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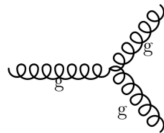
# A TMD-oriented analysis of $\pi^+\pi^-$ pairs in $e^+e^-$ collisions

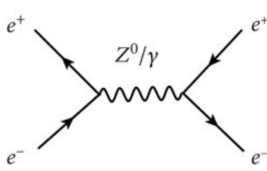
# QCD in one equation: the Lagrangian

$$\mathcal{L}_{\text{QCD}} = \underbrace{\bar{\psi} i \gamma^\mu \partial_\mu \psi}_{\text{kinetic}} + \underbrace{g_s \bar{\psi} \gamma^\mu t^a A_\mu^a \psi}_{\text{interaction}} - \underbrace{m_q \bar{\psi} \psi}_{\text{mass}} - \underbrace{\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu}}_{\text{gluon}} \quad [4]$$

Non-Abelian  $\Rightarrow$  gluon self-coupling.

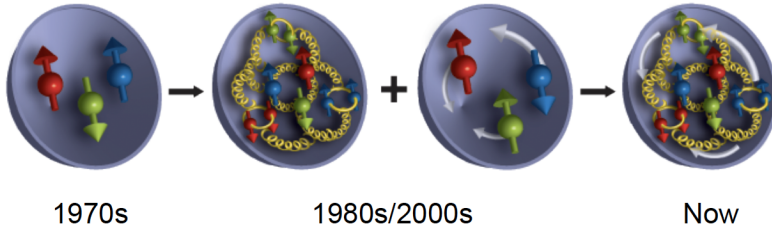
[4]





$Q$  sets resolution.

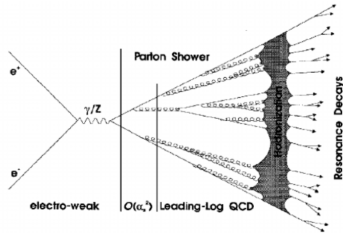
In  $e^+e^-$ :  $Q = \sqrt{s}$  (typically  $\gamma^*/Z$  exchange).



high  $Q \Rightarrow \alpha_s$  small; low  $Q \Rightarrow$  non-perturbative.

[4]

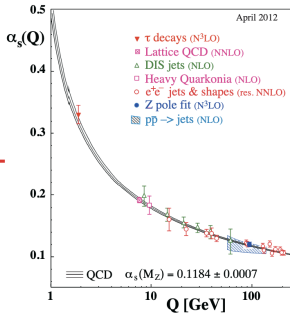
$$\alpha_s(Q^2) \sim \frac{1}{b_0 \ln(Q^2/\Lambda_{\text{QCD}}^2)} \Rightarrow Q \rightarrow \Lambda_{\text{QCD}} \text{ (Landau pole)}$$



Shower cutoff  $Q_0 \sim \mathcal{O}(1 \text{ GeV}) \Rightarrow$  hadronisation model.

[4]

Landau Pole  
& Confinement



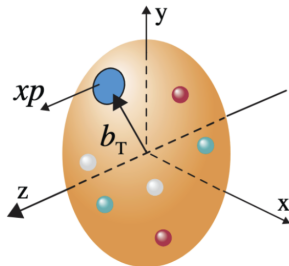
Asymptotic Freedom  
& Grand Unification?

## PDFs

$$f_i(x, Q)$$

- **Initial state:** parton  $i$  inside a hadron.
- $x$  = longitudinal momentum fraction of the hadron.

[1,3,4]



## FFs

$$D_i^h(z, Q)$$

- **Final state:** parton  $i \rightarrow$  hadron  $h$ .
- $z$  = fraction of parton momentum carried by  $h$ .

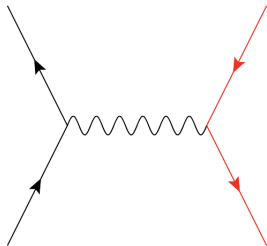
[1,2,3]

## TMD

TMD PDFs/FFs add transverse momentum:  $f_i(x, \mathbf{k}_T, Q)$ ,  $D_i^h(z, \mathbf{k}_T, Q)$  ("3D": longitudinal + transverse). [3,5]

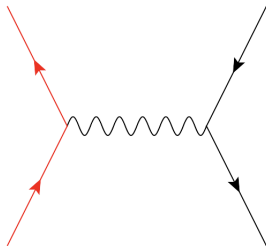
 $e^+e^-$  motivation (clean fragmentation)

In  $e^+e^-$  there is **no hadronic initial state**  $\Rightarrow$  no PDFs; focus is on the final-state fragmentation (FFs/TMD FFs). [1,2,3]



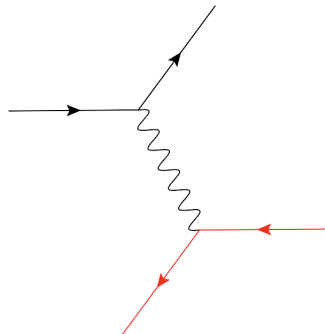
Hadronic  $Z$  decay

$$e^- e^+ \rightarrow \gamma^*/Z^0 \rightarrow q \bar{q}$$



Drell-Yan

$$q \bar{q} \rightarrow \gamma^*/Z^0 \rightarrow \ell^+ \ell^-$$



DIS

$$\ell \bar{q} \xrightarrow{\gamma^*/Z^*} \ell \bar{q}$$

# When PDFs are enough vs when you need TMDs

## Integrated over $q_T$ (collinear)

$$\frac{d\sigma}{dQ^2 dY} = \sum_{i,j} \int \frac{d\xi_a}{\xi_a} \frac{d\xi_b}{\xi_b} f_{i/H_a}(\xi_a) f_{j/H_b}(\xi_b) \frac{d\hat{\sigma}_{ij}}{dQ^2 dY} \left[ 1 + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{Q^2}\right) \right].$$

Inclusive in  $q_T$ : PDFs capture the dominant structure.

[5]

## Small $q_T$ (TMD-sensitive)

$$\frac{d\sigma}{d^4q} = \frac{1}{s} \sum_i \hat{\sigma}_{i\bar{i}}^{\text{TMD}}(Q) \int d^2\mathbf{k}_T f_{i/H_a}(x_a, \mathbf{k}_T) f_{\bar{i}/H_b}(x_b, \mathbf{q}_T - \mathbf{k}_T) \left[ 1 + \mathcal{O}\left(\frac{q_T^2}{Q^2}\right) \right].$$

Differential in  $q_T$ : the *shape* probes transverse momentum.

[5]

Often use Fourier/ $b_T$  space:  $\tilde{f}(x, \mathbf{b}_T) = \int d^2\mathbf{k}_T e^{-i\mathbf{b}_T \cdot \mathbf{k}_T} f(x, \mathbf{k}_T)$  (convolutions  $\rightarrow$  products).

[5]

# PYTHIA vs HERWIG (what is being tested)

**How to extract?** → need an observable **sensitive to TMDs**.

## Plan (in this project)

- Use back-to-back jets in  $e^+e^-$  and study  $\pi\pi$  pairs *within the two jets*. [3,5]
- Define the pion in each jet either as: (i) **first/closest** to the parent parton, or (ii) **highest-momentum**. [2]

hadronisation is **non-perturbative** → need **phenomenological models**.

PYTHIA = Lund/string      HERWIG = cluster

[2,4]

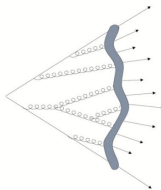


# PYTHIA vs HERWIG (hadronisation models)

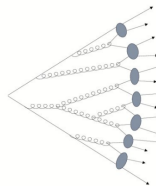
## PYTHIA

Lund/string: colour flux tube  
breaks  $\rightarrow q\bar{q} \rightarrow$  hadrons. [2,4]

STRING Hadronization



CLUSTER Hadronization



## HERWIG

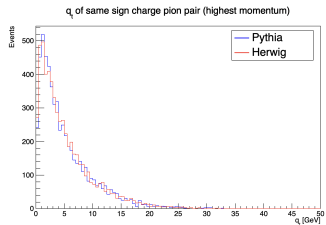
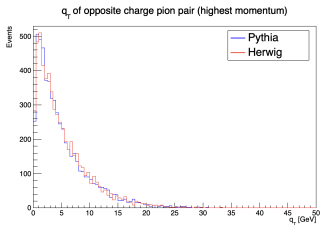
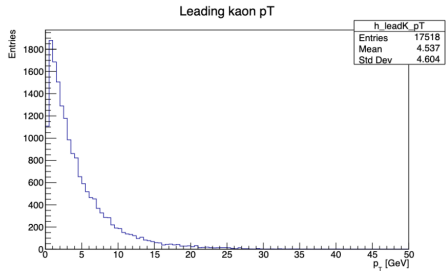
Cluster: colour-singlet clusters at  
shower cutoff  $\rightarrow$  hadron decays. [2,4]

row	event	size	no	id	name	st	m1	m2	d1	d2	px	py	pz	E
0	13	134	0	190	system	-11	0	0	0	0	0	0.000	0.000	0.000
1	13	134	11	-11	e+	-12	0	0	13	0	0	0.000	0.000	45.600
2	13	134	12	11	e-	-12	0	0	14	0	0	0.000	0.000	-45.600
3	13	134	13	-11	e+	-21	11	0	15	0	0	0.000	0.000	45.600
4	13	134	14	11	e-	-21	12	0	15	0	0	0.000	0.000	-45.600
5	13	134	15	23	Z0	-22	13	4	16	7	0	0.000	0.000	91.200
6	13	134	16	1	d	-23	15	0	18	9	0	-44.829	7.850	45.600
7	13	134	17	-1	dbar	-23	15	0	10	10	0	44.829	-7.850	45.600
8	13	134	18	1	d	-51	16	0	11	12	0	-41.505	0.505	4.102
9	13	134	19	121	g	-51	16	0	13	13	0	-0.289	6.815	-1.471

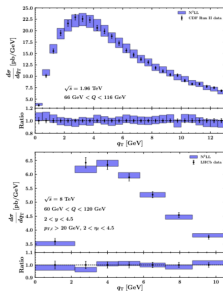
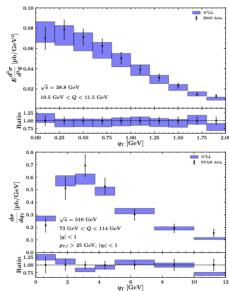
(showing rows 0 to 9 of 134 total)  
More [Enter=10, nnext event, qquit]:

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81	13	134	81	-3112	1d-3112	83	57	73	116	117	0	-3.258	0.221	-0.778
82	13	134	82	-211	pi-	83	57	73	0	0	0	-0.689	-0.776	0.356
83	13	134	83	211	pi+	83	57	73	0	0	0	-0.643	0.377	0.356
84	13	134	84	113	rho0	84	57	73	197	98	0	0.230	0.721	1.962
85	13	134	85	111	pi0	84	57	73	118	119	0	0.980	0.129	-0.747
86	13	134	86	-123	rho-	84	57	73	199	100	0	0.512	0.385	0.404
87	13	134	87	2212	pi+	84	57	73	0	0	0	1.593	1.236	0.164
88	13	134	88	-211	pi-	84	57	73	0	0	0	0.647	-0.002	0.279
89	13	134	89	-2114	1d-2114	84	57	73	101	102	0	3.813	0.916	0.066

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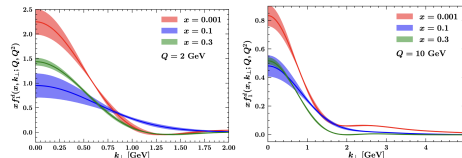


# $q_T$ spectra $\rightarrow$ TMD extraction (Drell–Yan example)



$q_T$  data  
constrains  
TMD  
shape/parameters.

Target: summarise shapes with few parameters. [6]



Low- $q_T$  shape discriminates global fits. [6]

**Goal:** use an analogous  $q_T$ -sensitive observable in  $e^+e^-$  to extract/compare TMD information. [6]

- 1 O. Biebel, D. de Florian, D. Milstead, W. Vogelsang (PDG), *Fragmentation Functions in  $e^+e^-$ ,  $ep$ , and  $pp$  Collisions (review)*, Prog. Theor. Exp. Phys. 2020, 083C01.
- 2 O. Biebel, P. Nason, B. R. Webber, *Jet fragmentation in  $e^+e^-$  annihilation*, arXiv:hep-ph/0109282v2 (2001).
- 3 A. Vossen, *Parton Fragmentation Functions*, arXiv:1702.01329 (2017).
- 4 P. Z. Skands, *Introduction to QCD*, arXiv:1207.2389 (2012).
- 5 R. Boussarie et al. (TMD Collaboration), *TMD Handbook*, arXiv:2304.03302 (2023).
- 6 A. Bacchetta et al., *Transverse-momentum-dependent parton distributions up to  $N^3LL$  from Drell-Yan data*, JHEP 07 (2020) 117, arXiv:1912.07550.