

PHYS 5319-001: Math Methods in Physics III

Differentiation & ODE

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Differentiation

For a function $f(x)$, the derivative is defined as

$$\frac{df}{dx} \equiv \lim_{h \rightarrow 0} \frac{f(x+h) - f(x)}{h}$$

Problem for numeric solution: subtractive cancellation

number precision: e.g. xx.xxxxxxxxxxxxxx???

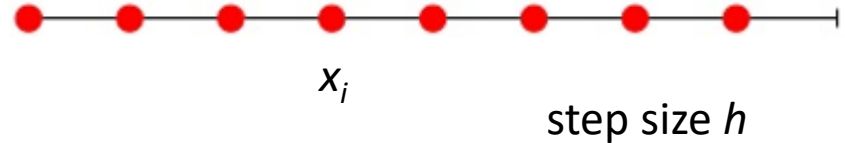
If $f(x+h)$ only differs $f(x)$ at the 9th decimal, \Rightarrow only 6 digits left!

When the step size $h \rightarrow 0 \Rightarrow$ the difference will reach the machine precision ϵ_m !

Also, too small h will increase the number of grid.



-Forward Difference



Taylor expansion

$$f(x + h) = f(x) + hf'(x) + \frac{h^2}{2!}f''(x) + \dots$$

So, the computed expression

$$\begin{aligned}f'_c(x) &\approx \frac{f(x + h) - f(x)}{h} \\ &\approx f'(x) + O(h)\end{aligned}$$

*Only need the
forward step info*

e.g. $f(x) = a + bx^2$ Exact: $f'(x) = 2bx$

$$f'_c(x) \approx \frac{f(x + h) - f(x)}{h} = 2bx + bh$$

*Only good when
 $h \ll 2x$*

-Central Difference

$$f'_c(x) \approx \frac{f\left(x + \frac{h}{2}\right) - f\left(x - \frac{h}{2}\right)}{h}$$

*Stepping forward by $h/2$
and backward by $h/2$*

$$\equiv D_c f(x, h)$$

From Taylor expansion of $f\left(x \pm \frac{h}{2}\right)$:

$$f'_c(x) \approx f'(x) + \frac{1}{3!} \left(\frac{h}{2}\right)^2 f^{(3)}(x) + \dots$$

$O(h^2)$

All the odd terms of h cancelled!

e.g. $f(x) = a + bx^2$

$$f'_c(x) \approx 2bx$$

-Extrapolated Difference

To manipulate $O(h^2)$ vanish since we know the analytic form of errors

instead of $D_c f(x, h)$, we introduce

$$D_c f(x, \frac{h}{2}) \equiv \frac{f\left(x + \frac{h}{4}\right) - f\left(x - \frac{h}{4}\right)}{h/2} \approx f'(x) + \frac{1}{3!} \left(\frac{h}{4}\right)^2 f^{(3)}(x) + \dots$$

$$f'_c(x) \approx f'(x) + \frac{1}{3!} \left(\frac{h}{2}\right)^2 f^{(3)}(x)$$

Mixing $D_c f(x, h)$ & $D_c f(x, h/2)$ to vanish $O(h^2)$:

$$f'_c(x) \approx \frac{4D_c\left(x, \frac{h}{2}\right) - D_c(x, h)}{3}$$

$$\approx f'(x) + \frac{1}{3} \left\{ 4 \frac{1}{5!} \left(\frac{h}{4}\right)^4 f^{(5)}(x) - \frac{1}{5!} \left(\frac{h}{2}\right)^4 f^{(5)}(x) \right\}$$

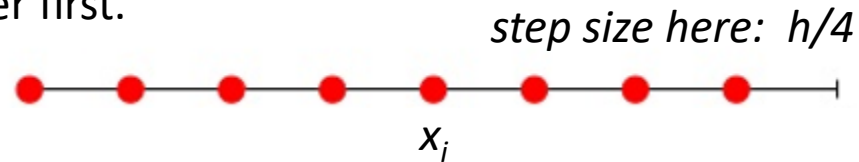
$$- \frac{h^4}{120 \times 4 \times 16} f^{(5)}(x) \quad O(h^4)$$

A better way to compute $f'_c(x) \approx \frac{4D_c\left(x, \frac{h}{2}\right) - D_c(x, h)}{3}$

→
$$f'_c(x) = \frac{1}{3h} \left\{ 8 \left[f\left(x + \frac{h}{4}\right) - f\left(x - \frac{h}{4}\right) \right] - \left[f\left(x + \frac{h}{2}\right) - f\left(x - \frac{h}{2}\right) \right] \right\}$$

Why?

Avoid a large number +/- a small number. Numbers in the same order of magnitude subtract each other first.



Cost:

Not only the nearest neighbor points, but also the 2nd n.n. points info is needed.

Question: more neighboring points, more accurate?

Fourier Transform

For a periodic (in L) function $f(x)$

FT

$$g(k) = \frac{1}{L} \int_0^L f(x) e^{ikx} dx$$

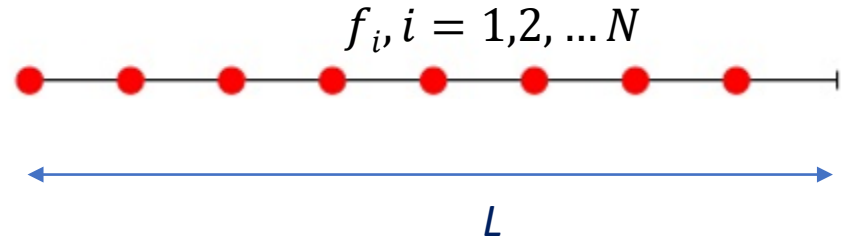
$$= \sum_{j=1}^N f_j e^{ikx_i}$$

FT⁻¹

$$f(x) = \sum_{l=1}^N g_l e^{-ik_l x}$$

Now, $\frac{d}{dx}$ on $f(x)$:

$$f'(x) = \sum_{l=1}^N (-ik_l) g_l e^{-ik_l x} = \sum_l G_l e^{-ik_l x}$$



Since $f(x+L)=f(x) \Rightarrow e^{ikL}=1$,
 $kL = 2\pi n, n = 0, 1, 2, \dots, N-1$



We have N discrete k -points:

$$k = \frac{2\pi n}{L}, n = 0, 1, 2, \dots, N-1$$

$g_l \equiv g(k_l)$, represented in Fourier space

Need info from N points!

-second derivatives

$$f''(x) \equiv \frac{d^2 f}{dx^2} \equiv \frac{d}{dx} \left(\frac{df}{dx} \right)$$

Central-Difference

$$f''(x) \approx \frac{f' \left(x + \frac{h}{2} \right) - f' \left(x - \frac{h}{2} \right)}{h}$$

apply Forward-Difference to the numerators:

$$f''(x) \approx \frac{[f(x+h) - f(x)] - [f(x) - f(x-h)]}{2h^2}$$

But for FT case, just a multiplication of $(-ik_l)$ again

$$f''(x) = \sum_{l=1}^N (-ik_l)^2 g_l e^{-ik_l x}$$

Differential Equations


A real physics problem often involves differential equation(s).
Could be a partial differential equation (PDE),
Or an ordinary differential equation (ODE).

Analytically, a PDE \Rightarrow ODEs by separation of variables *e.g. Hydrogen atom*

Numerically, you may directly solve some. e.g. Poisson equation:

$$\nabla^2 \varphi(\vec{r}) = -\rho(\vec{r})/4\pi\epsilon_0$$

In Fourier space (transform needed) $-k^2 \varphi(\vec{k}) = -\rho(\vec{k})/4\pi\epsilon_0$


$$\varphi(\vec{k}) = \rho(\vec{k})/(4\pi\epsilon_0 k^2)$$

Then an FT⁻¹ back to the real space, $\varphi(\vec{r})$.

But in most cases, we still need to dealing with ODEs

-Types of ODE

order \Rightarrow degree of derivatives

e.g. 1st-order: $\frac{dy}{dt} = f(t, y)$

More generally, 2nd-order: $\frac{d^2y}{dt^2} = f(t, \frac{dy}{dt}, y)$

“Linear” vs. “nonlinear”

Linear \Rightarrow law of linear superposition

Nonlinear \Rightarrow contains higher power of $y(t)$. A linear super position not working.

Standard form:

$$\frac{d\vec{y}}{dt} = \vec{f}(t, \vec{y})$$

Or a set of coupled 1st order ODEs:

$$\frac{dy_i(x)}{dx} = f_i(x, y_1, y_2, \dots, y_N)$$

-Numerical solution: (integration):

1. Start from the known initial value;
2. Use the derivative function to advance one step
3. Repeat till to the end



ODEs with initial values

- Starting with 1st order ODE

$$\frac{d}{dt}x(t) = f(x, t)$$

with initial condition $x(0) = x_0$.

- The Taylor series expansion of x around t gives

$$x(t + \Delta t) = x(t) + \frac{dx(t)}{dt} \Delta t + \frac{1}{2!} \frac{d^2x(t)}{dt^2} (\Delta t)^2 + \dots \quad (1)$$

Assume a reasonably smooth function $x(t)$ and a small interval Δt , we can propagate $x(t)$ to $x(t + \Delta t)$ to any accuracy desired, as long as we know the derivatives of $x(t)$.

- However, often the (higher order) derivatives for a general t is not known, and we want a propagation for a time interval $\gg \Delta t$.
- The problem then becomes: how to estimate the derivatives of an unknown function and the suitable repetitions of propagation by Δt .
- Several numerical methods here:
 - Euler method
 - Runge-Kutta approximation
 - Euler-Cromer method
 - Verlet method
 - ...



Numerical solutions for 1st order ODE

- When Δt is small, it is a good approximation by simply ignoring the second and higher orders of Δt in Eq. (1) as

$$x(t + \Delta t) \approx x(t) + \frac{dx}{dt} \Delta t = x(t) + f(x(t), t) \Delta t \quad (2)$$

- This is the First-order Approximation – **Euler method**.
- Obtaining $x(t + \Delta t)$ from $x(t)$ with Eq. (2), we can estimate $x(t + 2\Delta t)$ from $x(t + \Delta t)$ as
$$x(t + \Delta t + \Delta t) = x(t + \Delta t) + f(x(t + \Delta t), t + \Delta t) \Delta t$$
- Iteratively, we can obtain $x(t + N\Delta t)$. ($N\Delta t \rightarrow$ the interested)



- *e.g.* for radioactive decay

$$\frac{dN(t)}{dt} = -\lambda N(t)$$

- Applying the Euler method, we can estimate the number of active nuclei $N(t + \Delta t)$ at $t + \Delta t$ as

$$N(t + \Delta t) \approx N(t) + \frac{dN(t)}{dt} \Delta t = N(t) - \lambda N(t) \Delta t$$

- And $N(t + 2\Delta t)$ at $t + 2\Delta t$ as

$$N(t + 2\Delta t) \approx N(t + \Delta t) + \frac{dN(t)}{dt} \Delta t = N(t + \Delta t) - \lambda N(t + \Delta t) \Delta t$$

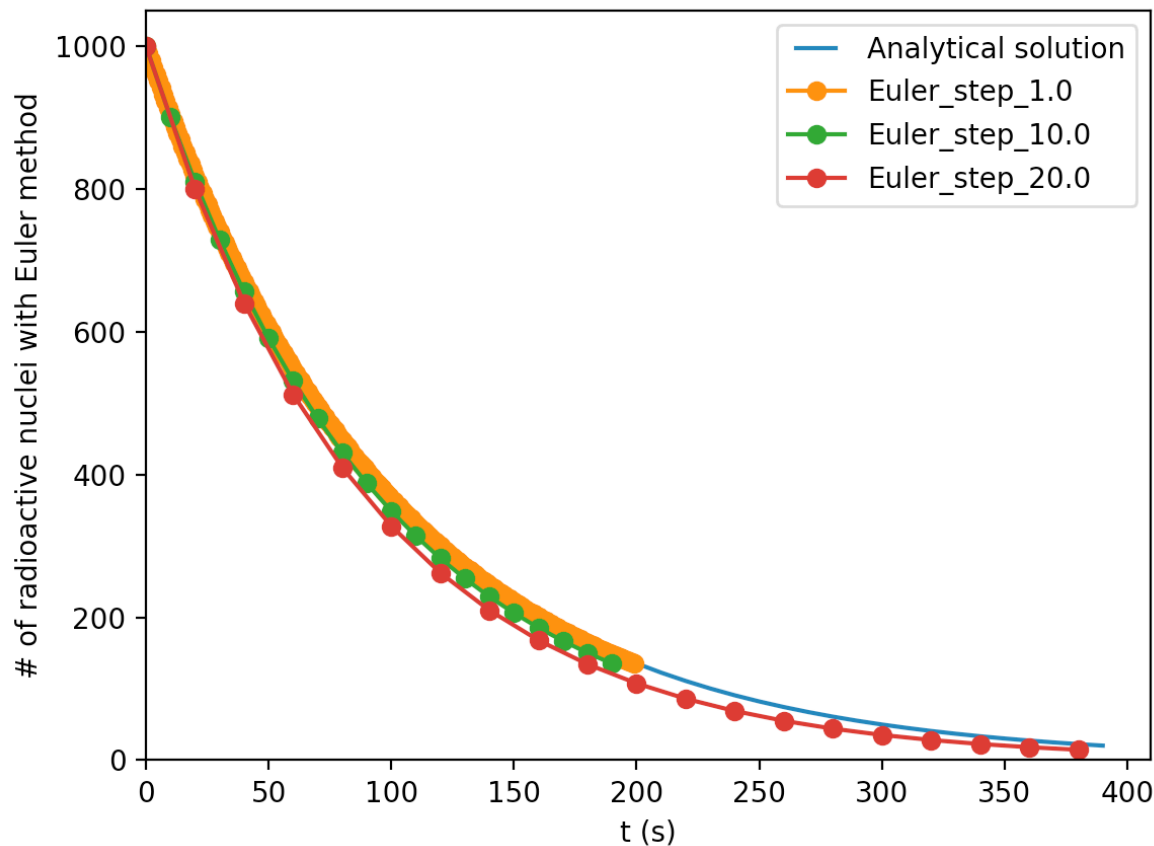
- ...

- And $N(t + i\Delta t)$ at $t + i\Delta t$ as

$$N(t + i\Delta t) \approx N(t + (i - 1)\Delta t) - \frac{dN(t)}{dt} \Delta t = N(t + (i - 1)\Delta t) - \lambda N(t + (i - 1)\Delta t) \Delta t$$

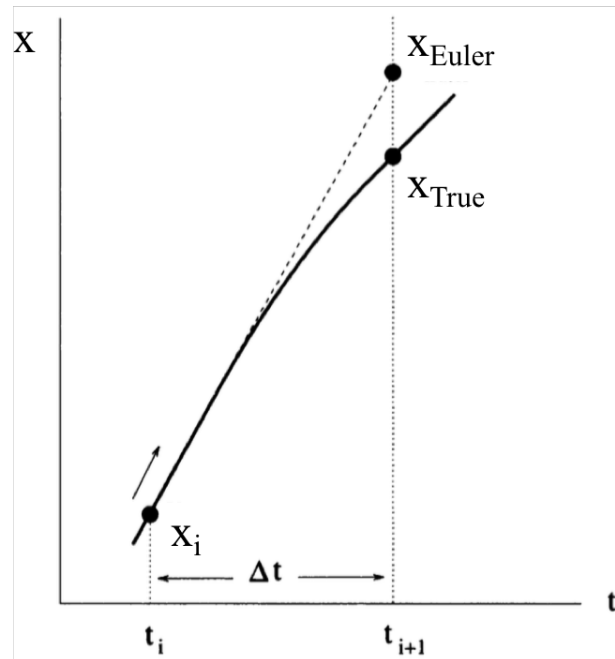
Apply Euler approximation to a specific decay example: $\frac{dN(t)}{dt} = -\lambda N(t)$
with $\lambda = 0.01$, $N_0 = 1000$ when $t_0 = 0$.

- The Euler approximation was found inaccurate compared to the analytical solution;
- The inaccuracy level is time step size affected.
- The computational time is also associated to the time step size.



- How to interpret Euler approximation geometrically?

$$x(t + \Delta t) \approx x(t) + \frac{dx}{dt} \Delta t = x(t) + f(x(t), t) \Delta t$$



Graphically, the Euler method corresponds to the linear extrapolation up to $(t + \Delta t)$ by using a line tangent to $x(t)$ at t .

A larger interval Δt , would cause a larger discrepancy between x_{true} and x_{Euler} .

How to systematically improve it?

2nd order ODE: e.g. Newton's 2nd law

$$\frac{d^2x}{dt^2} = F(x, t)/m \quad \text{with } x(0) \& v(0) \text{ given.}$$

In molecular dynamics: $m_k \ddot{\vec{x}}_k(t) = F_k(\vec{x}(t)) = -\nabla_{\vec{x}_k} V(\vec{x}(t))$

Split it into two coupled 1st
order ODE:

$$\frac{dx}{dt} = v(t)$$

$$\frac{dv}{dt} = F(x, t)$$

Assuming $m=1$ here

Euler method:

$$x(t + \Delta t) \approx x(t) + v(t)\Delta t$$

$$v(t + \Delta t) \approx v(t) + F(x(t), t)\Delta t$$

Error $O(\Delta t^2)$

Verlet method

$$x(t + \Delta t) = 2x(t) - x(t - \Delta t) + \left[\frac{F(x(t))}{m} \right] \Delta t^2$$

Error $O(\Delta t^4)$