Minicourse on information, complexity and organisation in multidimensional symbolic dynamics

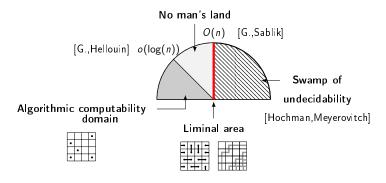
Exact computations of entropy for multidimensional SFT

Silvere Gangloff

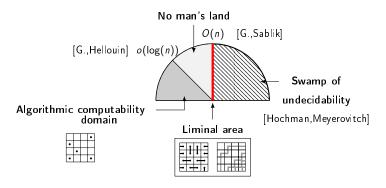
April 30, 2021

sgangloff@agh.edu.pl; silvere.gangloff@gmx.com

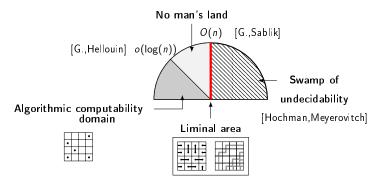
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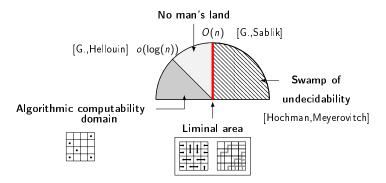


Reminder (third lecture):



In practice, formula for entropy: development of tools for analysis of many variables systems.

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Question: what makes the entropy of subshifts in the liminal area computable ?

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(Called Catalan constant)





w



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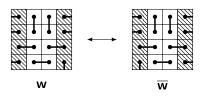


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Examples:





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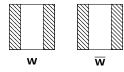
Proof: we have that:

$$\sum_{m} \left(N_n^{\mathbf{w}}(X_0)\right)^2 \leq N_{2n,n}^{c}(X_0) \leq \left(\sum_{m} N_n^{\mathbf{w}}(X_0)\right)^2.$$

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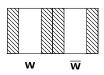
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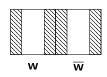
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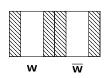


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Thus $h_c(X_0) = h(X_0)$.

In a similar way $h_t(X_0) = h(X_0)$.

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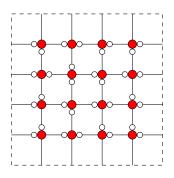
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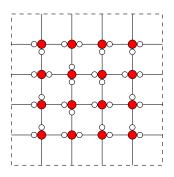


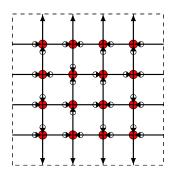
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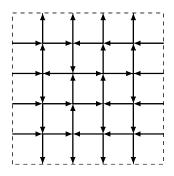
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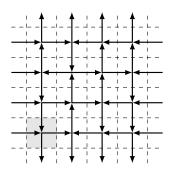
Diagonalisation of $K^{(n)} \to \text{formula for } N_n^t(X_0)$ as sum of trigonometric functions.

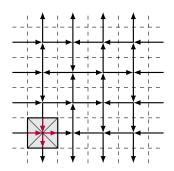




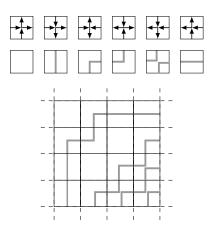




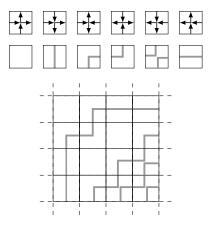




Square ice: Subshift X^s :

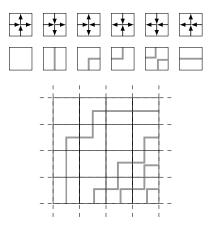


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E.H. Lieb, Residual entropy of square ice, Physical Review, 1967.

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- S. Gangloff, A proof that square ice entropy is $\frac{3}{2}\log_2(4/3)$, 2019 (based on the work of R.Baxter, K.Kozlowski).

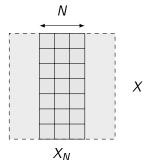
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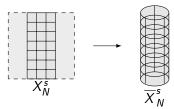
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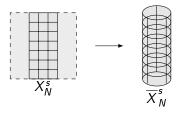
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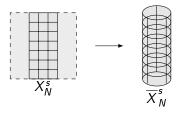
X

$$h(X) = \lim_{N} \frac{h(X_N)}{N}$$

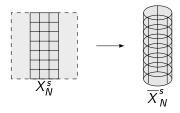


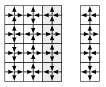


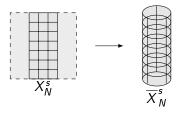


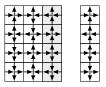


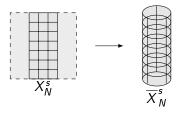


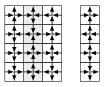


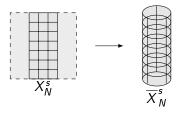


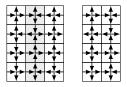


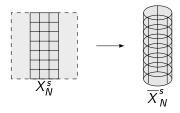


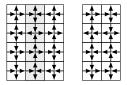


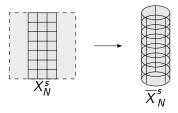


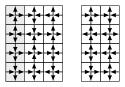


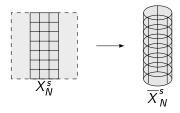


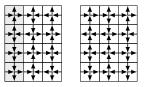


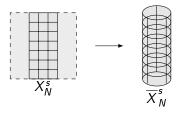


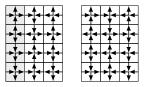


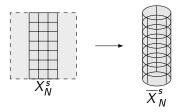


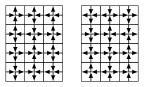


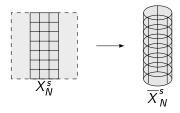


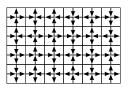


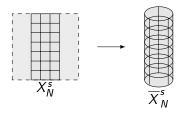




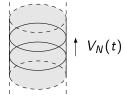


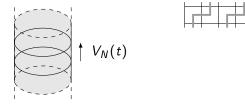


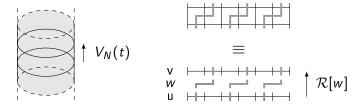


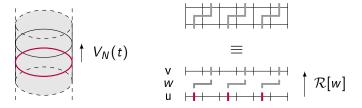


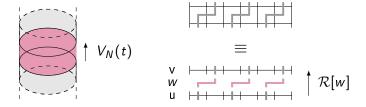
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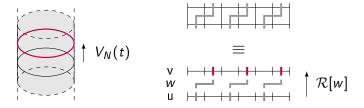


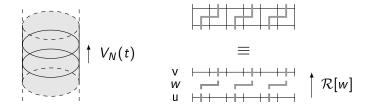








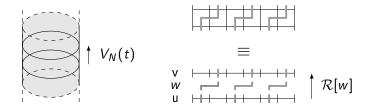




$$V_N(t)[\mathsf{u},\mathsf{v}] = \sum_{\mathsf{u} \mathcal{R}[w]\mathsf{v}} t^{|w|}.$$

where |w| = # of \square and \square

Square ice: Lieb transfer matrices:



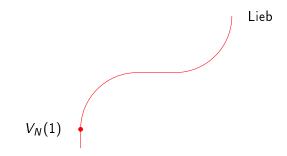
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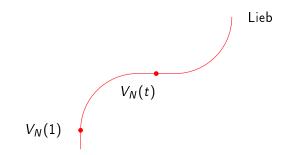
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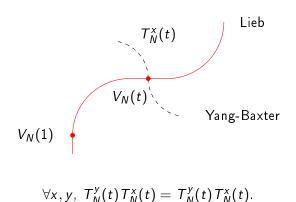
$$h(X^s) = \lim_{N} \frac{\log_2(\lambda_{\max}(V_N(1)))}{N}$$

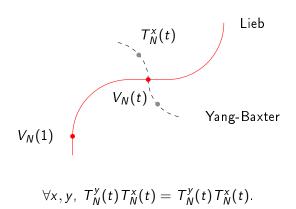
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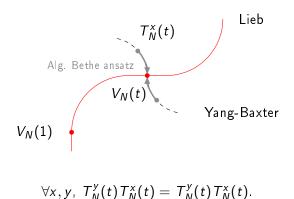
Square ice: Computing maximal eigenvalue of $V_N(1)$, strategy:

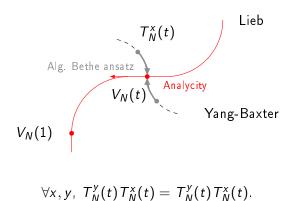












Square ice: R-matrices and monodromy matrices:





$$R(0,1)=\left(\begin{array}{c} \end{array}\right)$$



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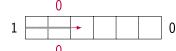
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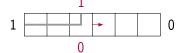


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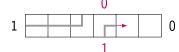


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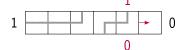


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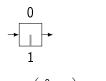




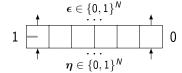
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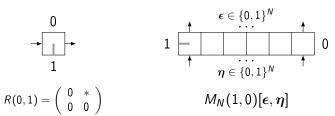


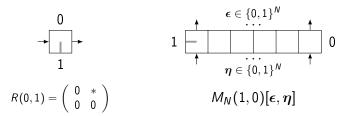
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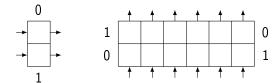


Yang-Baxter transfer matrices:

$$T_N[\epsilon, \eta] = \sum_{u \in \{0,1\}} M_N(u, u)[\epsilon, \eta].$$

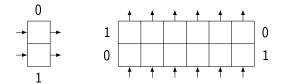
Square ice: Commutation of Yang-Baxter matrices:

Composition of R-matrices and monodromy matrices:

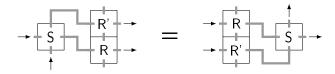


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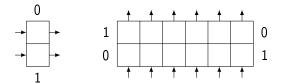


Yang-Baxter equation:

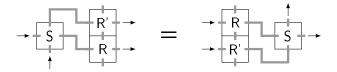


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Composition of R-matrices and monodromy matrices:



Yang-Baxter equation:



Yang-Baxter equation \Rightarrow transfer matrices commute.

Denote $\mu_t \in \left(0, \frac{\pi}{2}\right)$ s.t. $\cos(\mu_t) = \frac{2-t^2}{2}$.

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$$R_{\mu_t}^{x} = \frac{1}{\sin(\mu_t/2)} \left(\begin{array}{cccc} \sin(\mu_t - x) & 0 & 0 & 0 \\ 0 & \sin(x) & \sin(\mu_t) & 0 \\ 0 & \sin(\mu_t) & \sin(x) & 0 \\ 0 & 0 & 0 & \sin(\mu_t - x) \end{array} \right).$$

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where $R_{u_t}^{x}(0,0)$ is the up-left 2×2 part of this matrix, etc.

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If there exists $(p_j(t))_{j=1}^n$ solution of:

$$Np_{j} = 2\pi j - (n+1)\pi - \sum_{k=1}^{n} \Theta_{t}(p_{j}, p_{k})$$
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construction of a candidate eigenvector for the transfer matrix for value:

$$\prod_{k=1}^{n} L_{t}(e^{ip_{k}}) + \prod_{k=1}^{n} M_{t}(e^{ip_{k}}).$$

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Lemma: for all t there exists $(p_j(t))_{j=1}^n$ solution of (E_t) and it is an analytic function in t.

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- 4. Indentification of maximal eigenvalue and eigeinvector around $\sqrt{2}$ (positive coordinates), then on $(0,\sqrt{2})$ by analycity.

Square ice: Asymptotics:

Denote $(\mathbf{p}_j^{(k))})_j$ solutions of equations (E_t) for N_k, n_k .

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Change of variable $(\mathbf{p}_j^{(k))})_j = (\kappa_t(\pmb{lpha}_j^{(k)}))_j$.

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with $f(x) = \log_2(2|\sin(\kappa_t(x))/2|)$.

Lemma:

$$\lim_{N} \frac{\log_2(\lambda_{\max}(V_N(1)))}{N} = \lim_{k} \frac{1}{N_k} \sum_{i=1}^{n_k} f(\alpha_j^{(k)}),$$

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Theorem: there exists ρ_t s.t. for all $f \in L^1$:

$$\lim_{k} \frac{1}{N_k} \sum_{i=1}^{n_k} f(\alpha_j^{(k)}) = \int_{\mathbb{R}} f(\alpha) \rho_t(\alpha) d\alpha.$$

Square ice: Counting functions:

$$\xi_t^{(k)}: \alpha \mapsto \frac{1}{2\pi} \kappa_t(\alpha) + \frac{n_k + 1}{2N_k} + \frac{1}{2\pi N_k} \sum_{j=1}^{n_k} \theta_t(\alpha, \alpha_j^{(k)})$$

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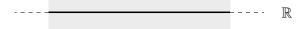
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Equality with a Riemann sum:

$$\lim_{k} \frac{1}{N_{k}} \sum_{j=1}^{n_{k}} f(\alpha_{j}^{(k)}) = \lim_{k} \frac{1}{N_{k}} \sum_{j=1}^{n_{k}} \left(\alpha_{j+1}^{(k)} - \alpha_{j}^{(k)}\right) \frac{\left(\xi_{t}^{(k)}(\alpha_{j+1}^{(k)}) - \xi_{t}^{(k)}(\alpha_{j}^{(k)})\right)}{\left(\alpha_{j+1}^{(k)} - \alpha_{j}^{(k)}\right)} f(\alpha_{j}^{(k)})$$

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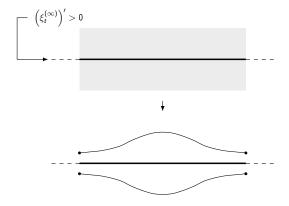
- 2. Assume $((\xi_t)^{\nu(k)})' \to F_t$ on any compact K.
- 3. F_t satisfies an integral equation with unique solution.

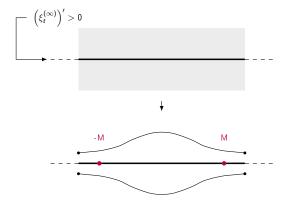
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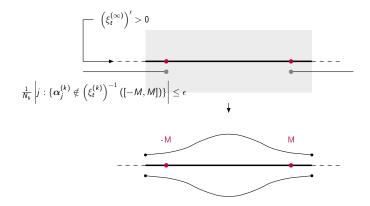
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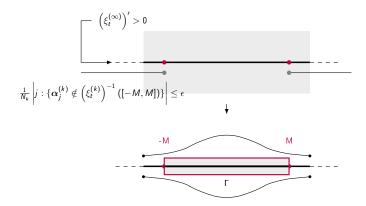
- 2. Assume $((\xi_t)^{\nu(k)})' \to F_t$ on any compact K.
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- 4. Thus, $(\xi_t^{(k)})'$ converges to this solution.

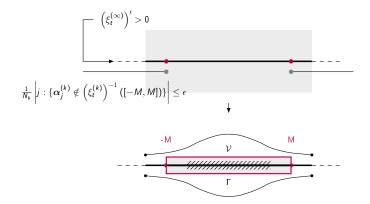


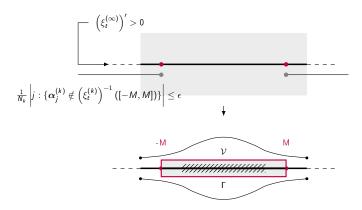






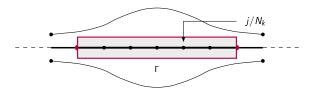




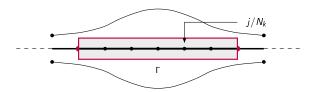


The functions $\xi_t^{(k)}$ have distinct values on $\mathcal V$ and Γ . Thus they are bihilomorphisms onto $\mathcal V$ (Cauchy formula).

Square ice: Lace integral expression of $\xi_t^{(k)}$:



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By residues theorem:

$$\xi_t^{(k)}(\alpha) = \frac{1}{2\pi} \kappa_t(\alpha) + \frac{n_k + 1}{2N_k} + \oint_{\Gamma} \theta_t \left(\left(\xi_t^{(k)} \right)^{-1}(\alpha) \right) \frac{e^{2i\pi s N_k}}{e^{2i\pi s N_k} - 1} ds + O(\epsilon).$$

Square ice: Fredholm integral equation: Limit and change of variable:

$$\xi_t^{(\infty)}(\alpha) = \frac{1}{2\pi} \kappa_t(\alpha) + \frac{1}{4} + \int_0^{+\infty} \theta_t(\alpha) \left(\xi_t^{(\infty)}\right)'(\alpha) d\alpha.$$

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Unique solution by Fourier transforms.

$$\lim_{N} \frac{\log_2(\lambda_{\max}(V_N(1)))}{N} = \lim_{k} \frac{1}{N_k} \sum_{j=1}^{n_k} f(\alpha_j^{(k)}),$$

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$$h(X^s) = \int_{\mathbb{D}} \log_2(2|\sin(\kappa_t(\alpha))/2|).\rho_t(\alpha)d\alpha.$$

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Expression of $ho_t = \left(\xi_t^{(\infty)}\right)'$ and lace integrals computations:

$$h(X^s) = \frac{3}{2}\log_2(4/3).$$

Friedland's theorem:

Theorem[Friedland(1967)]: if the set of forbidden patterns \mathcal{F} is stable by symmetry, $h(X_{\mathcal{F}})$ is a computable number.

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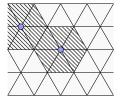
Examples: dimers, square ice, hard squares.

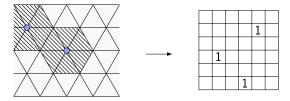
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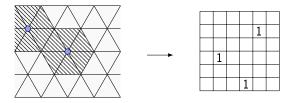
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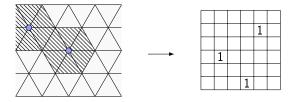
Examples: dimers, square ice, hard squares.

Question: what are the possible values of entropy for symmetric bidimensional subshifts of finite type ?



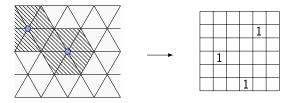




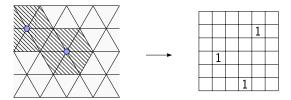


Formula for entropy as sum of a series:

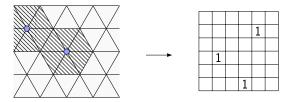
1. Transfer matrices → diagonal transfer matrices;



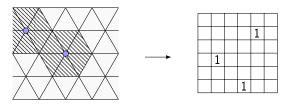
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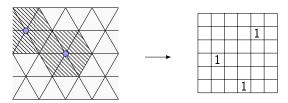
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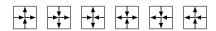
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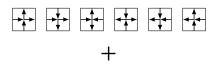


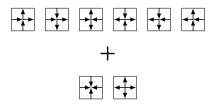
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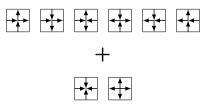
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Main problems: points 3, 4.









Entropy computation: similar to square ice; analytical part non verified.

Subsidiary questions:

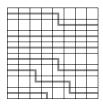
Question: can we use similar methods to talk about invariant measures (for instance $\times 2, \times 3$ conjecture) ?

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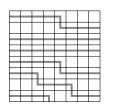
Question: can we use similar methods to talk about invariant measures (for instance $\times 2, \times 3$ conjecture)?

Question: can we find solutions of Yang-Baxter equations for other subshifts of finite type? *Example*: Kari-Culik tilings (know: positive entropy [Durand, Gamard, Grandjean (2017)]).

Definition subshift Δ_r ; ex for r = 3:

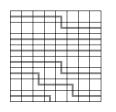


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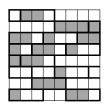


	3	1	2	3				
ĺ					1	2	3	1
	1	2	3	1	2	3	1	2
	2	3	1	2	3	1	2	3
	3	1	2	3				
	2	3			1	2	3	1
Ì			1	2	3			
	1	2	3			1	2	3

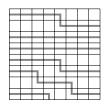
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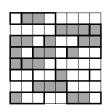
3	1	1	2				
)	1	2	3	_		_	
				1	2	3	1
1	2	3	1	2	3	1	2
2	3	1	2	3	1	2	3
3	1	2	3				
2	3			1	2	3	1
		1	2	3			
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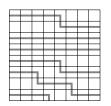


3	1	2	3				
				1	2	3	1
1	2	3	1	2	3	1	2
2	3	1	2	3	1	2	3
3	1	2	3				
2	3			1	2	3	1
		1	2	3			
1	2	3			1	2	3

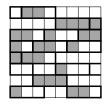


Theorem: [G., Sablik] $h(\Delta_r) = \frac{\log_2(r+1)}{r}$.

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3	1	2	3				
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3	1	2	3				
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1	2	3	_	<u>ა</u>	1	2	3



Theorem: [G., Sablik] $h(\Delta_r) = \frac{\log_2(r+1)}{r}$.

Question: for what kind of subshifts can we compute entropy with similar methods?

Second layer is trivial: we consider only first and third.

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Lower bound:

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Lower bound: $\mathcal{L}_n(\Delta_1)$ contains the following patterns:

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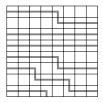
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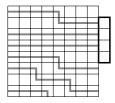
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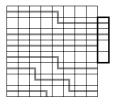


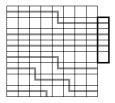


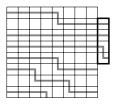
$$N_n(\Delta_1) \geq 2^{n^2}$$
.



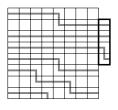






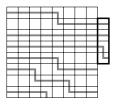


Upper bound: Consider a pattern in $\mathcal{L}_n(\Delta_1)$ for some n:

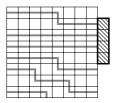


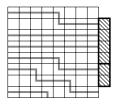
 2^3 choices in 2nd layer

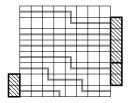
Upper bound: Consider a pattern in $\mathcal{L}_n(\Delta_1)$ for some n:

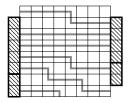


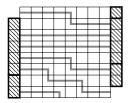
2³ choices in 2nd layer 2⁴ choices in total

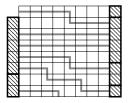


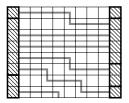


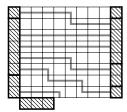


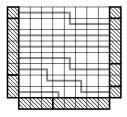


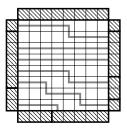


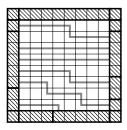




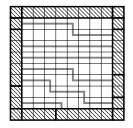






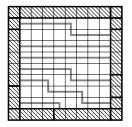


Upper bound: Consider a pattern in $\mathcal{L}_n(\Delta_1)$ for some n:



Number of possible extensions into a pattern of $\mathcal{L}_{n+1}(\Delta_1)$: $\leq 2^{4n}2^{3*3*4} = 2^{4n+4}2^{C}$.

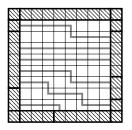
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$$h(\Delta_1)=1=\frac{\log_2(2)}{1}.$$

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- 1. For instance: use Hochman-Meyerovitch's theorem to realize square ice entropy; then simplify the construction without changing the entropy; find a subshift which is isomorphic to square ice.
- 2. Find smaller and smaller subshifts of square ice with same entropy.
- 3. Try Baxter's method for square ice; more precisely: do both transfer matrices of square ice have same maximal eigenvalues or not?