

# Project 3

Kandidatnummer

October 28, 2020

## 1 Introduction

Though simulating a solar system is interesting and fun enough on its own, it is naturally also quite useful for the study of astrodynamics. In addition, being able to simulate a solar system provides a good set of tools applicable to many other scientific areas. These tools include having a good understanding of different numerical integration methods, and being able to write a structured and fast code. In this project we wish to explore a model of our own solar system, beginning with simulating the simple two-body system including the Earth and the Sun. We will use this system to compare two different methods of numerical integration, the forward Euler method and the velocity Verlet method. This simple system also makes a good testing ground for exploring whether our model is consistent with known physical laws such as energy conservation and Kepler's laws of planetary motion. We will also test the stability of the velocity Verlet method by including Jupiter and playing around with its mass. From there we will include the rest of the planets in our solar system, and blah blah blah

general relativity.

## 2 Theory

### 2.1 Forward Euler

The forward Euler method is an algorithm to estimate the solution of a differential equation. The Forward Euler method wants to find the next point  $\mathbf{r}_{i+1}$ . To find the next point, it uses the point it is at,  $\mathbf{r}_n$ , a small time step,  $dt$ , and the derivative of its position. Which can be expressed like this

$$\mathbf{r}_{n+1} = \mathbf{r}_n + \mathbf{r}'_n \cdot dt \quad (1)$$

Where  $\mathbf{r}_{n+1}$  is the next step,  $\mathbf{r}_n$  is the current step,  $\mathbf{r}'_n$  is the derivative of the current step and  $dt$  is the time step. This algorithm is really based on an abbreviated version of a Taylor expansion, where we only expand the series one step at a time. But by only taking one step at a time, we will also get a local truncation error, which causes an er-

ror for each step we take.

$$\begin{aligned} \mathbf{r}(t_n + dt) &= \mathbf{r}_{n+1} \\ &= \mathbf{r}(t_n) + dt \left( \mathbf{r}'(t_n) + \frac{\mathbf{r}''(t_n)}{2!} dt \right. \\ &\quad \left. + \dots + \frac{\mathbf{r}^{(p)}(t_n)}{p!} dt^{p-1} \right) + O(dt^{p+1}) \end{aligned} \quad (2)$$

Where  $O(dt^2)$  is the local truncation error. Since the Forward Euler method is a first-order method, will the local truncation error be proportional to the square of the step size.

In our case  $\mathbf{r}(t_n) \rightarrow \mathbf{r}_n$  is the position,  $\mathbf{r}'(t_n) = \mathbf{v}(t_n) \rightarrow \mathbf{r}'_n = \mathbf{v}_n$  is the velocity and  $\mathbf{v}(t_n) = \mathbf{a}(t_n, r_n) \rightarrow \mathbf{v}_n = \mathbf{a}_n$  is the acceleration.

The update of our position and velocity with a time step is then given as

$$\mathbf{r}(t_{n+1}) = \mathbf{r}(t_n) + \mathbf{v}(t_n)dt + O(dt^2) \quad (3)$$

$$\mathbf{r}_{n+1} = \mathbf{r}_n + \mathbf{v}_n dt + O(dt^2) \quad (4)$$

$$\mathbf{v}(t_{n+1}) = \mathbf{v}(t_n) + \mathbf{a}_n dt + O(dt^2) \quad (5)$$

$$\mathbf{v}_{n+1} = \mathbf{v}_n + \mathbf{a}_n dt + O(dt^2) \quad (6)$$

Hence, will forward Euler algorithm look like this

$$\mathbf{r}_{n+1} = \mathbf{r}_n + \mathbf{v}_n dt \quad (7)$$

$$\mathbf{v}_{n+1} = \mathbf{v}_n + \mathbf{a}_n dt \quad (8)$$

$$t_{n+1} = t_n + dt \quad (9)$$

The number of FLOPS in the Forward Euler method will be 4 FLOPS for each time step, when we look at the addition and multiplication for this algorithm. Since the loop is looping over N times will the number of FLOPS be  $4N$ .

There is also important to notice since

the forward Euler method updates the position before the velocity leads to the energy not being conserved

## 2.2 Velocity Verlet

The velocity Verlet is based on the kinematic equations for an moving object, which in our case is the earth's orbit around the sun. If we want to find the next time step for the velocity and position we do a approximation and uses Taylor-expansion.

$$\mathbf{r}(t \pm dt) = \mathbf{r}(t) \pm \mathbf{r}'(t)dt + \frac{\mathbf{r}''(t)}{2}dt^2 \quad (10)$$

$$+ O(dt^3) \quad (11)$$

This gives us the position and velocity

$$\mathbf{r}_{n+1} = \mathbf{r}_n + \mathbf{v}_n dt + \frac{\mathbf{a}_n}{2} dt^2 + O(dt^3) \quad (12)$$

$$\mathbf{v}_{n+1} = \mathbf{v}_n + \mathbf{a}_n dt + \frac{\mathbf{v}''_n}{2} dt^2 + O(dt^3) \quad (13)$$

We rewrite expression 19 to define  $\mathbf{v}''_n$ , we do this by by defining the next step in acceleration  $\mathbf{a}_{n+1} = \mathbf{v}'_n$

$$\mathbf{a}_{n+1} = \mathbf{v}'_{n+1} = \mathbf{v}'_n + \mathbf{v}''_n dt + O(h^2) \quad (14)$$

$$\mathbf{v}''_n = \frac{\mathbf{v}'_{n+1} - \mathbf{v}_n}{dt} \quad (15)$$

$$= \frac{\mathbf{a}_{n+1} - \mathbf{a}_n}{dt} \quad (16)$$

Plugging equation 22 into 19 we get the general algorithm for Velocity Ver-

let.

$$\mathbf{r}_{n+1} = \mathbf{r}_n + \mathbf{v}_n dt + \frac{\mathbf{a}_n}{2} dt^2 + O(dt^2) \quad (17)$$

$$\mathbf{v}_{n+1} = \mathbf{v}_n + \frac{\mathbf{a}_{n+1} - \mathbf{a}_n}{2} dt + O(dt^3) \quad (18)$$

NUMBER OF FLOPS

### 3 Method

#### 3.1 The Earth-Sun system

We assume that the earths orbit around the sun is circular. And from circal motion and Newtons gravita-tional force we have that

$$F_G = G \frac{M_E M_\odot}{r^2} = \frac{M_E v^2}{r} \quad (19)$$

Where  $G$  is the gravitational constant,  $M_\odot$  is the mass of the sun,  $M_E$  is the mass of the earth,  $r$  is the distance between the earth and the sun, and  $v$  is the velocity of the moving object, earth.

We now want to show that the veloc-ity,  $v$ , of the earth can be written as

$$v^2 r = G M_\odot = 4\pi^2 \frac{AU^3}{yr^2} \quad (20)$$

Since we say it is circular orbits can we use the centripetal force to rewrite equation X

$$M_E \omega^2 r = G \frac{M_E M_\odot}{r^2} \quad (21)$$

where  $\omega^2$  is the angular velocity of the earth. Equation X can now be rewrit-ten as

$$M_E \left( \frac{2\pi}{P} \right)^2 r = G \frac{M_E M_\odot}{r^2} \quad (22)$$

$$(23)$$

Where  $P$  is a period of the earth around the sun, which is 1 yr. Now can we use Keepler's third law, which says that the square of an orbital pe-riod  $P^2$  equals the semi-major axis as it orbits the sun,  $a$ , OPPHØYD I TREDJE  $a^3$ . Where 1 a is equal to 1 AU. We will replace  $r$  with  $a$  in equation X.

$$P^2 = \frac{4\pi^2}{G M_\odot} a^3 \quad (24)$$

Where we substitute the circular ra-dius  $r$  with the semi-major axis  $a$ .

#### 3.2 Conservation of angular momentum

The definition of Keeplers law is that a area, where a connecting line be-tween a planet and the sun, is always the same for an equal time interval. Which is illustrated in the figure be-low:

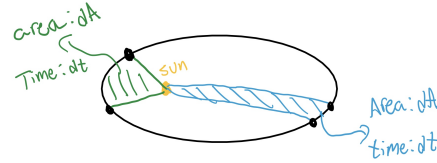


Figure 1: Keeplers second law. Where the area is the same for the same time interval.

The best way to show that the an-gular momentum is conserved by us-ing Keeplers second law is to make a drawing.

The best way to show that the an-gular momentum is conserved by us-ing Keeplers second law is to make a drawing:

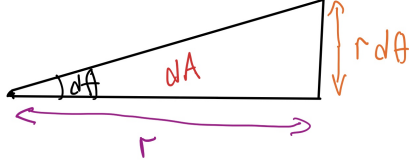


Figure 2: This square is the infinitesimal area  $dA$  that the planet has moved by a infinitesimal time interval  $dt$

Figure 2 shows us that the area is

$$dA = \frac{1}{2}r^2 d\theta \quad (25)$$

We now want to find an infinitesimal area where the planet has moved around the sun, over an infinitesimal time step, where we use equation 12.

$$\frac{dA}{dt} = \frac{1}{2}r^2 \frac{d\theta}{dt} \quad (26)$$

$$= \frac{1}{2}rv_\theta \quad (27)$$

And the definition of angular momentum is given as

$$L = mrv_\theta \quad (28)$$

With equation 14 and 15 we can show that the angular momentum is conserved.

### 3.3 Testing forms of the force

Until now we have looked at the movement for the planet with the gravitational force with the square of  $r$ . We will now try to change the gravitational force where the power of  $r$  is in the range  $\beta \in [2, 3]$  The gravitational force is then given as

$$\mathbf{F}_G = -\frac{GM_\odot M_{planet}}{r^\beta} \frac{\mathbf{r}}{r} \quad (29)$$

### 3.4 Escape velocity

We will now look at a planet which begins at a distance 1 AU from the sun, and see how fast the initial velocity must be for the planet to be able to escape the sun. We find the analytical solution to compare it to our numerical solution.

The analytical expression is given as

$$|\mathbf{V}_{esc}| = \sqrt{\frac{2GM_\odot}{r}} \quad (30)$$

By plugging equation 2 into 16, and use that our initial position of the planet is  $r = AU$  we get

$$|\mathbf{V}_{esc}^2| = \frac{2 \cdot 4\pi^2 AU^3}{yr^2} \cdot \frac{1}{r} \quad (31)$$

$$= \frac{2 \cdot 4\pi^2 AU^3}{yr^2} \cdot \frac{1}{AU} \quad (32)$$

$$|\mathbf{V}_{esc}| = 2\sqrt{2}\pi \cdot \frac{AU}{yr} \quad (33)$$

### 3.5 Perihelion precision

LES GJENNOM DETTE MARCUS

The definition of perihelion precision is that the perihelion point of an objects, in our case Mercury, changes its position because of the gravitational field deflection from light. RIKTIG? The perihelion point for Mercury changes with 43'' arcseconds for each century.

Mercury's orbit has a small precision for each century. This comes from that Mercury is the nearest planet to the sun, and therefore its acceleration and velocity will be large, which means that the space time will warp and we

have to obey the rules of General Relativity. Hence, we rewrite the expression for Newton's gravitational force to a relativistic gravitational force.

$$\mathbf{F} = -G \frac{M_{\odot} M_{\text{Mercury}}}{r^2} \left[ 1 + \left( \frac{3l^2}{r^2 c^2} \right) \right] \frac{\mathbf{r}}{r} \quad (34)$$

## 4 Results

### 4.1 The Earth-Sun system

From Kepler's third law we have

$$P^2 = \frac{4\pi^2}{GM_{\odot}} a^3 \quad (35)$$

$$GM_{\odot} = 4\pi^2 \frac{AU^3}{yr^2} \quad (36)$$

### 4.2 Conservation of angular momentum

If we substitute equation 15 into equation 14

$$\frac{dA}{dt} = \frac{1}{2} \frac{L}{m} \quad (37)$$

Since  $\frac{dA}{dt} = \text{constant}$  and the mass  $m$  is constant means that the angular momentum also is constant and therefore also conserved.

### 4.3 Escape velocity

In the figure below we can see the behavior for the Sun-Earth system with different initial velocities.

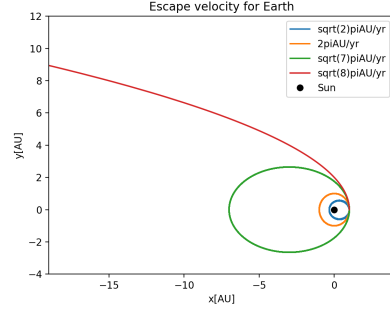


Figure 3: Escape velocity for the Sun-Earth system for different initial velocity

## 5 Discussion

### 5.1 The Earth-Sun system

It is important to note that Kepler's third law is given as  $P^2 = a^3$ .

$$P^2 = a^3 \quad (38)$$

### 5.2 Testing forms of the force

When we let  $\beta$  varies between  $[2, 3]$  we get different results depending on the initial values for position and velocity. If we for example set the distance between the sun and earth  $|\mathbf{r}| > AU$  larger than 1 AU, will the gravitational force decrease when we increase  $\beta$ , and the orbit will get more and more elliptical.

If we say that the distance  $|\mathbf{r}| < AU$ , is larger than 1 AU, the Earth will move towards the sun since the gravitational force will increase when we increase  $\beta$ .

### 5.3 Escape velocity

If we increase the initial velocity we see that the orbit of the earth changes from circular orbit to an elliptical orbit, and when the elliptical orbit increase will the position between the two object increase. Hence, will the force of gravity decrease. And when

the initial velocity to the planet make it to the escape velocity it will be tear loose from the gravitational field.

It is also important to add that we say the center of mass does not move, which means that the center of mass will be constant and it will be very close to the Sun.

## **A   Calculations**

### **References**

[1] Skriv inn kilde her.