Section 2.2, Equation 2.18 of thesis

$$rac{1}{H'-H_{e}\hat{n}_{N}}c_{N}^{\dagger}T=c_{N}^{\dagger}Trac{1}{H'-H_{h}(1-\hat{n}_{N})} \ \Longrightarrow \ H_{e}\hat{n}_{N}c_{N}^{\dagger}T=c_{N}^{\dagger}TH_{h}(1-\hat{n}_{N})$$

This seems to **require** H' **commuting with** T, because

$$c_N^{\dagger}TH'-c_N^{\dagger}TH_h(1-\hat{n}_N)=H'c_N^{\dagger}T-H_e\hat{n}_Nc_N^{\dagger}T$$

Why should H' commute with T?

(where 
$$H_e = Tr(H\hat{n}_N)$$
,  $H_h = Tr[H(1 - \hat{n}_N)]$  and  $T = Tr(Hc_N)$ )

Section 2.2, Equation 2.19 of thesis

$$\eta_N H \eta_N^{\dagger} = H_h (1 - n_N)$$

If I try to derive this using the result on the previous slide:

$$egin{aligned} \eta H \eta^\dagger &= \eta H_e \eta^\dagger = \eta H_e c^\dagger T G = \eta c^\dagger T H_h G \ &= \eta c^\dagger T G H_h = \eta \eta^\dagger H_h = H_h (1-\hat{n}) \end{aligned}$$

That required  $[G, H_h] = 0$ . How does that work out?

(where 
$$H_e = Tr(H\hat{n}_N)$$
,  $H_h = Tr[H(1 - \hat{n}_N)]$  and  $T = Tr(Hc_N)$ )

# Kondo Model appendix, Equation 9.61 of thesis

$$\begin{split} &\Delta \hat{H}_{(j)} = \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_j} \frac{(J^{(j)})^2}{2} \frac{\tau_{j,\hat{\mathbf{s}}_m,\beta}}{(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta}-\epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta}-J^{(j)}S^zs_{j,\hat{\mathbf{s}}_m}^z)} \\ &\times \left[ S^a S^b \sigma_{\alpha\beta}^a \sigma_{\beta\gamma}^b \sum_{\substack{(j_1,j_2< j),\\n,o}} c_{j_1,\hat{\mathbf{s}}_n,\alpha}^\dagger c_{j_2,\hat{\mathbf{s}}_o,\gamma} (1-\hat{n}_{j,\hat{\mathbf{s}}_m,\beta}) + ... \right. \\ &+ \sum_{m=1,}^{n_j} \frac{(J^{(j)})^2}{2(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta}-\epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta}-J^{(j)}S^zs_{j,\hat{\mathbf{s}}_m}^z)} \left[ S^x S^y \sigma_{\alpha\beta}^x \sigma_{\beta\alpha}^y c_{j,\hat{\mathbf{s}}_m,\alpha}^\dagger c_{j,\hat{\mathbf{s}}_m,\beta} c_{j,\hat{\mathbf{s}}_m,\beta} c_{j,\hat{\mathbf{s}}_m,\alpha} + ... \right. \end{split}$$

$$\begin{split} &\Delta \hat{H}_{(j)} = \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_j} \frac{(J^{(j)})^2}{2} \frac{\tau_{j,\hat{\mathbf{s}}_m,\beta}}{(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta}-\epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta}-J^{(j)}S^zs_{j,\hat{\mathbf{s}}_m}^z)} \\ &\times \left[ S^a S^b \sigma^a_{\alpha\beta} \sigma^b_{\beta\gamma} \sum_{\substack{(j_1,j_2 < j),\\n,o}} c^\dagger_{j_1,\hat{\mathbf{s}}_n,\alpha} c_{j_2,\hat{\mathbf{s}}_o,\gamma} (1-\hat{n}_{j,\hat{\mathbf{s}}_m,\beta}) + ... \right. \\ &+ \sum_{m=1,}^{n_j} \frac{(J^{(j)})^2}{2(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta}-\epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta}-J^{(j)}S^zs_{j,\hat{\mathbf{s}}_m}^z)} \left[ S^x S^y \sigma^x_{\alpha\beta} \sigma^y_{\beta\alpha} c^\dagger_{j,\hat{\mathbf{s}}_m,\alpha} c_{j,\hat{\mathbf{s}}_m,\beta} c_{j,\hat{\mathbf{s}}_m,\alpha} + ... \right. \end{split}$$

▶ The  $\tau$  should not be there in numerator i presume?

$$\begin{split} & \Delta \hat{H}_{(j)} = \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_{j}} \frac{(J^{(j)})^{2}}{2} \frac{\tau_{j,\hat{\mathbf{s}}_{m},\beta}}{(2\omega\tau_{j,\hat{\mathbf{s}}_{m},\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_{m},\beta} - J^{(j)}S^{z}S^{z}_{j,\hat{\mathbf{s}}_{m}})} \\ & \times \left[ S^{a}S^{b}\sigma^{a}_{\alpha\beta}\sigma^{b}_{\beta\gamma} \sum_{(j_{1},j_{2}< j),} c^{\dagger}_{j_{1},\hat{\mathbf{s}}_{n},\alpha}c_{j_{2},\hat{\mathbf{s}}_{o},\gamma}(1-\hat{n}_{j,\hat{\mathbf{s}}_{m},\beta}) + \dots \right. \\ & + \sum_{m=1,}^{n_{j}} \frac{(J^{(j)})^{2}}{2(2\omega\tau_{j,\hat{\mathbf{s}}_{m},\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_{m},\beta} - J^{(j)}S^{z}S^{z}_{j,\hat{\mathbf{s}}_{m}})} \left[ S^{x}S^{y}\sigma^{x}_{\alpha\beta}\sigma^{y}_{\beta\alpha}c^{\dagger}_{j,\hat{\mathbf{s}}_{m},\alpha}c_{j,\hat{\mathbf{s}}_{m},\beta}c^{\dagger}_{j,\hat{\mathbf{s}}_{m},\beta}c_{j,\hat{\mathbf{s}}_{m},\alpha} + \dots \right] \end{split}$$

Since coupling is  $\frac{J}{2}$ , shouldn't the thing be  $\frac{J^2}{4}$  instead of  $\frac{J^2}{2}$ ?

$$\Delta \hat{H}_{(j)} = \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_j} \frac{(J^{(j)})^2}{2} \frac{\tau_{j,\hat{\mathbf{s}}_m,\beta}}{(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta} - J^{(j)}S^z s_{j,\hat{\mathbf{s}}_m}^z)}$$

$$\times \left[ S^{a}S^{b}\sigma_{\alpha\beta}^{a}\sigma_{\beta\gamma}^{b} \sum_{\substack{(j_{1},j_{2}

$$+ \sum_{m=1,}^{n_{j}} \frac{(J^{(j)})^{2}}{2(2\omega\tau_{j,\hat{s}_{m},\beta}-\epsilon_{j,l}\tau_{j,\hat{s}_{m},\beta}-J^{(j)}S^{z}s_{j,\hat{s}_{m}}^{z})} \left[ S^{x}S^{y}\sigma_{\alpha\beta}^{x}\sigma_{\beta\alpha}^{y}c_{j,\hat{s}_{m},\alpha}^{\dagger}c_{j,\hat{s}_{m},\beta}c_{j,\hat{s}_{m},\beta}c_{j,\hat{s}_{m},\alpha} + ... \right.$$$$

➤ You mentioned the following in the google document- "interchange sigma\_a and sigma\_b (you get -1 sign)". But these are matrix elements (numbers). So why the minus sign?

$$\Delta \hat{H}_{(j)} = \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_j} \frac{(J^{(j)})^2}{2} \frac{\tau_{j,\hat{\mathbf{s}}_m,\beta}}{(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta} - J^{(j)}S^zs_{j,\hat{\mathbf{s}}_m}^z)}$$

$$imes \left[ S^a S^b \sigma^a_{lphaeta} \sigma^b_{eta\gamma} \sum_{\substack{(j_1,j_2 < j), \ n,o}} c^\dagger_{j_1,\hat{\mathbf{s}}_n,lpha} c_{j_2,\hat{\mathbf{s}}_o,\gamma} (1-\hat{n}_{j,\hat{\mathbf{s}}_m,eta}) + ... 
ight.$$

$$+\sum_{\substack{m=1,\\\beta=\pm/1}}^{\eta}\frac{(J^{(j)})^2}{2(2\omega\tau_{j,\hat{\mathbf{s}}_{m},\beta}-\epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_{m},\beta}-J^{(j)}S^zs_{j,\hat{\mathbf{s}}_{m}}^z)}\bigg[S^xS^y\sigma_{\alpha\beta}^x\sigma_{\beta\alpha}^yc_{j,\hat{\mathbf{s}}_{m},\alpha}^\dagger c_{j,\hat{\mathbf{s}}_{m},\beta}c_{j,\hat{\mathbf{s}}_{m},\beta}^\dagger c_{j,\hat{\mathbf{s}}_{m},\alpha}+...$$

► How do you combine the product of two sigmas ( $\sigma_{\alpha\beta}^a \sigma_{\beta\gamma}^b$ ) into a single  $\sigma_{\alpha\gamma}^c$ ?

Kondo URG coupling equation for J (equation 9.65):

$$\Delta J^{(j)} = n_j (J^{(j)})^2 \left[ \omega - \frac{\epsilon_{j,l}}{2} \right] \left[ (\frac{\epsilon_{j,l}}{2} - \omega)^2 - \frac{\left(J^{(j)}\right)^2}{16} \right]^{-1}$$

One-loop form (after setting  $\omega = \epsilon_{j,l}$ ):

$$\Delta J^{(j)} = \frac{n_j(J^{(j)})^2}{\omega - \frac{\epsilon_{j,l}}{2}} = 2\frac{n_j(J^{(j)})^2}{\epsilon_{j,l}} \rightarrow \frac{2\rho|\Delta D|J^2}{D} \quad [n_j, \rho \rightarrow \mathsf{DOS} \; \mathsf{per} \; \mathsf{spin}]$$

One-loop form in Coleman (Introduction to Many-Body Physics) ( $\tilde{J} = J/2$ ):

$$\Delta \tilde{J} = \frac{2\rho |\Delta D| \tilde{J}^2}{D} \implies \Delta J = \frac{\rho |\Delta D| J^2}{D}$$

Is there any reason for this difference?

▶ In the Kondo URG, are you considering two electrons on the shell  $\Lambda_N$ , one that we are decoupling  $(q\beta)$  and another with the same momentum but opposite spin  $(q\overline{\beta})$ ?

▶ If so, why does that kinetic energy piece  $(\epsilon_q \tau_{q\overline{\beta}})$  not come down in the denominator?

► Is that what gives rise to the second RG equation and hence the S<sup>z</sup>s<sup>z</sup> term in the effective Hamiltonian?

$$\begin{split} \Delta H_{(j)}^2 &= \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_j} \frac{(J^{(j)})^2}{(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta} - J^{(j)}S^z s_{j,\hat{\mathbf{s}}_m}^z)} \bigg[ S^x S^y \sigma_{\alpha\beta}^x \sigma_{\beta\alpha}^y c_{j,\hat{\mathbf{s}}_m,\alpha}^\dagger c_{j,\hat{\mathbf{s}}_m,\beta} c_{j,\hat{\mathbf{s}}_m,\beta} c_{j,\hat{\mathbf{s}}_m,\alpha} \\ &+ S^y S^x \sigma_{\alpha\beta}^x \sigma_{\beta\alpha}^y c_{j,\hat{\mathbf{s}}_m,\beta}^\dagger c_{j,\hat{\mathbf{s}}_m,\alpha} c_{j,\hat{\mathbf{s}}_m,\alpha} c_{j,\hat{\mathbf{s}}_m,\alpha} c_{j,\hat{\mathbf{s}}_m,\beta} \bigg] \end{split}$$

 $=\sum_{m=1,}^{n_j}\frac{(J^{(j)})^2}{(2\omega\tau_{j,\hat{\mathbf{s}}_m,\beta}-\epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_m,\beta}-J^{(j)}S^zs_{j,\hat{\mathbf{s}}_m}^z)}S^z\frac{\sigma_{\alpha\alpha}^z}{2}\bigg[\hat{\mathbf{n}}_{j,\hat{\mathbf{s}}_m,\alpha}(1-\hat{\mathbf{n}}_{j,\hat{\mathbf{s}}_m,\beta})-$ 

$$\hat{n}_{j,\hat{\mathbf{s}}_m,\beta}(1-\hat{n}_{j,\hat{\mathbf{s}}_m,lpha})$$

$$\begin{split} \Delta H_{(j)}^{2} &= \sum_{\substack{m=1,\\\beta=\uparrow/\downarrow}}^{n_{j}} \frac{(J^{(j)})^{2}}{(2\omega\tau_{j,\hat{\mathbf{s}}_{m},\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_{m},\beta} - J^{(j)}S^{z}S_{j,\hat{\mathbf{s}}_{m}}^{z})} \bigg[ S^{x}S^{y}\sigma_{\alpha\beta}^{x}\sigma_{\beta\alpha}^{y}C_{j,\hat{\mathbf{s}}_{m},\alpha}^{\dagger}C_{j,\hat{\mathbf{s}}_{m},\beta}C_{j,\hat{\mathbf{s}}_{m},\alpha}^{\dagger}C_{j,\hat{\mathbf{s}}_{m},\beta}C_{j,\hat{\mathbf{s}}_{m},\alpha}C_{j,\hat{\mathbf{s}}_{m},\alpha}C_{j,\hat{\mathbf{s}}_{m},\beta} \bigg] \\ &+ S^{y}S^{x}\sigma_{\alpha\beta}^{x}\sigma_{\beta\alpha}^{y}C_{j,\hat{\mathbf{s}}_{m},\beta}^{\dagger}C_{j,\hat{\mathbf{s}}_{m},\alpha}C_{j,\hat{\mathbf{s}}_{m},\alpha}C_{j,\hat{\mathbf{s}}_{m},\alpha}C_{j,\hat{\mathbf{s}}_{m},\beta} \bigg] \\ &= \sum_{m=1,}^{n_{j}} \frac{(J^{(j)})^{2}}{(2\omega\tau_{j,\hat{\mathbf{s}}_{m},\beta} - \epsilon_{j,l}\tau_{j,\hat{\mathbf{s}}_{m},\beta} - J^{(j)}S^{z}S_{j,\hat{\mathbf{s}}_{m}}^{z})} S^{z}\frac{\sigma_{\alpha\alpha}^{z}}{2} \bigg[ \hat{n}_{j,\hat{\mathbf{s}}_{m},\alpha}(1 - \hat{n}_{j,\hat{\mathbf{s}}_{m},\beta}) - \dots \bigg] \end{split}$$

# What I got:

$$S^{x}S^{y}\sigma_{\alpha\beta}^{x}\sigma_{\beta\alpha}^{y}c_{j,\hat{\mathbf{s}}_{m},\alpha}^{\dagger}c_{j,\hat{\mathbf{s}}_{m},\beta}c_{j,\hat{\mathbf{s}}_{m},\beta}^{\dagger}c_{j,\hat{\mathbf{s}}_{m},\alpha}=\underline{i^{2}}S^{z}\sigma_{\alpha\alpha}^{z}\hat{\eta}_{j,\hat{\mathbf{s}}_{m},\alpha}\left(1-\hat{\eta}_{j,\hat{\mathbf{s}}_{m},\beta}\right)$$

In eq. 2.21 of thesis,

$$UHU^{\dagger} = \frac{1}{2}Tr(H) + \tau Tr(H\tau) + \tau \{c^{\dagger}T, \eta\}$$

so the renormalization is

$$au\{m{c}^\daggerm{T},\eta\} = rac{1}{2} \left[ egin{array}{cccc} ext{particle sector} \ m{c}^\daggerm{T}\eta & - rac{\etam{c}^\daggerm{T}}{ ext{hole sector}} 
ight] = ext{difference of the 2 sectors}$$

Yet in most RG equations ( $\triangle H_F$  of 2d Hubbard,  $\triangle H_j$  of Kondo), you have added the two sectors. How/Why?

In the Kondo URG, you simplify the  $\hat{\omega}$  as

$$\hat{\omega} = \omega \tau$$

What is the formal way of doing this? Shouldn't it be

$$\hat{\omega} = \omega_1 \hat{\boldsymbol{n}} + \omega_1 (1 - \hat{\boldsymbol{n}})$$

Is this just an assumption?

In the RG equation for BCS instability (eq. 8.130 of thesis), you use

$$G^{-1} = \omega - \epsilon_1 \tau_1 - \epsilon_2 \tau_2$$

How is this choice of  $\hat{\omega}$  consistent with what was done in Kondo URG?