}^*)^{k'} \times (\overline{W})^{l'} \rightarrow \mathbb{C}$$

■ Space of type  $(0,2;0,0)$  antisymmetric tensors is 1-dimensional

# Spinors, Spinorial Tensors, and Spinor Space

■ Let  $W \cong \mathbb{C}^2$

▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$

▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$

▶ *Conjugate space*  $\overline{W}$ : the dual space of  $\overline{W}^*$

■ Tensor of type  $(k, l; k', l')$ :

$$T_{B_1 \dots B_l B'_1 \dots B'_{l'}}^{A_1 \dots A_k A'_1 \dots A'_{k'}} : (W^*)^k \times (W)^l \times (\overline{W}^*)^{k'} \times (\overline{W})^{l'} \rightarrow \mathbb{C}$$

■ Space of type  $(0,2;0,0)$  antisymmetric tensors is 1-dimensional

■ Spinor space:  $(W, \epsilon_{AB})$

# Spinors, Spinorial Tensors, and Spinor Space

■ Let  $W \cong \mathbb{C}^2$

▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$

▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$

▶ *Conjugate space*  $\overline{W}$ : the dual space of  $\overline{W}^*$

■ Tensor of type  $(k, l; k', l')$ :

$$T_{B_1 \dots B_l B'_1 \dots B'_{l'}}^{A_1 \dots A_k A'_1 \dots A'_{k'}} : (W^*)^k \times (W)^l \times (\overline{W}^*)^{k'} \times (\overline{W})^{l'} \rightarrow \mathbb{C}$$

■ Space of type  $(0,2;0,0)$  antisymmetric tensors is 1-dimensional

■ Spinor space:  $(W, \epsilon_{AB})$

▶  $\lambda^A \in W$  is called a *spinor*

# Spinors, Spinorial Tensors, and Spinor Space

## ■ Let $W \cong \mathbb{C}^2$

- ▶ *Dual space*  $W^*$ : linear maps  $\lambda_A : W \rightarrow \mathbb{C}$
- ▶ *Conjugate dual space*  $\overline{W}^*$ : anti-linear maps  $\lambda_{A'} : W \rightarrow \mathbb{C}$
- ▶ *Conjugate space*  $\overline{W}$ : the dual space of  $\overline{W}^*$

## ■ Tensor of type $(k, l; k', l')$ :

$$T_{B_1 \dots B_l B'_1 \dots B'_{l'}}^{A_1 \dots A_k A'_1 \dots A'_{k'}} : (W^*)^k \times (W)^l \times (\overline{W}^*)^{k'} \times (\overline{W})^{l'} \rightarrow \mathbb{C}$$

- Space of type  $(0,2;0,0)$  antisymmetric tensors is 1-dimensional
- Spinor space:  $(W, \epsilon_{AB})$ 
  - ▶  $\lambda^A \in W$  is called a *spinor*
  - ▶ Tensors over  $W$  are called *spinorial tensors*

# Spinor Conventions

# Spinor Conventions

- (Un)primed index order irrelevant:  $T^{AD'B}{}_C \leftrightarrow T^{AB}{}_C{}^{D'}$

# Spinor Conventions

- (Un)primed index order irrelevant:  $T^{AD'B}{}_C \leftrightarrow T^{AB}{}_C{}^{D'}$
- Conjugation maps  $(k, l; k', l')$  tensors to  $(k', l'; k, l)$  tensors:

$$T^A{}_{BC} \leftrightarrow \overline{T}^{A'}{}_{B'C'}$$

- $\epsilon^{AB}, \epsilon_{AB}$  raise/lower unprimed indices;  $\bar{\epsilon}_{A'B'}$  for primed



# Spinor Conventions

- (Un)primed index order irrelevant:  $T^{AD'B}{}_C \leftrightarrow T^{AB}{}_C{}^{D'}$
- Conjugation maps  $(k, l; k', l')$  tensors to  $(k', l'; k, l)$  tensors:

$$T^A{}_{BC} \leftrightarrow \overline{T}^{A'}{}_{B'C'}$$

- $\epsilon^{AB}, \epsilon_{AB}$  raise/lower unprimed indices;  $\bar{\epsilon}_{A'B'}$  for primed
- Contraction occurs over *first* index of  $\epsilon$ :

$$\phi_A = \epsilon_{BA} \phi^B = -\epsilon_{AB} \phi^B$$

# Spinor Conventions

- (Un)primed index order irrelevant:  $T^{AD'B}{}_C \leftrightarrow T^{AB}{}_C{}^{D'}$
- Conjugation maps  $(k, l; k', l')$  tensors to  $(k', l'; k, l)$  tensors:

$$T^A{}_{BC} \leftrightarrow \overline{T}^{A'}{}_{B'C'}$$

- $\epsilon^{AB}, \epsilon_{AB}$  raise/lower unprimed indices;  $\bar{\epsilon}_{A'B'}$  for primed
- Contraction occurs over *first* index of  $\epsilon$ :

$$\phi_A = \epsilon_{BA}\phi^B = -\epsilon_{AB}\phi^B$$

►  $\implies \phi_A\phi^A = 0$

# Spinor Conventions

- (Un)primed index order irrelevant:  $T^{AD'B}{}_C \leftrightarrow T^{AB}{}_C{}^{D'}$
- Conjugation maps  $(k, l; k', l')$  tensors to  $(k', l'; k, l)$  tensors:

$$T^A{}_{BC} \leftrightarrow \overline{T}^{A'}{}_{B'C'}$$

- $\epsilon^{AB}, \epsilon_{AB}$  raise/lower unprimed indices;  $\bar{\epsilon}_{A'B'}$  for primed
- Contraction occurs over *first* index of  $\epsilon$ :

$$\phi_A = \epsilon_{BA}\phi^B = -\epsilon_{AB}\phi^B$$

►  $\implies \phi_A\phi^A = 0$

- $\delta^A{}_B = \mathbb{I}_W$  differs by a sign from  $\delta_C{}^D = \mathbb{I}_{W^*}$

# Spinor Conventions

- (Un)primed index order irrelevant:  $T^{AD'B}{}_C \leftrightarrow T^{AB}{}_C{}^{D'}$
- Conjugation maps  $(k, l; k', l')$  tensors to  $(k', l'; k, l)$  tensors:

$$T^A{}_{BC} \leftrightarrow \overline{T}^{A'}{}_{B'C'}$$

- $\epsilon^{AB}, \epsilon_{AB}$  raise/lower unprimed indices;  $\bar{\epsilon}_{A'B'}$  for primed
- Contraction occurs over *first* index of  $\epsilon$ :

$$\phi_A = \epsilon_{BA}\phi^B = -\epsilon_{AB}\phi^B$$

►  $\implies \phi_A\phi^A = 0$

- $\delta^A{}_B = \mathbb{I}_W$  differs by a sign from  $\delta_C{}^D = \mathbb{I}_{W^*}$

►  $\rightarrow$  use  $\epsilon^A{}_B, \epsilon_C{}^D$  and their conjugates to avoid confusion

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Let  $L^A_B : W \rightarrow W$  be a linear transformation

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Let  $L^A_B : W \rightarrow W$  be a linear transformation
  - ▶  $\det(L) := \frac{1}{2} \epsilon_{AB} \epsilon_{CD} L^A_C L^B_D$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Let  $L^A_B : W \rightarrow W$  be a linear transformation
  - ▶  $\det(L) := \frac{1}{2} \epsilon_{AB} \epsilon_{CD} L^A_C L^B_D$
- $SL(2, \mathbb{C})$  is simply all  $L$  with  $\det(L) = 1$



# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Let  $L^A_B : W \rightarrow W$  be a linear transformation
  - ▶  $\det(L) := \frac{1}{2} \epsilon_{AB} \epsilon_{CD} L^A_C L^B_D$
- $SL(2, \mathbb{C})$  is simply all  $L$  with  $\det(L) = 1$ 
  - ▶ Polar decomp:  $L = UH \rightarrow 6$  real d.o.f.

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Let  $L^A_B : W \rightarrow W$  be a linear transformation
  - ▶  $\det(L) := \frac{1}{2} \epsilon_{AB} \epsilon_{CD} L^A_C L^B_D$
- $SL(2, \mathbb{C})$  is simply all  $L$  with  $\det(L) = 1$ 
  - ▶ Polar decomp:  $L = UH \rightarrow$  6 real d.o.f.
  - ▶ Simply connected Lie group  $\cong S^3 \times \mathbb{R}^3$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Let  $L^A_B : W \rightarrow W$  be a linear transformation
  - ▶  $\det(L) := \frac{1}{2} \epsilon_{AB} \epsilon_{CD} L^A_C L^B_D$
- $SL(2, \mathbb{C})$  is simply all  $L$  with  $\det(L) = 1$ 
  - ▶ Polar decomp:  $L = UH \rightarrow$  6 real d.o.f.
  - ▶ Simply connected Lie group  $\cong S^3 \times \mathbb{R}^3$
  - ▶  $\det(L) = 1 \iff L^A_C L^B_D \epsilon_{AB} = \epsilon_{CD}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Tensors  $\phi^{AA'} \in W_{1,0;1,0}$  comprise a  $\mathbb{C}^4$  vector space

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Tensors  $\phi^{AA'} \in W_{1,0;1,0}$  comprise a  $\mathbb{C}^4$  vector space
  - ▶ Let  $\{o^A, \iota^A\}$  be a basis for  $W$  with  $o_A \iota^A = 1$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Tensors  $\phi^{AA'} \in W_{1,0;1,0}$  comprise a  $\mathbb{C}^4$  vector space
  - ▶ Let  $\{o^A, \iota^A\}$  be a basis for  $W$  with  $o_A \iota^A = 1$
  - ▶ A basis for  $W_{1,0;1,0}$  can be given by

$$\begin{aligned}t^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{o}^{A'} + \iota^A \bar{\iota}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\x^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{\iota}^{A'} + \iota^A \bar{o}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\y^{AA'} &= \frac{i}{\sqrt{2}} \left( o^A \bar{\iota}^{A'} - \iota^A \bar{o}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix} \\z^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{o}^{A'} - \iota^A \bar{\iota}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\end{aligned}$$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Tensors  $\phi^{AA'} \in W_{1,0;1,0}$  comprise a  $\mathbb{C}^4$  vector space
  - ▶ Let  $\{o^A, \iota^A\}$  be a basis for  $W$  with  $o_A \iota^A = 1$
  - ▶ A basis for  $W_{1,0;1,0}$  can be given by

$$\begin{aligned} t^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{o}^{A'} + \iota^A \bar{\iota}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\ x^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{\iota}^{A'} + \iota^A \bar{o}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ y^{AA'} &= \frac{i}{\sqrt{2}} \left( o^A \bar{\iota}^{A'} - \iota^A \bar{o}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix} \\ z^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{o}^{A'} - \iota^A \bar{\iota}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{aligned}$$

- Under conjugation,  $\overline{W}_{1,0;1,0} = W_{1,0;1,0}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- Tensors  $\phi^{AA'} \in W_{1,0;1,0}$  comprise a  $\mathbb{C}^4$  vector space

- ▶ Let  $\{o^A, \iota^A\}$  be a basis for  $W$  with  $o_A \iota^A = 1$

- ▶ A basis for  $W_{1,0;1,0}$  can be given by

$$\begin{aligned} t^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{o}^{A'} + \iota^A \bar{\iota}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\ x^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{\iota}^{A'} + \iota^A \bar{o}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ y^{AA'} &= \frac{i}{\sqrt{2}} \left( o^A \bar{\iota}^{A'} - \iota^A \bar{o}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix} \\ z^{AA'} &= \frac{1}{\sqrt{2}} \left( o^A \bar{o}^{A'} - \iota^A \bar{\iota}^{A'} \right) &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{aligned}$$

- Under conjugation,  $\overline{W}_{1,0;1,0} = W_{1,0;1,0}$

- ▶  $\phi^{AA'} \in W_{1,0;1,0}$  s.t.  $\overline{\phi}^{AA'} = \phi^{AA'}$  are called *real*



# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$ 
  - ▶ Automatically,  $\lambda^{AA'}_{CC'} \lambda^{BB'}_{DD'} g_{AA'BB'} = g_{CC'DD'}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$ 
  - ▶ Automatically,  $\lambda^{AA'}_{CC'} \lambda^{BB'}_{DD'} g_{AA'BB'} = g_{CC'DD'}$
  - ▶ But this means  $\lambda \in O(3, 1)$  (in fact,  $SO(3, 1)^+$ )!!

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$ 
  - ▶ Automatically,  $\lambda^{AA'}_{CC'} \lambda^{BB'}_{DD'} g_{AA'BB'} = g_{CC'DD'}$
  - ▶ But this means  $\lambda \in O(3, 1)$  (in fact,  $SO(3, 1)^+$ )!!
  - ▶  $L_1, L_2 \rightarrow \lambda \implies L_1 = \pm L_2$  (a double cover)



# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$ 
  - ▶ Automatically,  $\lambda^{AA'}_{CC'} \lambda^{BB'}_{DD'} g_{AA'BB'} = g_{CC'DD'}$
  - ▶ But this means  $\lambda \in O(3, 1)$  (in fact,  $SO(3, 1)^+$ )!!
  - ▶  $L_1, L_2 \rightarrow \lambda \implies L_1 = \pm L_2$  (a double cover)
- Let  $\{t^a, x^a, y^a, z^a\}$  be a basis for  $\mathbb{R}^{3,1}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$ 
  - ▶ Automatically,  $\lambda^{AA'}_{CC'} \lambda^{BB'}_{DD'} g_{AA'BB'} = g_{CC'DD'}$
  - ▶ But this means  $\lambda \in O(3, 1)$  (in fact,  $SO(3, 1)^+$ )!!
  - ▶  $L_1, L_2 \rightarrow \lambda \implies L_1 = \pm L_2$  (a double cover)
- Let  $\{t^a, x^a, y^a, z^a\}$  be a basis for  $\mathbb{R}^{3,1}$
- Define  $\sigma^a_{AA'} := t^a t_{AA'} - x^a x_{AA'} - y^a y_{AA'} - z^a z_{AA'}$

# $SL(2, \mathbb{C})$ and $SO(3, 1)^+$

- $\{t^{AA'}, x^{AA'}, y^{AA'}, z^{AA'}\}$  defined above are clearly real
- They span a 4-real dimensional space  $V \subset W_{1,0;1,0}$
- Define  $g : V \times V \rightarrow \mathbb{R}$  by  $g_{AA'BB'} := \epsilon_{AB}\bar{\epsilon}_{A'B'}$ 
  - ▶  $g$  is nondegenerate with signature  $(+, -, -, -)$ ; a Lorentz metric!
- Define  $\lambda : V \rightarrow V$  by  $\lambda^{AA'}_{BB'} := L^A_B \bar{L}^{A'}_{B'}$ 
  - ▶ Automatically,  $\lambda^{AA'}_{CC'} \lambda^{BB'}_{DD'} g_{AA'BB'} = g_{CC'DD'}$
  - ▶ But this means  $\lambda \in O(3, 1)$  (in fact,  $SO(3, 1)^+$ )!!
  - ▶  $L_1, L_2 \rightarrow \lambda \implies L_1 = \pm L_2$  (a double cover)
- Let  $\{t^a, x^a, y^a, z^a\}$  be a basis for  $\mathbb{R}^{3,1}$
- Define  $\sigma^a_{AA'} := t^a t_{AA'} - x^a x_{AA'} - y^a y_{AA'} - z^a z_{AA'}$ 
  - ▶  $\sigma$  is an isomorphism between  $\text{Re}[W_{1,0;1,0}]$  and  $\mathbb{R}^{3,1}$

# Spinors and Null Vectors

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$

# Spinors and Null Vectors

■ Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$

►  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$



# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation
- Def.  $\epsilon_{AA' BB' CC' DD'} := \epsilon_{AB} \epsilon_{CD} \bar{\epsilon}_{A' C'} \bar{\epsilon}_{B' D'} - \epsilon_{AC} \epsilon_{BD} \bar{\epsilon}_{A' B'} \bar{\epsilon}_{C' D'}$

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation
- Def.  $\epsilon_{AA' BB' CC' DD'} := \epsilon_{AB} \epsilon_{CD} \bar{\epsilon}_{A' C'} \bar{\epsilon}_{B' D'} - \epsilon_{AC} \epsilon_{BD} \bar{\epsilon}_{A' B'} \bar{\epsilon}_{C' D'}$ 
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural orientation

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation
- Def.  $\epsilon_{AA' BB' CC' DD'} := \epsilon_{AB} \epsilon_{CD} \bar{\epsilon}_{A' C'} \bar{\epsilon}_{B' D'} - \epsilon_{AC} \epsilon_{BD} \bar{\epsilon}_{A' B'} \bar{\epsilon}_{C' D'}$ 
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural orientation
- *Null flag*:  $F^{AA' BB'} := \psi^A \psi^B \bar{\epsilon}^{A' B'} + \bar{\psi}^{A'} \bar{\psi}^{B'} \epsilon^{AB}$

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation
- Def.  $\epsilon_{AA' BB' CC' DD'} := \epsilon_{AB} \epsilon_{CD} \bar{\epsilon}_{A' C'} \bar{\epsilon}_{B' D'} - \epsilon_{AC} \epsilon_{BD} \bar{\epsilon}_{A' B'} \bar{\epsilon}_{C' D'}$ 
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural orientation
- Null flag:  $F^{AA' BB'} := \psi^A \psi^B \bar{\epsilon}^{A' B'} + \bar{\psi}^{A'} \bar{\psi}^{B'} \epsilon^{AB}$ 
  - ▶  $F^{AA' BB'} = -F^{BB' AA'}$ ,  $F_{AA' BB'} F^{AA' BB'} = F_{AA' BB'} \psi^B \bar{\psi}^{B'} = 0$

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_{\psi}^{AA'}, k_{\phi}^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation
- Def.  $\epsilon_{AA' BB' CC' DD'} := \epsilon_{AB} \epsilon_{CD} \bar{\epsilon}_{A' C'} \bar{\epsilon}_{B' D'} - \epsilon_{AC} \epsilon_{BD} \bar{\epsilon}_{A' B'} \bar{\epsilon}_{C' D'}$ 
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural orientation
- Null flag:  $F^{AA' BB'} := \psi^A \psi^B \bar{\epsilon}^{A' B'} + \bar{\psi}^{A'} \bar{\psi}^{B'} \epsilon^{AB}$ 
  - ▶  $F^{AA' BB'} = -F^{BB' AA'}$ ,  $F_{AA' BB'} F^{AA' BB'} = F_{AA' BB'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $\implies F^{AA' BB'} = k^{AA'} m^{BB'} - k^{BB'} m^{AA'}$  for some  $m^{AA'}$

# Spinors and Null Vectors

- Let  $\psi^A \in W$ , then  $\psi^A \bar{\psi}^{A'} =: k^{AA'} \in \text{Re}[W_{1,0;1,0}]$ 
  - ▶  $k_{AA'} k^{AA'} = g_{AA' BB'} k^{AA'} k^{BB'} = \epsilon_{AB} \bar{\epsilon}_{A' B'} \psi^A \bar{\psi}^{A'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $k^{AA'}$  is thus a null vector ( $\psi^A$  “square root” of a null vector?)
- Let  $\psi^A, \phi^A \in W$ , then  $\psi_A \bar{\psi}_{A'} \phi^A \bar{\phi}^{A'} = |\psi_A \phi^A|^2 > 0$ 
  - ▶  $k_\psi^{AA'}, k_\phi^{AA'}$  on the same side of the light cone: future direction
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural time orientation
- Def.  $\epsilon_{AA' BB' CC' DD'} := \epsilon_{AB} \epsilon_{CD} \bar{\epsilon}_{A' C'} \bar{\epsilon}_{B' D'} - \epsilon_{AC} \epsilon_{BD} \bar{\epsilon}_{A' B'} \bar{\epsilon}_{C' D'}$ 
  - ▶  $\text{Re}[W_{1,0;1,0}]$  has a natural orientation
- Null flag:  $F^{AA' BB'} := \psi^A \psi^B \bar{\epsilon}^{A' B'} + \bar{\psi}^{A'} \bar{\psi}^{B'} \epsilon^{AB}$ 
  - ▶  $F^{AA' BB'} = -F^{BB' AA'}$ ,  $F_{AA' BB'} F^{AA' BB'} = F_{AA' BB'} \psi^B \bar{\psi}^{B'} = 0$
  - ▶  $\implies F^{AA' BB'} = k^{AA'} m^{BB'} - k^{BB'} m^{AA'}$  for some  $m^{AA'}$
  - ▶  $\psi, \psi' \rightarrow F \iff \psi = \pm \psi'$



# A Couple Neat Identities

# A Couple Neat Identities

- Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )

# A Couple Neat Identities

- Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )
- Antisymmetrization:

# A Couple Neat Identities

- Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )
- Antisymmetrization:
  - ▶  $T_{[ab]} = T_{(AB)[A'B']} + T_{[AB](A'B')} = \phi_{AB}\bar{\epsilon}_{A'B'} + \bar{\phi}_{A'B'}\epsilon_{AB}$   
with  $\phi_{AB} = \frac{1}{2}T_{(AB)A'}{}^{A'}$  symmetric

# A Couple Neat Identities

- Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )
- Antisymmetrization:
  - ▶  $T_{[ab]} = T_{(AB)[A'B']} + T_{[AB](A'B')} = \phi_{AB}\bar{\epsilon}_{A'B'} + \bar{\phi}_{A'B'}\epsilon_{AB}$   
with  $\phi_{AB} = \frac{1}{2} T_{(AB)A'}{}^{A'}$  symmetric
- Symmetrization:

# A Couple Neat Identities

■ Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )

■ Antisymmetrization:

►  $T_{[ab]} = T_{(AB)[A'B']} + T_{[AB](A'B')} = \phi_{AB}\bar{\epsilon}_{A'B'} + \bar{\phi}_{A'B'}\epsilon_{AB}$   
with  $\phi_{AB} = \frac{1}{2} T_{(AB)A'}{}^{A'}$  symmetric

■ Symmetrization:

►  $T_{(ab)} = T_{(AB)(A'B')} + T_{[AB][A'B']} = T_{(AB)(A'B')} + \frac{1}{4}\epsilon_{AB}\bar{\epsilon}_{A'B'} T$   
where  $T = T_A{}^A{}_A{}^A = T_a T^a$

# A Couple Neat Identities

■ Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )

■ Antisymmetrization:

►  $T_{[ab]} = T_{(AB)[A'B']} + T_{[AB](A'B')} = \phi_{AB}\bar{\epsilon}_{A'B'} + \bar{\phi}_{A'B'}\epsilon_{AB}$   
with  $\phi_{AB} = \frac{1}{2}T_{(AB)A'}{}^{A'}$  symmetric

■ Symmetrization:

►  $T_{(ab)} = T_{(AB)(A'B')} + T_{[AB][A'B']} = T_{(AB)(A'B')} + \frac{1}{4}\epsilon_{AB}\bar{\epsilon}_{A'B'}T$   
where  $T = T_A{}^A{}_A{}^A = T_a T^a$

► Note also that  $T_{[AB]A'}{}^{A'} = \frac{1}{2}\epsilon_{AB}T$

# A Couple Neat Identities

- Let  $T_{ab}$  ( $T_{AA'BB'}$ ) be a tensor on  $\mathbb{R}^{3,1}$  ( $\text{Re}[W_{0,1;0,1}]$ )
- Antisymmetrization:
  - ▶  $T_{[ab]} = T_{(AB)[A'B']} + T_{[AB](A'B')} = \phi_{AB}\bar{\epsilon}_{A'B'} + \bar{\phi}_{A'B'}\epsilon_{AB}$   
with  $\phi_{AB} = \frac{1}{2}T_{(AB)A'}{}^{A'}$  symmetric
- Symmetrization:
  - ▶  $T_{(ab)} = T_{(AB)(A'B')} + T_{[AB][A'B']} = T_{(AB)(A'B')} + \frac{1}{4}\epsilon_{AB}\bar{\epsilon}_{A'B'}T$   
where  $T = T_A{}^A{}_A{}^A = T_a{}^a$
  - ▶ Note also that  $T_{[AB]A'}{}^{A'} = \frac{1}{2}\epsilon_{AB}T$
- $\partial_{AA'}\partial_B{}^{A'} = \frac{1}{2}\epsilon_{AB}\square$ , where  $\square = \partial_{AA'}\partial^{AA'}$



# The Universal Enveloping Algebra

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$
  - ▶ Fully lifted,  $[\cdot, \cdot]$  obeys Leibniz's law:  $T(\mathfrak{g})$  is a *poisson algebra*

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$
  - ▶ Fully lifted,  $[\cdot, \cdot]$  obeys Leibniz's law:  $T(\mathfrak{g})$  is a *poisson* algebra
  - ▶  $\mathcal{U}(\mathfrak{g})$  is what remains after “modding out” the poisson structure, i.e.  $\mathcal{U}(\mathfrak{g}) = T(\mathfrak{g}) / \sim$  where the equivalence relation is given by  $[a, b] = a \otimes b - b \otimes a$

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$
  - ▶ Fully lifted,  $[\cdot, \cdot]$  obeys Leibniz's law:  $T(\mathfrak{g})$  is a *poisson* algebra
  - ▶  $\mathcal{U}(\mathfrak{g})$  is what remains after “modding out” the poisson structure, i.e.  $\mathcal{U}(\mathfrak{g}) = T(\mathfrak{g}) / \sim$  where the equivalence relation is given by  $[a, b] = a \otimes b - b \otimes a$
- Casimir elements



# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$
  - ▶ Fully lifted,  $[\cdot, \cdot]$  obeys Leibniz's law:  $T(\mathfrak{g})$  is a *poisson algebra*
  - ▶  $\mathcal{U}(\mathfrak{g})$  is what remains after “modding out” the poisson structure, i.e.  $\mathcal{U}(\mathfrak{g}) = T(\mathfrak{g}) / \sim$  where the equivalence relation is given by  $[a, b] = a \otimes b - b \otimes a$
- Casimir elements
  - ▶ *Center*  $Z(\mathcal{U}(\mathfrak{g}))$ : all elements that commute with all of  $\mathcal{U}(\mathfrak{g})$

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$
  - ▶ Fully lifted,  $[\cdot, \cdot]$  obeys Leibniz's law:  $T(\mathfrak{g})$  is a *poisson algebra*
  - ▶  $\mathcal{U}(\mathfrak{g})$  is what remains after “modding out” the poisson structure, i.e.  $\mathcal{U}(\mathfrak{g}) = T(\mathfrak{g}) / \sim$  where the equivalence relation is given by  $[a, b] = a \otimes b - b \otimes a$
- Casimir elements
  - ▶ *Center*  $Z(\mathcal{U}(\mathfrak{g}))$ : all elements that commute with all of  $\mathcal{U}(\mathfrak{g})$
  - ▶ *Casimir elements* form a basis of  $Z(\mathcal{U}(\mathfrak{g}))$

# The Universal Enveloping Algebra

- Given a Lie algebra  $\mathfrak{g}$ , the *universal enveloping algebra*  $\mathcal{U}(\mathfrak{g})$  is the unique unital associative algebra whose representations correspond exactly to the representations of  $\mathfrak{g}$
- Construction
  - ▶ Form the *Tensor algebra*  $T(\mathfrak{g}) := \mathbb{C} \oplus \mathfrak{g} \oplus (\mathfrak{g} \otimes \mathfrak{g}) \cdots$
  - ▶ Recursively lift  $[\cdot, \cdot]$  from  $\mathfrak{g}$  to  $T(\mathfrak{g})$
  - ▶ Fully lifted,  $[\cdot, \cdot]$  obeys Leibniz's law:  $T(\mathfrak{g})$  is a *poisson algebra*
  - ▶  $\mathcal{U}(\mathfrak{g})$  is what remains after “modding out” the poisson structure, i.e.  $\mathcal{U}(\mathfrak{g}) = T(\mathfrak{g}) / \sim$  where the equivalence relation is given by  $[a, b] = a \otimes b - b \otimes a$
- Casimir elements
  - ▶ *Center*  $Z(\mathcal{U}(\mathfrak{g}))$ : all elements that commute with all of  $\mathcal{U}(\mathfrak{g})$
  - ▶ *Casimir elements* form a basis of  $Z(\mathcal{U}(\mathfrak{g}))$
  - ▶ Casimir representatives proportional to identity

# Representations and Hilbert Space

# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s + 1)$ , respectively.

# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s + 1)$ , respectively.
- Cases:

# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s+1)$ , respectively.
- Cases:
  1.  $m^2 > 0$

# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s+1)$ , respectively.
- Cases:
  1.  $m^2 > 0$
  2.  $m^2 = 0$  with nontrivial translations



# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s+1)$ , respectively.
- Cases:
  1.  $m^2 > 0$
  2.  $m^2 = 0$  with nontrivial translations
    - 2.1 Helicity parameterization:  $s = 0, \pm\frac{1}{2}, \pm 1, \pm\frac{3}{2} \dots$
  3.  $m^2 = 0$ ; trivial translations

# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s+1)$ , respectively.
- Cases:
  1.  $m^2 > 0$
  2.  $m^2 = 0$  with nontrivial translations
    - 2.1 Helicity parameterization:  $s = 0, \pm\frac{1}{2}, \pm 1, \pm\frac{3}{2} \dots$
    - 2.2 "Continuous spin"
  3.  $m^2 = 0$ ; trivial translations
  4.  $m^2 < 0$  (tachyons!)

# Representations and Hilbert Space

- $\mathcal{U}(ISL(2, \mathbb{C}))$  has two independent Casimir elements:  $P^2$  and  $S^2$  with eigenvalues  $m^2$  and  $s(s+1)$ , respectively.
- Cases:
  1.  $m^2 > 0$
  2.  $m^2 = 0$  with nontrivial translations
    - 2.1 Helicity parameterization:  $s = 0, \pm\frac{1}{2}, \pm 1, \pm\frac{3}{2} \dots$
    - 2.2 "Continuous spin"
  3.  $m^2 = 0$ ; trivial translations
  4.  $m^2 < 0$  (tachyons!)
- Physically relevant cases are 1 and 2.1

# Representations and Hilbert Space

- Specify representatives for class 1 via

$()$

# Summary

# Summary

