fluid

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Citing

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Contents

1	Flow past a cylinder	1
2	Conformal mapping	2
3	Computation of flux-barrier base points 3.1 Magnet insertion	3 4 4
4	Outer base points	5
5	Flux-barrier sideline points	6
6	Example of Matlab/Octave plot	7

1 Flow past a cylinder

Let ρ_0 be the radius of the cylinder, ρ, ξ the polar coordinate system in use. One possible solution of this problem have these potential and streamline

functions:

$$\phi(\rho,\xi) = \left(\rho + \frac{\rho_0^2}{\rho}\right)\cos\xi\tag{1}$$

$$\psi(\rho,\xi) = \left(\rho - \frac{\rho_0^2}{\rho}\right)\sin\xi\tag{2}$$

Although these equations are deeply coupled, the radius ρ and the phase ξ can be obtained as a function of the other quantities. For our purposes, we use ψ .

$$\rho(\psi,\xi) = \frac{\psi + \sqrt{\psi^2 + 4\rho_0^2 \sin^2 \xi}}{2\sin \xi}$$
 (3)

$$\xi(\psi, \rho) = \arcsin\left(\frac{\rho \,\psi}{\rho^2 - \rho_0^2}\right) \tag{4}$$

The velocity field can also be derived through

$$v_{\rho}(\rho,\xi) = \frac{\partial \phi}{\partial \rho} = \left(1 - \frac{\rho_0^2}{\rho^2}\right) \cos \xi$$

$$v_{\xi}(\rho,\xi) = \frac{1}{\rho} \frac{\partial \phi}{\partial \xi} = -\left(1 + \frac{\rho_0^2}{\rho^2}\right) \sin \xi$$
(5)

2 Conformal mapping

From the reference plane, which is equivalent to a two-pole machine, we use a complex map to obtain the quantities in the actual plane. Let p be the number of pole pairs. Then:

$$\zeta \xrightarrow{\mathcal{M}} z = \sqrt[p]{\zeta}$$

$$\rho e^{j\xi} \xrightarrow{\mathcal{M}} r e^{j\vartheta} = \sqrt[p]{\rho} e^{j\xi/p}$$

$$\chi + j\eta \xrightarrow{\mathcal{M}} x + jy$$
(6)

It is easy to find the inverse map:

$$\mathcal{M} \colon \sqrt[p]{\cdot} \qquad \mathcal{M}^{-1} \colon (.)^p \tag{7}$$

In the transformed plane, the velocities have a different expression:

$$v_r(r,\vartheta) = p \left(r^{p-1} - \frac{R_0^{2p}}{r^{p+1}} \right) \cos p\vartheta$$

$$v_{\vartheta}(r,\vartheta) = -p \left(r^{p-1} + \frac{R_0^{2p}}{r^{p+1}} \right) \sin p\vartheta$$
(8)

This vector field is tangent to the streamlines in every point in the transformed plane. In order to work with this field in x, y coordinates, we need a rotational map:

$$v_x(r,\vartheta) = v_r \cos \vartheta - v_\vartheta \sin \vartheta$$

$$v_\vartheta(r,\vartheta) = v_r \sin \vartheta + v_\vartheta \cos \vartheta$$
(9)

3 Computation of flux-barrier base points

Refer to Figure 1 for the points naming scheme. Keep in mind that A' is not simply the projection of A onto the q-axis, but it represent the original starting point for the barrier sideline, so it lies on the flux-barrier streamline. The same is true for points B', B, C', C, and D', D.

Let the flux-barrier and flux-carrier thicknesses be given. Then the base points for the flux-barriers can be computed easily. Let $D_{\rm r}$ be the rotor outer diameter, $w_{\rm rib,t}$ the tangential iron rib width, $w_{\rm c,k}$ the k-th flux-carrier width, and $t_{\rm c,k}$ the k-th flux-barrier thickness.¹ Then

$$R_{\rm r} = \frac{D_{\rm r}}{2} - w_{\rm rib,t}$$

$$R_{\rm A'_{1}} = R_{\rm r} - w_{\rm c,1}$$

$$R_{\rm B'_{1}} = R_{\rm A'_{1}} - t_{\rm c,1}$$

$$\vdots$$
(10)

¹You may wonder why the main dimensions of the flux-carrier and flux-barrier differ in the name (width versus thickness). This is due to a choice of mine, because I prefer to refer to width when the flux flows perpendicularly to the dimension, and to thickness when it flows in parallel.

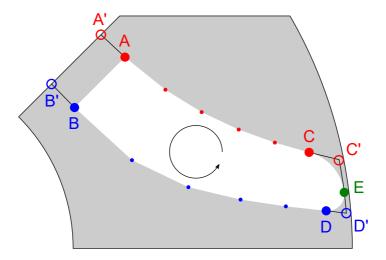


Figure 1: Flux-barrier base points description.

where R represents the radius from the origin. So, in general:

$$R_{\mathsf{A}'_{\mathsf{k}}} = R_{\mathsf{B}'_{\mathsf{k}-1}} - w_{\mathsf{c},k}$$

$$R_{\mathsf{B}'_{\mathsf{k}}} = R_{\mathsf{A}'_{\mathsf{k}-1}} - t_{\mathsf{c},k}$$
(11)

with the exception $R_{\mathsf{B}_0'} = R_{\mathrm{r}}$.

Now we know both the radii and the angle – always $\pi/(2p)$ – of the flux-barrier internal points. So we can compute their respective streamline value.

3.1 Magnet insertion

$$w_{\mathrm{rib},k} = w_{\mathrm{rib},k} + w_{\mathrm{m},k}$$

where $w_{m,k}$ is the k-th magnet width.

3.2 Central base points

We refer to points A and B. If the rib width is zero $A \equiv A'$ and $B \equiv B'$.

The line describing the q-axis is

$$y = mx + q$$

$$m = \tan \frac{\pi}{2p}$$

$$q = \frac{w_{\text{rib}}}{2\cos \frac{\pi}{2p}}$$
(12)

$$\begin{cases} y_{\mathsf{A}} - mx_{\mathsf{A}} - q = 0 \\ x_{\mathsf{A}} - r_{\mathsf{A}}(\psi_{\mathsf{A}'}, \vartheta_{\mathsf{A}})\cos\vartheta_{\mathsf{A}} = 0 \\ y_{\mathsf{A}} - r_{\mathsf{A}}(\psi_{\mathsf{A}'}, \vartheta_{\mathsf{A}})\sin\vartheta_{\mathsf{A}} = 0 \end{cases}$$
(13)

where ϑ_{A} is used as the third degree of freedom and r_{A} is then a function of it. The solution of such system can be determined solving the single equation

$$r_{\mathsf{A}}(\psi_{\mathsf{A}'}, \vartheta_{\mathsf{A}}) \left(\sin \vartheta_{\mathsf{A}} - m \cos \vartheta_{\mathsf{A}} \right) - q = 0 \tag{14}$$

in the unknown ϑ_A . The function $r(\psi, \vartheta)$ is simply

$$r(\psi, \vartheta) = \sqrt[p]{\rho(\psi, \vartheta/p)}$$

The same equation can be written for point B with the proper substitution and repeated for all the flux-barriers.

4 Outer base points

We refer to points C, D, and E. If the flux-barrier angle, $\alpha_{\rm b}$, is given, then

$$x_{\mathsf{E}} = R_r \cos(\frac{\pi}{2p} - \alpha_{\mathsf{b}})$$

$$y_{\mathsf{E}} = R_r \sin(\frac{\pi}{2p} - \alpha_{\mathsf{b}})$$
(15)

Points C and D results from the connection of the flux-barrier sidelines and point E. This connection should be as smooth as possible in order to avoid dangerous mechanical stress concentrations. We are going to use circular arcs to make this connection. So we impose the tangency between the flux-barrier sideline and the arc, between the arc and the radius through point E. The tangent to the sideline can be obtained through the velocity field described above.

Then we want point C to lay on the flux-barrier sideline. These conditions represent a nonlinear system of 4 equations, in 6 unknowns. So we need two more equations, which are that points C and E belong to the fillet circle with radius R.

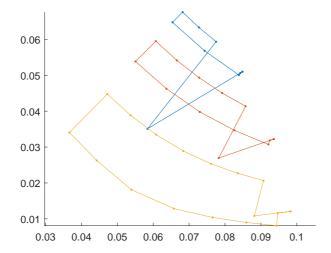
$$\begin{cases} x_{\mathsf{C}} - r_{\mathsf{C}}(\psi_{\mathsf{C}}, \vartheta_{\mathsf{C}}) \cos \vartheta_{\mathsf{C}} = 0 \\ y_{\mathsf{C}} - r_{\mathsf{C}}(\psi_{\mathsf{C}}, \vartheta_{\mathsf{C}}) \sin \vartheta_{\mathsf{C}} = 0 \\ (x_{\mathsf{C}} - x_{\mathsf{O}_{\mathsf{C}}})^{2} + (y_{\mathsf{C}} - y_{\mathsf{O}_{\mathsf{C}}})^{2} - R_{\mathsf{EC}}^{2} = 0 \\ (x_{\mathsf{E}} - x_{\mathsf{O}_{\mathsf{C}}})^{2} + (y_{\mathsf{E}} - y_{\mathsf{O}_{\mathsf{C}}})^{2} - R_{\mathsf{EC}}^{2} = 0 \\ (x_{\mathsf{O}_{\mathsf{C}}} - x_{\mathsf{E}})y_{\mathsf{E}} - (y_{\mathsf{O}_{\mathsf{C}}} - y_{\mathsf{E}})x_{\mathsf{E}} = 0 \\ (x_{\mathsf{O}} - x_{\mathsf{C}})v_{x}(r_{\mathsf{C}}, \vartheta_{\mathsf{C}}) + (y_{\mathsf{O}} - y_{\mathsf{C}})v_{y}(r_{\mathsf{C}}, \vartheta_{\mathsf{C}}) \end{cases}$$
(16)

The very same system can be written and solved for point D.

5 Flux-barrier sideline points

to be continued...

6 Example of Matlab/Octave plot



Here is an example of a Matlab/Octave output plot. The V-shaped lines represent the radii of the fillet arcs, which were not worth to be shown in Matlab/Octave.