

Using the commutator of H_0^* with s^+ to bring H_0^* to the left, and using $s^+s^- = s^z + \frac{1}{2} = 1$, we get

$$\frac{J^2}{4\left(E_g + \frac{J}{4}\right)} \left[\frac{H_0^*}{E_g + \frac{J}{4}} + \left(\frac{H_0^*}{E_g + \frac{J}{4}} \right)^2 - \sum_{kk'qq'} \left\{ \frac{\epsilon_k - \epsilon_{k'}}{E_g + \frac{J}{4}} + \left(\frac{\epsilon_k - \epsilon_{k'}}{E_g + \frac{J}{4}} \right)^2 \right\} c_{k\uparrow}^\dagger c_{k'\downarrow} c_{q\downarrow}^\dagger c_{q'\uparrow} \right] \quad (9.135)$$

The full effective Hamiltonian, for $K = 0$, up to quartic interactions, is

$$H_0^* + \frac{J}{4} \left(\frac{J}{E_g + \frac{J}{4}} - 1 \right) - \frac{2V^2}{E_g} - \frac{2V^2}{E_g} \left[\frac{H_0^*}{E_g} + \left(\frac{H_0^*}{E_g} \right)^2 \right] + \frac{J^2}{4\left(E_g + \frac{J}{4}\right)} \left[\frac{H_0^*}{E_g + \frac{J}{4}} + \left(\frac{H_0^*}{E_g + \frac{J}{4}} \right)^2 \right] + \sum_{kk'qq'} F_{kk'qq'} c_{k\uparrow}^\dagger c_{k'\downarrow} c_{q\downarrow}^\dagger c_{q'\uparrow} \quad (9.136)$$

The coefficient $F_{kk'qq'}$ is

$$F_{kk'qq'} = \frac{V^2}{E_g N^* (E_g + \frac{J}{4})} \left[\frac{V^2}{E_g N^*} (\xi_{k'} + 2 - \xi_k) (\xi_q + \xi_{q'}) + \frac{J}{2} (\xi_{k'} + 2 - \xi_k + \xi_q + \xi_{q'}) \right] - \frac{J^2}{4\left(E_g + \frac{J}{4}\right)} \left[\frac{\epsilon_k - \epsilon_{k'}}{E_g + \frac{J}{4}} + \left(\frac{\epsilon_k - \epsilon_{k'}}{E_g + \frac{J}{4}} \right)^2 \right] \quad (9.137)$$

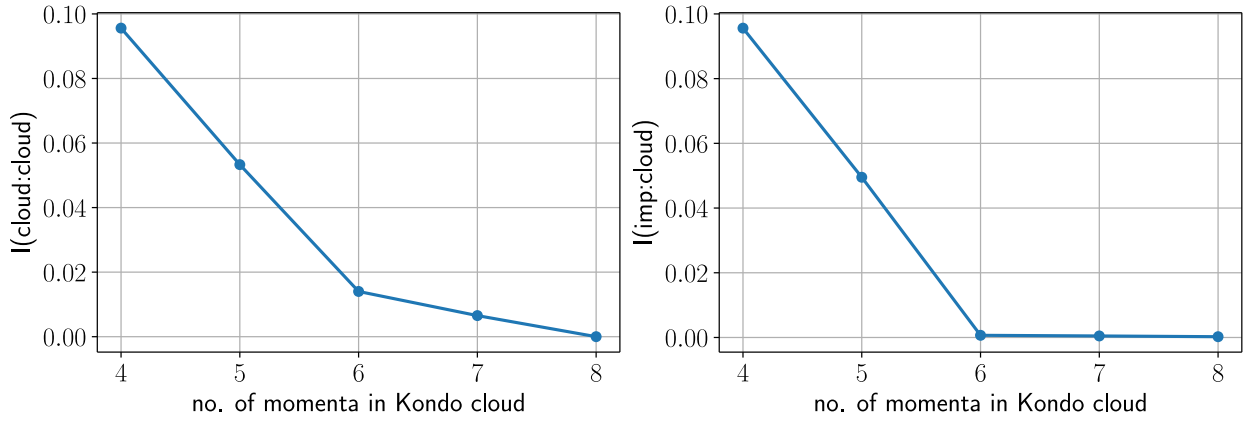


Figure 27: *Left:* Mutual information between two conduction electrons inside the cloud. *Right:* Mutual information between a conduction electron inside the cloud and an impurity electron.

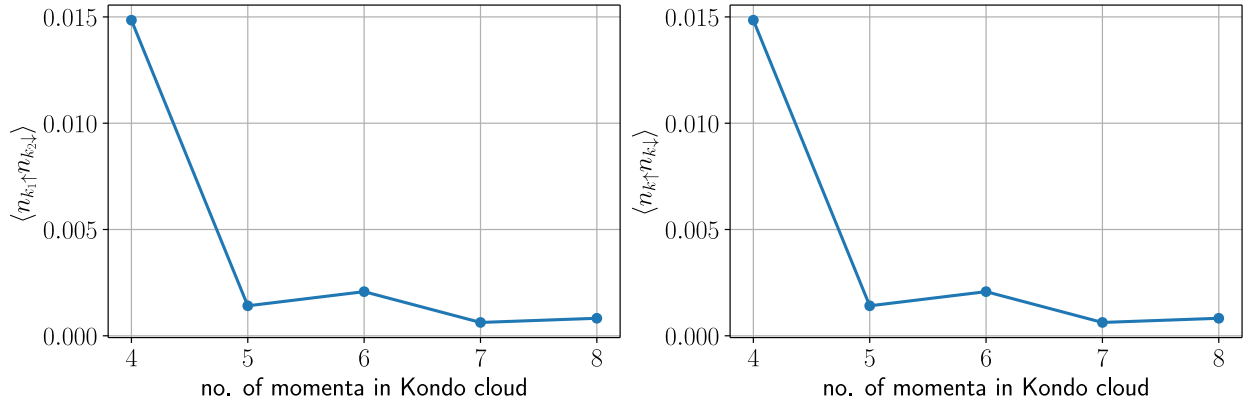


Figure 28: Diagonal correlation functions between cloud electrons

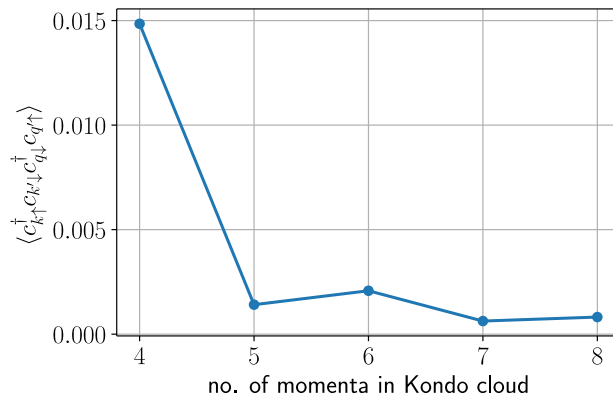


Figure 29: off-diagonal correlation function