Large-Scale Numerical Investigations into the Dynamics of Nonlinear Classical Systems

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Outline

Acknowledgements

- I would like to thank A.I. Nicolin, and V. Băran for helping and motivating me.
- The author has been supported by the research project PN-III-P4-ID-PCE-2016-0792 funded by the Romanian Minisrty of Research and Inovation.
- All numerical simulations were performed on the computing cluster of Department of Computational Physics and Information Technologies, “Horia Hulubei” National Institute for Physics and Nuclear Engineering.

Introduction

The model

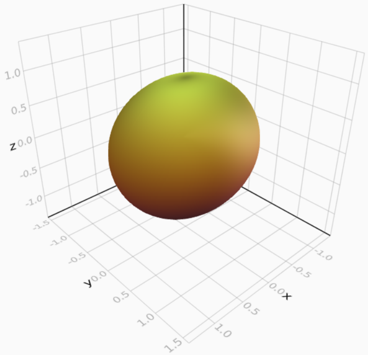
- The physical system that we model is the surface of heavy nuclei.
- The Hamiltonian describes the constrained motion of the vibrational quadrupole degrees of freedom of nuclear surface.

The model

The Hamiltonian of the system

$$H = \frac{A}{2}\left(p_0^2 + p_2^2\right) + \frac{A}{2}\left(q_0^2 + q_2^2\right) + \frac{B}{\sqrt{2}} q_0 \left(3q_2^2 - q_0^2\right) + \frac{D}{4}\left(q_0^2 + q_2^2\right)^2$$

- Harmonic oscillator part
- Integrable part
- Non-integrable term



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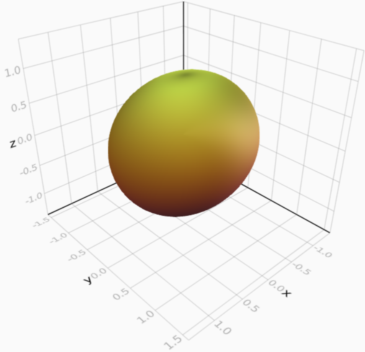

The model

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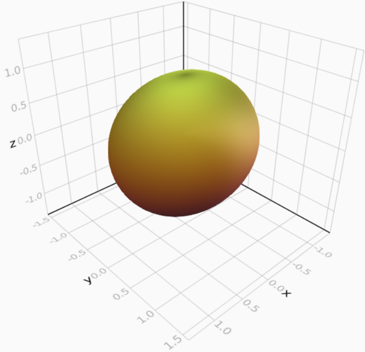

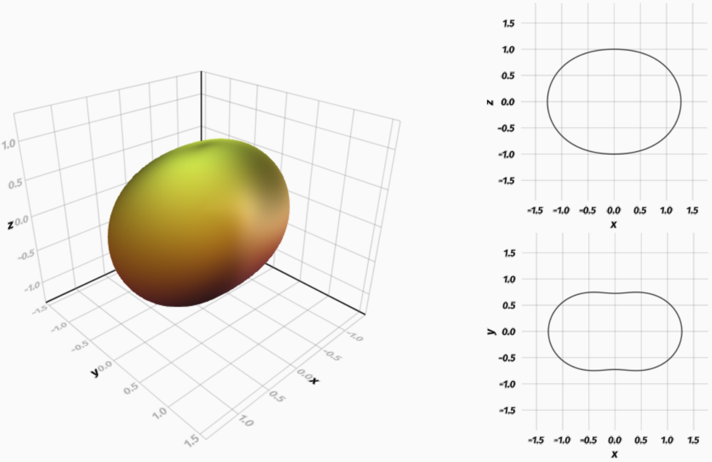

Figure 1: The nuclear surface and its sections

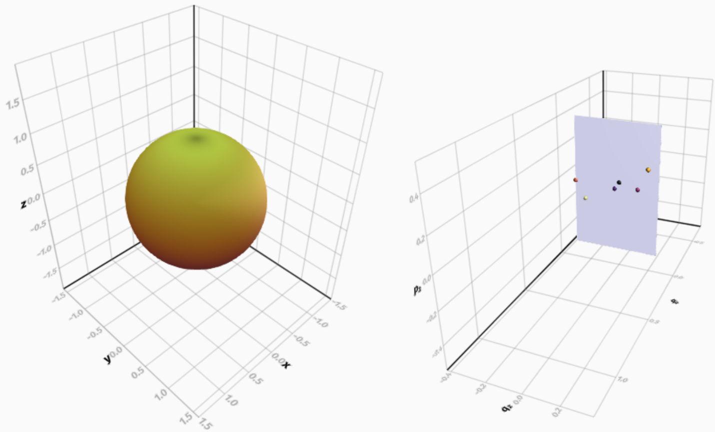

Figure 2: The nucleus and the corresponding trajectory in the phase space for a chaotic trajectory with $B = 0.5, E = 0.3$

Numerical simulations

Julia

- Numerical simulations and the visualizations of the results was done in Julia (using DifferentialEquations.jl and DynamicalSystems.jl).
- Having access to the implementations of a large number of integrators helps us taking an informed decision for the choiche of the integration algorithm.

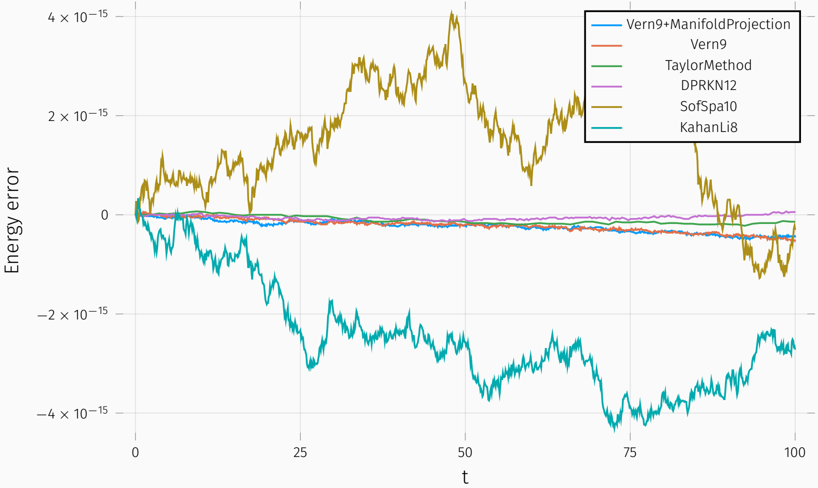

Figure 3: Energy error benchmark for short integration time

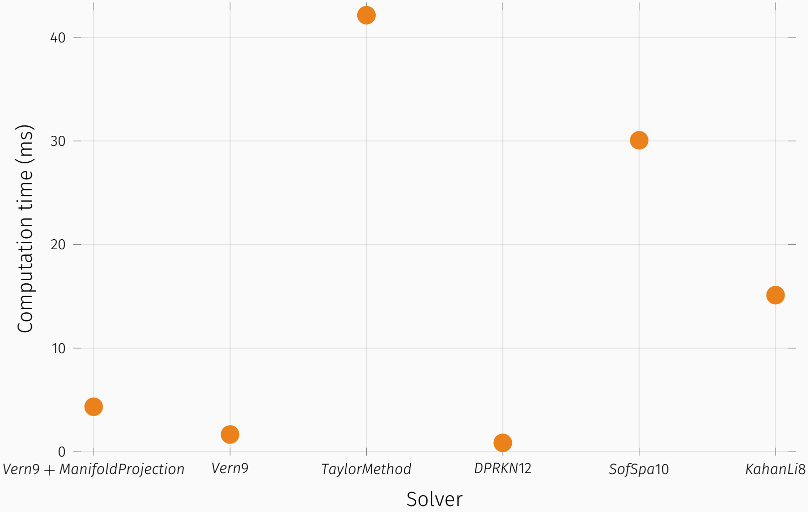

Figure 4: Computational time benchmark for short integration time

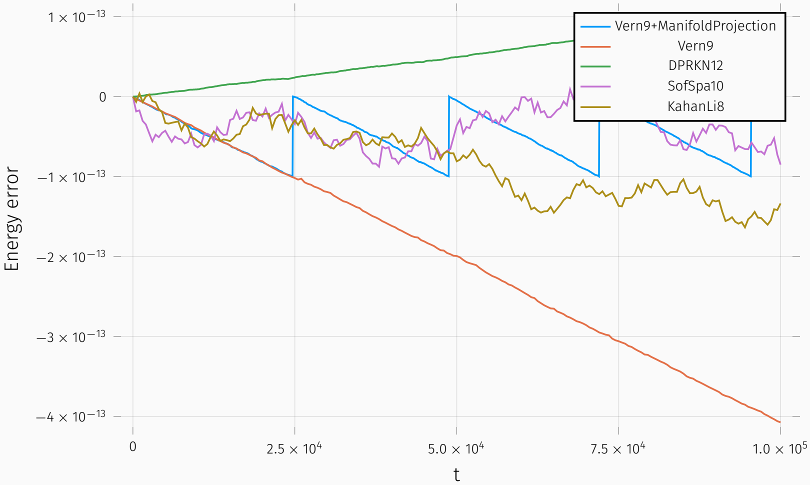

Figure 5: Energy error benchmark for long integration time

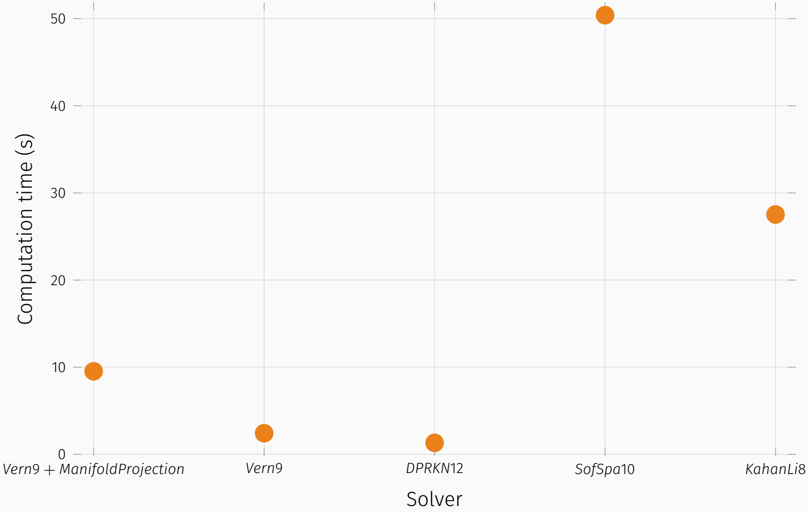

Figure 6: Computational time benchmark for long integration time

Quantifying chaoticity

- The maximal Lyapunov exponent is defined as

$$\lambda = \lim_{t \to \infty} \frac{1}{t} \ln \frac{\mathrm{d}x\,(t)}{\mathrm{d}x\,(0)}$$

- One of the signature characheristics of chaos is the sensitivity to the initial conditions ($\lambda > 0$).
- We can get some intuitive insight for the sensitivity to the initial conditions by following the evolution of two nearby initial conditions.

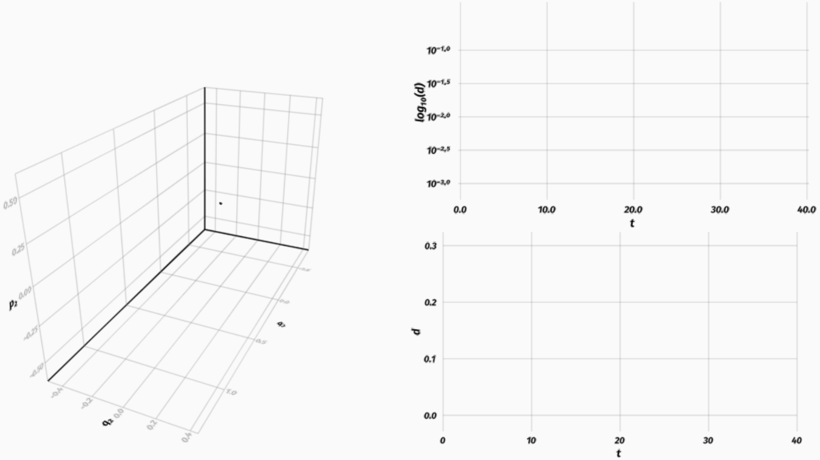

Figure 7: The distance between nearby trajectories for $B = 0.5, E = 0.3$

Interactive exploratinos

- For a given set of parameters (B and E) we have a set of compatible initial conditions.
- Poincaré sections give us a global picture of the dynamics.
- For a better (visual) understanding of the dynamics we can show the Lyapunov exponents on the Poincaré map.

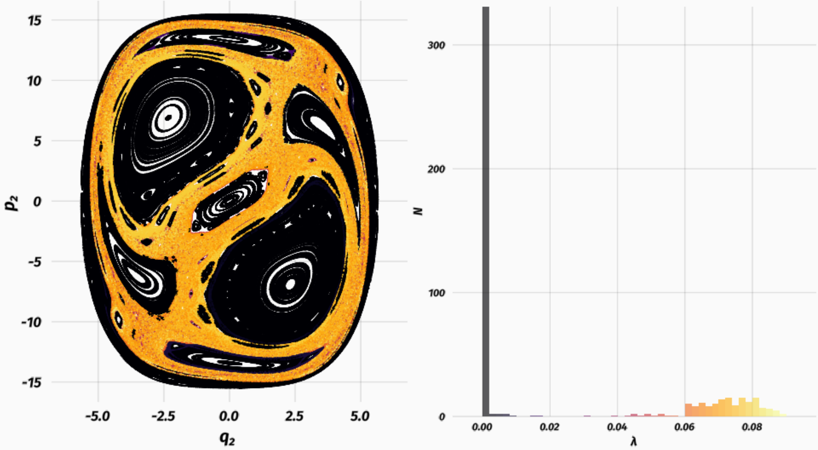

Figure 8: A Poincaré section at $B = 0.5, E = 120$

A note regarding the λ histogram

- Theoretically the Lyapunov exponent histogram should have two sharp peaks: one for the regular part and one for the chaotic one.
- The spread in the chaotic part is given by finite time effects.
- To better understand this we will take a look at how the integration time affects the results.

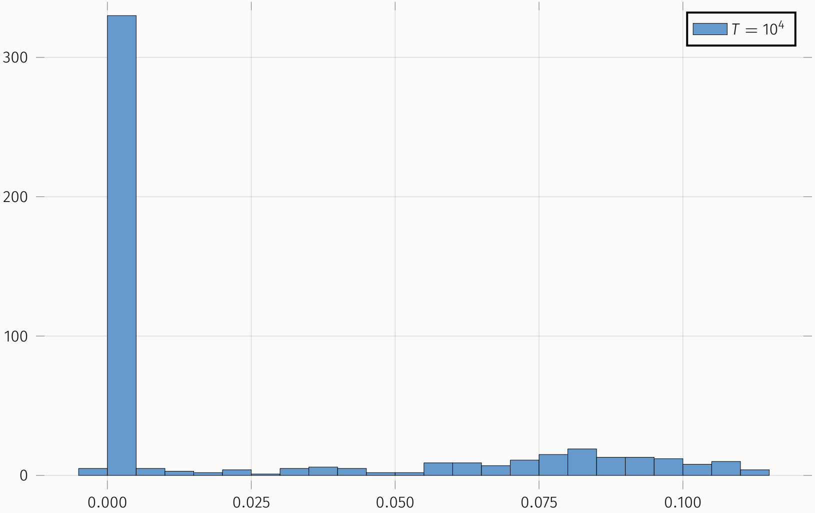

Figure 9: Maximal Lyapunov coefficient histogram for $B = 0.5$, $E = 120$.

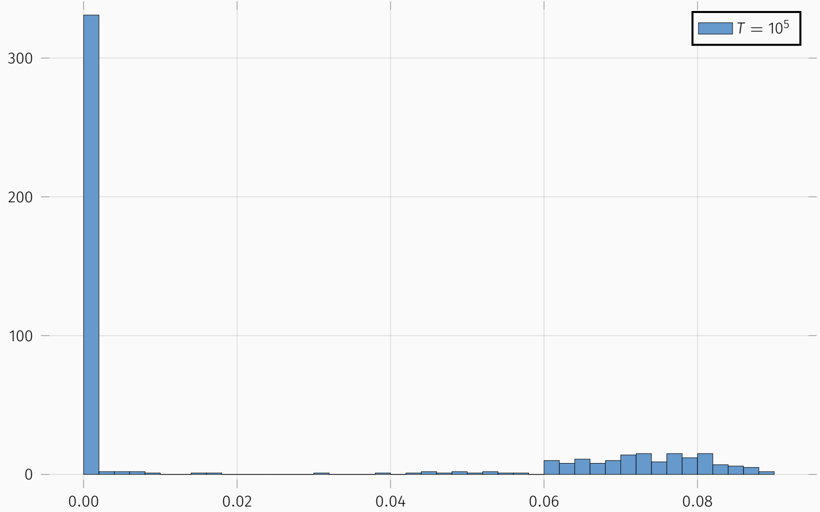

Figure 10: Maximal Lyapunov coefficient histogram for $B = 0.5, E = 120$.

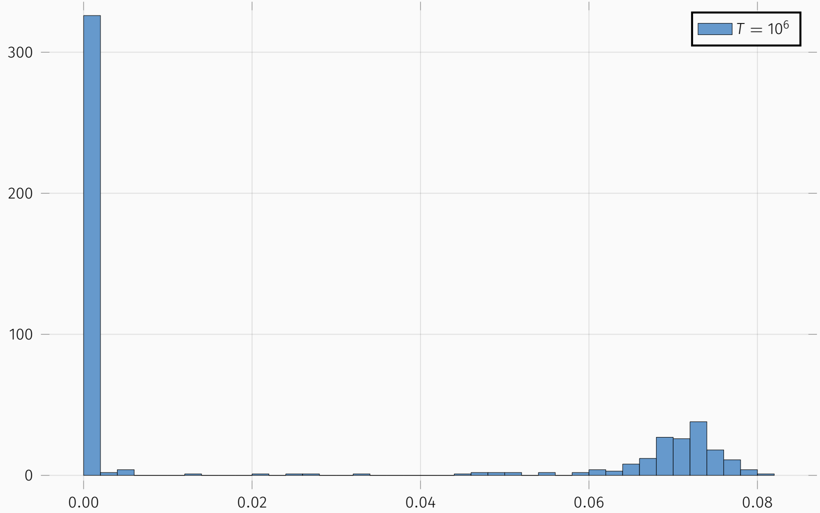

Figure 11: Maximal Lyapunov coefficient histogram for $B = 0.5, E = 120$.

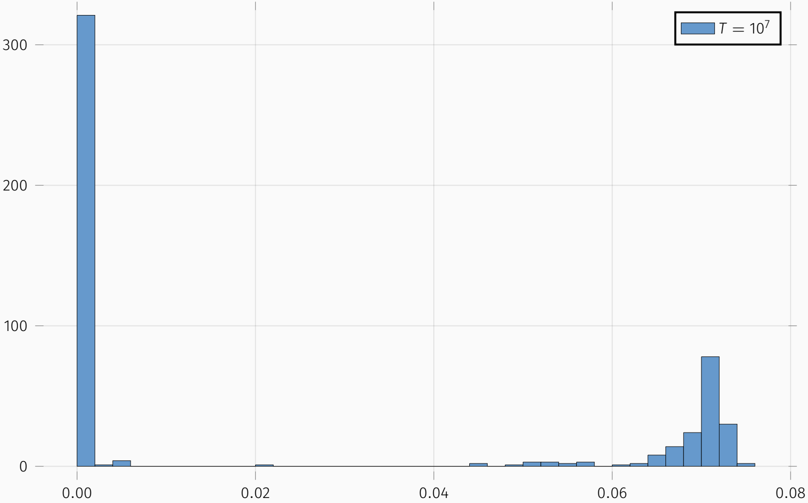

Figure 12: Maximal Lyapunov coefficient histogram for $B = 0.5, E = 120$.

Averaging λ over the initial conditions

- We define the averaged Lyapunov coefficient as the mean of the maximal Lyapunov exponents in the chaotic region.
- We need a sufficiently robust method of selecting the λs in the chaotic region.
- We consider as chaotic everything after the first local maxima in the histogram.

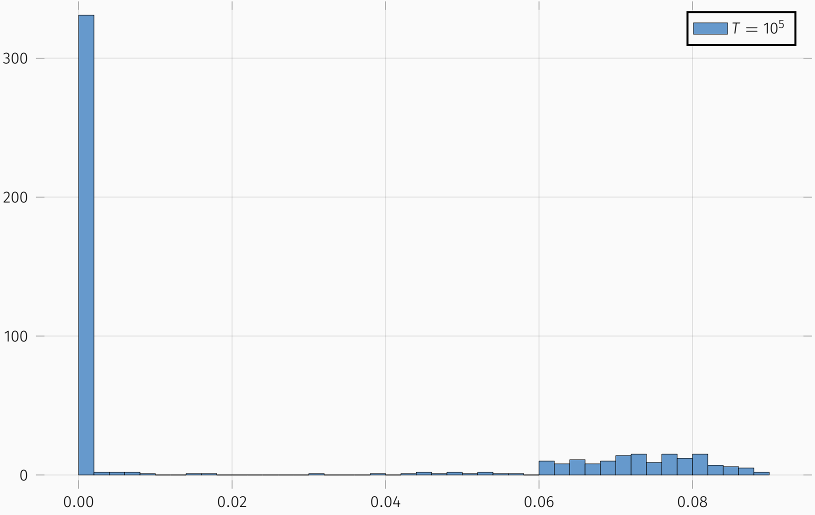

Figure 13: Selecting the chaotic trajectories for $B = 0.5, E = 120$.

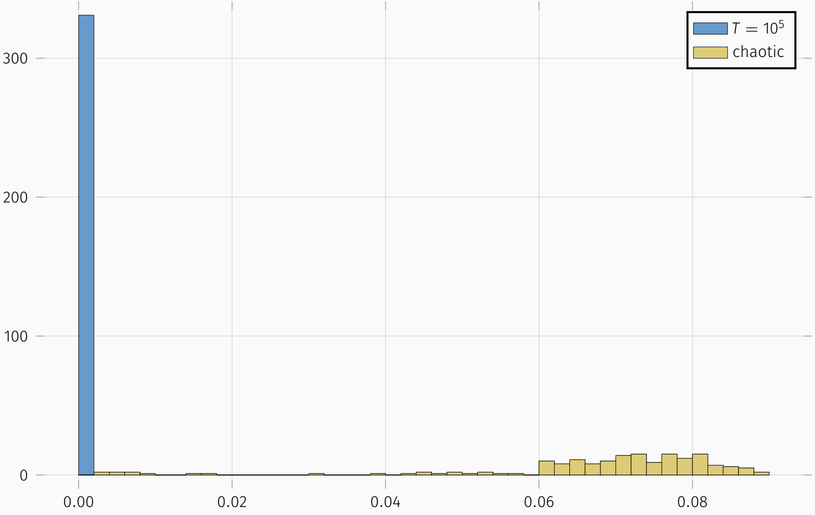

Figure 14: Selecting the chaotic trajectories for $B = 0.5$, $E = 120$.

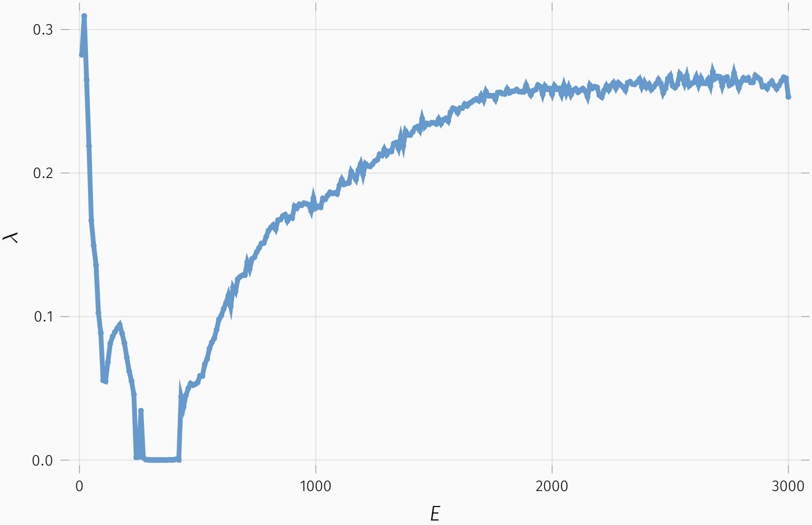

Figure 15: Averaged λ for $B = 0.5, E \in (10, 3000)$.

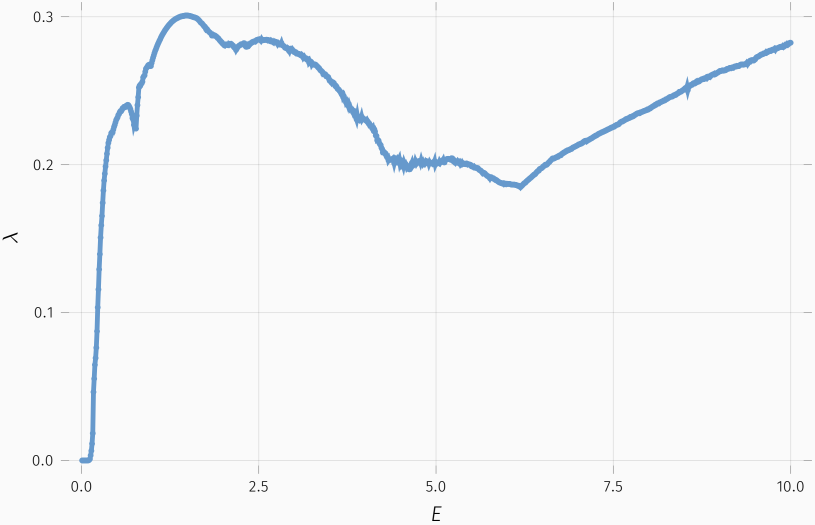

Figure 16: Averaged λ for $B = 0.5, E \in (0.01, 10)$.

Other indicators

- We can look at the distance between two nearby trajectories in the limit of $T \to \infty$ in order to get an estimate of the phase space volume.
- We define Γ as

$$\Gamma = \frac{e^{\lambda} - 1}{d_{\infty}}$$

- We can apply similar averaging techniques as for λ.

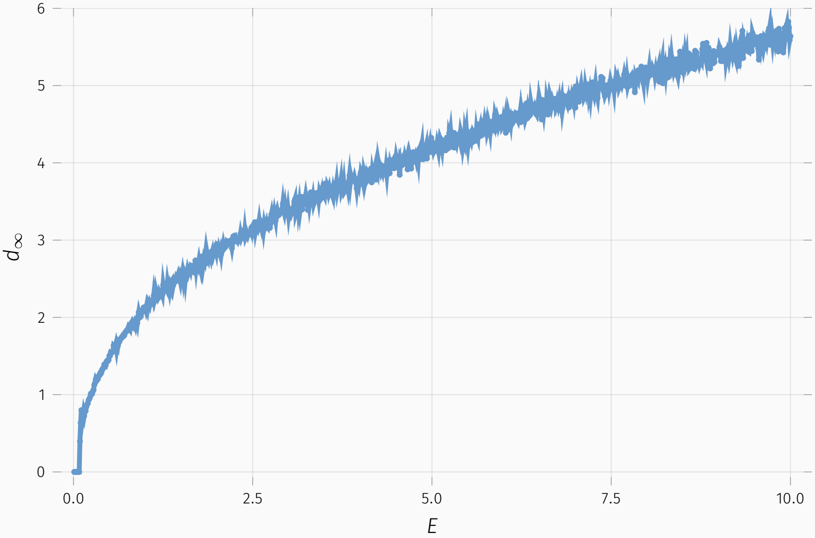

Figure 17: Averaged d_∞ for $B = 0.5, E \in (0.01, 10)$.

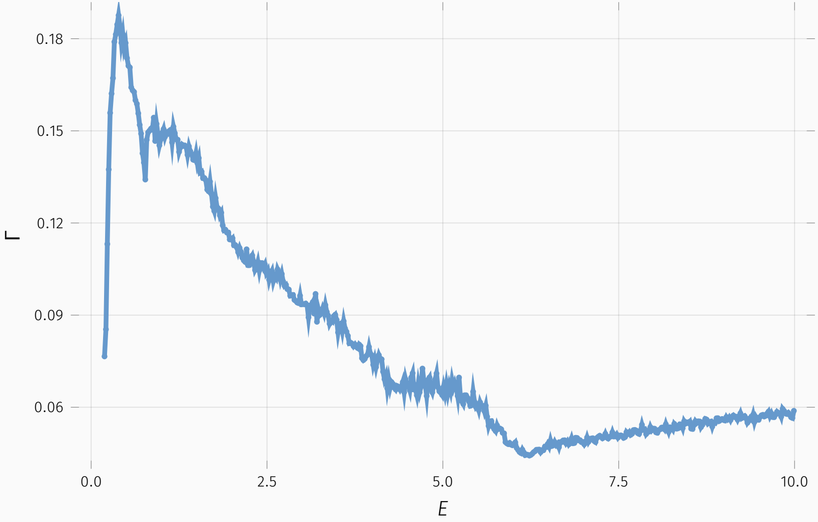

Figure 18: Averaged Γ for $B = 0.5, E \in (0.01, 10)$.

Averaging λ over the initial conditions

- We can get a global picture of the sistem by integrating the averaged Lyapunov coefficient over an energy interval.
- For a large energy interval we have a non-trivial dependence of the averaged Lyapunov exponent with respect to the non-integrability parameter *B*.

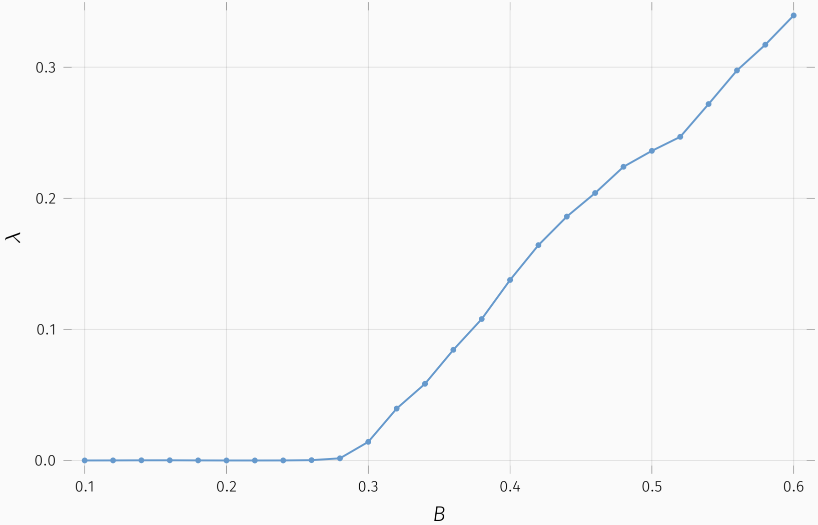

Figure 19: Averaged λ for $B = 0.5, E \in (10, 3000)$ and $B \in (0.1, 0.6)$.

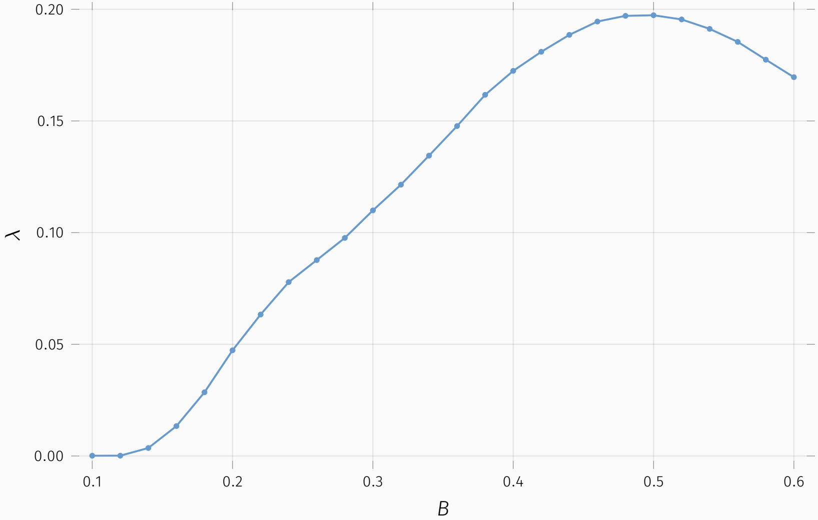

Figure 20: Averaged λ for $B = 0.5$, $E \in (0.01, 10)$ and $B \in (0.1, 0.6)$.

Conclusions

Thank you!

References i

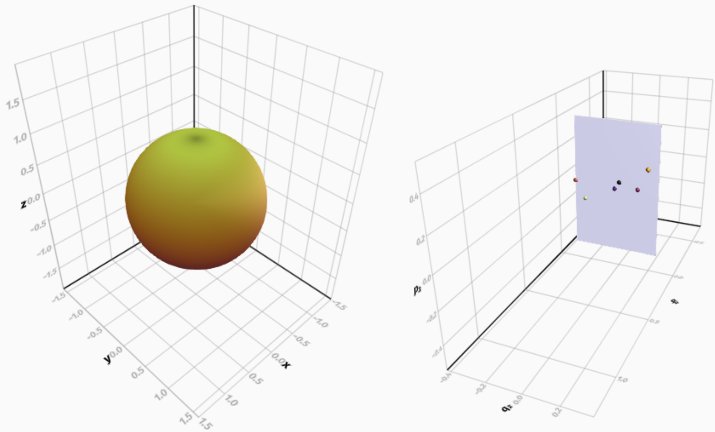

Figure 21: The nucleus and the corresponding trajectory in the phase space for a regular trajectory with $B = 0.5, E = 0.3$

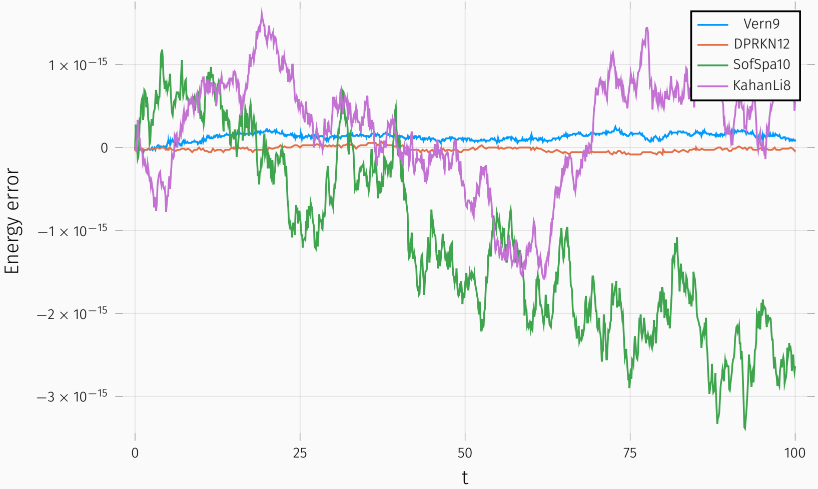

Figure 22: Energy error benchmark for short integration time with rescaling

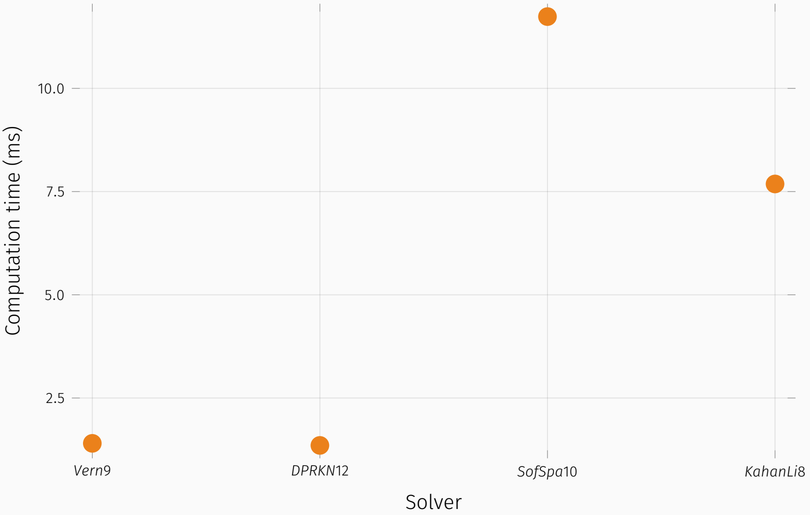

Figure 23: Computational time benchmark for short integration time with rescaling

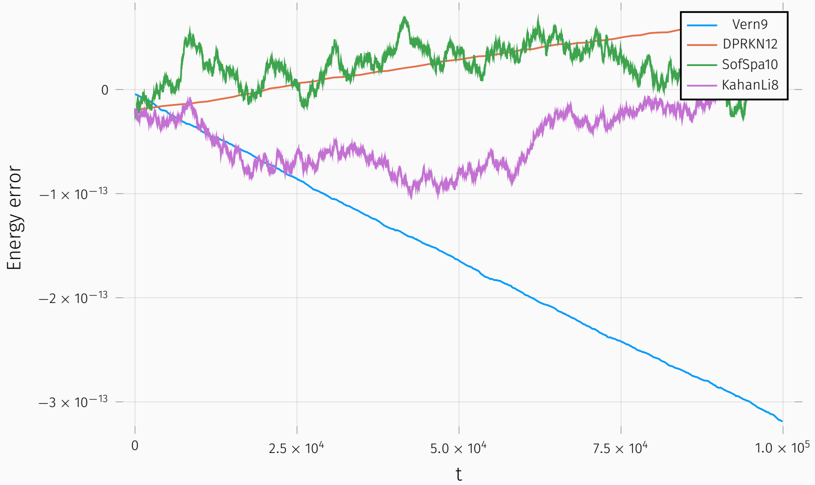

Figure 24: Energy error benchmark for long integration time with rescaling

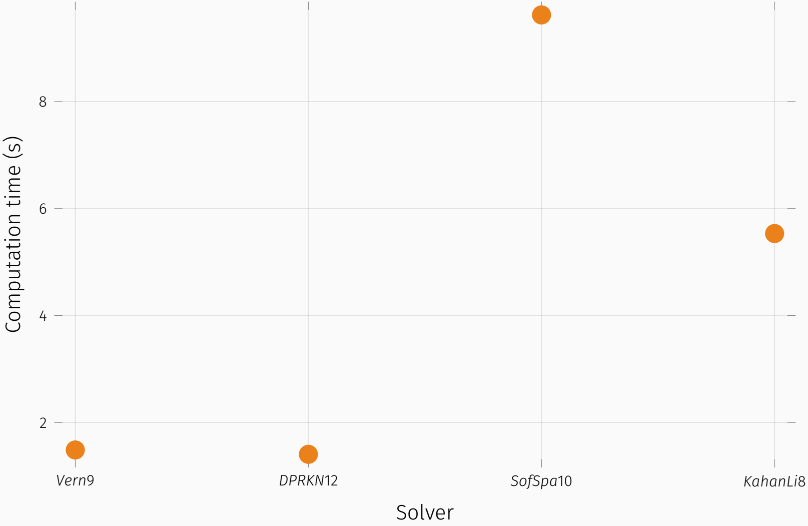

Figure 25: Computational time benchmark for long integration time with rescaling