

Test 1:

Planck function: $B_\nu(T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT} - 1}, \quad B_\lambda(T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{hc/\lambda kT} - 1}$

with units $\text{erg s}^{-1}\text{Hz}^{-1}\text{cm}^{-2}\text{sr}^{-1}$ and $\text{erg s}^{-1}\text{cm}^{-1}\text{cm}^{-2}\text{sr}^{-1}$, respectively. We also have $h = 6.63 \times 10^{-27} \text{ erg s}$, $c = 3.00 \times 10^{10} \text{ cm s}^{-1}$, $k = 1.38 \times 10^{-16} \text{ erg K}^{-1}$.

Note that $B_\nu(T) = I_\nu(T)$ at thermal equilibrium (when collisions dominate).

Rayleigh-Jeans law: $B_\lambda(T) \simeq \frac{2ckT}{\lambda^4}$, obtained when $hc \ll \lambda kT$

Wien's displacement law: $\lambda_{\text{max}} = \frac{0.29}{T} \text{ cm}$

Stefan-Boltzmann law: $F = \sigma T^4$, $\sigma = \frac{2\pi^5 k^4}{15h^3 c^2} = 5.67 \text{ erg s}^{-1}\text{cm}^{-2}\text{K}^{-4}$.

Also note that $F = \pi \int_0^\infty B_\lambda(T) d\lambda$

Luminosity: $L_* = 4\pi R_*^2 \sigma T_{\text{eff}}^4$

Magnitude: $m_1 - m_2 = 2.5 \log \left(\frac{F_1}{F_2} \right)$

Distance modulus: $m - M = 5 \log \left(\frac{d}{10 \text{ pc}} \right)$

Energy of hydrogen electron: $E_n = 13.6(1 - 1/n^2)$

Degeneracy for hydrogen: $g_n = 2n^2$

Boltzmann equation: $\frac{n_i}{n_j} = \frac{g_i}{g_j} e^{-(E_i - E_j)/kT}$, $\frac{n_i}{n_{\text{ion}}} = \frac{g_i}{Z} e^{-E_i/kT}$

Partition function: $Z = \sum_i g_i e^{-E_i/kT}$

Saha equation: $\frac{n_{i+1}n_e}{n_i} = \left(\frac{2\pi m_e kT}{h^2} \right)^{3/2} \frac{2Z_{i+1}}{Z_i} e^{-\chi_i/kT}$

where $m_e = 9.109 \times 10^{-28}$, and χ_i is the ionization potential.

Fractional ionization: $f_i = \frac{n_i}{n_1 + n_2 + n_3 + \dots} = \frac{\left(\frac{n_i}{n_{i-1}} \right) \left(\frac{n_{i-1}}{n_{i-2}} \right) \dots \left(\frac{n_2}{n_1} \right)}{1 + \left(\frac{n_2}{n_1} \right) + \left(\frac{n_3}{n_2} \right) \left(\frac{n_2}{n_1} \right) + \left(\frac{n_4}{n_3} \right) \left(\frac{n_3}{n_2} \right) \left(\frac{n_2}{n_1} \right) + \dots}$

Mean molecular weight: $\mu = \frac{1}{n_{\text{tot}} m_H} \sum_i n_i m_i$

Maxwellian velocity distribution: $N(v) = N 4\pi v^2 \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-mv^2/2kT}$

Hydrostatic equilibrium: $\frac{dP(r)}{dr} = -\rho(r)g(r)$

Universal law of gravitation: $g(r) = \frac{GM(r)}{r^2}$, $G = 6.67 \times 10^{-8} \text{ dyn cm}^2 \text{ g}^{-2}$

Gravitational potential energy of sphere: $\Omega = - \int_0^{M_*} \frac{GM(r)}{r} dM$

For a sphere: $V = \frac{4}{3}\pi r^3$, $A_{\text{surface}} = 4\pi r^2$

Equipartition theorem: $K = \frac{3}{2}kT$

Ideal gas law: $P(r) = n(r)kT = \frac{\rho(r)}{\mu m_H} kT$

Column mass: $m(r) \equiv M(r)/dA$, where $M(r) = dA \int \rho(r) dr$

Pressure of an ideal gas: $P = \frac{2}{3}\varepsilon$ for energy per unit volume $\varepsilon = \frac{NK}{V} = \frac{3}{2}nkT$.

Virial theorem: $2U + \Omega = 0$ for a system in thermal equilibrium.

Jeans mass: $M_J = \left(\frac{kT}{\mu m_H G} \right)^{3/2} \left(\frac{3}{4\pi\rho} \right)^{1/2}$

Jeans length: $\lambda_J = \left(\frac{3M_J}{4\pi\rho} \right)^{1/3}$

Free-fall time: $t_{\text{ff}} = \sqrt{\frac{3\pi}{32G\rho}}$

Test 2:

Specific intensity change: $dI_\nu = -(\alpha_\nu + \varsigma_\nu)I_\nu ds$

where α_ν and ς_ν are the absorption and scattering coefficients, respectively.

Specific intensity profile: $I_\nu = I_\nu(0)e^{-(\alpha_\nu + \varsigma_\nu)s}$ for constant extinction and no emissivity.

Note I_ν has units of $\text{erg s}^{-1}\text{Hz}^{-1}\text{cm}^{-2}\text{sr}^{-1}$.

Mean free path as the length-scale: $\ell_\nu = (\alpha_\nu + \varsigma_\nu)^{-1}$

Probability of absorption: $\epsilon_\nu = \frac{\alpha_\nu}{\alpha_\nu + \varsigma_\nu}$

Mean displacement: $\ell_\nu^{*2} = \frac{\ell_\nu}{\sqrt{\epsilon_\nu}} = \frac{1}{\sqrt{\alpha_\nu(\alpha_\nu + \varsigma_\nu)}}$

Opacity: $k_\nu = \frac{1}{\rho}(\alpha_\nu + \varsigma_\nu) = \kappa_\nu + \sigma_\nu$ [cm^2g^{-1}]

Mean intensity: $J_\nu(z) = \frac{1}{4\pi} \int I_\nu(z, \theta) d\Omega = \frac{1}{2} \int_{-1}^1 I_\nu(z, u) du$ where $u = \cos(\theta)$.

Net flux: $F_\nu(z) = \int I_\nu(z, \theta)(\mathbf{n} \cdot \mathbf{e}_z) d\Omega = 2\pi \int_{-1}^1 I_\nu(z, u)u du$ [$\text{erg s}^{-1}\text{Hz}^{-1}\text{cm}^{-2}$]

Radiative moments: $M_\nu(z, n) = \frac{1}{2} \int_{-1}^1 I_\nu(z, u)u^n du$,

where $M_\nu(z, 0) = J_\nu(z)$, $M_\nu(z, 1) = H_\nu(z) = F_\nu(z)/4\pi$, and $M_\nu(z, 2) = K_\nu(z)$.

Eddington's approximation: $J_\nu(z) = 3K_\nu(z)$

Eddington flux: $H_\nu(\tau_\nu) \simeq -\frac{1}{3k_\nu\rho} \frac{dB_\nu(T)}{dT} \frac{dT}{dz}$

Total Eddington flux: $H(z) = -\frac{1}{3k_R\rho} \frac{dB(T)}{dT} \frac{dT}{dz}$

Rosseland mean opacity: $\frac{1}{k_R} = \frac{\pi}{4\sigma T^3} \int_0^\infty \frac{1}{k_\nu} \frac{dB_\nu(T)}{dT} d\nu$

Radiative transfer equation: $\frac{1}{\rho} \frac{dI_\nu}{ds} = -k_\nu I_\nu + j_\nu$,

where ρ is mass density and j_ν is the emissivity in units of $\text{erg s}^{-1}\text{Hz}^{-1}\text{sr}^{-1}\text{g}^{-1}$.

Radiative transfer in terms of optical depth: $u \frac{dI_\nu(\tau_\nu, u)}{d\tau_\nu} = I_\nu(\tau_\nu, u) - S_\nu(\tau_\nu)$,

for optical depth $d\tau_\nu \equiv -k_\nu \rho dz$ and source function $S_\nu(\tau_\nu) \equiv \frac{j_\nu}{k_\nu}$.

General solution: $I_\nu(\tau_\nu, u) = I_\nu(0, u)e^{\tau_\nu/u} - \frac{e^{\tau_\nu/u}}{u} \int_0^{\tau_\nu} S(\lambda)e^{-\lambda/u} d\lambda$

For constant source: $I_\nu(\tau_\nu, u) = I_\nu(0, u)e^{\tau_\nu/u} + S_\nu(1 - e^{\tau_\nu/u})$

Optically thick / LTE conditions: $I_\nu(|\tau_\nu| > 1, u) \approx S_\nu$

Optically thin / non-LTE conditions: $I_\nu(|\tau_\nu| \ll 1, u) \simeq I_\nu(0, u) + S_\nu \frac{|\tau_\nu|}{u}$

Grey atmosphere temperature profile: $T(\tau) = T_{\text{eff}} \left[\frac{3}{4} \left(\tau + \frac{2}{3} \right) \right]$

Heisenberg inequality: $\Delta E \Delta t \geq \frac{\hbar}{2}, \quad \Delta p_x \Delta x \geq \frac{\hbar}{2}$

Lorentz profile for natural broadening: $\varphi_\nu = \frac{\Gamma}{4\pi^2} \frac{1}{(\nu - \nu_0)^2 + (\Gamma/4\pi)^2},$

where $\Gamma/2\pi$ is the full-width-half-magnitude of the curve.

Opacity redefined: $k_\nu = \alpha(\nu)/\rho$, where now $\alpha(\nu) = \frac{\pi e^2}{mc^2} f_{ij} \varphi_\nu$

Radial velocity: $n(v_z) = n \sqrt{\frac{m}{2\pi kT}} e^{-mv_z^2/2kT}$

Non-relativistic doppler effect: $\nu' = \nu \left(1 - \frac{U}{c} \right), \quad \frac{\Delta \nu}{\nu} = \frac{\Delta V}{c}$

Voigt profile: $\varphi_\nu = \frac{U(a, v)}{\Delta v_D}$

Final:

Equation of hydrostatic equilibrium: $\rho \frac{d^2 r}{dt^2} = -\frac{\rho(r)GM(r)}{r^2} - \frac{dP}{dr}$

Equation of mass conservation: $\frac{dM(r)}{dr} = 4\pi r^2 \rho(r)$

Energy-transport equation: $\frac{dT(r)}{dr} = -\frac{3k_R \rho}{64\pi r^2 \sigma T^3} L(r)$

Equation of energy conservation: $\frac{dL(r)}{dr} = 4\pi r^2 \rho(r) \varepsilon(r),$

where $\varepsilon(r)$ is the rate of thermonuclear energy generated per unit mass in $\text{erg s}^{-1} \text{g}^{-1}$.

Temperature gradient: $\nabla \equiv \frac{d \ln(T)}{d \ln(P)}$

For pure radiative transport and hydrostatic equilibrium: $\nabla_{\text{rad}} = \frac{3k_R}{64\pi r^2 g} \frac{P}{\sigma T^4} L(r).$

Note that $\frac{dy}{dx} = y \frac{d \ln y}{dx}.$

Monochromatic Eddington flux: $H_\nu(r) = \frac{k^3}{8h^2 c^2} \frac{k_R}{k_\nu} \frac{T_{\text{eff}}^4}{T} \left(\frac{R}{r} \right)^2 \mathcal{P}(u),$ with $\mathcal{P}(u) = \frac{u^4 e^u}{(e^u - 1)^2}, \quad u = \frac{h\nu}{kT}$

Volume of a cell in phase space: $\Delta x \Delta y \Delta z \Delta p_x \Delta p_y \Delta p_z = h^3$

Non-degenerate: $n_e(p) \leq \frac{8\pi p^2 dp}{h^3}$ **Degenerate:** $n_e(p) \gtrsim \frac{8\pi p^2 dp}{h^3}$

Conductive Eddington flux: $H_{\text{cond}} = \frac{4\sigma T^3}{3\pi k_{\text{cond}} \rho} \frac{dT}{dr}$

Radiative Eddington flux: $H_{\text{rad}} = \frac{4\sigma T^3}{3\pi k_{\text{rad}} \rho} \frac{dT}{dr}$

Total Eddington flux: $H_{\text{tot}} = H_{\text{rad}} + H_{\text{cond}} = \frac{4\sigma T^3}{3\pi k_{\text{tot}} \rho} \frac{dT}{dr},$ where $\frac{1}{k_{\text{tot}}} = \frac{1}{k_{\text{rad}}} + \frac{1}{k_{\text{cond}}}$

Bouyant force: $ma = (\rho_f - \rho_c)g\Delta V$

Condition for convection: $\left(\frac{d\rho}{dr} \right)_{\text{adi}} < \left(\frac{d\rho}{dr} \right)_{\text{rad}}$

Energy-transport in convective equilibrium: $\frac{dT}{dr} = - \left(\frac{\gamma - 1}{\gamma} \right) \frac{\rho g T}{P}$

Schwarzschild criterion for convection: $\frac{3k_R}{64\pi r^2 g} \frac{P}{\sigma T^4} L(r) > \frac{\gamma - 1}{\gamma},$ where $\gamma = c_P/c_V$

Poisson equation: $\frac{1}{r^2} \frac{d}{dr} \left[\frac{r^2}{\rho(r) \frac{dP(r)}{dr}} \right] = -4\pi G \rho(r)$

Polytropic equation of state: $P = K \rho^{(n+1)/n}$

Lane-Emden equation: $\frac{d^2\theta(\xi)}{d\xi^2} + \frac{2}{\xi} \frac{d\theta(\xi)}{d\xi} + \theta^n(\xi) = 0$

Equation of state for an ideal gas: $P = P_{\text{ions}} + P_{\text{e}} = (n_{\text{ions}} + n_{\text{e}})kT = \left(\frac{1}{\mu_{\text{ions}}} + \frac{1}{\mu_{\text{e}}} \right) \frac{\rho kT}{m_{\text{H}}}$

Mass fraction of a species: $X_i = \frac{\rho_i}{\rho} = \frac{n_i m_i}{\sum_j n_j m_j}$

Total mean molecular weight: $\frac{1}{\mu} = \frac{1}{\mu_{\text{ions}}} + \frac{1}{\mu_{\text{e}}},$

where $\frac{1}{\mu_{\text{ions}}} = \sum_i \frac{X_i}{A_i}$ and $\frac{1}{\mu_{\text{e}}} = \sum_i \frac{z_i X_i}{A_i} = \frac{n_{\text{e}}}{\rho/m_{\text{H}}} = \frac{\text{number of free electrons}}{\text{number of nucleons}}$

For full ionization: $\mu \simeq \frac{2}{3X + Y/2 + 1},$

For neutral: $\mu \simeq \frac{1}{X + Y/4}$

with mass fractions X and Y for hydrogen and helium, respectively.

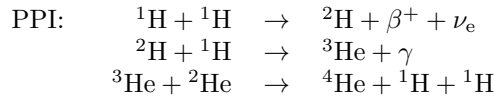
Equation of state for degenerate gas: $P_{\text{e}} \propto n_{\text{e}}^\gamma \propto \left(\frac{\rho}{\mu_{\text{e}}} \right)^\gamma$

Radiation pressure: $P_{\text{rad}} = \frac{4\pi}{c} K_\nu$

If radiation pressure dominates over gas pressure: $P_{\text{rad}} \approx \frac{4\sigma T^4}{3c}$

Atomic species: ${}_Z^AX$, A = atomic mass, Z = atomic number, X = element symbol

Proton-proton chains:



Energy generated: $Q = [M_{\text{initial}} - M_{\text{final}}]c^2 = [4M({}^1\text{H}) - M({}^4\text{He})]c^2 = 26.73 \text{ MeV}$

Energy conversion rate: $\varepsilon = \frac{Q}{M({}^AX)c^2}$

Energy produced through fusion: $E_{\text{ms}} \sim L_* t_{\text{ms}}$

Time on main sequence: $t_{\text{ms}} \sim \varepsilon \frac{M_* c^2}{L_*}$