汽车动力学 2018秋 作业1

肖飞宇 2018210441

1 PROBLEM

某轮胎额定载荷 F_z = 8000N, 在此载荷系数下附着系数为 μ_y = 0.8,侧偏刚度为K = 81000N/rad, 转折系数 E_y = 0.1. 该轮胎半径为R = 0.36m,接地印记长度为l = 0.3m,载荷在印记上的分布为抛物线(P = a + bx^2 , $\frac{1}{2}$ $\leq x \leq \frac{1}{2}$),沿宽度分布为常数。设侧向力-侧偏角的关系为:

$$F_{y} = \mu_{y} \times F_{z} \times \left[1 - e^{-(\psi + E_{y} \times \psi^{3})}\right]$$

其中, $\psi = \frac{K \times tg\partial}{\mu_y \times F_z}$, ∂ 为侧偏角 忽略轮胎侧向变形所产生的附加回正力矩的情况下,求:

- 1. 回正力矩-侧偏角特性的解析解和数值解,并绘制曲线
- 2. 设轮胎的滚动阻力系数为f = 0.01,粗事故垂直压力沿印记方向的分布为 $P = a + bx^2 c\sin(\frac{2\pi x}{c})$,求解此时的回正力矩-侧偏角特性的数值解,并绘制曲线。

1.1 (1)

首先求解垂直力的分布, 首先由

$$\int_{-\frac{l}{2}}^{\frac{l}{2}} P dx = F_z \tag{1.1}$$

得

$$F_z = al + \frac{b}{12}l^3 \tag{1.2}$$

同时, 在边缘处应该为力的零值边界条件, 有

$$P(-\frac{l}{2}) = P(\frac{l}{2}) = a + \frac{b}{4}l^2 = 0$$
 (1.3)

联立方程1.2和方程1.3可得

$$a = \frac{3F_z}{2l}$$
$$b = -\frac{6F_z}{l^3}$$

那么载荷的垂直分布为

$$P = \frac{3F_z}{2I} - \frac{6F_z}{I^3}x^2$$

同时注意到,最大侧向力轮廓为:

$$F_y = \mu_y P = \mu_y \frac{3F_z}{2I} - \mu_y \frac{6F_z}{I^3} x^2$$

为了求解侧向力的分布,假设其为简单的线性分布,即假设其方程为

$$f_{y} = k_{y}(x - \frac{l}{2}) \tag{1.4}$$

接下来就是求出这一分布和最大侧向力轮廓的交点(显然,从这一交点往印记后方,侧向李将会沿着最大侧向力轮廓分布),不妨设交点的横坐标为 x_0 ,有

$$k_{y}(x_{0} - \frac{l}{2}) = \mu_{y} \frac{3F_{z}}{2l} - \mu_{y} \frac{6F_{z}}{l^{3}} x_{0}^{2}$$
(1.5)

容易求得

$$k_y = \mu_y \frac{F_z(3l^2 - 12x_0^2)}{l^3(2x_0 - l)} \tag{1.6}$$

从而可以得到侧向力的分布

$$f_{y}(x) = \begin{cases} k_{y}(x - \frac{1}{2}) & x_{0} \le x \le \frac{1}{2} \\ \mu_{y} \frac{3F_{z}}{2I} - \mu_{y} \frac{6F_{z}}{I^{3}} x^{2} & -\frac{1}{2} \le x \le x_{0} \end{cases}$$

$$(1.7)$$

对侧向力的分布进行积分,可以得到总的侧向力

$$F_{y} = \int_{-\frac{1}{2}}^{\frac{1}{2}} f_{y} dx \tag{1.8}$$

代入求得

$$F_{y} = -\frac{k_{y}}{2}(x_{0} - \frac{l}{2})^{2} + \frac{\mu_{y}F_{z}}{2l^{3}}(l^{3} + 3l^{2}x_{0} - 4x_{0}^{3})$$
(1.9)

而题干中已知

$$F_{y} = \mu_{y} \times F_{z} \times \left[1 - e^{-(\psi + E_{y} \times \psi^{3})}\right]$$

$$(1.10)$$

得到一元三次方程

$$8x_0^3 - 12lx_0^2 + 6l^2x_0 + l^3(8e^{-(\psi + E_y \times \psi^3)} - 1) = 0$$
(1.11)

注意到

$$\psi = \frac{K \times tg\partial}{\mu_y \times F_z} \tag{1.12}$$

结合方程1.11和1.28,可以将 x_0 表示为只和侧偏角 ∂ 相关的函数,不妨记为 $x_0 = \Psi(\partial)$ 需要特别指出的是,之所以此处求出 x_0 和侧偏角 ∂ 之间的关系,是为了后面绘制回正力矩和侧偏角之间的关系。

基于之前求得的侧向力公式1.23,可以求出回正力矩为

$$M_z = \int_{-\frac{l}{2}}^{\frac{l}{2}} f_y dx = \frac{u_y F_z}{2l^3} \left(x_0^4 - 10l x_0^3 - \frac{9}{2} x_0^2 + \frac{5}{2} l^3 x_0 + \frac{17}{16} l^4 \right)$$
 (1.13)

编写**Python**程序(见附录2.1)求解出临界侧向滑移点位置和回正力矩分别和侧偏角的关系如图1.1和图1.2所示

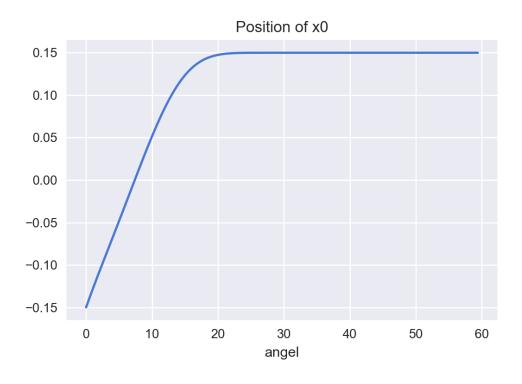


Figure 1.1: 临界侧向滑移点位置和侧偏角关系

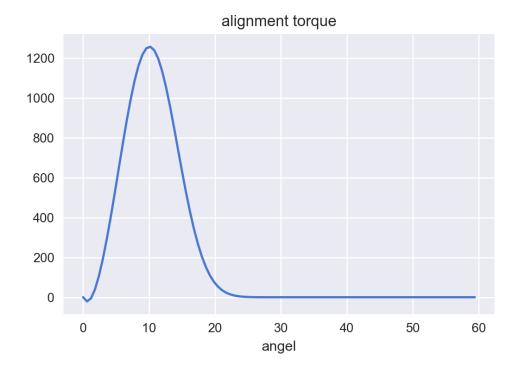


Figure 1.2: 回正力矩和侧偏角关系

1.2 (2)

首先求解垂直力的分布,和(1)类似首先由

$$\int_{-\frac{1}{2}}^{\frac{1}{2}} P dx = F_z \tag{1.14}$$

得

$$F_z = al + \frac{b}{12}l^3 - c\sin(\frac{2\pi x}{l})$$
(1.15)

同时,在边缘处应该为力的零值边界条件,有

$$P(-\frac{l}{2}) = P(\frac{l}{2}) = 0 \tag{1.16}$$

联立方程1.15和方程1.16可得

$$a = \frac{3F_z}{2l}$$
$$b = -\frac{6F_z}{l^3}$$

那么载荷的垂直分布为

$$P = \frac{3F_z}{2l} - \frac{6F_z}{l^3}x^2 - c\sin(\frac{2\pi x}{l})$$

有滚动阻力特性

$$M_{y} = fRF_{z} = \int_{-\frac{1}{2}}^{\frac{1}{2}} Px dx \tag{1.17}$$

将方程1.17代入方程1.2中可以求出

$$c = -\frac{2\pi f R F_z}{I^2} \tag{1.18}$$

那么载荷的垂直分布求得

$$P = \frac{F_z}{2l^3}(3l^2 - 12x^2) + \frac{2\pi f R F_z}{l^2} \sin(\frac{2\pi x}{l})$$
 (1.19)

为了求解侧向力的分布,假设其为简单的线性分布,即假设其方程为

$$f_{y} = k_{y}(x - \frac{l}{2}) \tag{1.20}$$

接下来就是求出这一分布和最大侧向力轮廓的交点(显然,从这一交点往印记后方,侧向 李将会沿着最大侧向力轮廓分布),不妨设交点的横坐标为 x_0 ,有

$$k_{y}(x_{0} - \frac{l}{2}) = \mu_{y} \left(\frac{3F_{z}}{2l} - y \frac{6F_{z}}{l^{3}} x_{0}^{2} + \frac{2\pi f RF_{z}}{l^{2}} sin(\frac{2\pi x}{l}) \right)$$
 (1.21)

容易求得

$$k_{y} = \frac{\mu_{y}}{x_{0} + l/2} \left(\frac{F_{z}}{2l^{3}} (3l^{2} - 12x_{0}^{2}) + \frac{2\pi f R F_{z}}{l^{2}} sin(\frac{2\pi x}{l}) \right)$$
(1.22)

从而可以得到侧向力的分布

$$f_{y}(x) = \begin{cases} k_{y}(x - \frac{1}{2}) & x_{0} \le x \le \frac{1}{2} \\ \mu_{y}P & -\frac{1}{2} \le x \le x_{0} \end{cases}$$
 (1.23)

对侧向力的分布进行积分,可以得到总的侧向力

$$F_{y} = \int_{-\frac{1}{2}}^{\frac{1}{2}} f_{y} dx \tag{1.24}$$

代入求得

$$F_{y} = -\frac{k_{y}}{2}(x_{0} - \frac{l}{2})^{2} + \frac{\mu_{y}F_{z}}{2l^{3}}(l^{3} + 3l^{2}x_{0} - 4x_{0}^{3})$$

$$-\frac{\mu_{y}fRF_{z}}{l}\left(1 + \cos(\frac{2\pi x_{0}}{l})\right)$$

$$-\frac{F_{z}\mu_{y}\pi fRl}{2l^{3}}(2x_{0} - l)sin(\frac{2\pi x_{0}}{l})$$
(1.25)

而题干中已知

$$F_y = \mu_y \times F_z \times \left[1 - e^{-(\psi + E_y \times \psi^3)} \right]$$
 (1.26)

得到方程

$$8x_0^3 - 12lx_0^2 + 6l^2x_0 + l^3(8e^{-(\psi + E_y \times \psi^3)} - 1)$$

$$-8fRl^2 \left(1 + \cos(\frac{2\pi x_0}{l}) \right)$$

$$-4\pi fRl(2x_0 - l)sin(\frac{2\pi x_0}{l}) = 0$$
(1.27)

注意到

$$\psi = \frac{K \times tg\partial}{\mu_v \times F_z} \tag{1.28}$$

结合方程1.11和1.28,可以将 x_0 表示为只和侧偏角 ∂ 相关的函数,不妨记为 $x_0 = \Psi(\partial)$ 基于之前求得的侧向力公式1.23,可以求出回正力矩为

$$M_{z} = \int_{-\frac{l}{2}}^{\frac{l}{2}} f_{y} dx = \frac{u_{y} F_{z}}{2l^{3}} \left(x_{0}^{4} - 10 l x_{0}^{3} - \frac{9}{2} x_{0}^{2} + \frac{5}{2} l^{3} x_{0} + \frac{17}{16} l^{4} \right)$$

$$+ \frac{\mu_{y} f R F_{z}}{l} \left(-\frac{l}{2} + \frac{l}{2\pi} sin(\frac{2\pi x_{0}}{l}) + x_{0} cos(\frac{2\pi x_{0}}{l}) \right)$$

$$- \frac{\mu_{y} \pi f F_{z} R}{12 l^{2}} (8x_{0}^{2} - 2l x_{0} - l^{2}) sin(\frac{2\pi x_{0}}{l})$$

$$(1.29)$$

编写**Python**程序(见附录2.2),特别地,采用牛顿迭代法进行方程求解,求解出临界侧向滑移点位置和回正力矩分别和侧偏角的关系如图1.3和图1.4所示

2 APPENDIXS

2.1 APPENDIXA

```
# -*- coding: utf-8 -*-
"""

@author: feiyuxiao
"""

from sympy.solvers import solve
from sympy import Symbol
import math
import numpy as np
import matplotlib.pyplot as plt
import seaborn as sns

sns.set(style="darkgrid", palette="muted", color_codes=True)
x=Symbol('x')
```

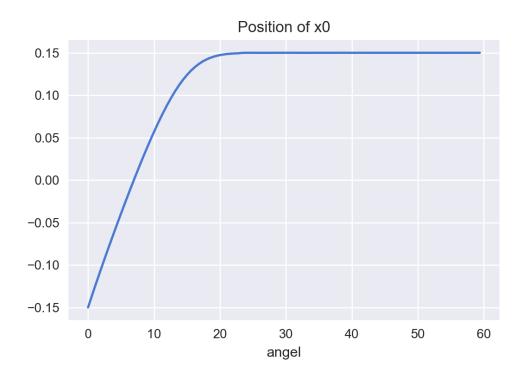


Figure 1.3: 临界侧向滑移点位置和侧偏角关系

```
k = 81000/(0.8*8000)
Ey = 0.1
1 = 0.3
def g(rad):
    theta = k * math.tan(rad)
    return math.exp(-(theta+Ey*theta**3))

size = 100

kk = 0.8*8000/(2*1**3)
def M(x):
    return kk*(x**4 - 10*1*x**3 -4.5*1**2*x**2 + 2.5*1**3*x + 17*1**4/16)

angle = np.zeros(size)
angle_theta = np.zeros(size)
x0 = np.zeros(size)
M0 = np.zeros(size)
```

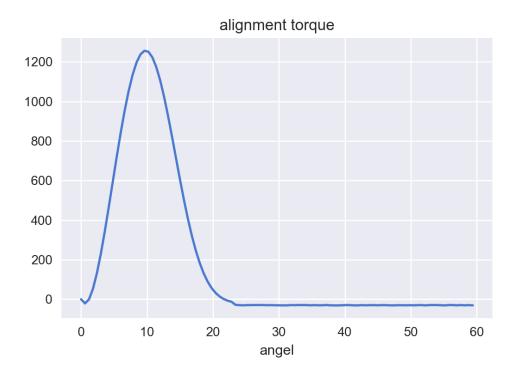


Figure 1.4: 回正力矩和侧偏角关系

```
for i in range(size):
   angle[i] = i*(math.pi/(3*size))
   angle_theta[i] = 60*i/size
for i in range(size):
   f = 8 * x**3 - 12*1*x**2 + 6*1*1*x + 1*1*1*(g(angle[i])-1)
   s=solve(f, x)
   x0[i] = s[0]
   MO[i] = M(xO[i])
plt.figure()
plt.plot(angle_theta,x0)
plt.xlabel("angel")
plt.title("Position of x0")
plt.savefig("1-1.png",dpi=200)
plt.figure()
plt.plot(angle_theta,MO)
plt.xlabel("angel(rad)")
```

```
plt.title("alignment torque")
plt.savefig("1-2.png",dpi=200)
```

2.2 APPENDIXB

```
# -*- coding: utf-8 -*-
@author: feiyuxiao
import sympy
from sympy import Symbol
import math
import numpy as np
import matplotlib.pyplot as plt
import seaborn as sns
import random
pi = 3.1416
sns.set(style="darkgrid", palette="muted", color_codes=True)
x=Symbol('x')
k = 81000/(0.8*8000)
Ey = 0.1
1 = 0.3
def g(rad):
   theta = k * math.tan(rad)
   return math.exp(-(theta+Ey*theta**3))
size = 100
kk1 = 0.8*8000/(2*1**3)
kk2 = 0.8*0.01*0.36*8000/0.3
kk3 = -0.8*pi*0.01*8000*0.36/(12*1*1)
def M(x):
   return kk1*(x**4 - 10*1*x**3 -4.5*1**2*x**2 + 2.5*1**3*x + 17*1**4/16) \
+ kk2*(-0.5*1 + 1*math.sin(2*pi*x/1)/(2*pi)+x*math.cos(2*pi*x/1)) 
+kk3*(8*x**2 -2*1*x -1*1)*math.sin(2*pi*x/1)
```

```
angle = np.zeros(size)
angle_theta = np.zeros(size)
x0 = np.zeros(size)
MO = np.zeros(size)
for i in range(size):
   angle[i] = i*(pi/(3*size))
   angle_theta[i] = 60*i/size
F1 = -8*0.01*0.36*1*1
F2 = -4*pi*0.01*0.36*1
F0 = 2*pi/1
for i in range(size):
   f = 8 * x**3 - 12*1*x**2 + 6*1*1*x + 1*1*1*(8*g(angle[i])-1) \
   + F1*(1+sympy.cos(F0*x)) + F2*(2*x-1)*sympy.sin(F0*x)
   ffunc = sympy.diff(f, x)
   begin = 1
   end = 2
   MAXSTEP = 100
   step_count = 0
   xx0 = random.uniform(begin, end)
   temp = f.subs(x, xx0)
   while step_count < MAXSTEP and abs(temp) > 1e-10:
       xx0 = xx0 - (temp / (ffunc.subs(x, xx0)))
       temp = f.subs(x, xx0)
       step_count += 1
   x0[i] = xx0
   #print(step_count)
   MO[i] = M(xO[i])
plt.figure()
plt.plot(angle_theta,x0)
plt.xlabel("angel")
plt.title("Position of x0")
plt.savefig("2-1.png",dpi=200)
```

```
plt.figure()
plt.plot(angle_theta,M0)
plt.xlabel("angel")
plt.title("alignment torque")
plt.savefig("2-2.png",dpi=200)
```