

Transport Equation: $u_t + cu_x = 0 \implies u(x, t) = f(x - ct)$ Solved via method of characteristics: $\frac{dx}{dt} = c \implies x = at + C$ and function is constant along the characteristic lines.

FOL PDE: $a(x, y)u_x + b(x, y)u_y = 0$ or $\langle \nabla u, (a, b) \rangle = 0$. if a, b constant $u(x, y) = f(bx - ay)$ Solved via method of characteristics: $\frac{dy}{dx} = \frac{b}{a} \implies y = \frac{b}{a}x + C$ and function is constant along the characteristic lines. OR by Change of Var for $x' = ax + by, y' = bx - ay$

If non constant a, b then solve for $\frac{dy}{dx} = \frac{b(x, y)}{a(x, y)} \implies y = \text{something}$. Solve for constant and that is that the function is constant along.

Well Posed: 1. Existence 2. Uniqueness 3. Stable/Continuous

Waves: $u_{tt} = c^2 u_{xx} \implies u(x, t) = f(x - ct) + g(x + ct)$. Product of 2 transports. **With IC:** $u(x, 0) = \phi(x), u_t(x, 0) = \psi(x)$. On infinite x: $u(x, t) = \frac{\phi(x+ct) - \phi(x-ct)}{2} - \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds$ aka D'Alembert's Formula.

Energy For Wave: $E(t) = \frac{1}{2} \int_{-\infty}^{\infty} (u_t^2 + c^2 u_x^2) dx$ can make $c^2 = T/\rho$. Energy is conserved. Finite speed of propagation.

Diffusion: $u_t = k u_{xx}$

Max Principle Weak: max on all is on boundary. Strong: max of all can only be on boundary.

Uniqueness Heat is unique proven by either max principle with difference of 2 solutions = 0 or by energy method.

Energy for Heat: $E(t) = \int_{-\infty}^{\infty} u^2 dx$

Stability: $\int |u(x, t) - v(x, t)|^2 dx \leq \int |u(x, 0) - v(x, 0)|^2 dx$

Invariance on Whole Line: 1. Translation, 2. Derivative, 3. Linear Combination, 4. Integral of sol with anything. 5. Scaling by $u(x, t) \rightarrow u(\alpha x, \alpha^2 t)$

Fundamental Solution: $u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy$. for IC $u(x, 0) = \phi(x)$ $S = \frac{1}{\sqrt{4\pi kt}} e^{-\frac{x^2}{4kt}}$ is the fundamental solution. S is a green's function.

Separation of Variables: $u(x, t) = X(x)T(t)$. Create ratio of X and T and set equal to constant. Solve for X and T separately with the BC/IC that are homogenous.

Dirichlet BC: $u(0, t) = u(L, t) = 0$: $X(x) = \sin(\frac{n\pi x}{L})$. Waves: $T(t) = A_n \cos(\frac{n\pi ct}{L}) + B_n \sin(\frac{n\pi ct}{L})$ Diffusion:

$T(t) = A_n e^{-\frac{n^2 \pi^2 kt}{L^2}}$ Laplace: $Y(y) = A_n \cosh(\frac{n\pi y}{L}) + B_n \sinh(\frac{n\pi y}{L})$ Connected to a fourier sine series.

Neumann BC: $u_x(0, t) = u_x(L, t) = 0$: $X(x) = \cos(\frac{n\pi x}{L})$. Waves: $T(t) = A_n \cos(\frac{n\pi ct}{L}) + B_n \sin(\frac{n\pi ct}{L})$ Diffusion:

$T(t) = A_n e^{-\frac{n^2 \pi^2 kt}{L^2}}$ Laplace: $Y(y) = A_n \cosh(\frac{n\pi y}{L}) + B_n \sinh(\frac{n\pi y}{L})$ Connected to a fourier cosine series.

Robin BC: $u_x(0, t) - \alpha_0 u(0, t) = u_x(L, t) + \alpha_1 u(L, t) = 0$: Waves: $T(t) = A_n \cos(\sqrt{\lambda_n} t) + B_n \sin(\sqrt{\lambda_n} t)$ Diffusion: $T(t) = A_n e^{-\lambda_n kt}$. They take the form of a fourier series.

Fourier Sine Series: $f(x) = \sum_{n=1}^{\infty} B_n \sin(\frac{n\pi x}{L}) \implies B_n = \frac{2}{L} \int_0^L f(x) \sin(\frac{n\pi x}{L}) dx$ Relates to Odd functions.

Fourier Cosine Series: $f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} A_n \cos(\frac{n\pi x}{L}) \implies A_n = \frac{2}{L} \int_0^L f(x) \cos(\frac{n\pi x}{L}) dx$ Relates to Even functions.

Fourier Series: $f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} A_n \cos(\frac{n\pi x}{L}) + B_n \sin(\frac{n\pi x}{L}) \implies A_n = \frac{1}{L} \int_{-L}^L f(x) \cos(\frac{n\pi x}{L}) dx, B_n = \frac{1}{L} \int_{-L}^L f(x) \sin(\frac{n\pi x}{L}) dx$
Note that our interval is $2L$.

Generally $c_n = \frac{1}{2L} \langle f(x), e^{inx} \rangle$ $c_0 = \frac{a_0}{2}$ $c_n = \frac{a_n - ib_n}{2}$ $c_{-n} = \frac{a_n + ib_n}{2}$ $a_n = c_n + c_{-n}$ $b_n = i(c_n - c_{-n})$

Converges: Each point of continuity converges to the average of the left and right limits. All L^2 functions converge to the function. **Bessel's Inequality:** $\sum_{n \in \mathbb{N}} |c_n|^2 \leq \|f\|^2$ $c_n \rightarrow 0$ **Plancherel's Theorem:** $\|f\|^2 = \sum_{n \in \mathbb{N}} |c_n|^2$

Laplace's Equation: $\Delta u = 0$ $u_{xx} + u_{yy} = 0$ $u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} = 0$

Fundamental solution in 2D is $u = \frac{1}{2\pi} \log(r)$ and 3d is $u = \frac{1}{4\pi r}$.

Max Principle: The max of u is found on the boundary.

Uniqueness: Given sufficient BC (on all sides) the solution is unique.

Invariance Invariant under Rigid Motion. This implies that the solution doesn't care about rotation/direction hence the solution is a function of r only.

Rectangle and Cube: Split each inhomogeneous side into its own solution as $u = u_1 + u_2 + u_3 + u_4$ where u_i is the solution to the i th side with inhomogeneous rest homogenous. Note that most solutions will be in form of $(\cos + \sin)(\cosh + \sinh)$ where terms die out based on IC

Poisson's Formula: $\Delta u = 0$ for $x^2 + y^2 < a^2$ $u = h(\theta)$ for $x^2 + y^2 = a^2$

For $\mathbf{x} = (r, \theta)$ $\mathbf{x}' = (a, \phi)$ we have $r = |\mathbf{x}|$ $a = |\mathbf{x}'|$ $|\mathbf{x} - \mathbf{x}'|^2 = a^2 - 2ar \cos(\theta - \phi) + r^2$

$$u(r, \theta) = \frac{(a^2 - r^2)}{2\pi} \int_0^{2\pi} \frac{h(\phi)}{a^2 - 2ar \cos(\theta - \phi) + r^2} d\phi \text{ and } u(\mathbf{x}) = \frac{a^2 - |\mathbf{x}|^2}{2\pi} \int_{|\mathbf{x}'|=a} \frac{u(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|^2}$$

Mean Value Property: $u(0) = \frac{a^2}{2\pi a} \int_{|x|=a} \frac{u(\mathbf{x}')}{a^2} dS$ in other words the value at the center is the average of the boundary.

Wedges: $\{0 < \theta < \theta_0, 0 < r < a\}$. **Annulus:** $\{a < r < b, 0 < \theta < 2\pi\}$. **Exterior of Circle:** $\{r > a, 0 < \theta < 2\pi\}$. Make sure to consider **ALL** eigenvalues.

Divergence Theorem: $\int_D \nabla \cdot F dx = \int_{\partial D} F \cdot n dS$

Green's First Identity: $\int_{\partial D} v \frac{\partial u}{\partial n} dS = \int_D \nabla v \cdot \nabla u dx + \int_D v \Delta u dx$

Mean Value Principle: $u(0) = \frac{1}{\text{Area of } D} \int_D u dx$. Derived by recognizing the value of the integral in the sphere is not dependent on the radius.

Dirichlet Principle Among all w that solve $w = h$ on ∂D , the one with lowest energy is the harmonic solution.

Where $E[w] = \frac{1}{2} \int_D |\nabla w|^2 dx$

Green's Second Identity: $\int_D u \Delta v - v \Delta u dx = \int_{\partial D} u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} dS$

Representation Formula: $u(x_0) = \int_{\partial D} [-u(x)] \frac{\partial}{\partial n} \left(\frac{1}{|\mathbf{x} - \mathbf{x}_0|} \right) + \frac{1}{|\mathbf{x} - \mathbf{x}_0|} \frac{\partial u}{\partial n} \frac{dS}{4\pi}$

Green's Functions: the Green's Function $G(x)$ for the operator $-\Delta$ and the domain D at the point $x_0 \in D$ is a function defined for $x \in D$ st.

(i) $G(x)$ has cont. 2nd Derivatives and $\Delta G = 0$ in D except at $x = x_0$ where $\Delta G = -\delta(x - x_0)$

(ii) $G(x) = 0$ on ∂D

(iii) The functions $G(x) + \frac{1}{4\pi|\mathbf{x} - \mathbf{x}_0|}$ is finite at x_0 and has continuous 2nd Derivatives and is harmonic at x_0

We can also denote it as $G(x, x_0)$

Green's Function for Laplace's Equation with Dirichlet BC: $u(x) = \int_{\partial D} u \frac{\partial G}{\partial n} dS$

Symmetric: Green's function is Symmetric $G(x, x_0) = G(x_0, x)$

Half Space: $G(x, x_0) = -\frac{1}{4\pi|\mathbf{x} - \mathbf{x}_0|} + \frac{1}{4\pi|\mathbf{x} - \mathbf{x}_0^*|}$ where $x_0^* = (x_0, y_0, -z_0)$ Note that x_0^* is the reflection of x_0 across the plane $z = 0$ and provides the "opposite" energy field to make the field 0 on the boundary.

Half Space with Dirichlet BC: $u = h \in \partial D$ $u(x_0) = \frac{z_0}{2\pi} \int_{\partial D} \frac{h(x)}{|\mathbf{x} - \mathbf{x}_0|^3} dS$

Sphere For a sphere of radius a centered at 0, $x_0^* = \frac{a^2 x_0}{|x_0|^2}$ we denote $\rho = |\mathbf{x} - \mathbf{x}_0|$ $\rho^* = |\mathbf{x} - \mathbf{x}_0^*|$ Then the green's function is $G(x, x_0) = -\frac{1}{4\pi\rho} + \frac{a}{|x_0|} \frac{1}{4\pi\rho^*}$ and for $G(x, 0) = -\frac{1}{4\pi|\mathbf{x}|} + \frac{1}{4\pi a}$ Note that in application this becomes the solution using Poisson's Formula.

Sum/Difference $\sin(A) + \sin(B) = 2\sin(\frac{A+B}{2})\cos(\frac{A-B}{2})$ $\sin(A) - \sin(B) = 2\cos(\frac{A+B}{2})\sin(\frac{A-B}{2})$

$\cos(A) + \cos(B) = 2\cos(\frac{A+B}{2})\cos(\frac{A-B}{2})$ $\cos(A) - \cos(B) = -2\sin(\frac{A+B}{2})\sin(\frac{A-B}{2})$

Add/Subtract $\sin(A \pm B) = \sin(A)\cos(B) \pm \cos(A)\sin(B)$ $\cos(A \pm B) = \cos(A)\cos(B) \mp \sin(A)\sin(B)$

Product: $\sin(A)\sin(B) = \frac{1}{2}[\cos(A-B) - \cos(A+B)]$ $\cos(A)\cos(B) = \frac{1}{2}[\cos(A-B) + \cos(A+B)]$

$\sin(A)\cos(B) = \frac{1}{2}[\sin(A+B) + \sin(A-B)]$

Squared: $\sin^2(A) = \frac{1}{2}[1 - \cos(2A)]$ $\cos^2(A) = \frac{1}{2}[1 + \cos(2A)]$

Double Angle: $\sin(2A) = 2\sin(A)\cos(A)$ $\cos(2A) = \cos^2(A) - \sin^2(A) = 1 - 2\sin^2(A) = 2\cos^2(A) - 1$