

Emergence in free and correlated fermions: from impurity models to the bulk

JRF-to-SRF Presentation

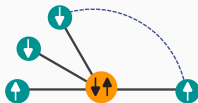
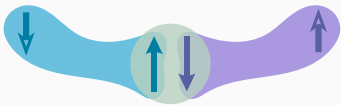
August 11, 2022

Abhirup Mukherjee

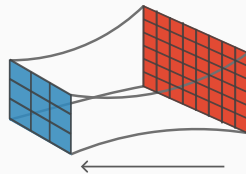
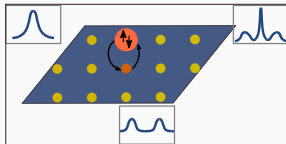
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Summary of Work



Completed Projects

- ✓ Unveiling the Kondo cloud: Unitary renormalization-group study of the Kondo model

Phys. Rev. B 105, 085119, arXiv:2111.10580v3

A. Mukherjee, Abhirup Mukherjee, N. S. Vidhyadhiraja, A. Taraphder, and S. Lal

- ✓ Frustration shapes multi-channel Kondo physics: A star graph perspective

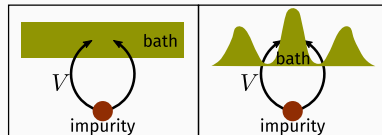
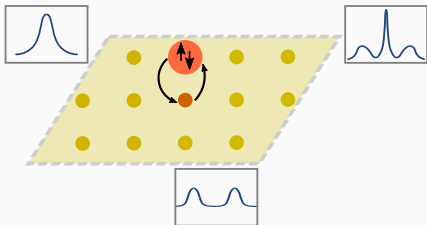
under review at PRB, arXiv:2205.00790

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Ongoing Projects

- ✓ Metal-insulator transition in an extended Anderson impurity model (*manuscript in preparation*)
- ✓ Holography and topology of entanglement scaling in free fermions (*manuscript in preparation*)
- ✓ URG-based auxiliary model approach to correlated systems (*ongoing*)

Local MIT in an extended Anderson impurity model

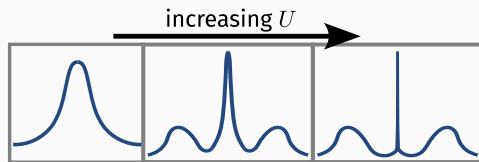


Introducing the extended Anderson impurity model

Introducing the extended Anderson impurity model

Standard Anderson impurity model

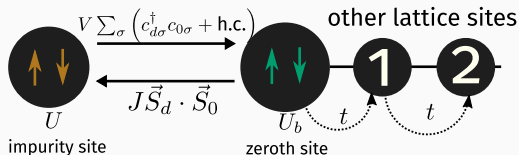
- no local-moment phase, $A(\omega)$ gapless
- cannot explain insulating phase of DMFT



Gap in spectral function requires additional physics!

Extended Anderson impurity model

- impurity-bath spin correlation: J
- bath zeroth site local correlation: U_b



Phase Diagram & Ground-States

Nature of the transition

Universal theory near the transition

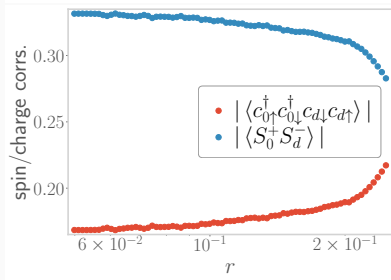
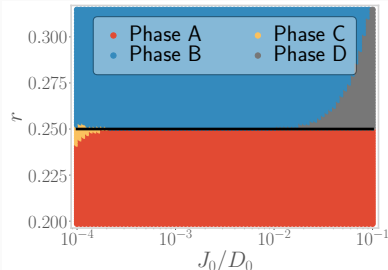
Insights into DMFT

Phase Diagram & Nature of Ground-States

RG Phase diagram:

Local MIT in an extended Anderson impurity model

- Competition between J and U_b leads to phase transition from screened singlet phase at $|U_b| \leq 4J$ to unscreened local moment phase at $|U_b| > 4J$.
- Impurity spectral function becomes gapped beyond the critical point.
- Decoupling the impurity model leads to an effective model with the zeroth site as the correlated impurity, demonstrating the symmetry between the impurity and zeroth site.
- Geometric entanglement and mutual information track the transition by vanishing beyond the critical point.
- Subdominant pairing tendencies are observed near the quantum critical point.

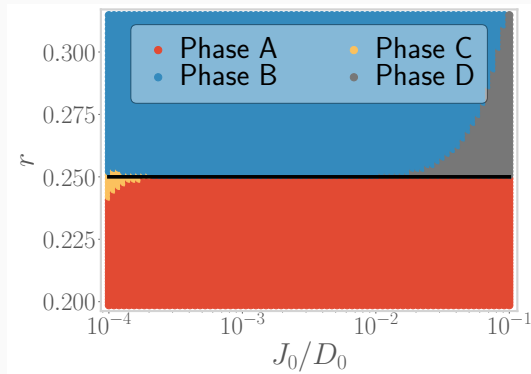


RG Phase Diagram

RG equations reveal critical point where J, V become irrelevant

$$r = -U_b/J$$

1. orange phase: J is relevant: strong-coupling
2. blue phase: J is irrelevant: local moment
3. yellow phase: spin+charge liquid
4. gray phase: all couplings irrelevant



Presence of a phase transition

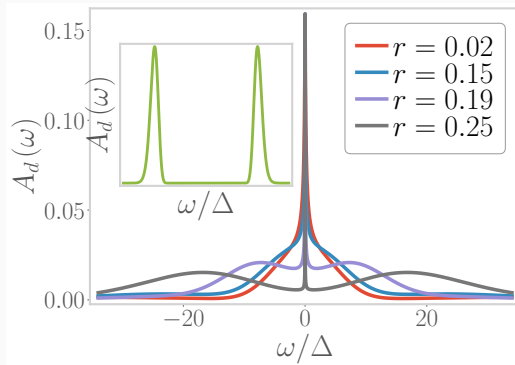
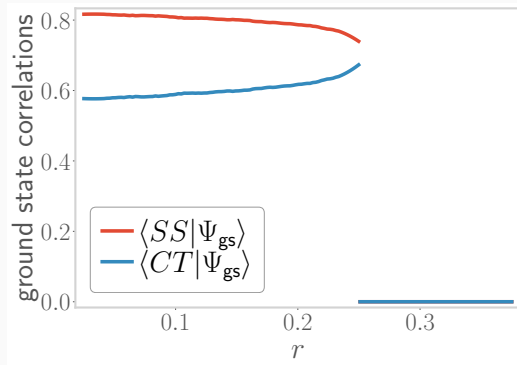
singlet \longrightarrow spin+charge liquid \longrightarrow local moment

impurity spectral function gaps out

$$r = -U_b/J$$

$$|SS\rangle = |\uparrow, \downarrow\rangle - |\downarrow, \uparrow\rangle$$

$$|CT\rangle = |2, 0\rangle + |0, 2\rangle$$



Bath spectral function: towards self-consistency

- Decoupling the impurity site leads to an Anderson impurity model

$$H_{0+\text{rest}} = \underbrace{- (U_0 + U_b) (\hat{n}_{0\uparrow} - \hat{n}_{0\downarrow})^2}_{\text{new correlated impurity}} \underbrace{- t \sum_{\substack{j \in \text{n.n. of } 0, \\ \sigma}} (c_{0\sigma}^\dagger c_{j\sigma} + \text{h.c.})}_{\text{hopping between new impurity \& new bath}} \underbrace{- t \sum_{\langle i,j \rangle} (c_{i\sigma}^\dagger c_{j\sigma} + \text{h.c.})}_{\text{K.E. of new bath}}$$

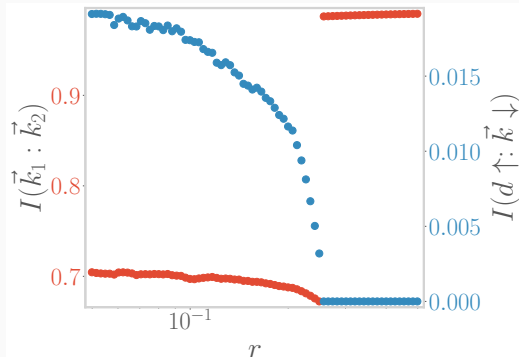
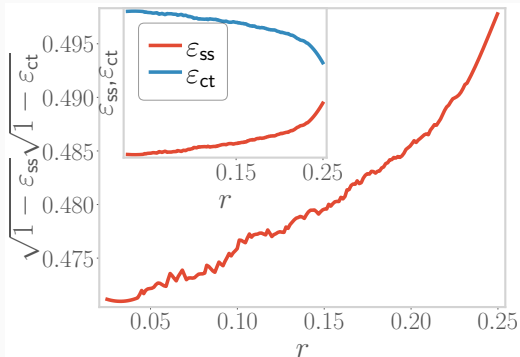
- correlated, dominant spin-flip processes lead to repulsive $U_{\text{eff}} = U_0 + U_b \sim J^*/8$
- J symmetrises the two sites, leading to similar spectral functions \rightarrow essence of self-consistency

Entanglement as a probe for the transition

Geometric entanglement: $\varepsilon(\psi_1, \psi_2) = 1 - |\langle \psi_1 | \psi_2 \rangle|^2$

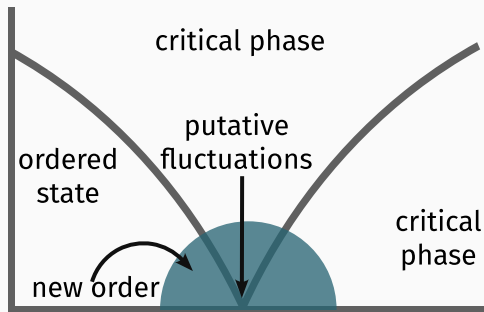
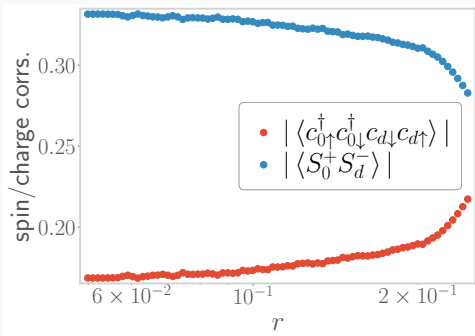
$\rightarrow \sqrt{1 - \varepsilon_{ss}} \sqrt{1 - \varepsilon_{ct}}$ is maximised, then vanishes

Mutual information between impurity and cloud vanishes

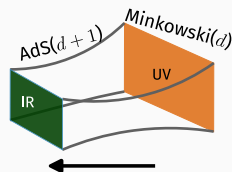
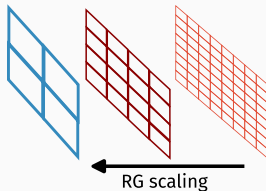


Presence of subdominant pair fluctuations

- **pairing tendencies** observed near the quantum critical point
- might lead to **superconductivity** with doping
- seen in cuprates, heavy-fermions materials, pnictides, etc



Entanglement scaling in free fermions: holography & topology



Creating subsystems

Free Dirac fermions on torus: $k_x^n = \frac{2\pi}{L_x} n$, $n \in \mathbb{Z}$; define **sparsity** $= \Delta n = 1$

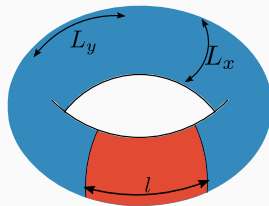
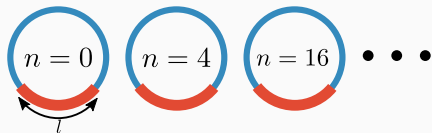
Simplest choice: the entire set

sparsity = 1 $\rightarrow n \in \{-N, -(N-1), -(N-2), \dots, -1, 0, 1, \dots, N-2, N-1, N\}$

Coarser choices: increase sparsity

sparsity = 2 $\rightarrow n \in \{-N, -(N-2), -(N-4), \dots, -2, 0, 2, \dots, N-4, N-2, N\}$

sparsity = 4 $\rightarrow n \in \{-N, -(N-4), -(N-8), \dots, -4, 0, 4, \dots, N-8, N-4, N\}$



Subsystem entanglement entropy: Entanglement hierarchy

$$S_{A_z(j)} = f_z(j) c \alpha L_x - c \log |2 \sin(\pi f_z(j) \phi)|$$

$$i < j, S_{i \cup j} = \begin{cases} S_i, & z > 0 \\ S_j, & z < 0 \end{cases}$$



- presents a **hierarchy** of entanglement \longrightarrow EE distributed across RG steps
RG transformation \longrightarrow reveals entanglement
- distribution of entanglement also present in **multipartite** entanglement

Mutual information = distance

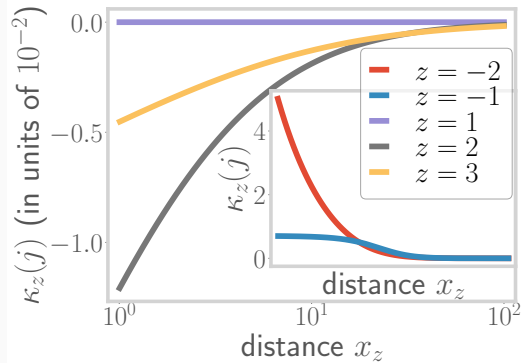
Mutual information: $I^2(A : B) \equiv S(A) + S(B) - S(A \cup B)$ (non-negative)

Define distances using mut. info.

$$x_z(j) = \log t_z(j), \quad y_z(j) = \log t_z(j+1)$$

$$v_z(j) \equiv \Delta y_z(j) / \Delta x_z(j), \quad v' = \Delta v_z(j) / \Delta x_z(j)$$

Curvature as well: $\kappa_z(j) = \frac{v'_z(j)}{[1 + v_z(j)^2]^{\frac{3}{2}}}$



RG evolution = emergent distance

- Distances and curvature can be related to an RG **beta function**
- Amounts to an **explicit demonstration** of the holographic principle
- Sign of curvature is **topological**, can be written in terms of winding numbers

Topological nature of geometry-independent term

$$S_{A_z(j)} = f_z(j) c \alpha L_x - \underbrace{c \log |2 \sin(\pi f_z(j) \phi)|}_{=Q(\phi), \text{geometry-independent term}}$$

- $Q(\phi)$ is periodic in the flux ϕ , $\phi = 1$ transports a charge across Fermi surface
- pole structure of $\left(\sin \frac{\pi}{4} - |\sin(\pi f_z(j) \phi)|\right)^{-1}$ counts number of states \rightarrow tracks Luttinger volume
- Luttinger volume is topological, so is $Q(\phi)$; $Q(\phi)$ can be expressed in terms of winding numbers

Future Prospects

Future Prospects

- Better model can be obtained by taking multiple impurities and general impurity filling
- novel auxiliary model method can be used for studying other models of strong-correlations as well as topologically active or flat band systems
- The URG can be applied to heavy-fermion materials towards a study of phase diagram and unconventional superconductivity, as well as Kondo insulators
- Interacting systems in a magnetic field is also a potential area of study, specifically fractional Chern insulators (e.g. the fractional quantum hall effects)

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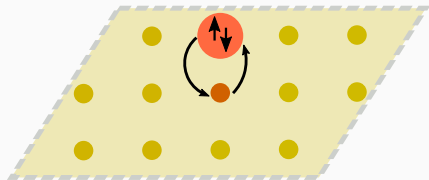
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Further Details

Theory for the single-channel Kondo cloud

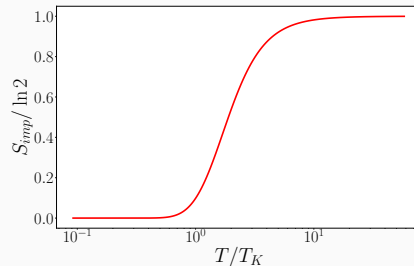
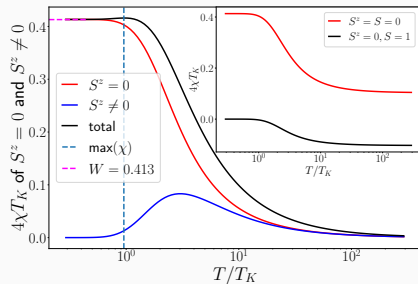
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Theory for the single-channel Kondo cloud

✓ spectral function & magnetic susceptibility

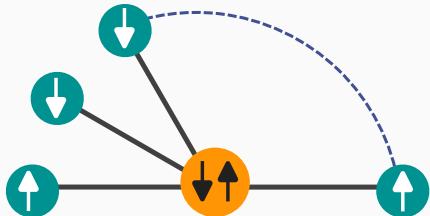


- ✓ local Fermi liq. & orthogonality catastrophe
- ✓ thermal entropy

Role of degeneracy in the multi-channel Kondo problem

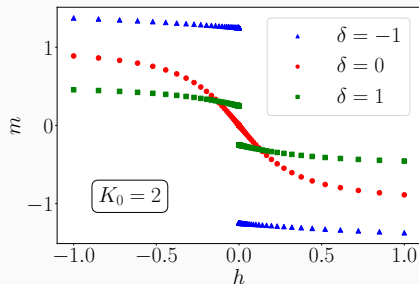
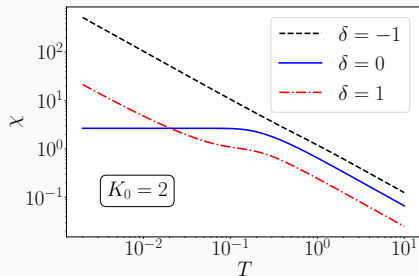
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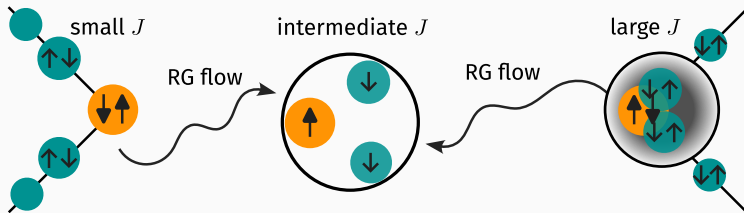
Role of degeneracy in the multi-channel Kondo problem

- ✓ Intermediate-coupling RG fixed point Hamiltonian and **degenerate** ground states
- ✓ Degree of compensation, magnetization and susceptibility show **incomplete screening**



Role of degeneracy in the multi-channel Kondo problem

- ✓ Local **marginal Fermi liquid** within the low-energy excitations of the bath
- ✓ **Duality** relations constrain the RG flows of the MCK model



Holography and topology of entanglement scaling in free fermions

Future Prospects

Improvements to the auxiliary model

- Better model can be obtained by using multiple impurities
- Allows entangled liquid-like insulating phases
- Might also provide k -space resolution
 - partial gapping of Fermi surface?
 - pseudogap phases
- Introducing general impurity filling
 - new phases?
 - dominant pair fluctuations?

A novel auxiliary model approach

- Using local impurity models to create bulk lattice models (Bloch's theorem)

$$H_{\text{bulk}} = \sum_i H_{\text{local}}(i), \quad \Psi_{\text{bulk}}(\vec{k}) \sim \sum_i e^{i\vec{k} \cdot \vec{r}_i} \Psi_{\text{local}}(i)$$

- Relates bulk correlation functions to those of the auxiliary model
- phase transition in the extended AIM \longrightarrow phase transition in the bulk model, **metal-insulator transition** in Hubbard-Heisenberg model

A novel auxiliary model approach

- Should be useful for studying other models of strong-correlations
 - periodic Anderson/Kondo models
 - Heisenberg models
- Another potential application: topologically active systems:
 - Fractional quantum hall systems
- Extend the formalism towards higher order Greens functions
 - two-particle Greens functions, doublon-holon correlations
 - can provide more info on the MIT

- Materials with very high quasiparticle masses
- Outstanding questions exist about the nature of phases and phase transitions
 - microscopic justification of certain phases
 - theory for the strange metal excitations
 - microscopic justification for the origin of unconventional superconductivity
- the URG, MERG and auxiliary model methods should prove useful