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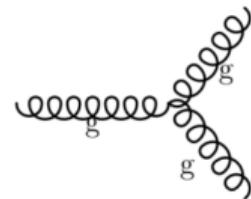
A TMD-oriented analysis of $\pi^+\pi^-$ pairs in e^+e^- collisions

QCD in one equation: the Lagrangian

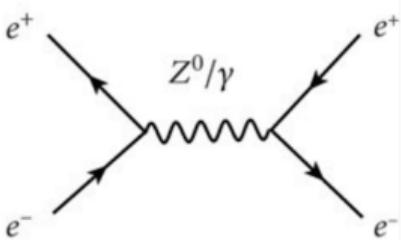
$$\mathcal{L}_{\text{QCD}} = \underbrace{\bar{\psi} i\gamma^\mu \partial_\mu \psi}_{\text{kinetic}} + \underbrace{g_s \bar{\psi} \gamma^\mu t^a A_\mu^a \psi}_{\text{interaction}} - \underbrace{m_q \bar{\psi} \psi}_{\text{mass}} - \underbrace{\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu}}_{\text{gluon}} \quad [4]$$

Non-Abelian \Rightarrow gluon self-coupling.

[4]

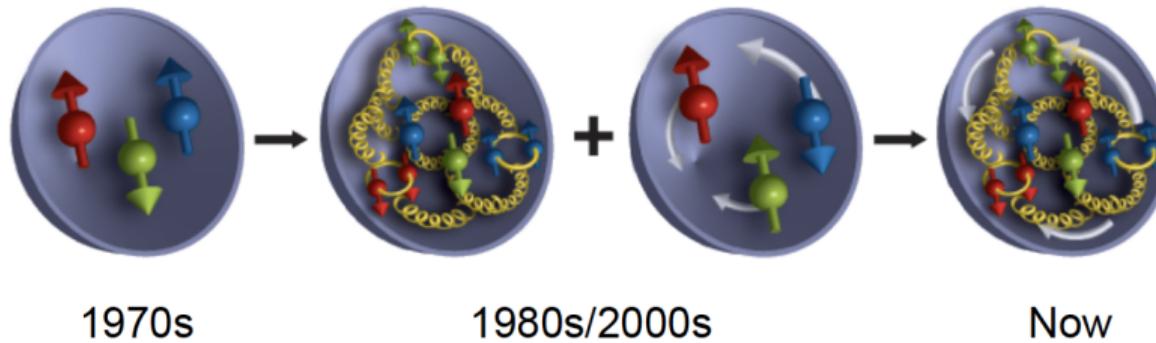


Proton model



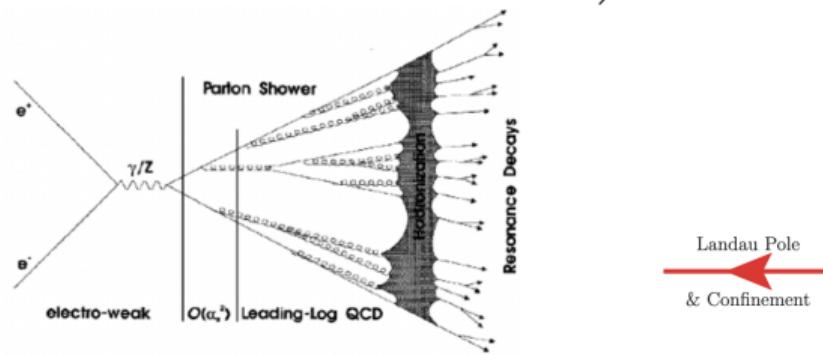
Q sets resolution.

In e^+e^- : $Q = \sqrt{s}$ (typically γ^*/Z exchange).



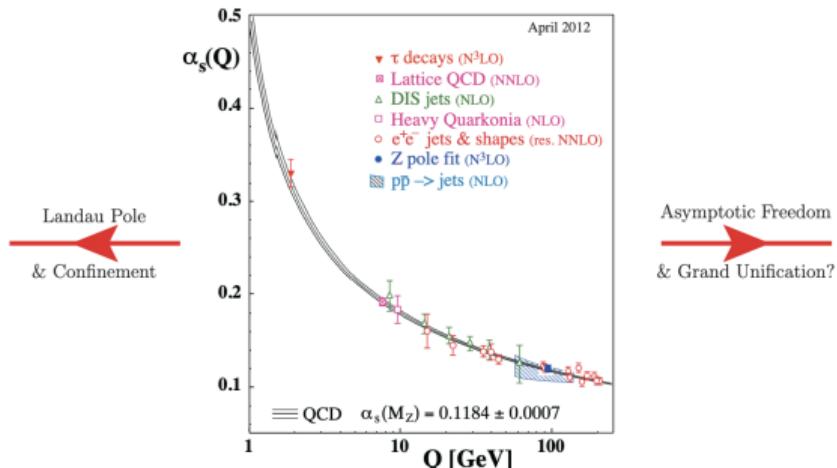
high $Q \Rightarrow \alpha_s$ small; low $Q \Rightarrow$ non-perturbative.

$$\alpha_s(Q^2) \sim \frac{1}{b_0 \ln(Q^2/\Lambda_{\text{QCD}}^2)} \quad \Rightarrow \quad Q \rightarrow \Lambda_{\text{QCD}} \text{ (Landau pole)}$$



Shower cutoff $Q_0 \sim \mathcal{O}(1 \text{ GeV}) \Rightarrow$ hadronisation model.

[4]

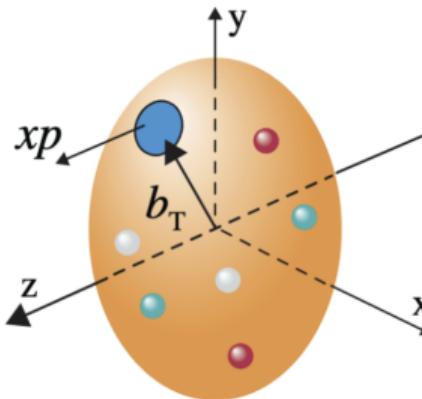


PDFs

$$f_i(x, Q)$$

- Initial state: parton i inside a hadron.
- $x =$ longitudinal momentum fraction of the hadron.

[1,3,4]



FFs

$$D_i^h(z, Q)$$

- Final state: parton $i \rightarrow$ hadron h .
- $z =$ fraction of parton momentum carried by h .

[1,2,3]

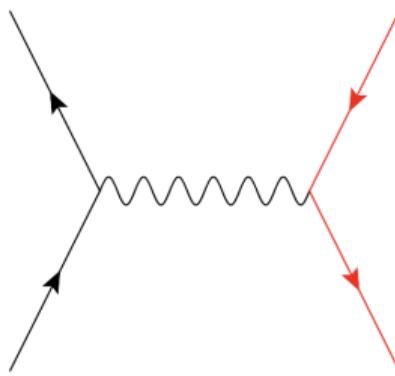
TMD

TMD PDFs/FFs add transverse momentum: $f_i(x, k_T, Q)$, $D_i^h(z, k_T, Q)$ ("3D": longitudinal + transverse). [3,5]

 e^+e^- motivation (clean fragmentation)

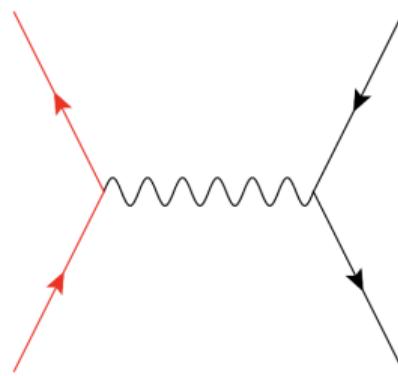
In e^+e^- there is no hadronic initial state \Rightarrow no PDFs; focus is on the final-state fragmentation (FFs/TMD FFs). [1,2,3]

TMDs across reactions



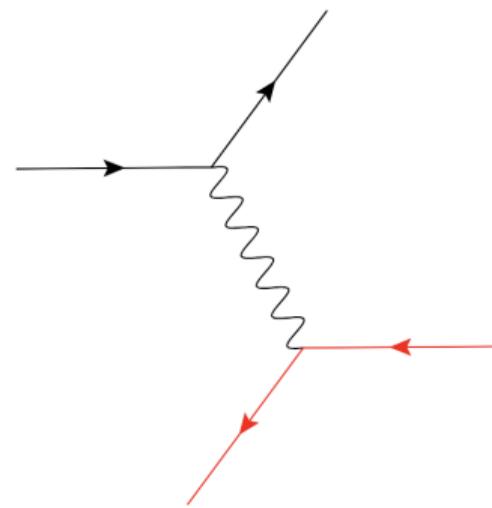
Hadronic Z decay

$$e^- e^+ \rightarrow \gamma^*/Z^0 \rightarrow q\bar{q}$$



Drell-Yan

$$q\bar{q} \rightarrow \gamma^*/Z^0 \rightarrow \ell^+ \ell^-$$



DIS

$$\ell\bar{q} \xrightarrow{\gamma^*/Z^*} \ell\bar{q}$$

When PDFs are enough vs when you need TMDs

Integrated over q_T (collinear)

$$\frac{d\sigma}{dQ^2 dY} = \sum_{i,j} \int \frac{d\xi_a}{\xi_a} \frac{d\xi_b}{\xi_b} f_{i/H_a}(\xi_a) f_{j/H_b}(\xi_b) \frac{d\hat{\sigma}_{ij}}{dQ^2 dY} \left[1 + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{Q^2}\right) \right].$$

Inclusive in q_T : PDFs capture the dominant structure.

[5]

Small q_T (TMD-sensitive)

$$\frac{d\sigma}{d^4 q} = \frac{1}{s} \sum_i \hat{\sigma}_{i\bar{i}}^{\text{TMD}}(Q) \int d^2 \mathbf{k}_T f_{i/H_a}(x_a, \mathbf{k}_T) f_{\bar{i}/H_b}(x_b, \mathbf{q}_T - \mathbf{k}_T) \left[1 + \mathcal{O}\left(\frac{q_T^2}{Q^2}\right) \right].$$

Differential in q_T : the *shape* probes transverse momentum.

[5]

Often use Fourier/ b_T space: $\tilde{f}(x, \mathbf{b}_T) = \int d^2 \mathbf{k}_T e^{-i\mathbf{b}_T \cdot \mathbf{k}_T} f(x, \mathbf{k}_T)$ (convolutions \rightarrow products).

[5]

PYTHIA vs HERWIG (what is being tested)

How to extract? → need an observable **sensitive to TMDs**.

Plan (in this project)

- Use back-to-back jets in e^+e^- and study $\pi\pi$ pairs *within the two jets.* [3,5]
- Define the pion in each jet either as: (i) **first/closest** to the parent parton, or (ii) **highest-momentum.** [2]

hadronisation is **non-perturbative** → need **phenomenological models.**

PYTHIA = Lund/string HERWIG = cluster

[2,4]

PYTHIA vs HERWIG (hadronisation models)

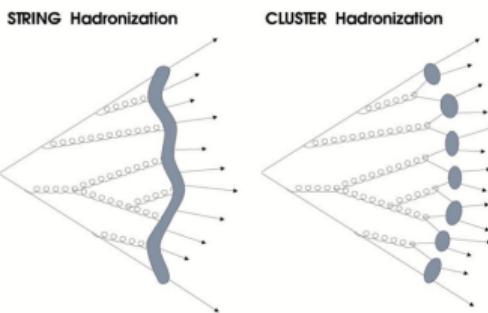
PYTHIA

Lund/string: colour flux tube
 breaks $\rightarrow q\bar{q} \rightarrow$ hadrons. [2,4]

row	event	size	no	id	name	st	m1	m2	d1	d2	px	py	pz	E
0	[3	134	9	98	system	-11	9	0	0	0	0.000	0.000	0.000	91.200
1	[3	134	1	-12	[e+	-12	0	0	3	0	0.000	0.000	0.000	45.600
2	[3	134	12	14	[e-	-12	0	0	4	0	0.000	0.000	0.000	45.600
3	[3	134	3	-31	[e+	-21	1	0	5	0	0.000	0.000	0.000	45.600
4	[3	134	4	11	[e-	-21	12	0	5	0	0.000	0.000	0.000	45.600
5	[3	134	16	23	Z0	-22	3	4	6	7	0.000	0.000	0.000	91.200
6	[3	134	16	1	[d	-23	5	0	8	9	-44.829	7.856	2.821	45.600
7	[3	134	17	-1	[dbar	-23	5	0	10	18	-44.829	-2.821	45.600	87
8	[3	134	8	1	[u	-51	6	0	11	12	-1.595	0.505	-4.502	41.711
9	[3	134	9	21	[g	-51	6	0	13	13	-0.289	0.613	-1.471	0.976

(showing rows 0 to 9 of 134 total)

More [Enter=+10, n=next event, q=quit]:



HERWIG

Cluster: colour-singlet clusters at shower cutoff \rightarrow hadron decays.

[2,4]

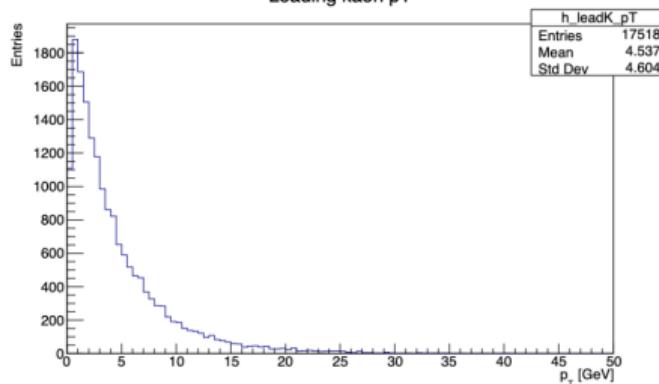
row	event	size	no	id	name	st	m1	m2	d1	d2	px	py	pz	E
00	[3	134	80	-322	K-	83	57	73	19	0	-0.031	-0.152	-0.157	8.543
01	[3	134	81	-3112	[pi-3112	83	57	73	116	117	-3.258	0.231	-0.778	3.564
02	[3	134	82	-2112	[pi-2112	83	57	73	19	0	-0.449	-0.776	0.356	1.051
03	[3	134	83	[211	[pi+	83	57	73	0	0	-0.643	0.377	0.396	0.056
04	[3	134	84	[113	[rho0	84	57	73	97	98	0.238	0.721	-1.855	1.962
05	[3	134	85	[111	[pi0	85	57	73	116	118	0.319	0.79	-0.175	1.175
06	[3	134	86	-2110	[rho-2110	84	57	73	19	140	0.512	0.365	0.484	1.136
07	[3	134	87	[2212	[pi+	84	57	73	19	0	1.593	1.236	0.164	2.229
08	[3	134	88	-2115	[pi-2115	84	57	73	0	0	0.647	-0.002	0.279	0.728
09	[3	134	89	-2116	[id-2116	84	57	73	101	102	3.613	0.916	0.066	4.158

(showing rows 00 to 09 of 134 total)

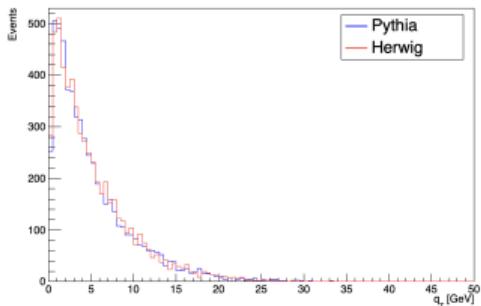
More [Enter=+10, n=next event, q=quit]:

Current plots

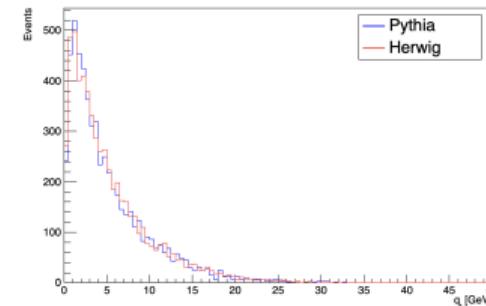
Leading kaon pT



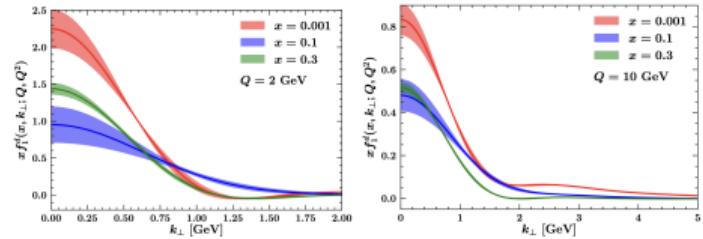
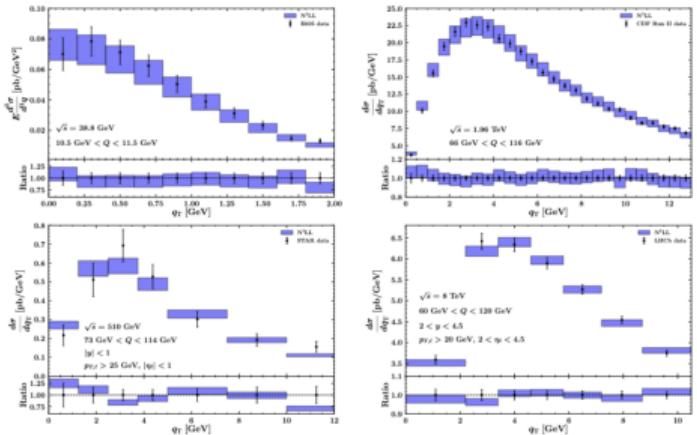
q_T of opposite charge pion pair (highest momentum)



q_T of same sign charge pion pair (highest momentum)



q_T spectra → TMD extraction (Drell–Yan example)



q_T data constrains TMD shape/parameters. [6]

Target: summarise shapes with few parameters. [6]

TMD

shape/parameters.

Low- q_T shape discriminates global fits.

[6]

Goal: use an analogous q_T -sensitive observable in e^+e^- to extract/compare TMD information. [6]

- 1 O. Biebel, D. de Florian, D. Milstead, W. Vogelsang (PDG), *Fragmentation Functions in e^+e^- , ep , and pp Collisions (review)*, Prog. Theor. Exp. Phys. 2020, 083C01.
- 2 O. Biebel, P. Nason, B. R. Webber, *Jet fragmentation in e^+e^- annihilation*, arXiv:hep-ph/0109282v2 (2001).
- 3 A. Vossen, *Parton Fragmentation Functions*, arXiv:1702.01329 (2017).
- 4 P. Z. Skands, *Introduction to QCD*, arXiv:1207.2389 (2012).
- 5 R. Boussarie et al. (TMD Collaboration), *TMD Handbook*, arXiv:2304.03302 (2023).
- 6 A. Bacchetta et al., *Transverse-momentum-dependent parton distributions up to N^3LL from Drell-Yan data*, JHEP 07 (2020) 117, arXiv:1912.07550.