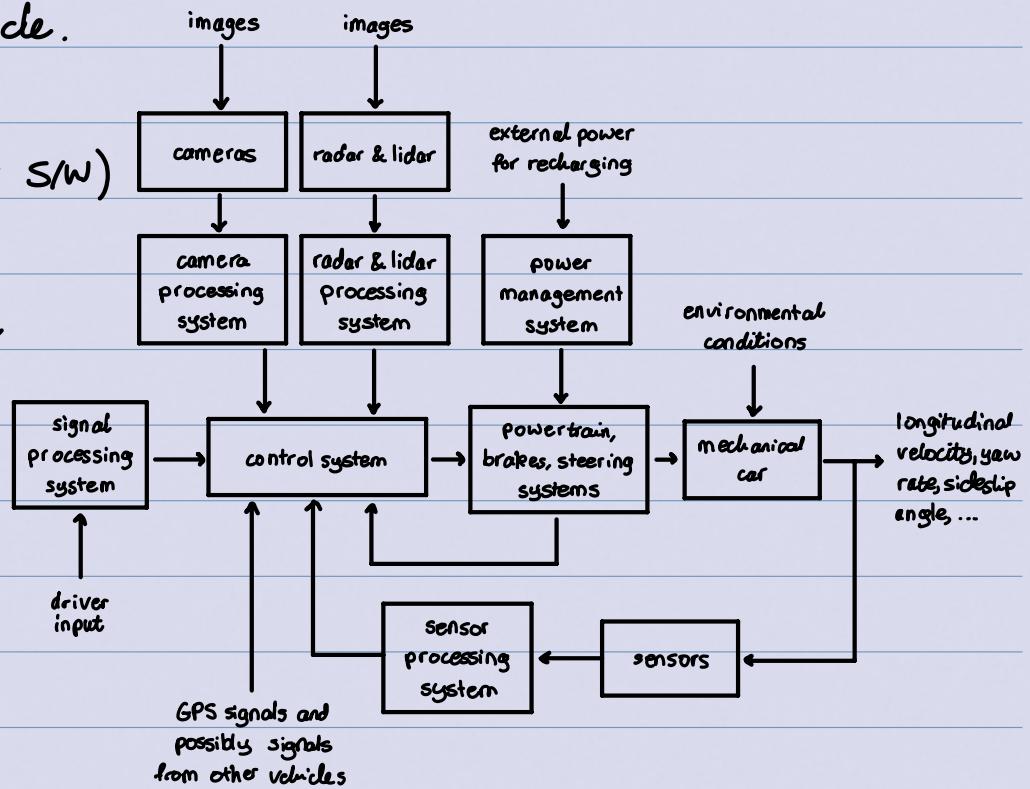


Motivating Example: design of software and hardware used to control an electric vehicle.

boxes = systems (H/W or S/W)

arrows = signals

⇒ this is an example of a block diagram



Can decompose this into 3 types of tasks:

### 1. Modelling

- we need some way to describe how systems process to generate outputs

### 2. Analysis

- we need tools to determine and study the behavior of the various systems
- e.g., is the control system stable? is it fast or slow? How does it respond in windy conditions?

### 3. Design

- we need to have a systematic way to create and tune the various systems (control system, image processing system, etc.)

Signal: a function of one or more independent variables, generally containing information about the behavior of some phenomenon of interest.

⇒ Example: a signal may represent a force, a torque, an angle, a

Speed, a stock price, available SSD memory, etc.

- We will deal with only the situation where there is one independent variable, namely time:

- If time is varying consistently, it's a continuous-time signal.

↳ we denote time by  $t$  and continuous-time signals as  $x(t)$ ,  $u(t)$ ,  $y(t)$ , etc

- If time jumps from one value to the next, it's a discrete-time signal.

↳ we denote time by  $k$  and discrete-time signals as  $x[k]$ ,  $u[k]$ ,  $y[k]$ , etc

**System:** a device, process, or algorithm that takes one or more input signals and generates one or more output signals.

⇒ Example: each of the blocks in the electric vehicle system, a rocket, a heart, a phone, a planet, etc

It's traditional to denote a generic input signal by  $u$  (ie, either  $u(t)$  or  $u[k]$ ) and a generic output signal by  $y$  (ie, either  $y(t)$  or  $y[k]$ )

Systems that have one input signal and one output signal are called single-input single-output (SISO). Systems that have multiple inputs and multiple outputs are called multi-input multi-output (MIMO).

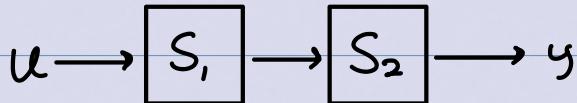
The output of the system is also called the response of the system.

If both the input signal(s) and output signal(s) are continuous-time signals, then we say the system is a continuous-time system.

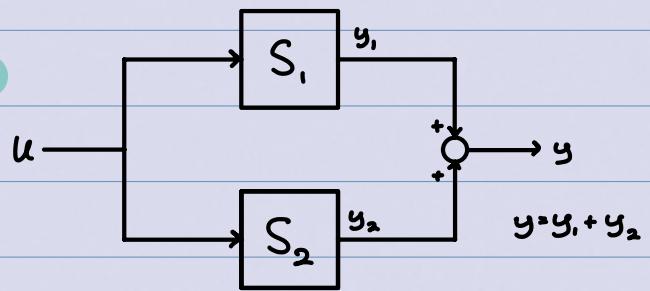
Similarly, if both are discrete-time signals, then we say the system is a discrete-time system. Any other combination results in a hybrid system.

In a block diagram, blocks can be connected in **Series** (aka a **Cascade connection**) or in **parallel** (with the help of a **Summer**):

**Series:**



**parallel:**



**Differential Equation:** any math equation that, in contrast to a purely algebraic equation, includes the derivatives of one or more dependent variable with respect to one or more independent variables.

**Ordinary Differential Equation (ODE):** a differential equation with only one independent variable.

**Partial Differential Equation (PDE):** a differential equation with more than one independent variable.

**Order of a Differential Equation:** the order of the highest derivative in the equation.

⇒ Example: Are the following algebraic, ODEs, or PDEs?

$$\cdot \frac{d^3y}{dt^3} + 4y = \frac{du}{dt} + 2u$$

ODE, 3<sup>rd</sup> order

$$\cdot F = ma$$

algebraic

$$\cdot F = m \frac{d^2y}{dt^2}$$

ODE, 2<sup>nd</sup> order

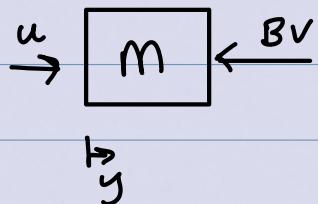
$$\cdot \ddot{y} + 2(2 - \dot{y}^3)\dot{y} + 4y = u$$

ODE, 3<sup>rd</sup> order

$$\cdot \frac{\partial y(x,t)}{\partial t} - K \frac{\partial^2 y^2(x,t)}{\partial^2 x} = u(x,t)$$

PDE, 2<sup>nd</sup> order

**Example:** consider the dynamics of a vehicle moving in a straight line. The system is affected mainly by the force applied by the engine and air resistance (friction). Let  $u$  = input force due to engine and  $v$  = output velocity ( $= \dot{y}$ )

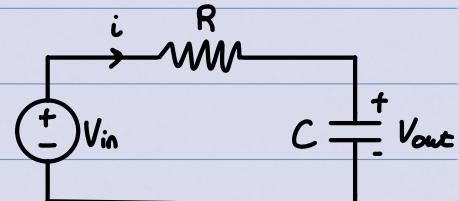


$$F = ma \rightarrow u - Bv = m\ddot{y}$$

$$\rightarrow u - Bv = m\dot{v}$$

$$\rightarrow u = m\dot{v} + Bv \quad \therefore, 1^{\text{st}} \text{ order ODE}$$

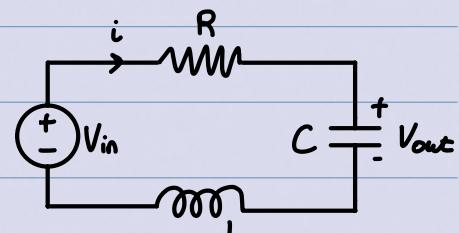
**Example:** consider this RC circuit. We desire to know the dynamic relationship between the output voltage  $V_{\text{out}}$  and the input voltage  $V_{\text{in}}$ . Let  $i$  be the current in the loop.



$$Ri - V_{\text{in}} + V_{\text{out}} = 0, \text{ and } i = C \frac{dV_{\text{out}}}{dt}.$$

$$\hookrightarrow V_{\text{out}} + RC \frac{dV_{\text{out}}}{dt} = V_{\text{in}} \quad \therefore, 1^{\text{st}} \text{ order ODE}$$

**Example:** Same as previous example, but now with an inductor included.



$$Ri - V_{\text{in}} + V_{\text{out}} + L \frac{di}{dt} = 0 \text{ and } i = C \frac{dV_{\text{out}}}{dt}$$

$$\hookrightarrow V_{\text{out}} + RC \frac{dV_{\text{out}}}{dt} + LC \frac{d^2V_{\text{out}}}{dt^2} = V_{\text{in}} \quad \therefore, 2^{\text{nd}} \text{ order ODE}$$

**Static / Memoryless System:** at each time instant, each possible output doesn't depend on any value of the input except perhaps for the input at the same time instant.

↳ else, the system is said to be Dynamic or to have Memory.

Example: a resistor ( $V(t) = i(t)R$ ) is a static system

Example: a capacitor ( $C \frac{dv(t)}{dt} = i(t)$  or  $v(t) = \frac{1}{C} \int_{-\infty}^t i(\tau) d\tau$ ) is a dynamic system.

Causal / Non-Anticipative System: at each time instant, each possible output does not depend on the future values of the input.

↳ else, the system is said to be non-causal or acausal.

Example: the discrete-time system  $y[k] = u[k] + 2u[k] + 3u[k-1]$  is causal.

Example: the discrete-time system  $y[k] = u[k] + 2u[k] + 3u[k+1]$  is noncausal.

Example: the system  $y(t) = K \int_{-\infty}^t u(\tau) d\tau$  is causal

⇒ A static system is causal!

Linear System: satisfies the superposition property, that is, for any input signals  $u_1$  (w/ output  $y_1$ ) and  $u_2$  (w/ output  $y_2$ ) and any constants  $\alpha_1$  and  $\alpha_2$ , a response to the input signal  $u = \alpha_1 u_1 + \alpha_2 u_2$  is  $y = \alpha_1 y_1 + \alpha_2 y_2$ .

↳ Else, the system is said to be nonlinear.

**Homogeneity Property:** for any input signal  $u$ , (w/ output  $y_1$ ) and any constant  $\alpha_1$ , a response to the input signal  $u = \alpha_1 u_1$  is  $y = \alpha_1 y_1$ .

**Additivity Property:** for any input signals  $u_1$ , (w/ output  $y_1$ ) and  $u_2$  (w/ output  $y_2$ ), a response to the input signal  $u = u_1 + u_2$  is  $y = y_1 + y_2$ .

A system satisfies the Superposition property (ie, the system is linear) if and only if it satisfies both the homogeneity property and the additivity property.

Proof:

( $\Rightarrow$ ) If superposition holds:

- set  $\alpha_1 = \alpha_2 = 1$  to conclude additivity holds
- set  $\alpha_2 = 0$  to conclude homogeneity holds

( $\Leftarrow$ ) If both homogeneity and additivity hold, then let  $y_1 = S(u_1)$  and  $y_2 = S(u_2)$ . Then a response to input  $\alpha_1 u_1 + \alpha_2 u_2$  is:  
 $S(\alpha_1 u_1 + \alpha_2 u_2)$   
=  $S(\alpha_1 u_1) + S(\alpha_2 u_2)$  by additivity  
=  $\alpha_1 S(u_1) + \alpha_2 S(u_2)$  by homogeneity  
=  $\alpha_1 y_1 + \alpha_2 y_2$ .

$\therefore$  superposition property is satisfied.

$\Rightarrow$  it's often faster to check for both homogeneity and additivity instead of superposition directly!

Example: are the following systems linear or nonlinear?

a)  $y(t) = Ku(t)$

• Satisfies superposition:

Apply input  $u_1$  to get output  $y_1 = Ku_1$

Apply input  $u_2$  to get output  $y_2 = Ku_2$

Apply input  $a_1u_1 + a_2u_2$  to get output  $K(a_1u_1 + a_2u_2)$   
 $= a_1(Ku_1) + a_2(Ku_2) = a_1y_1 + a_2y_2$

∴, the system is linear.

b)  $y(t) = Ku(t) + 1$

Fails both homogeneity and additivity. Eg:

Apply input  $u_1$  to get output  $y_1 = Ku_1 + 1$

Apply input  $u_2$  to get output  $y_2 = Ku_2 + 1$

Apply input  $u_1 + u_2$  to get output  $K(u_1 + u_2) + 1 \neq y_1 + y_2$ .

∴, the system is nonlinear.

c)  $y(t) = au(t) + bu^2(t)$

- Fails additivity. Eg:

Apply input  $u_1$  to get output  $y_1 = au_1 + bu_1^2$

Apply input  $u_2$  to get output  $y_2 = au_2 + bu_2^2$

Apply input  $u_1 + u_2$  to get output  $a(u_1 + u_2) + b(u_1 + u_2)^2 \neq y_1 + y_2$

∴, the system is nonlinear (if  $b \neq 0$ )

d)  $y(t) = \sin(u(t))$

- Fails homogeneity. Eg:

Apply input  $u_1$  to get output  $y_1 = \sin(u_1)$

Apply input  $a_1u_1$  to get output  $y = \sin(a_1u_1) \neq a_1u_1$

∴, the system is nonlinear

e)  $y(t) = \sin(t)u(t)$

• Satisfies superposition!

Apply input  $u_1$  to get output  $y_1(t) = \sin(t) u_1(t)$

Apply input  $u_2$  to get output  $y_2(t) = \sin(t) u_2(t)$

Apply input  $\alpha_1 u_1 + \alpha_2 u_2$  to get output

$$y(t) = \sin(t)(\alpha_1 u_1(t) + \alpha_2 u_2(t)) = \alpha_1 y_1(t) + \alpha_2 y_2(t).$$

∴ the system is linear!

f)  $y(t) = 5u(t) + u(t)y(t)$

Solving for  $y$ :  $y = 5u + uy \rightarrow y = \frac{5u}{1-u}$

Fails homogeneity:

Apply input  $u_1$  to get output  $y_1 = \frac{5u_1}{1-u_1}$  ( $u_1 \neq 1$ )

Apply input  $\alpha_1 u_1$  to get output  $y = \frac{5\alpha_1 u_1}{1-\alpha_1 u_1} \neq \alpha_1 y_1$

∴ the system is nonlinear!

g)  $y(t) = K \frac{du(t)}{dt}$

• Satisfies superposition:

Apply input  $u_1$  to get output  $y_1 = K \frac{du_1}{dt}$

Apply input  $u_2$  to get output  $y_2 = K \frac{du_2}{dt}$

Apply input  $\alpha_1 u_1 + \alpha_2 u_2$  to get output  $y = K \frac{d}{dt}(\alpha_1 u_1 + \alpha_2 u_2) = \alpha_1 y_1 + \alpha_2 y_2$

∴ the system is linear!

h)  $y(t) = K \int_{-\infty}^t u(\tau) d\tau$

• Satisfies superposition:

Apply input  $u_1$  to get output  $y_1 = K \int_{-\infty}^t u_1(\tau) d\tau$

Apply input  $u_2$  to get output  $y_2 = K \int_{-\infty}^t u_2(\tau) d\tau$

Apply input  $\alpha_1 u_1 + \alpha_2 u_2$  to get output  $y = K \int_{-\infty}^t \alpha_1 u_1(\tau) + \alpha_2 u_2(\tau) d\tau = \alpha_1 y_1 + \alpha_2 y_2$

∴ the system is linear!

$$i) M \frac{d^2y(t)}{dt^2} + B \frac{dy(t)}{dt} + K_s y(t) = u(t) \quad (\text{mass-spring damper system})$$

• Satisfies Superposition

Apply input  $u_1$  to get output  $y_1$ , satisfying  $M \frac{d^2y_1}{dt^2} + B \frac{dy_1}{dt} + K_s y_1 = u_1$  ①

Apply input  $u_2$  to get output  $y_2$  satisfying  $M \frac{d^2y_2}{dt^2} + B \frac{dy_2}{dt} + K_s y_2 = u_2$  ②

Compute  $\alpha_1 \cdot ① + \alpha_2 \cdot ②$ :

$$M \frac{d^2}{dt^2} (\alpha_1 y_1 + \alpha_2 y_2) + B \frac{d}{dt} (\alpha_1 y_1 + \alpha_2 y_2) + K_s (\alpha_1 y_1 + \alpha_2 y_2) = \alpha_1 u_1 + \alpha_2 u_2$$

$\therefore \alpha_1 y_1 + \alpha_2 y_2$  is a solution to the ODE when  $u = \alpha_1 u_1 + \alpha_2 u_2$ .

So, the system is linear!

$$j) \frac{dP(t)}{dt} = \alpha P(t) + b P^2(t) + u(t)$$

• Fails homogeneity.

Apply input  $u_1$  to get output satisfying  $\frac{dP_1}{dt} = \alpha P_1 + b P_1^2 + u_1$  ①

Suppose the system satisfies homogeneity. Then,  $\frac{da_i P_i}{dt} = \alpha a_i P_i + b(a_i P_i)^2 + u_i$   
 $\Rightarrow \frac{dP_i}{dt} = \alpha P_i + b a_i P_i^2 + u_i$  ②

If ① and ② both hold, then so does ① - ②:

$$0 = bP_1^2 - b\alpha_1 P_1^2 \Rightarrow b(\alpha_1 - 1)P_1^2 = 0 \quad ③$$

But ③ doesn't hold in general (except if  $b=0$  or  $\alpha_1=1$  or  $P_1=0$  which are not of interest).

$\therefore$  the system does not satisfy homogeneity by contradiction, and therefore, the system is nonlinear!

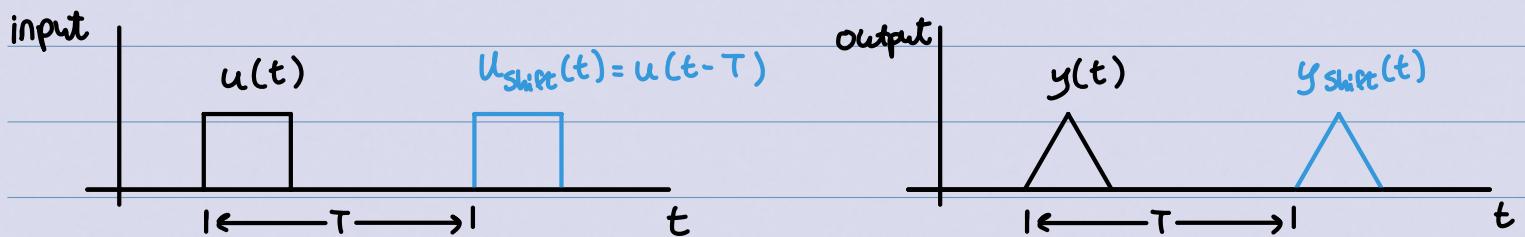
## Time-Invariant Systems vs Time-Varying Systems

A system is said to be time-invariant if a time shift in the input signal always causes the same time shift (but no other distortion) in the output signal. More formally:

Assume input  $u$  is applied to the system with an associated output  $y$ . For a constant  $T$  (with  $-\infty < T < \infty$ ), let  $u_{\text{shift}}(t) = u(t-T)$  denote the shifted output. The system is said to be time-invariant if, for all  $u$  and all  $T$ , there exists an output associated with the input  $u_{\text{shift}}$ , denoted  $y_{\text{shift}}$ , such that:

$$y_{\text{shift}}(t) = y(t-T) \text{ for } -\infty < T < \infty.$$

A system that is not time-invariant is said to be time-varying.



↳ system is time-invariant if  $y_{\text{shift}}(T) = y(t-T)$ .

Example: indicate if the following systems are time-invariant or time-varying:

a)  $y(t) = Ku(t)$

Apply input  $u$  to get output  $y(t) = Ku(t)$

Apply input  $u_{\text{shift}}$  to get output  $y_{\text{shift}}(t) = Ku_{\text{shift}}(t)$

$\therefore$  the system is time-invariant!  $= Ku(t-T) = y(t-T)$

b)  $y(t) = Ku(t) + 1$

Apply input  $u$  to get output  $y(t) = Ku(t) + 1$

Apply input  $u_{\text{shift}}$  to get output  $y_{\text{shift}}(t) = Ku_{\text{shift}}(t) + 1$

$\therefore$  the system is time-invariant!  $= K(u(t-T)) + 1 = y(t-T)$

c)  $y(t) = Ku(t) + t$

Apply input  $u$  to get output  $y(t) = Ku(t) + t$

Apply input  $u_{\text{shift}}$  to get output  $y_{\text{shift}}(t) = Ku_{\text{shift}}(t) + t$

$\therefore$  the system is time-varying.

$$= Ku(t-T) + t \neq Ku(t-T) + t - T = y(t-T)$$

d)  $y(t) = \sin(t)u(t)$

Apply input  $u$  to get output  $y(t) = \sin(t)u(t)$

Apply input  $u_{shift}$  to get output  $y_{shift}(t) = \sin(t)u_{shift}(t)$

$\therefore$  the system is time-varying.

$$\begin{aligned} &= \sin(t)u(t-T) \neq \sin(t-T)u(t-T) \\ &= y(t-T) \end{aligned}$$

e)  $y(t) = \sin(u(t))$

Apply input  $u$  to get output  $y(t) = \sin(u(t))$

Apply input  $u_{shift}$  to get output  $y_{shift}(t) = \sin(u_{shift}(t))$

$\therefore$  the system is time-invariant.

$$= \sin(u(t-T)) = y(t-T).$$

f)  $y(t) = K \frac{du(t)}{dt}$

Apply input  $u$  to get output  $y(t) = K \frac{du(t)}{dt}$

Apply input  $u_{shift}$  to get output  $y_{shift}(t) = K \frac{du_{shift}(t)}{dt}$

$\therefore$  the system is time-invariant.

$$= K \frac{du(t-T)}{dt} = y(t-T).$$

g)  $y(t) = K \int_{-\infty}^t u(\tau) d\tau$

Apply input  $u$  to get output  $y(t) = K \int_{-\infty}^t u(\tau) d\tau$

Apply input  $u_{shift}$  to get output  $y_{shift}(t) = K \int_{-\infty}^t u_{shift}(\tau) d\tau$

$\therefore$  the system is time-invariant!

$$\begin{aligned} &= K \int_{-\infty}^t u(t-\tau) d\tau = K \int_{-\infty}^t u(\tau) d\tau \\ &= y(t-T) \end{aligned}$$

h)  $M \frac{d^2y(t)}{dt^2} + B \frac{dy(t)}{dt} + K_s y(t) = u(t)$

Apply input  $u$  to get output  $y$  satisfying:

$$M \frac{d^2y(t)}{dt^2} \Big|_t + B \frac{dy}{dt} \Big|_t + K_s y(t) = u(t) \quad (1)$$

Apply input  $u_{shift}(t)$  to get output  $y_{shift}$  satisfying:

$$M \frac{d^2y_{shift}(t)}{dt^2} \Big|_t + B \frac{dy_{shift}}{dt} \Big|_t + K_s y_{shift}(t) = u_{shift}(t) = u(t-T)$$

let  $x = t-T \rightarrow dx = dt$

$$\rightarrow M \frac{d^2 y_{\text{shift}}(t)}{dx^2} \Big|_{x+T} + B \frac{dy_{\text{shift}}}{dx} \Big|_{x+T} + K_S y_{\text{shift}}(t) = u(x) \quad (2)$$

① and ② represent the same ODE with the same input.

Therefore, the outputs match:  $y_{\text{shift}}(x+T) = y(x)$

$$\Leftrightarrow y_{\text{shift}}(t) = y(t-T).$$

∴, the system is time-invariant.

### Definition of a solution to an ODE

A general  $n^{\text{th}}$ -order ODE with input  $u$  and output  $y$  can be written as:  $F(t, y, y', y'', \dots, y^{(n)}, u, u', u'', \dots, u^{(m)}) = 0 \quad (n \geq m)$

For a given input  $u$ , any output  $y$  which satisfies the equation is considered to be a solution to the ODE. More formally:

Assume the input  $u(t)$  is defined on some time interval  $t_0 < t < t_1$ , (possibly with  $t_0 = -\infty$  and/or  $t_1 = \infty$ ) and that the first  $m$  derivatives of  $u$  are defined for  $t_0 < t < t_1$ . Then a function  $y(t) = \phi(t)$  is said to be a solution to the ODE if the first  $n$  derivatives of  $\phi$  exist for  $t_0 < t < t_1$ , and if  $F(t, y, y', y'', \dots, y^{(n)}, u, u', u'', \dots, u^{(m)}) = 0$  for  $t_0 < t < t_1$ .

Example: this is the logistic equation which models the dynamics of population growth:  $\frac{dp}{dt} = \alpha p - bP^2 + u$ . Consider the special case where  $u=0$  and  $b=0$ , resulting in  $\frac{dp}{dt} = \alpha p$ . Show that  $P(t) = 3e^{\alpha t}$  and  $P(t) = 10e^{\alpha t}$  (for  $-\infty < t < \infty$ ) are both solutions:

$$P(t) = 3e^{\alpha t} \rightarrow \frac{dp}{dt} = 3\alpha e^{\alpha t} = \alpha P(t). \quad \therefore \frac{dp}{dt} = \alpha p \quad \checkmark$$

$$P(t) = 10e^{\alpha t} \rightarrow \frac{dp}{dt} = 10\alpha e^{\alpha t} = \alpha P(t). \quad \therefore \frac{dp}{dt} = \alpha p \quad \checkmark$$

$\therefore$ , both  $P(t) = 3e^{at}$  and  $P(t) = 10e^{at}$  are solutions.

↳ This shows that there need not be only one solution to an ODE. Usually, there exists a whole family of solutions, namely  $P(t) = ce^{at}$  (for  $-\infty < t < \infty$ ) for some constant  $c$ . Formally:

It's normally (but not always) true that there exists a family of solutions to  $F(t, y, y', y'', \dots, y^{(n)}, u, u', u'', \dots, u^{(m)}) = 0$ , and the family is parameterised by  $n$  constants  $c_1, c_2, \dots, c_n$ . The family of solutions is said to be the general solution to the ODE. If specific values are chosen for the constants  $c_1, c_2, \dots, c_n$ , we say that the resulting  $y(t)$  is a particular solution to the ODE.

Example: imagine a mass-spring damper system modeled by the ODE  $M\ddot{y} + B\dot{y} + K_S y = u$ . Consider the case where  $u=0$ ,  $M=1$ ,  $B=3$ , and  $K_S=2$ :  $\ddot{y} + 3\dot{y} + 2y = 0$ .

a) verify that  $y(t) = -c_1 e^{-2t} + 2c_1 e^{-t} - c_2 e^{-2t} + c_2 e^{-t}$  (for  $-\infty < t < \infty$ ), parameterised by 2 arbitrary constants  $c_1$  &  $c_2$ , is a family of solutions.

$$y = -c_1 e^{-2t} + 2c_1 e^{-t} - c_2 e^{-2t} + c_2 e^{-t}$$

$$\dot{y} = 2c_1 e^{-2t} - 2c_1 e^{-t} + 2c_2 e^{-2t} - c_2 e^{-t}$$

$$\ddot{y} = -4c_1 e^{-2t} + 2c_1 e^{-t} - 4c_2 e^{-2t} + c_2 e^{-t}$$

$$\therefore \ddot{y} + 3\dot{y} + 2y = 0 \Rightarrow -4c_1 e^{-2t} + 2c_1 e^{-t} - 4c_2 e^{-2t} + c_2 e^{-t} + 3(2c_1 e^{-2t} - 2c_1 e^{-t} + 2c_2 e^{-2t} - c_2 e^{-t})$$

$$+ 2(-C_1 e^{-2t} + 2C_1 e^{-t} - C_2 e^{-2t} + C_2 e^{-t}) = 0$$

$$\rightarrow \underbrace{-4C_1 e^{-2t}}_{-3C_2 e^{-t}} + \underbrace{2C_1 e^{-t}}_{-2C_1 e^{-2t}} - \underbrace{4C_2 e^{-2t}}_{4C_1 e^{-t}} + \underbrace{C_2 e^{-t}}_{-2C_2 e^{-2t}} + \underbrace{6C_1 e^{-2t}}_{2C_2 e^{-t}} - \underbrace{6C_1 e^{-t}}_{6C_2 e^{-2t}} = 0$$

→ everything cancels, get  $0=0$  ✓

b) find a particular solution to  $\ddot{y} + 3\dot{y} + 2y = 0$

arbitrarily,  $C_1 = C_2 = 1$  :  $y(t) = -e^{-2t} + 2e^{-t} - e^{-2t} + e^{-t}$   
 $\hookrightarrow y(t) = -2e^{-2t} + 3e^{-t}$

equally arbitrarily,  $C_1 = C_2 = 0$ :  $y(t) = 0 + 0 + 0 + 0$   
 $\hookrightarrow y(t) = 0$ .

We should get used to the idea that there's a family of solutions to an ODE. However, there are exceptions:

example: consider the system described by the following

nonlinear ODE (for  $-\infty < t < \infty$ ):  $(\frac{dy}{dt})^2 + y^2 = u$

a) if  $u=-1$ , there are no solutions (why?)

the LHS cannot be negative!

b) if  $u=0$ , there is only one solution (why?)

only solution is  $\frac{dy}{dt} = 0$  and  $y = 0$ .

c) if  $u=2$ , there exist multiple solutions (why?)

One possible solution:  $y(t) = \sqrt{2}$ , so  $\frac{dy}{dt} = 0$ .

$$\therefore (\frac{dy}{dt})^2 + y^2 = \sqrt{2}^2 + 0 = 2.$$

Another possible solution:  $y(t) = \sin(t) + \cos(t)$ , so  $\frac{dy}{dt} = \cos(t) - \sin(t)$ .

$$\begin{aligned}\therefore \left(\frac{dy}{dt}\right)^2 + y^2 &= (\cos t - \sin t)^2 + (\sin t + \cos t)^2 \\ &= \cos^2 t - 2\cos t \sin t + \sin^2 t + \sin^2 t + 2\cos t \sin t + \cos^2 t \\ &= 2\sin^2 t + 2\cos^2 t = 2(\sin^2 t + \cos^2 t) = 2\end{aligned}$$

## The Initial-Value-Problem (IVP)

Consider the ODE  $F(t, y, y', y'', \dots, y^{(n)}, u, u', u'', \dots, u^{(m)}) = 0$ . Assume the input  $u$  is given. The initial-value-problem involves finding a solution (if one exists) of the ODE for  $t_0 \leq t \leq t_1$ , (or  $t_0 \leq t \leq t_1$ ), possibly with  $t_1 = \infty$ , subject to the following initial conditions:

$$y(t_0) = y_0, y'(t_0) = y_1, \dots, y^{(n-1)}(t_0) = y_{n-1}$$

(assuming that  $y_0, y_1, \dots, y_{n-1}$  are given constants)

note: initial conditions are all specified at the same time

instant,  $t = t_0$ . If different times are used (eg  $y(0) = 2$  and  $y'(3) = -4$ ), the problem is no longer an IVP, but a boundary-value-problem (BVP). We will not deal with BVPs.

note 2: an IVP usually has a unique solution (with exceptions)

Example: consider again the mass-spring-damper system. Find the particular solution to the following IVP:

Solve  $\ddot{y} + 3\dot{y} + 2y = 0$  subject to  $y(0) = 2$  and  $\dot{y}(0) = 3$

We saw previously that  $y(t) = -C_1 e^{-2t} + 2C_1 e^{-t} - C_2 e^{-2t} + C_2 e^{-t}$  is a family of solutions to  $\ddot{y} + 3\dot{y} + 2y = 0$ . at initial  $t = 0$ :

$$y(0) = -C_1 e^{-2(0)} + 2C_1 e^{-(0)} - C_2 e^{-2(0)} + C_2 e^{-(0)} = -C_1 + 2C_1 - C_2 + C_2 = C_1$$

$$\dot{y}(0) = 2C_1 e^{-2(0)} - 2C_1 e^{-(0)} + 2C_2 e^{-2(0)} - C_2 e^{-(0)} = 2C_1 - 2C_1 + 2C_2 - C_2 = C_2$$

$\therefore y(t) = y(0)(-e^{-2t} + 2e^{-t}) + \dot{y}(0)(-e^{-2t} + e^{-t})$  for  $t \geq 0$  is

another way of writing the family of solutions.

We were given now that  $y(0)=2$  and  $\dot{y}(0)=3$ , so:

$$y(t) = 2(-e^{-2t} + 2e^{-t}) + 3(-e^{-2t} + e^{-t}) \\ = -2e^{-2t} + 4e^{-t} - 3e^{-2t} + 3e^{-t} = -5e^{-2t} + 7e^{-t}$$

$\therefore y(t) = 7e^{-t} - 5e^{-2t}$  is the particular solution to the IVP subject to the given initial conditions.

However, it's not always true that an IVP has a unique solution. Non-uniqueness arises in the following example because of the specific form of nonlinearity.

Example: the IVP  $\ddot{y}(t) - 4t\sqrt{y(t)} = 0$ , subject to  $y(0)=0$  has 2 solutions.

$$1. y(t) = 0 \quad (t \geq 0)$$

$$2. y(t) = t^4 \quad (t \geq 0) \rightarrow 4t^3 - 4t\sqrt{t^4} = 4t^3 - 4t^3 = 0$$

### Factors that affect the solution to an IVP

There are two distinct factors that affect the solution to an IVP: the input signal and the initial conditions.

Example: Consider again the mass-spring damper system, but now with  $u=2$ . As before, we'll have  $M=1$ ,  $B=3$ ,  $K_s=2$ . Using tools we'll later learn, we can show that the solution to the IVP

$\ddot{y} + 3\dot{y} + 2y = 2$  subject to  $y(0)=2$  and  $\dot{y}(0)=3$  is, for  $t \geq 0$ :

$$y(t) = \underbrace{y(0)(-e^{-2t} + 2e^{-t}) + \dot{y}(0)(-e^{-2t} + e^{-t})}_{\text{response due to initial conditions}} + \underbrace{1 - 2e^{-t} + e^{-2t}}_{\text{response due to input signal}} \\ = -4e^{-2t} + 5e^{-t} + 1. \rightarrow \text{same transient but different steady-state response}$$

transient response: response for finite  $t$

steady-state response: response as  $t \rightarrow \infty$

the three approaches that we will consider to solve first-order ODEs are:

1. phase portrait sketch
2. separation of variables
3. exact differential approach

### Approach 1: Phase Portrait Sketch

Simple qualitative approach for determining solutions to an ODE that has the form:

$$\frac{dy}{dt} = f(y, u)$$

for constant input  $u$ . By considering values of  $y$  and  $u$  where  $f(y, u) = 0$ , and the sign of  $f(y, u)$  between those values, we can get a good idea (qualitatively) of how  $y(t)$  behaves, at least for the situation where  $u$  is a constant.

Example: consider the logistic equation with  $u = 0$ ,  $\frac{dP}{dt} = P(a - bP)$  we can determine the sign of the derivative for different values of  $P$ :

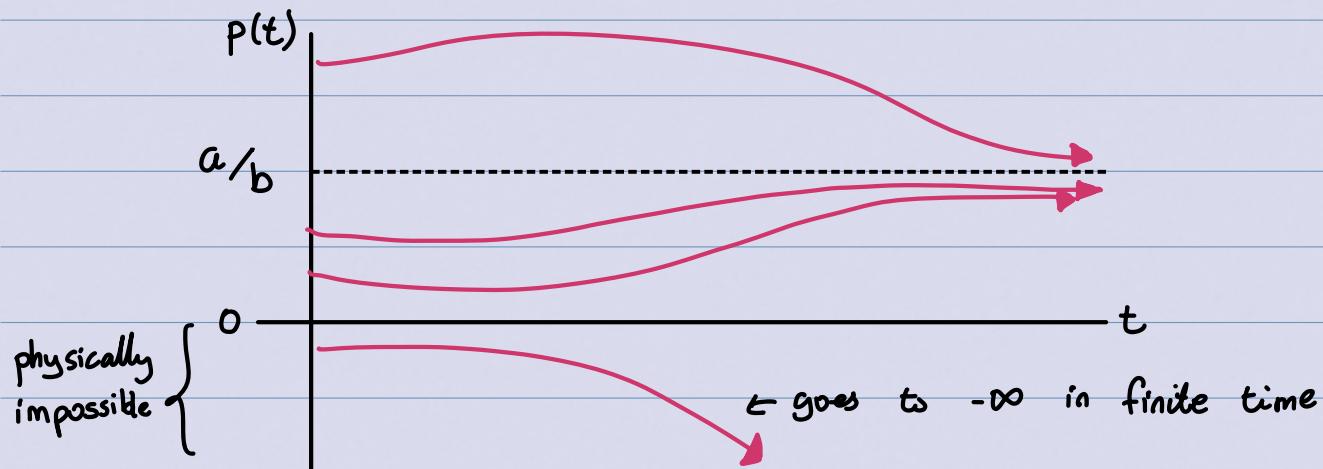
P value	Sign of $dP/dt$
$P < 0$	- $\rightarrow$ this case isn't possible
$P = 0$	0
$0 < P < a/b$	+
$P = a/b$	0
$P > a/b$	-

Using this information, we can sketch on an axis the

phase portrait of the ODE:



So, we can expect the solution  $P(t)$  to look something like:



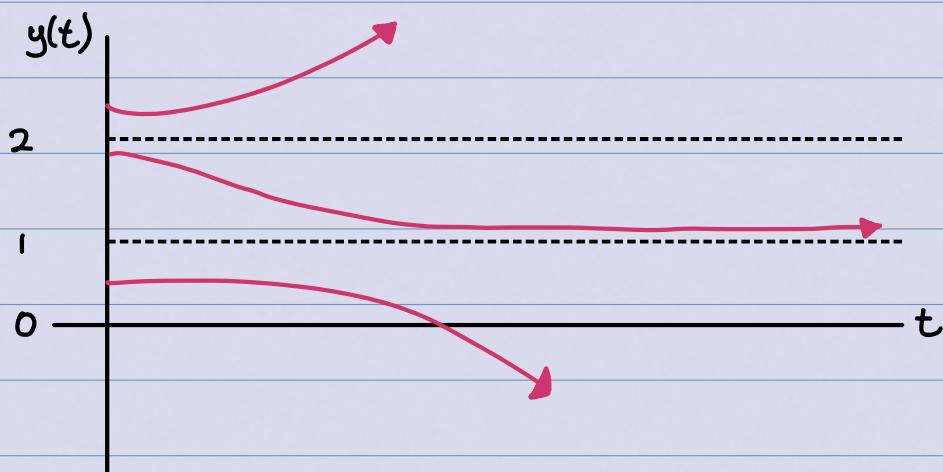
→ we expect it to always converge to  $a/b$  in steady-state

Example: consider the following ODE  $\frac{dy}{dt} = (y-1)^2(y-2)$  w/  $y(0)=y_0$ . Use the phase portrait approach to characterize the solutions.

$y$ value	sign of $dy/dt$
$y < 1$	-
$y = 1$	0
$1 < y < 2$	-
$y = 2$	0
$y > 2$	+



the expected solution for various values of  $y_0$ :



## Approach 2: Separation of Variables

Phase portrait method is simple to apply and usually insightful. However, if we want more detail, we need quantitative analysis.

Separation of Variables works for any first-order ODE that can be written (for nonzero  $u$ ) in the form:  $\frac{dy}{dt} = f(t) \cdot g(y)$ .

An ODE in this form is said to be separable.

↳ for nonzero  $u$ , the structure is:  $\frac{dy}{dt} = f(t, u) \cdot g(y)$  or  $\frac{dy}{dt} = f(t, u) \cdot g(y, u)$

This form is nice since we can rearrange into a form that can often be solved explicitly via integration.

Example: which of the following can be written in separable form?

a)  $\frac{dy}{dt} = e^t y$  yes! already in separate form

b)  $t \frac{dy}{dt} = ye^{y+2t} \Rightarrow \frac{dy}{dt} = \left(\frac{e^{2t}}{t}\right)(ye^y)$  yes!

c)  $\frac{dy}{dt} = y + e^t$  not separable

Example: a) use separation of variables to determine the general solution to  $\frac{dy}{dt} = e^t y$ . b) what is the particular solution associated with the initial condition  $y(0) = 2$ ?

a)  $\frac{dy}{dt} = e^t y$

↓ !! can split up  $\frac{dy}{dt}$   
if  $y \neq 0$

$$\Rightarrow \frac{dy}{y} = e^t dt \Rightarrow \int \frac{1}{y} dy = \int e^t dt \Rightarrow \ln|y| = e^t + C, \Rightarrow |y| = e^{e^t + C}$$

$$\therefore y(t) = C_2 e^{e^t}$$

special case  $y(t) = 0$ : this is a valid solution to the ODE, and the general solution addresses it when  $C_2 = 0$ . so  $C_2$  can be anything!

$$\therefore y(t) = C_2 e^{e^t} \quad (-\infty < t < \infty, -\infty < C_2 < \infty)$$

b)  $y(0) = 2 \Rightarrow C_2 e^{e^0} = 2 \Rightarrow C_2 e^1 = 2 \Rightarrow C_2 = 2e^{-1}$

$$\therefore y(t) = 2e^{e^{t-1}}, \quad t \geq 0.$$

Example: a) consider the logistic equation, simplified with  $b, u = 0$ :  
 $\frac{dP(t)}{dt} = \alpha P(t)$ . Use the separation of variables method to determine the general solution. b) what is the particular solution associated with the initial condition  $P(0) = P_0$ ?

a)  $\frac{dP}{dt} = \alpha P(t)$

$$\Rightarrow \int \frac{1}{P} dP = \int \alpha dt \Rightarrow \ln|P| = \alpha t + C_1 \Rightarrow |P| = e^{\alpha t + C_1} \Rightarrow P(t) = C_2 e^{\alpha t}$$

Special case:  $P(t) = 0 \Rightarrow$  a solution to the ODE and the general

Solution includes this case when  $C_2 = 0$ .

$$\therefore P(t) = C_2 e^{at} \quad (-\infty < t < \infty, C_2 \text{ arbitrary constant})$$

b)  $P(0) = P_0 : P_0 = C_2 e^{a(0)} = C_2$

$$\therefore P(t) = P_0 e^{at}, \quad t \geq 0.$$

### Approach 3: Exact Differential approach

(Review) Schwartz's Theorem: consider a function  $\phi(t, y)$  whose second-order (partial) derivatives  $\frac{\partial^2 \phi}{\partial t^2}, \frac{\partial^2 \phi}{\partial y^2}, \frac{\partial^2 \phi}{\partial t \partial y}, \frac{\partial^2 \phi}{\partial y \partial t}$  all exist and are continuous. Then it's a fact that:

$$\frac{\partial^2 \phi}{\partial t \partial y} = \frac{\partial^2 \phi}{\partial y \partial t}.$$

In addition, the differential:  $d\phi = \frac{\partial \phi}{\partial t} dt + \frac{\partial \phi}{\partial y} dy$  exists.

Example: consider  $\phi(t, y) = t + 2y + e^{ty}$ .

All the partial derivatives exist and are continuous:

$$\frac{\partial \phi}{\partial t} = 1 + ye^{ty}$$

$$\frac{\partial \phi}{\partial y} = 2 + te^{ty}$$

$$\hookrightarrow \frac{\partial^2 \phi}{\partial y \partial t} = e^{ty} + tye^{ty}$$

$$\hookrightarrow \frac{\partial^2 \phi}{\partial t \partial y} = e^{ty} + tye^{ty}$$

$$\therefore \frac{\partial^2 \phi}{\partial y \partial t} = \frac{\partial^2 \phi}{\partial t \partial y} !$$

& the differential of  $\phi(t, y)$  is  $d\phi = \frac{\partial \phi}{\partial t} dt + \frac{\partial \phi}{\partial y} dy = (1 + ye^{ty})dt + (2 + te^{ty})dy$

Consider a first-order ODE of the form (for zero u):  $\frac{dy}{dt} = -\frac{M(t, y)}{N(t, y)}$

and rewrite it as  $M(t, y)dt + N(t, y)dy = 0$ .

If there is a function  $\phi(t, y)$  with continuous second-order partial derivatives whose differentials happen to match the above equation, that is,  $d\phi(t, y) = \frac{\partial \phi}{\partial t} dt + \frac{\partial \phi}{\partial y} dy = M(t, y)dt + N(t, y)dy$ , then we say that  $M(t, y)dt + N(t, y)dy$  is an exact differential.

If  $M(t, y)dt + N(t, y)dy$  is an exact differential, then  $M(t, y)dt + N(t, y)dy = 0$  can be written equivalently as  $d\phi(t, y) = 0$ .

We conclude that the family of solutions is (for arbitrary constant  $c$ ):  $\phi(t, y) = c$ .

Example: Use the exact differential approach to find all solutions to the ODE  $\frac{dy}{dt} = -\left(\frac{1+ye^{ty}}{2+te^{ty}}\right)$

$$(1+ye^{ty})dt + (2+te^{ty})dy = 0$$

from the previous example, we saw that  $(1+ye^{ty})dt + (2+te^{ty})dy$  is an exact differential with  $\phi(t, y) = t + 2y + e^{ty}$ .

∴ general solution of the ODE is  $t + 2y + e^{ty} = c \quad (-\infty < t < \infty)$   
↳ can't solve for  $y \rightarrow$  an "implicit" solution

Theorem: Assume  $M(t, y)$  and  $N(t, y)$  are both continuous with continuous first-order partial derivatives. Then  $M(t, y)dt + N(t, y)dy$  is an exact differential if and only if  $\frac{\partial M(t, y)}{\partial y} = \frac{\partial N(t, y)}{\partial t}$ .

↳ Proof:

( $\Rightarrow$ ) if  $M(t, y)dt + N(t, y)dy$  is an exact differential, then (by definition)  $\exists$  a function  $\phi(t, y)$  st.  $d\phi(t, y) = M(t, y)dt + N(t, y)dy$  with  $\frac{\partial^2 \phi}{\partial t \partial y} = \frac{\partial^2 \phi}{\partial y \partial t}$ .

But,  $\frac{\partial^2 \phi}{\partial t \partial y} = \frac{\partial^2 \phi}{\partial y \partial t}$  is equivalent to  $\frac{\partial M(t, y)}{\partial y} = \frac{\partial N(t, y)}{\partial t}$  Since:

$$\frac{\partial^2 \phi}{\partial t \partial y} = \frac{\partial}{\partial t} \left( \frac{\partial \phi}{\partial y} \right) = \frac{\partial}{\partial t} (N(t, y)) \quad \text{and} \quad \frac{\partial^2 \phi}{\partial y \partial t} = \frac{\partial}{\partial y} \left( \frac{\partial \phi}{\partial t} \right) = \frac{\partial}{\partial y} (M(t, y)).$$

$\therefore \Rightarrow$  proved.

( $\Leftarrow$ ) suppose that  $\frac{\partial M(t, y)}{\partial y} = \frac{\partial N(t, y)}{\partial t}$  holds. we must construct a function  $\phi(t, y)$  such that  $d\phi(t, y) = M(t, y)dt + N(t, y)dy$ . This is hard so it'll be illustrated in the following examples:

Example) which of the following are exact differentials?

a)  $ydt + tdy$   $\frac{\partial(y)}{\partial y} - \frac{\partial(t)}{\partial t} = 1 \Rightarrow$  yes!

b)  $ydt - tdy$   $\frac{\partial(y)}{\partial y}(y) \neq \frac{\partial(-t)}{\partial y}(-t) ! \Rightarrow$  no

c)  $-ysin(t)dt + cos(t)dy$   $\frac{\partial}{\partial y}(-ysint) = \frac{\partial}{\partial t}(cost) = -sint \Rightarrow$  yes!

d)  $t^2dt + (1+y+y^2)dy$   $\frac{\partial}{\partial y}(t^2) = \frac{\partial}{\partial t}(1+y+y^2) = 0 \Rightarrow$  yes!

Example: use the exact differentials approach to find the solution to  $\frac{dy(t)}{dt} + u(t)y(t) = 0$  for the initial condition  $y(1) = 4$  with  $u(t) = \frac{1}{t}$  for  $t \geq 1$ .

$$\frac{dy}{dt} + \frac{1}{t}y = 0 \Rightarrow ydt + tdy = 0, \quad (\text{so } M=y \text{ and } N=t).$$

See that  $\frac{\partial M}{\partial y} = \frac{\partial N}{\partial t} = 1$ , so  $ydt + tdy$  is an exact differential.

Therefore, there exists a  $\phi(t, y)$  such that  $d\phi = ydt + tdy$ :

*don't assume +C, just do +f(t)!!*

$$\text{Step 1: } \frac{\partial \phi}{\partial y} = t \Rightarrow \phi(t, y) = ty + f(t), \text{ arbitrary function } f(t)$$

$$\text{Step 2: } \frac{\partial \phi}{\partial t} = y \Rightarrow \frac{\partial}{\partial t}(ty + f(t)) = y \Rightarrow y + f'(t) = y \Rightarrow f'(t) = 0.$$

Since the derivative of  $f(t) = 0$ ,  $f(t)$  must be a constant, say  $C$ ,  
 $\therefore f(t) = C_1$ .

$$\text{So, } \phi(t, y) = ty + C_1.$$

$\therefore$  the general solution is  $\phi(t, y) = C$ , ie  $ty = C_2$  or  $y(t) = \frac{C_2}{t}$

$$\text{now force } y(1) = 4 \text{ to get } C_2 = 4. \quad \therefore y(t) = \frac{4}{t}, \quad t \geq 1.$$

Example: Use the exact differential approach for  $\frac{dP}{dt} = aP$ .

$$dP - aPdt = 0 \quad (M=1, \quad N=aP).$$

$$\text{but, } \frac{\partial M}{\partial t} = 0 \neq \frac{\partial N}{\partial t} = aP'(t).$$

we can fix this by rearranging to  $\frac{1}{P}dP - adt = 0$ ,  
so  $\frac{\partial M}{\partial t} = 0 = \frac{\partial N}{\partial t}$  ! Now we must find  $\phi(t, y)$  st  $d\phi = \frac{1}{P}dP - adt$ .

$$\text{Step 1: } \frac{\partial \phi}{\partial P} = \frac{1}{P} \Rightarrow \int d\phi = \int \frac{1}{P}dP \Rightarrow \phi = \ln|P| + f(t)$$

Step 2:  $\frac{\partial \phi}{\partial t} = -\alpha \Rightarrow \frac{\partial}{\partial t} (\ln|P| + f(t)) = -\alpha \Rightarrow f'(t) = -\alpha$   
 so,  $f(t) = -\alpha t + C_1$ . P is function of t?

$\therefore \phi = \ln|P| - \alpha t + C_1$ , so, general solution is  $\phi(t, y) = C$ ,  
 ie,  $\ln|P| - \alpha t = C_1$ .

$$\hookrightarrow \text{so, } |P(t)| = e^{\alpha t + C_1} = C_2 e^{\alpha t}. \quad \therefore P(t) = C_3 e^{\alpha t}$$

remove absolute value

now we force  $P(0) = P_0$  to get  $P_0 = C_3 (1) = C_3$ .

Therefore,  $P(t) = P_0 e^{\alpha t}, t \geq 0$ .

Example: use the exact differential approach to find the solution to  $\frac{dy}{dt} = \frac{-ty^2}{2+t^2y}$  with initial condition  $y(0) = 3$ .

rewrite as  $ty^2 dt + (2+t^2y) dy = 0$ . ( $M = ty^2$ ,  $N = 2+t^2y$ )

$\frac{\partial}{\partial y}(ty^2) = 2ty = \frac{\partial}{\partial t}(2+t^2y)$ ! So, we must find  $\phi(t, y)$  such that  $d\phi(t, y) = ty^2 dt + (2+t^2y) dy$

Step 1:  $\frac{\partial \phi}{\partial t} = ty^2 \Rightarrow \int \partial \phi = y^2 \int t dt \Rightarrow \phi(t, y) = \frac{y^2 t^2}{2} + f(y)$

Step 2:  $\frac{\partial \phi}{\partial y} = 2+t^2y \Rightarrow \frac{\partial}{\partial y} (\frac{1}{2}y^2 t^2 + f(y)) = 2+t^2y \Rightarrow t^2y + f'(y) = 2+t^2y$   
 $\Rightarrow f'(y) = 2 \Rightarrow f(y) = 2y + C_1$ .

$\therefore \phi(t, y) = \frac{1}{2}t^2y^2 + 2y + C_1$ .

The family of solutions to the ODE is  $\frac{1}{2}t^2y^2 + 2y = C_2$ .

force  $y(0) = 3$  to get  $6 = C_2$ .

$$\therefore y(t) = \begin{cases} \frac{2}{t^2} \sqrt{1+3t^2} - 1, & t > 0 \\ 3 & t = 0 \end{cases}$$

## Use of Integrating Factors

Sometimes a re-arrangement is needed to make an ODE solvable by the exact differential approach.

Suppose that  $M(t,y)dt + N(t,y)dy = 0$  is NOT an exact differential. We modify this function by multiplying each side by a yet-to-be-determined  $\mu(t,y)$ :  $\mu(t,y)M(t,y)dt + \mu(t,y)N(t,y)dy = 0$

we know this is an exact differential if and only if:

$$\frac{\partial \mu(t,y)M(t,y)}{\partial y} = \frac{\partial \mu(t,y)N(t,y)}{\partial t}$$

$$\text{or: } \mu(t,y) \frac{\partial M(t,y)}{\partial y} + \frac{\partial \mu(t,y)}{\partial y} M(t,y) = \mu(t,y) \frac{\partial N(t,y)}{\partial t} + \frac{\partial \mu(t,y)}{\partial t} N(t,y)$$

This is super messy, but we can solve for  $\mu(t,y)$  under some following special cases:

- Suppose that we require that  $\mu(t,y)$  depends only on  $t$ . then it simplifies to an ODE:

$$\frac{d\mu(t)}{dt} = \Delta_1(t,y)\mu(t) \quad \text{where} \quad \Delta_1(t,y) = \frac{\left(\frac{\partial M(t,y)}{\partial y} - \frac{\partial N(t,y)}{\partial t}\right)}{N(t,y)}$$

if  $\Delta_1(t,y)$  also happens to depend only on  $t$ , then we can use separation of variables to solve for  $\mu(t)$ :  $\mu(t) = C e^{\int \Delta_1(t) dt}$

↳ can use  $C=1$

• or, suppose  $u(t,y)$  depends only on  $y$ . Then the messy equation simplifies to:

$$\frac{\partial u(y)}{\partial y} = \Delta_2(t,y) \text{ where } \Delta_2(t,y) = \frac{\left( \frac{\partial N(t,y)}{\partial t} - \frac{\partial M(t,y)}{\partial y} \right)}{M(t,y)}.$$

if  $\Delta_2(t,y)$  also happens to depend only on  $y$ , we can then use separation of variables to solve for  $u(y)$ :  $u(y) = Ce^{\int_{\Delta_2(y)} dy}$  ↪ can use  $C=1$

Example: use the exact differential approach to find the solution to  $6tydt + (4y + 9t^2)dy = 0$  with  $y(0)=1$ .

$$\text{let } ty = M \quad \& \quad 4y + 9t^2 = N.$$

then  $\frac{\partial M}{\partial t} = y \neq \frac{\partial N}{\partial y} = 4$ . so let's see if the special cases apply:

$$\Delta_1(t,y) = \frac{\left( \frac{\partial M(t,y)}{\partial y} - \frac{\partial N(t,y)}{\partial t} \right)}{N(t,y)} = \frac{6t - 18t}{4y + 9t^2} \times \text{depends on both } t \text{ & } y$$

$$\Delta_2(t,y) = \frac{\left( \frac{\partial N(t,y)}{\partial t} - \frac{\partial M(t,y)}{\partial y} \right)}{M(t,y)} = \frac{18t - 6t}{6ty} = \frac{12t}{6ty} = \frac{2}{y} \checkmark \text{depends only on } y!$$

so we can use the following integrating factor:  $u(y) = e^{\int_{\Delta_2(y)} dy}$

$$u(y) = e^{\int^y \Delta_2(y') dy'} = e^{2 \ln|y|} = e^{\ln y^2} = y^2.$$

the modified ODE is:  $6ty^3 dt + (4y^3 + 9t^2 y^2) dy = 0$

$$\hookrightarrow \frac{\partial}{\partial y}(6ty^3) = 18ty^2 = \frac{\partial}{\partial t}(4y^3 + 9t^2 y^2) \checkmark$$

$$\text{Step 1: } \frac{\partial \phi}{\partial t} = 6ty^3 \Rightarrow \phi(t, y) = 6y^3 \int t dt \Rightarrow \phi(t, y) = 3t^2 y^3 + f(y)$$

$$\text{Step 2: } \frac{\partial \phi}{\partial y} = 4y^3 + 9t^2 y^2 \Rightarrow \frac{\partial}{\partial y} (3t^2 y^3 + f(y)) = 4y^3 + 9t^2 y^2$$

$$\Rightarrow 9t^2 y^2 + f'(y) = 4y^3 + 9t^2 y^2 \Rightarrow f'(y) = 4y^3 \Rightarrow f(y) = y^4 + C_1$$

$$\therefore \phi(t, y) = 3t^2 y^3 + y^4 + C_1.$$

Solution is  $3t^2 y^3 + y^4 = C_2$ . Apply  $y(0)=1$  to get  $C_2=1$ .

$$\therefore 3t^2 y^3 + y^4 = 1.$$

### Solving Constant-Coefficient Linear ODEs

If we specialize to the case where the ODE is linear and time-invariant, there are 3 methods that can solve any linear constant-coefficient ODE of any order (we will only do the last method)

A general  $n^{\text{th}}$ -order constant coefficient linear ODE, with input  $u(t)$  and output  $y(t)$  has the following form (for  $a_n \neq 0$ ):

$$a_n \frac{d^n y}{dt^n} + a_{n-1} \frac{d^{n-1} y}{dt^{n-1}} + \dots + a_1 \frac{dy}{dt} + a_0 y = b_m \frac{d^m u}{dt^m} + b_{m-1} \frac{d^{m-1} u}{dt^{m-1}} + \dots + b_1 \frac{du}{dt} + b_0 u$$

We're usually interested in the associated IVP for time  $t \geq 0$  with the initial conditions:  $y(0) = y_0, y'(0) = y_1, y''(0) = y_2, \dots, y^{(n-1)}(0) = y_{n-1}$ .

Notes:

- we only deal with situations where  $n \geq m$
- there is no loss of generality in stating the initial conditions

at time  $t_0 = 0$  (because it's time-invariant!)

The presence of initial conditions somewhat complicates the type of linearity and time-invariance analysis that we did previously.

**Linearity:** the IVP is linear if and only if the initial conditions  $y_0, \dots, y_m$  are all zero. However, we will now allow for the possibility of nonzero initial conditions throughout.

**Time-invariance:** the presence of initial conditions slightly affects how we analyse time-invariance - the time at which the ICs are assessed must be shifted in the same manner that the input signal is shifted.

**Theorem:** Consider the IVP above. There exists a solution defined for all  $t \geq 0$ , and the solution is unique!

Also, for constant-coefficient linear ODEs, the system can be decomposed into the sum of two responses, one due to the initial conditions and the other due to the input signal:

$$y(t) = (\text{response due to initial conditions}) + (\text{response due to input signal})$$

## Laplace Transform Method

**Fundamental Idea:** ODEs and IVPs are challenging because of the presence of derivatives, so we seek a method of transforming the original time-domain IVP into a new domain where the IVP is represented purely by algebraic equations. We call the new domain the Laplace domain or the s-domain. ∴, calculations with derivatives now become routine algebraic calculations!

## Major Benefits of the Laplace Method

• Systematic and Straightforward!

- Can readily handle inputs that are discontinuous or that contain impulses
- Clearly distinguishes between the effect of the initial conditions and that of the input signal.

Example: Use the Laplace Transform Methods to find the solution to:

Solve  $\ddot{y} + 3\dot{y} + 2y = \alpha$  subject to  $y(0) = y_0$  and  $\dot{y}(0) = 0$ .

Step 1: apply the Laplace transform of each side of the ODE to map the problem to the S-domain:

$$(s^2 Y(s) - sy_0) + 3(sY(s) - y_0) + 2Y(s) = \frac{\alpha}{s}$$

↳  $Y(s)$  denotes the Laplace transform of  $y(t)$ .

Step 2: solve algebraically for  $Y(s)$ :

$$\begin{aligned} Y(s) &= \left[ \frac{1}{s(s^2 + 3s + 2)} \right] \alpha + \left[ \frac{s+3}{s^2 + 3s + 2} \right] y_0 \\ &= \left[ \frac{0.5}{s} - \frac{1}{s+1} + \frac{0.5}{s+2} \right] \alpha + \left[ \frac{2}{s+1} - \frac{1}{s+2} \right] y_0 \end{aligned}$$

Step 3: apply the inverse Laplace transform to map back to time:

$$y(t) = (0.5 - e^{-t} + 0.5e^{-2t}) \alpha + (2e^{-t} - e^{-2t}) y_0 \quad \text{for } t \geq 0.$$

The Laplace transform of a signal  $f(t)$  is denoted equivalently by  $\mathcal{L}\{f(t)\}$ , or  $\mathcal{L}\{f\}$ , or  $F(s)$

The inverse Laplace transform of  $F(s)$  recovers  $f(t)$ . We write

$$\mathcal{L}^{-1}\{F(s)\}$$

time domain

$$\xrightarrow{\text{Laplace Transform}} \quad F(s) = \mathcal{L}\{f(t)\}$$
  

$$\xleftarrow{\text{inverse Laplace transform}} \quad f(t) = \mathcal{L}^{-1}\{F(s)\}$$

Laplace domain

the Laplace transform of a signal  $f(t)$  defined for  $t \geq 0$  is the following function of  $s \in \mathbb{C}$ :

$$\mathcal{L}\{f(t)\} = \int_0^\infty f(t) e^{-st} dt \quad (= \lim_{\tau \rightarrow \infty} \int_0^\tau f(t) e^{-st} dt)$$

**Region of Convergence (ROC)** of the Laplace transform: the set of all  $s$  values for which the above integral converges  
 ↳ Depending on the signal  $f(t)$ , the ROC may be the empty set, the entire complex plane, or something in between.

Example: the unit step (aka Heaviside function) is defined to be the signal:  $U_{\text{step}}(t) = \begin{cases} 0, & t < 0 \\ 1, & t \geq 0 \end{cases}$ . Determine its Laplace transform.

$$\mathcal{L}\{U_{\text{step}}(t)\} = \int_0^\infty U_{\text{step}}(t) e^{-st} dt$$



$$= \lim_{\tau \rightarrow \infty} \int_0^\tau 1 \cdot e^{-st} dt = \lim_{\tau \rightarrow \infty} -\frac{1}{s} e^{-st} \Big|_0^\tau$$

$$= \lim_{\tau \rightarrow \infty} -\frac{1}{s} (e^{-s\tau} - 1) = \frac{1}{s} - \frac{1}{s} \lim_{\tau \rightarrow \infty} e^{-s\tau}$$

Analysing  $\lim_{T \rightarrow \infty} e^{-sT}$ : let  $s = \sigma + j\omega$ .

$$\therefore \lim_{T \rightarrow \infty} e^{-sT} = \lim_{T \rightarrow \infty} e^{-(\sigma+j\omega)T} = \lim_{T \rightarrow \infty} e^{-\sigma T} e^{-j\omega T}.$$

If  $\sigma < 0$ :  $|e^{-\sigma T} e^{-j\omega t}| = e^{-\sigma T} \rightarrow \infty \therefore$  limit DNE

If  $\sigma = 0$ :  $e^{-\sigma T} e^{-j\omega t} = e^{-j\omega t}$  a complex # that rotates as  $T \uparrow \therefore$  limit DNE

If  $\sigma > 0$ :  $|e^{-\sigma T} e^{-j\omega t}| = e^{-\sigma T} \rightarrow 0$  as  $T \rightarrow \infty \therefore$  limit converges

so,  $\mathcal{L}\{u_{\text{step}}(t)\} = \frac{1}{s}$  and the ROC is the half-plane  $\{s : \text{Re}(s) > 0\}$ .

Example: determine the Laplace transform of the unit ramp, defined to be

$u_{\text{ramp}}(t) = \begin{cases} 0, & t < 0 \\ t, & t \geq 0 \end{cases}$ . Note that  $u_{\text{ramp}}(t) = t u_{\text{step}}(t)$ .

$$\mathcal{L}\{u_{\text{ramp}}(t)\} = \int_{0^-}^{\infty} u_{\text{ramp}}(t) e^{-st} dt = \int_{0^-}^{\infty} t e^{-st} dt$$



$$= \lim_{T \rightarrow \infty} \int_0^T t e^{-st} dt = \lim_{T \rightarrow \infty} \left( -\frac{te^{-st}}{s} \Big|_0^T + \frac{1}{s} \int_0^T e^{-st} dt \right) \quad \begin{matrix} \leftarrow \text{by integration by} \\ \text{parts!} \end{matrix}$$

$$= \lim_{T \rightarrow \infty} \left( -\frac{te^{-st}}{s} \Big|_0^T - \frac{e^{-st}}{s^2} \Big|_0^T \right) = \lim_{T \rightarrow \infty} \left( -\frac{Te^{-st}}{s} - 0 - \frac{e^{-st}}{s^2} + \frac{1}{s^2} \right)$$

$$= \frac{1}{s^2} - \frac{1}{s^2} \lim_{T \rightarrow \infty} (sT+1)e^{-sT}$$

Analysing  $\lim_{T \rightarrow \infty} (sT+1)e^{-sT}$ : set  $s = \sigma + j\omega$

$$\lim_{T \rightarrow \infty} (sT+1)e^{-sT} = \lim_{T \rightarrow \infty} (\sigma T + j\omega T + 1) e^{-\sigma T} e^{-j\omega T}$$

If  $\sigma < 0$ , the limit DNE since  $|(\sigma T + j\omega T + 1) e^{-\sigma T} e^{-j\omega T}| = |\sigma T + j\omega T + 1| e^{-\sigma T} \geq e^{-\sigma T} \rightarrow \infty$  as  $T \rightarrow \infty$

If  $\sigma = 0$ , the limit DNE since  $|(\sigma T + j\omega T + 1) e^{-\sigma T} e^{-j\omega T}| = |j\omega T + 1| \rightarrow \infty$  as  $T \rightarrow \infty$

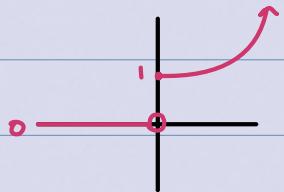
If  $\sigma > 0$ , the limit is 0 since  $|(\sigma T + j\omega T + 1) e^{-\sigma T} e^{-j\omega T}| = |\sigma T + j\omega T + 1| e^{-\sigma T}$

$$\leq (|\sigma T + j\omega T + 1|) e^{-\sigma T} = T |\sigma + j\omega| e^{-\sigma T} + e^{-\sigma T} \rightarrow 0 \text{ as } T \rightarrow \infty$$

so,  $\mathcal{L}\{u_{\text{ramp}}(t)\} = \frac{1}{s^2}$  and the ROC is the half-plane  $\{s : \text{Re}(s) > 0\}$

Example: determine the Laplace transform of the exponential function  $f(t) = e^{at} u_{\text{step}}(t)$ . Allow for the possibility that  $a$  is a complex number.

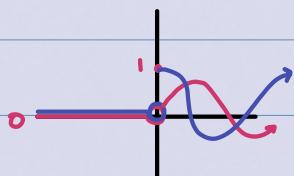
$$\begin{aligned}\mathcal{L}\{f(t)\} &= \int_0^\infty f(t)e^{-st} dt = \lim_{T \rightarrow \infty} \int_0^T e^{(a-s)t} dt \\ &= \lim_{T \rightarrow \infty} \frac{1}{a-s} e^{(a-s)t} \Big|_0^T = \lim_{T \rightarrow \infty} \left( \frac{e^{(a-s)T}}{a-s} - \frac{1}{a-s} \right) \\ &= \frac{1}{s-a} - \frac{1}{s-a} \lim_{T \rightarrow \infty} e^{(a-s)T}\end{aligned}$$



Similar analysis to the unit step example (unit step is identical, but with  $a=0$ )

so,  $\mathcal{L}\{f(t)\} = \frac{1}{s-a}$  and the ROC is  $\{s : \operatorname{Re}(s) > \operatorname{Re}(a)\}$ .

Example: determine the LT of the signals  $f(t) = \sin(at) u_{\text{step}}(t)$  and  $f(t) = \cos(at) u_{\text{step}}(t)$ .



$$\begin{aligned}\mathcal{L}\{\cos(at) u_{\text{step}}(at)\} &= \int_0^\infty \cos(at) e^{-st} dt \\ &= \lim_{T \rightarrow \infty} \int_0^T \cos(at) e^{-st} dt = \lim_{T \rightarrow \infty} \left( \frac{a \sin(at) - s \cos(at)}{s^2 + a^2} e^{-st} \right) \Big|_0^T \\ &= \lim_{T \rightarrow \infty} \left( \frac{a \sin(aT) - s \cos(aT)}{s^2 + a^2} e^{-sT} \right) - \left( \frac{-s}{s^2 + a^2} \right)\end{aligned}$$

$$= \frac{s}{s^2 + a^2} + \frac{1}{s^2 + a^2} \lim_{T \rightarrow \infty} \underbrace{[a \sin(aT) - s \cos(aT)] e^{-sT}}_{\text{let this be } z}$$

$$\text{let } s = \sigma + j\omega. \therefore \lim_{T \rightarrow \infty} z \cdot e^{-(\sigma+j\omega)T} = \lim_{T \rightarrow \infty} z \cdot e^{-\sigma T} e^{-j\omega T}$$

$z \cdot e^{-\sigma T} e^{-j\omega T}$  is a complex number with magnitude  $|z| \cdot e^{-\sigma T}$ .

if  $\sigma < 0$ : limit DNE since  $|z|$  varies with  $T$  but doesn't converge to 0

and  $e^{-\sigma T} \rightarrow \infty$  as  $T \rightarrow \infty$ .

if  $\sigma = 0$ : limit DNE since  $|z|$  varies with  $T$  but doesn't converge to 0  
and  $e^{-\sigma T} = 1$  and  $e^{-j\omega T}$  has magnitude 1 with rotating phase.

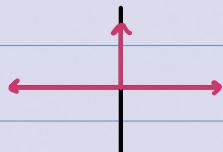
if  $\sigma > 0$ : limit is 0 since  $|z|$  is a bounded function of  $T$  (in fact,  $|z| \leq a + |s|$ ) and  $e^{-\sigma T} \rightarrow 0$  as  $T \rightarrow \infty$ .

So,  $\mathcal{L}\{\cos(at)\} = \frac{s}{s^2 + a^2}$  and the ROC is the half-plane  $\{s: \operatorname{Re}(s) > 0\}$

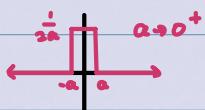
Similarly,  $\mathcal{L}\{\sin(at)\} = \frac{a}{s^2 + a^2}$  and the ROC is also the half-plane  $\{s: \operatorname{Re}(s) > 0\}$

## The Impulse Function $\delta(t)$ :

loosely:  $\delta(t) = \begin{cases} \infty, & t=0 \\ 0, & t \neq 0 \end{cases}$



better:  $\delta_a(t) = \begin{cases} \frac{1}{2a}, & -a \leq t \leq a \\ 0, & \text{otherwise} \end{cases}$



note: the area under the curve of  $\delta_a(t)$  is 1 for all  $a > 0$ .

we rely heavily on the following facts:

1.  $\int_{-\infty}^{\infty} \delta_a(t) dt = 1$

2.  $\int_{-\infty}^{\infty} \delta_a(t) g(t) dt = g(0)$

3.  $\int_{-\infty}^{\infty} \delta(t-t_0) g(t) dt = g(t_0)$

4.  $\int_{-\infty}^{\infty} \delta(t) g(t-t_0) dt = g(-t_0)$

} the sifting property of the impulse function!

Example:  $\int_{-\infty}^{\infty} \delta(t)(t+2) dt = 2$ .

Example:  $\int_{-\infty}^{\infty} \delta(t-5)(t^2) dt = 5^2 = 25$ .

Also,  $\delta(t)$  is <sup>kind of</sup> the derivative of  $U_{\text{step}}(t)$ :  $\frac{d}{dt} U_{\text{step}}(t) = \delta(t)$ .

Example: find the Laplace transform of the impulse function.

We need to handle the lower limit carefully. Let's use  $t = -\epsilon$  (for small  $\epsilon > 0$ ) instead of the vague "0-" notation, and consider only  $\alpha$  small enough that  $0 < \alpha < \epsilon$ .

$$\begin{aligned} \mathcal{L}\{f_\alpha(t)\} &= \int_{-\epsilon}^{\infty} f_\alpha(t) e^{-st} dt = \int_{-\alpha}^{\alpha} \frac{1}{2\alpha} e^{-st} dt \quad \text{Graph: } \begin{array}{c} \text{A rectangular pulse from } t=-\alpha \text{ to } t=\alpha \\ \text{height } \frac{1}{2\alpha} \end{array} \\ &= \left. \frac{1}{2\alpha} \left( -\frac{e^{-st}}{s} \right) \right|_{-\alpha}^{\alpha} = \frac{e^{as} - e^{-as}}{2as} \end{aligned}$$

Let's now take the limit  $\alpha \rightarrow 0^+$  as follows:

$$\begin{aligned} \mathcal{L}\{\delta(t)\} &= \lim_{\alpha \rightarrow 0^+} \mathcal{L}\{f_\alpha(t)\} = \lim_{\alpha \rightarrow 0^+} \frac{e^{as} - e^{-as}}{2as} \stackrel{\text{L'H}}{=} \lim_{\alpha \rightarrow 0^+} \frac{se^{as} + se^{-as}}{2s} \\ &= \frac{2s}{2s} = 1 \quad \text{::} \end{aligned}$$

Alternatively, the result quickly follows from the sifting property of the impulse function:

$$\text{LT of } \delta(t) = \int_{0^-}^{\infty} \delta(t) e^{-st} dt = e^{-st} \Big|_{t=0} = 1 // .$$

$\therefore \mathcal{L}\{\delta_\alpha(t)\} = 1$  and the ROC is the entire S-plane!

## Laplace Transform Pairs Table

### Time Domain

1.  $u_{\text{step}}(t)$

S-Domain,  $F(s) = \mathcal{L}\{f(t)\}$

$\frac{1}{s}$

2. $U_{ramp}(t)$ (ie, $tU_{step}(t)$ )	$\frac{1}{s^2}$
3. $e^{at} U_{step}(t)$	$\frac{1}{s-a}$
4. $\sin(at) U_{step}(t)$	$\frac{a}{s^2+a^2}$
5. $\cos(at) U_{step}(t)$	$\frac{s}{s^2+a^2}$
6. $\delta(t)$	1
7. $t^n U_{step}(t)$ for $n \geq 1$	$\frac{n!}{s^{n+1}}$
8. $\sin^2(at) U_{step}(t)$	$\frac{2a^2}{s(s^2+4a^2)}$
9. $\cos^2(at) U_{step}(t)$	$\frac{s^2+2a^2}{s(s^2+4a^2)}$
10. $t \sin(at) U_{step}(t)$	$\frac{2as}{(s^2+a^2)^2}$
11. $t \cos(at) U_{step}(t)$	$\frac{s^2-a^2}{(s^2+a^2)^2}$
12. $[\sin(at) - at \cos(at)] U_{step}(t)$	$\frac{2a^3}{(s^2+a^2)^2}$
13. $e^{bt} \sin(at) U_{step}(t)$	$\frac{a}{(s-b)^2+a^2}$
14. $e^{bt} \cos(at) U_{step}(t)$	$\frac{s-b}{(s-b)^2+a^2}$
15. $t e^{at} U_{step}(t)$	$\frac{1}{(s-a)^2}$
16. $t^2 e^{at} U_{step}(t)$	$\frac{2}{(s-a)^3}$
17. $t^n e^{at} U_{step}(t)$ for $n \geq 1$	$\frac{n!}{(s-a)^{n+1}}$

memorise  
these!