

# Low-field electron mobility analytical models

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The following analytical models for low-field electron mobility ( $\mu$ ) calculations from different scattering mechanism are based on the Reference-[1, 2] and the references therein.

$$\mu = \left( \frac{m^*}{e} \sum_i \frac{1}{\tau_i} \right)^{-1}$$

where  $\tau_i$ 's are the momentum relaxation rate from different scattering mechanisms presented below.

# 1 Parameters

Parameters name	Symbol	SI Unit
Electron elementary charge	$e$	C
Boltzmann constant	$k_B$	J K <sup>-1</sup>
Temperature	$T$	K
Mass density	$\rho$	kg m <sup>-3</sup>
Lattice parameter	$a_0$	Å
Lattice parameter	$c_0$	Å
Unit cell volume	$\Omega_0$	Å <sup>3</sup>
Bandgap energy	$E_g$	eV
Alloy-disordered scattering potential	$U_0$	eV
Electron effective mass	$m^*$	$m_0$
Electron rest mass	$m_0$	kg
Static dielectric constant	$\varepsilon_s$	$\varepsilon_0$
High frequency dielectric constant	$\varepsilon_h$	$\varepsilon_0$
Vacuum permittivity	$\varepsilon_0$	C V <sup>-1</sup> m <sup>-1</sup>
Longitudinal acoustic phonon velocity	$v_{LA}$	m s <sup>-1</sup>
Transversal acoustic phonon velocity	$v_{TA}$	m s <sup>-1</sup>
Deformation potential	$E_D$	eV
Polar optical phonon energy	$E_{POP}$	eV
Electromechanical coupling coefficient	$K^2$	-
Scattering wave vector	$k$	m <sup>-1</sup>
Fermi wave vector	$k_F$	m <sup>-1</sup>
Thomas-Fermi wave vector	$q_{TF}$	m <sup>-1</sup>
Fang-Howard variational WF parameter	$b$	m <sup>-1</sup>
Fang-Howard WF form factor	$G, F$	-
POP wave vector	$k_0$	m <sup>-1</sup>
2D electron gas density	$n_{2D}$	m <sup>-2</sup>
Interface RMS roughness	$\Delta$	m
Interface roughness correlation length	$L$	m
Dislocation density	$N_{DIS}$	m <sup>-2</sup>
Dislocation occupancy	$f_{DIS}$	-

**Note:** Proper unit conversions for all parameters have been implemented within the equations.

## 2 Scattering mechanisms

$$b = \left[ \frac{33m^*e^2n_{2D}}{8\varepsilon_s\hbar^2} \right]^{\frac{1}{3}} \quad (1)$$

$$k = 2k_F u \quad (2)$$

$$k_F = \sqrt{2\pi n_{2D}} \quad (3)$$

$$q_{TF} = \frac{m^*e^2}{2\pi\varepsilon_s\hbar^2} \quad (4)$$

### 2.1 Structural related scattering

#### 2.1.1 Interface roughness momentum relaxation rate

$$\frac{1}{\tau_{IFR}} = \frac{m^*e^4(\Delta L n_{2D})^2}{8\varepsilon_s^2\hbar^3} \times \int_0^1 \frac{u^4 \exp(-(Lk_F u)^2)}{\left(u + \frac{q_{TF}G(u)}{2k_F}\right)^2 \sqrt{1-u^2}} du \quad (5)$$

$$G(k) = \frac{2\eta(k)^3 + 3\eta(k)^2 + 3\eta(k)}{8} \quad (6)$$

$$\eta(k) = \frac{b}{b+k} \implies \eta(u) = \frac{b}{b+2k_F u} \quad (7)$$

#### 2.1.2 Dislocation momentum relaxation rate

$$\frac{1}{\tau_{DIS}} = \frac{m^*e^4 N_{DIS} f_{DIS}^2}{4\pi k_F^4 c_0^2 \varepsilon_s^2 \hbar^3} \times \int_0^1 \frac{1}{\left(u + \frac{q_{TF}}{2k_F}\right)^2 \sqrt{1-u^2}} du \quad (8)$$

#### 2.1.3 Alloy-disordered momentum relaxation rate

$$\frac{1}{\tau_{AD}} = \frac{m^*\Omega_0 U_0^2 x(1-x)}{\hbar^3} \times \int_{-\infty}^{\infty} \Psi(z)^4 dz = \frac{m^*\Omega_0 U_0^2 x(1-x)}{\hbar^3} \frac{3b}{16} \quad (9)$$

## 2.2 Phonon related scattering

### 2.2.1 Deformation potential momentum relaxation rate

$$\frac{1}{\tau_{DP}} = \frac{6m^*E_D^2k_BT k_F^2b}{2\pi\rho v_{LA}^2\hbar^3} \times \int_0^1 \frac{u^4}{(2k_Fu + q_{TF}F(u))^2\sqrt{1-u^2}} du \quad (10)$$

### 2.2.2 Piezoelectric effect momentum relaxation rate

$$\frac{1}{\tau_{PE}} = \frac{4m^*e^2K^2k_BT k_F}{\pi\varepsilon_s\hbar^3} \times \int_0^1 \frac{u^3F(u)}{(2k_Fu + q_{TF}F(u))^2\sqrt{1-u^2}} du \quad (11)$$

$$F(k) = \eta(k)^3 \quad (12)$$

$$\eta(k) = \frac{b}{b+k} \implies \eta(u) = \frac{b}{b+2k_Fu} \quad (13)$$

### 2.2.3 Acoustic phonon scattering relaxation rate

$$\frac{1}{\tau_{AP}} = \frac{1}{\tau_{DP}} + \frac{1}{\tau_{PE}} \quad (14)$$

### 2.2.4 Polar optical phonon momentum relaxation rate

$$\frac{1}{\tau_{POP}} = \frac{m^*e^2E_{POP}G(k_0)}{2\varepsilon^*k_0\hbar^3} \times \frac{y(1+y-e^{-y})^{-1}}{\exp\left(\frac{E_{POP}}{k_BT}\right) - 1} \quad (15)$$

$$G(k_0) = \frac{2\eta(k_0)^3 + 3\eta(k_0)^2 + 3\eta(k_0)}{8} \quad (16)$$

$$\eta(k_0) = \frac{b}{b+k_0} \quad (17)$$

$$k_0 = \sqrt{\frac{2m^*E_{POP}}{\hbar^2}} \quad (18)$$

$$\frac{1}{\varepsilon^*} = \frac{1}{\varepsilon_h} - \frac{1}{\varepsilon_s} \quad (19)$$

$$y = \frac{\pi\hbar^2n_{2D}}{m^*k_BT} \quad (20)$$

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## References

- [1] Julien Bassaler, Jash Mehta, Idriss Abid, Leszek Konczewicz, Sandrine Juillaguet, Sylvie Contreras, Stéphanie Rennesson, Sebastian Tamariz, Maud Nemoz, Fabrice Semon, Julien Pernot, Farid Medjdoub, Yvon Cordier, and Philippe Ferrandis. Al-rich algan channel high electron mobility transistors on silicon: A relevant approach for high temperature stability of electron mobility. *Advanced Electronic Materials*, page 2400069, 2024.
- [2] JinFeng Zhang, Yue Hao, JinCheng Zhang, and JinYu Ni. The mobility of two-dimensional electron gas in AlGaN/GaN heterostructures with varied al content. *Science in China Series F: Information Sciences*, 51(6):780–789, 2008.