

Project 4

Problem 1

Solve the wave equation

$$\frac{\partial^2 u(t, x)}{\partial t^2} = \frac{\partial^2 u(t, x)}{\partial x^2}$$

numerically. Use a 1 + 1-dimensional domain, and use Dirichlet boundary conditions $u(t, -L) = u(t, +L) = 0$. Use the spatial domain bounds $[-1; +1]$, and use the Gaussian-like function

$$f(t, x) = A \exp\left(-\frac{1}{2} \frac{\sin^2(2\pi(x - t))}{W^2}\right)$$

as initial condition. Use $L = 1$, $A = 1$, $W = 0.2$ as parameters. Evolve from $t = 0$ to $t = 4L$, i.e. for two crossing times. Plot the result.

Solution:

```
In [1]: using CairoMakie
        using LinearAlgebra
```

```
In [2]: A = 1
        W = 0.2
        L = 1;
```

```
In [3]: function coords(nn, ni, dt, dx)
        t = [n * dt for i in -ni:ni, n in 0:nn]
        x = [i * dx for i in -ni:ni, n in 0:nn]
        return t, x
    end;
```

The initial condition

$$f(t, x) = A \exp\left(-\frac{1}{2} \frac{\sin^2(2\pi(x - t))}{W^2}\right)$$

provided in the question is not consistent with the boundary condition, because at $x \rightarrow 1, t \rightarrow 0$ initial condition produces $u(0, 1) = 1$ although boundary condition is $u(0, 1) = 0$. To avoid inconsistency, we shifted the Gaussian initial condition to half wavelength. The initial wavefunction would like the following:

$$f(t, x) = A \exp\left(-\frac{1}{2} \frac{\sin^2(2\pi(x - t + 0.25))}{W^2}\right)$$

$$= A \exp\left(-\frac{1}{2} \frac{\cos^2(2\pi(x - t))}{W^2}\right)$$

```
In [4]: # Define initial conditions for a standing wave
function initial_condition(t, x)
    ni, nn = size(t)
    u = zeros(ni, nn)
    # Initial conditions for first two points in time
    # f(t,x) = A exp(-0.5*cos^2(2 π (x-t))/W^2)
    for n in 1:2, i in 1:ni
        u[i, n] = A * exp.(-0.5*(cos(2*π * (x[i, n]-t[i,n])))^2/W^2)
    end
    return u
end;
```

```
In [5]: function evolve_point!(u, dt, dx, n, i)
    ni, nn = size(u)
    # Finite difference for second derivative in x direction at point [i,n-1]
    uxx = (u[i-1,n-1] - 2*u[i,n-1] + u[i+1,n-1]) / dx^2
    # Time derivative
    u[i,n] = uxx * dt^2 - u[i,n-2] + 2*u[i,n-1]
    return
end;
```

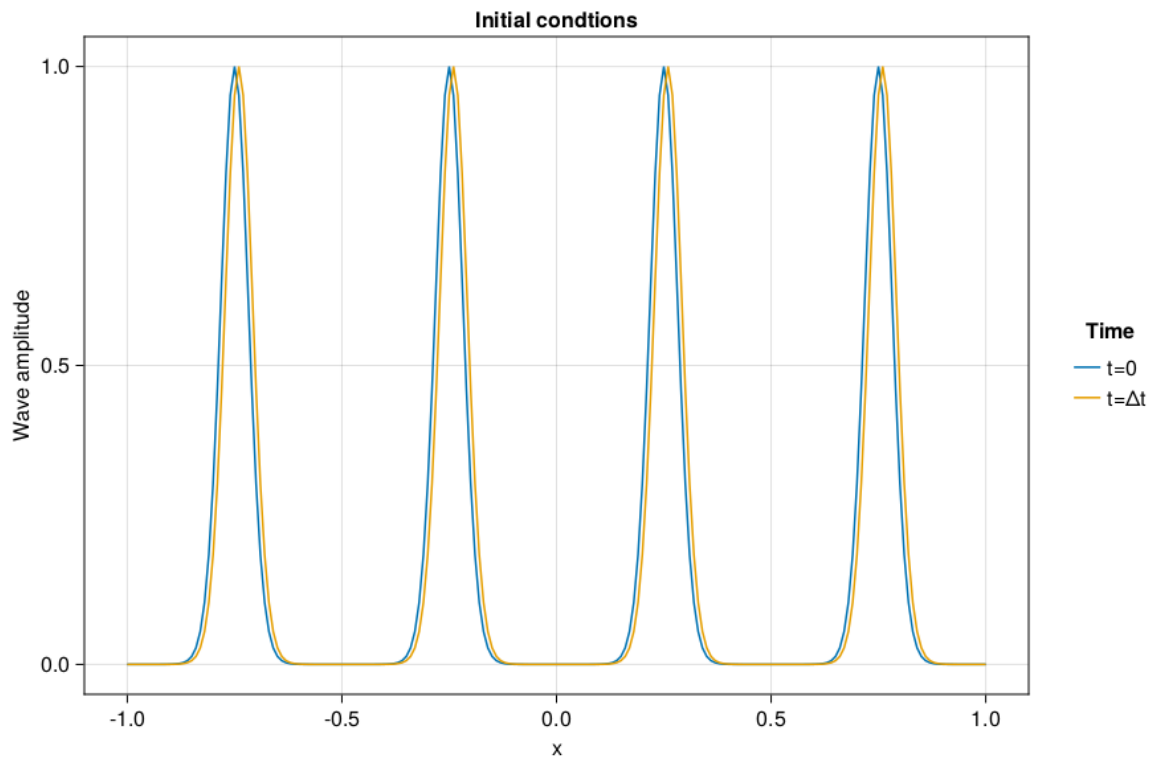
```
In [6]: function evolve_step!(u, dt, dx, n)
    ni, nn = size(u)
    # Apply boundary condition at left boundary
    u[1,n] = 0
    # Evolve interior
    for i in 2:ni-1
        evolve_point!(u, dt, dx, n, i)
    end
    # Apply boundary condition at right boundary
    u[ni,n] = 0
    return
end;
```

```
In [7]: function evolve!(u, dt, dx)
    ni, nn = size(u)
    # Loop over all times
    for n in 3:nn
        evolve_step!(u, dt, dx, n)
    end
    return
end;
```

```
In [8]: nn = 400
        ni = 100
        dt = 4 / nn
        dx = 1 / ni
        t, x = coords(nn, ni, dt, dx);
```

```
In [9]: u = initial_condition(t, x);
        fig = Figure(resolution = (900, 600))
        ax = Axis(fig[1, 1], xlabel="x", ylabel="Wave amplitude", title="Initial condtions")
        lines!(x[:,1],u[:,1],label="t=0")
        lines!(x[:,2],u[:,2],label="t=Δt")
        fig[1, 2] = Legend(fig, ax, "Time", framevisible = false)
        fig
```

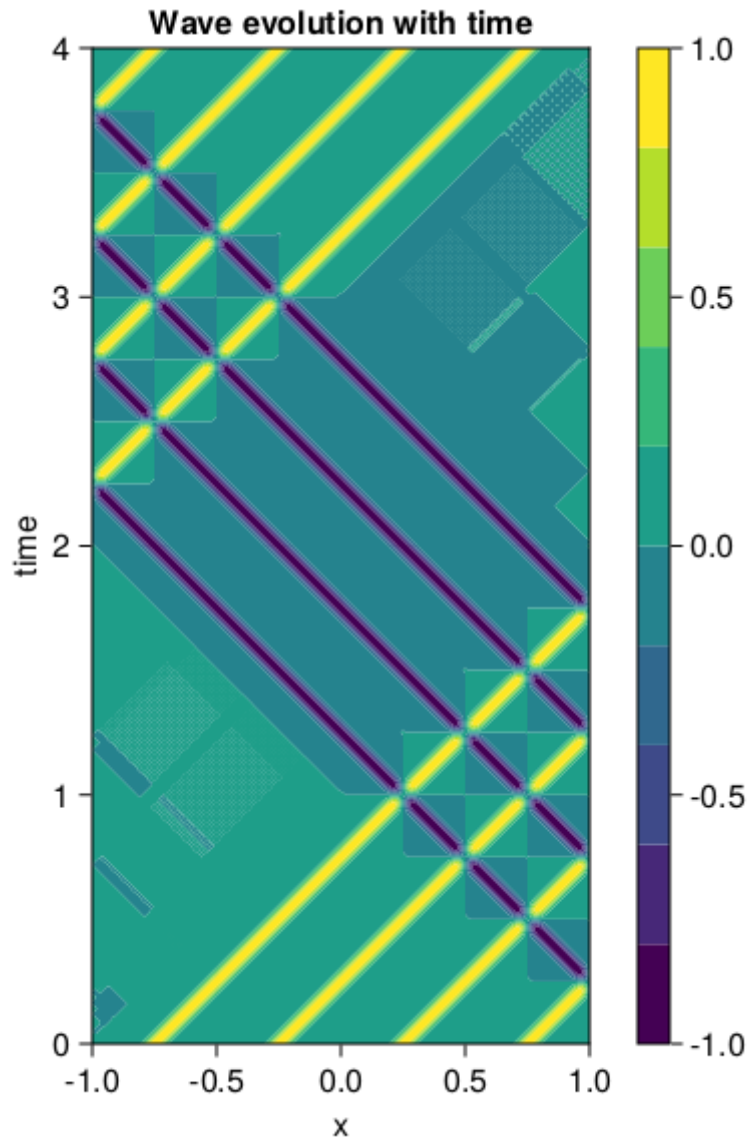
Out[9]:



```
In [10]: evolve!(u, dt, dx)
```

```
In [11]: fig = Figure(resolution = (400, 600))
        ax = Axis(fig[1, 1],title="Wave evolution with time",xlabel="x",ylabel="time")
        co = contourf!(fig[1,1], vec(x), vec(t), vec(u))
        Colorbar(fig[1,2],co)
        fig
```

Out[11]:



Remarks: This wave pattern makes sense because of the hard boundary the wave crest gets reflected to ± 1 to ∓ 1 at each boundary and bounced back. Moreover, when two opposite crest meet it causes destructive interference which is clear in the above picture.

Problem 2

Vary the time step size dt . If you choose dt too large, the discrete system will be unstable. What is the stability limit? Why does this happen? How could you avoid this limit?

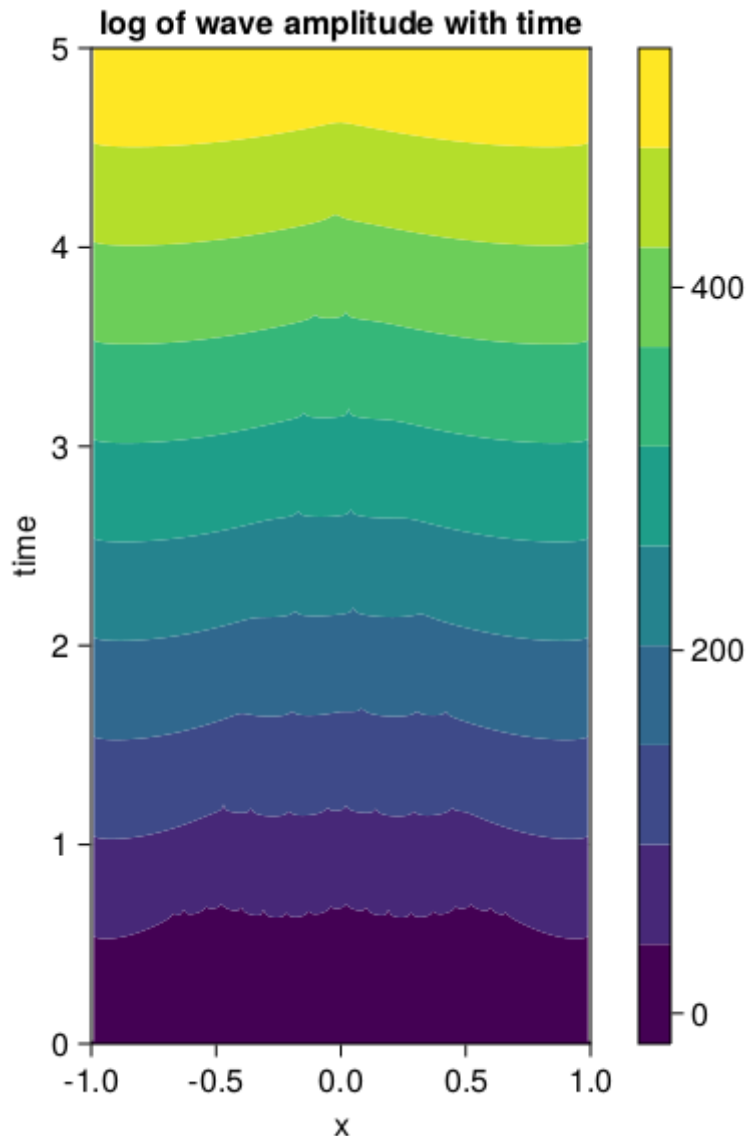
Solution:

Just for illustration purpose we have used $dt = 1.25 \, dx$. This gives a very unstable solution and the wave amplitude grows unbounded in this case. To demonstrate the effect here we have plotted $\log|u|$ here below.

```
In [12]: dt = 5 / nn #dt = 1.25 dx
t, x = coords(nn, ni, dt, dx);
u = initial_condition(t, x);
evolve!(u, dt, dx)
```

```
In [13]: fig = Figure(resolution = (400, 600))
ax = Axis(fig[1, 1], title="log of wave amplitude with time", xlabel="x", ylabel="time")
co = contourf!(fig[1,1], vec(x), vec(t), log.(abs.(vec(u))))
Colorbar(fig[1,2],co)
fig
```

Out[13]:



Stability condition: From the Courant–Friedrichs–Lewy condition

$$\frac{v\Delta t}{\Delta x} \leq C_{CFL}$$

we can estimate that for this particular problem. For our particular setup velocity of wave propagation $v = 1$ and Courant number $C_{CFL} = 1$. Hence, the discrete time difference must be less than equal to discrete location difference in order for the solution to be stable i.e. $\Delta t \leq \Delta x$.

Explanation: If the *stability condition* is violated we end up computing numerically the waveamplitude before the actual wave reaches the space-time point (with velocity $v = 1$). Hence, we are getting unphysical result from numerical solution.

Ways to avoid the limit:

- Increase the value of `dx` keeping `dt` fixed.
- Decrease the value of `dt` keeping `dx` fixed.
- Use different algorithm (e.g. implicit Euler method) to increase C_{CFL} .

Problem 3

The energy density of a scalar wave is

$$\epsilon(t, x) = \frac{1}{2} \left[\left(\frac{du(t, x)}{dt} \right)^2 + \left(\frac{du(t, x)}{dx} \right)^2 \right]$$

Calculate the energy density as a function of space and time for your solution. Plot the result.

Solution:

```
In [14]: dt = 4 / nn #dt = dx
t, x = coords(nn, ni, dt, dx);
u = initial_condition(t, x);
evolve!(u, dt, dx)
```

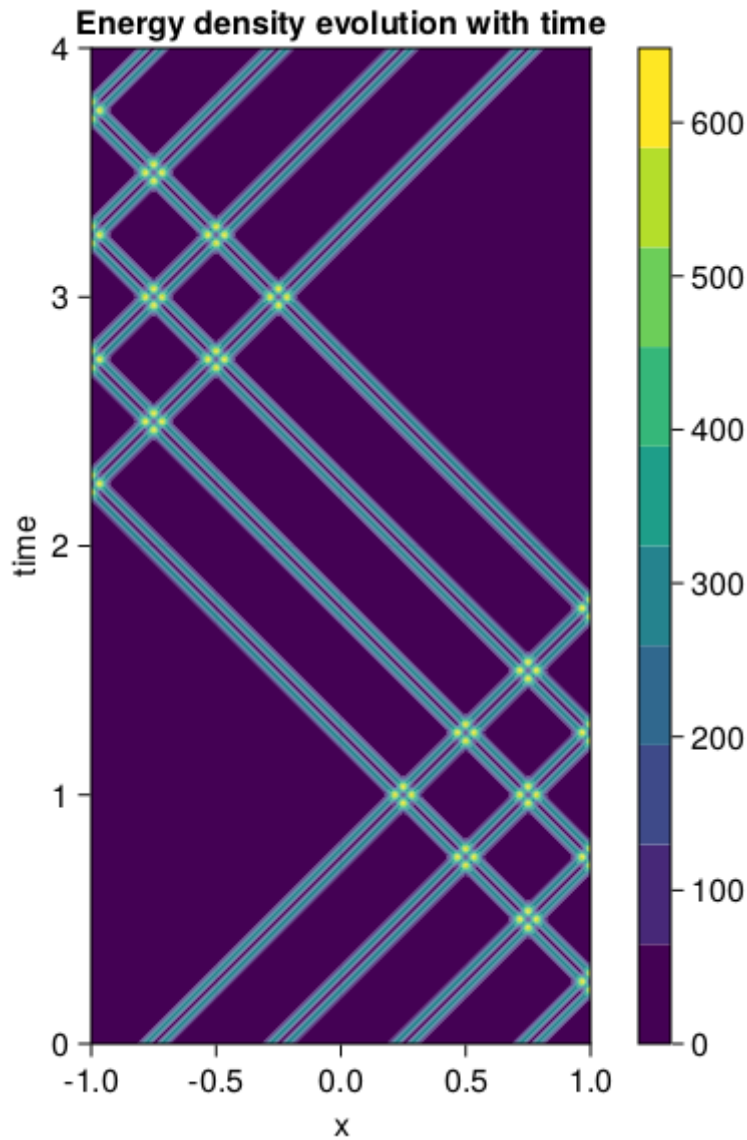
We will take center derivative to calculate bulk energy density here. We will use forward and backward derivative for all the peripheral points.

```
In [15]: function energy_density(u,dx,dt)
    dux = similar(u)
    dudt = similar(u)
    # center derivate in x direction only for the bulk region
    dux[2:end-1,:] = (u[3:end,:]-u[1:end-2,:])./(2*dx)
    # now we handle boundary with forward and backward derivative respectively
    dux[1,:] = (u[2,:]-u[1,:])./(dx)
    dux[end,:] = (u[end,:]-u[end-1,:])./(dx)
    # center derivate in t direction only for the bulk region
    dudt[:,2:end-1] = (u[:,3:end]-u[:,1:end-2])./(2*dt)
    # now we handle boundary with forward and backward derivative respectively
    dudt[:,1] = (u[:,2]-u[:,1])./(dt)
    dudt[:,end] = (u[:,end]-u[:,end-1])./(dt)
    # energy density formula for each location
    return ε = 0.5*(dux.^2 .+ dudt.^2)
end;
```

```
In [16]: ε = energy_density(u,dx,dt);
```

```
In [17]: fig = Figure(resolution = (400, 600))
ax = Axis(fig[1, 1],title="Energy density evolution with time",xlabel="x",ylabel="")
co = contourf!(fig[1,1], vec(x), vec(t), vec(ε))
Colorbar(fig[1,2],co)
fig
```

Out[17]:



Problem 4

To calculate the total energy $E(t)$ as a function of time, you need to integrate it over space. One way to do so is using the function below. Calculate and show the total energy $E(t)$ as a function of time. How well is the energy conserved? How does this depend on the grid spacings dx and dt ?

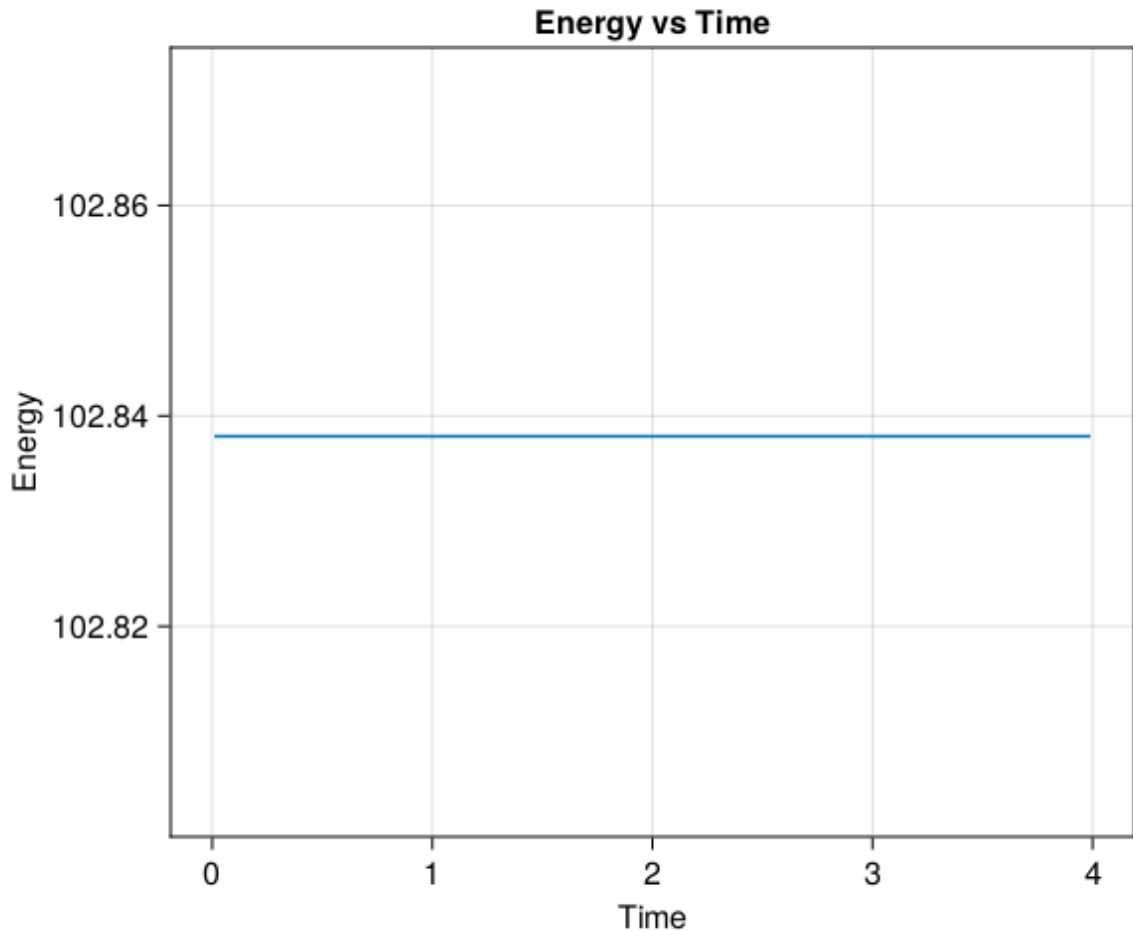
```
In [18]: # Integrate `epsilon` over space
function integrate(epsilon, n, dx)
    ni, nt = size(epsilon)
    # The boundary points are special; they have a weight of 1/2.
    # There are `ni` points overall, and the grid spacing is `dx = 2L/(ni-1)`.
    # Integrating the constant function `epsilon=1` thus gives exactly `2L`.
    E = epsilon[1, n] / 2
    for i in 2:ni-1
        E += epsilon[i, n]
    end
    E += epsilon[ni, n] / 2
    # Scale by the grid spacing
    E *= dx
    return E
end;
```

Solution:

```
In [19]: E = [integrate(ϵ,i,dx) for i in 1:(nn+1)];
```

```
In [20]: fig = Figure(resolution = (600, 500))
ax = Axis(fig[1, 1], xlabel="Time", ylabel="Energy", title="Energy vs Time")
lines!(t[1,2:end-1],E[2:end-1]) # energy in bulk region
ylims!(102.8, 102.875)
fig
```

Out[20]:



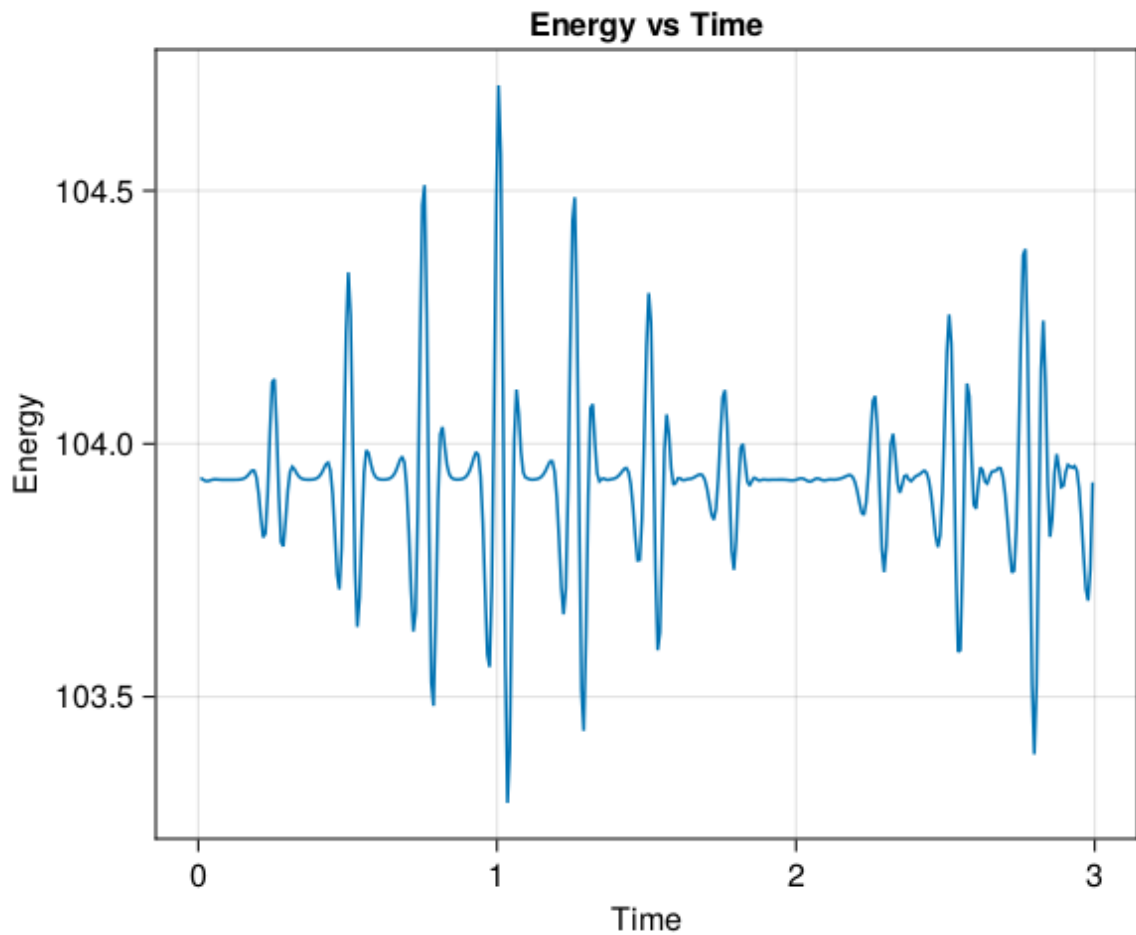
Remarks: From the plot it is apparent that here energy is well-consevered.

Case 1: $\frac{\Delta t}{\Delta x} = 0.75$

```
In [21]: dt = 3 / nn #dt = 0.75 dx
t, x = coords(nn, ni, dt, dx);
u = inital_condition(t, x);
evolve!(u, dt, dx)
ϵ = energy_density(u,dx,dt);
E = [integrate(ϵ,i,dx) for i in 1:(nn+1)];
```

```
In [22]: fig = Figure(resolution = (600, 500))
ax = Axis(fig[1, 1], xlabel="Time", ylabel="Energy", title="Energy vs Time")
lines!(t[1,2:end-1],(E[2:end-1])) # energy in bulk region
fig
```


Out[22]:



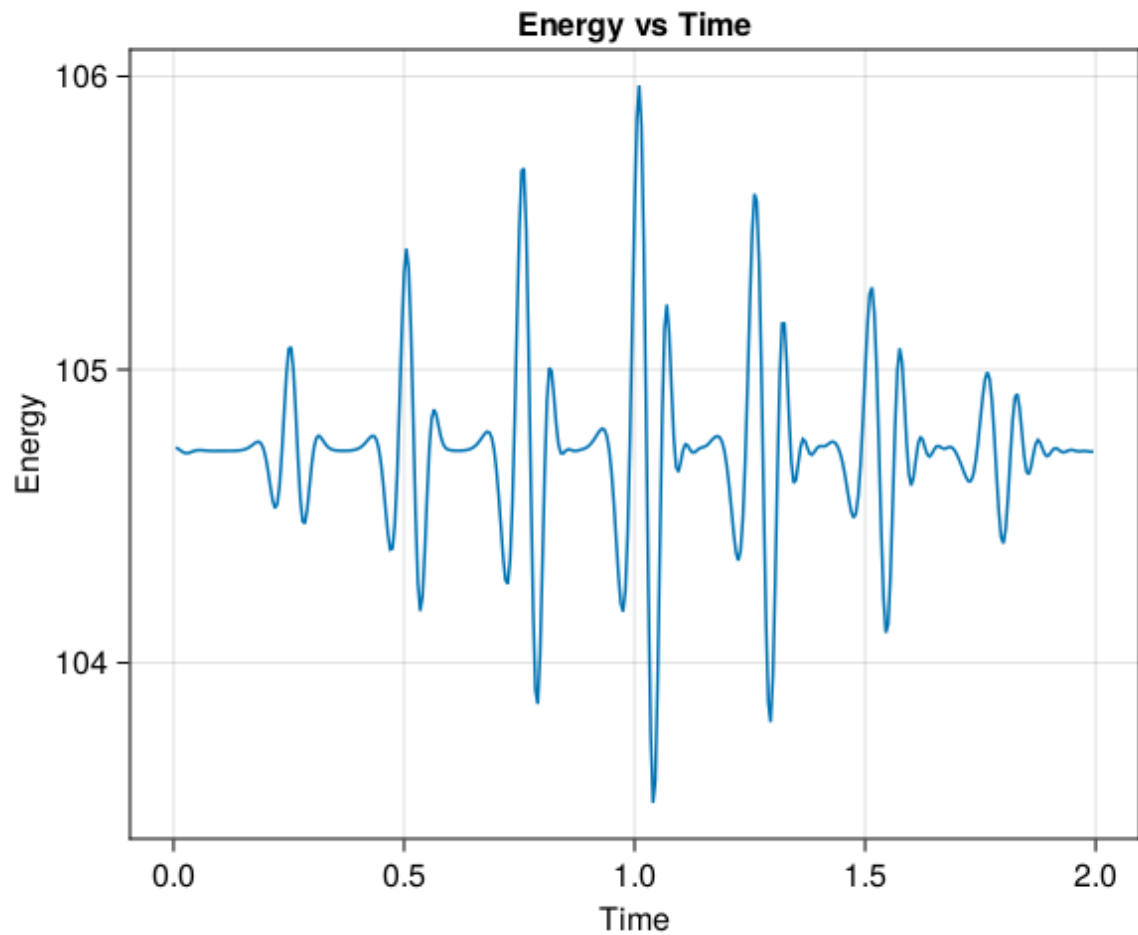
Remarks: Energy conservation is stable most of the time; however, it fluctuates when waves bounces at the boundary or collides with each other.

Case 2: $\frac{\Delta t}{\Delta x} = 0.5$

```
In [23]: dt = 2 / nn #dt = 0.5 dx
t, x = coords(nn, ni, dt, dx);
u = inital_condition(t, x);
evolve!(u, dt, dx)
ϵ = energy_density(u,dx,dt);
E = [integrate(ϵ,i,dx) for i in 1:(nn+1)];
```

```
In [24]: fig = Figure(resolution = (600, 500))
ax = Axis(fig[1, 1], xlabel="Time", ylabel="Energy", title="Energy vs Time")
lines!(t[1,2:end-1],(E[2:end-1])) # energy in bulk region
fig
```

Out[24]:



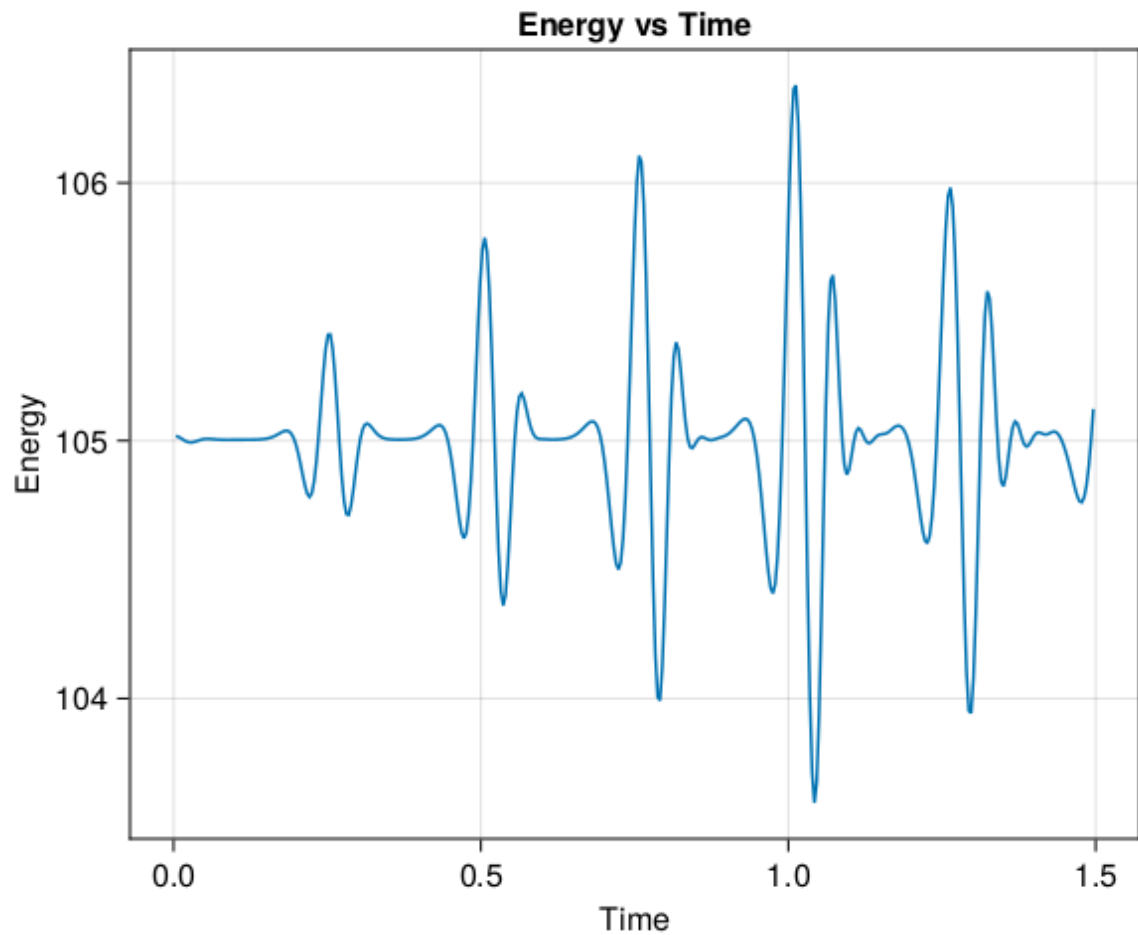
Remarks: Energy conservation is more or less fine, but the fluctuation due to numerical error is worse than previous case.

Case 3: $\frac{\Delta t}{\Delta x} = 0.375$

```
In [25]: dt = 1.5 / nn
t, x = coords(nn, ni, dt, dx);
u = initial_condition(t, x);
evolve!(u, dt, dx)
ϵ = energy_density(u,dx,dt);
E = [integrate(ϵ,i,dx) for i in 1:(nn+1)];
```

```
In [26]: fig = Figure(resolution = (600, 500))
ax = Axis(fig[1, 1], xlabel="Time", ylabel="Energy", title="Energy vs Time")
lines!(t[1,2:end-1],(E[2:end-1])) # energy in bulk region
fig
```

Out[26]:



Remarks: Energy fluctuation due to numerical error is worse than pervious case. We can estimate that numerical error will grow if reduce the `dt` and `dx`. However, if we ignore the periodic fluctuation ($< 2\%$) here, energy is pretty much conserved for all three cases,