

# 1 URG as a double-bracket flow

The difference RG equation for URG can be written in the form

$$\Delta\mathcal{H}(\omega, D) = \frac{1}{\omega_1 - \omega_0} [G[\mathcal{H}^d, \mathcal{H}^I], \mathcal{H}] \quad (1)$$

This is not in the standard double-bracket form, primarily because it takes into account the off-diagonal terms in the Hamiltonian inside the generator of the unitary transformation. It can be given a double-bracket form by taking some approximations, as was shown in eq. ??.

Just like the standard double-bracket flow equation, the URG equation acts as an optimizer - it minimizes the function

$$\chi_j = \text{Tr} \left[ (\mathcal{H}_j^I)^2 \right] \quad (2)$$

The definition of this function first requires a scheme to be defined. We can order the energy of the electrons as  $\epsilon_1 < \epsilon_2 < \dots < \epsilon_j < \dots < \epsilon_N$ . The URG consists of sequentially decoupling the states  $\epsilon_N$ , then  $\epsilon_{N-1}$ , and so on. At the  $j^{\text{th}}$  step, the Hamiltonian can be partitioned in the subspace of the electron being decoupled; the partitioning looks like

$$\mathcal{H}_j^0 + c_j^\dagger T_j + T_j^\dagger c_j \quad (3)$$

$\mathcal{H}_j^0$  is the part that *does not* scatter between  $|\hat{n}_j\rangle = 0, 1$ , while  $\mathcal{H}_j^I = c_j^\dagger T_j + T_j^\dagger c_j$  is the part that *does* scatter between states with a definite value of  $\hat{n}_j$ .

The first observation that we make is that  $\chi_j$  is semi-positive definite. This is because it can be expressed as the norm-squared of a state vector.

$$\chi_j = \sum_{i=1}^N \langle \psi_i | (\mathcal{H}_j^I)^2 | \psi_i \rangle = \sum_{i=1}^N \langle \phi_i | \phi_i \rangle \geq 0, \text{ [where } |\phi_i\rangle = \mathcal{H}_j^I |\psi_i\rangle] \quad (4)$$

The difference equation for  $\chi_j$  is

$$\Delta\chi_j = 2\text{Tr} [\mathcal{H}_j^I \Delta\mathcal{H}_j^I] = 2\text{Tr} [\mathcal{H}_j^I (\mathcal{H}_{j-1}^I - \mathcal{H}_j^I)] \quad (5)$$

The first part of the trace is zero. To see why, note that from the nature of URG, once  $j$  has been decoupled, it is diagonal in all the subsequent Hamiltonians. Hence,  $\mathcal{H}_{j-1}^I$  will be diagonal in  $j$ , while  $\mathcal{H}_j^I$  is, by definition, off-diagonal in  $j$ . The product  $\mathcal{H}_j^I \mathcal{H}_{j-1}^I$  will hence be off-diagonal and will change  $\hat{n}_j$ . Hence, it will vanish when taken inside a trace. What remains is

$$\Delta\chi_j = -2\text{Tr} \left[ (\mathcal{H}_j^I)^2 \right] = -2\chi_j \leq 0 \quad (6)$$

At the fixed point  $j^*$  of URG, the off-diagonal part of the Hamiltonian vanishes, so we can write  $\mathcal{H}_{j^*}^I = 0 \implies \Delta\chi^* = 0$ . Combining the three points:

$$\chi_j \geq 0, \quad \Delta\chi_j \leq 0, \quad \Delta\chi_{j^*} = 0 \quad (7)$$

we can say that URG starts from a non-minimal value of  $\chi$  and flows to its minimum  $\chi^* = 0$  at the fixed point.