

# A Test Particle Simulation for the Jovian Magnetospheric Electrons Precipitating into Europa's Oxygen Atmosphere

SPS 2023 • Appendix  
**Shinnosuke Satoh**

THIS STUDY IS NOW IN PREPARATION TO SUBMIT TO JGR.

## Jupiter's Dipole Field

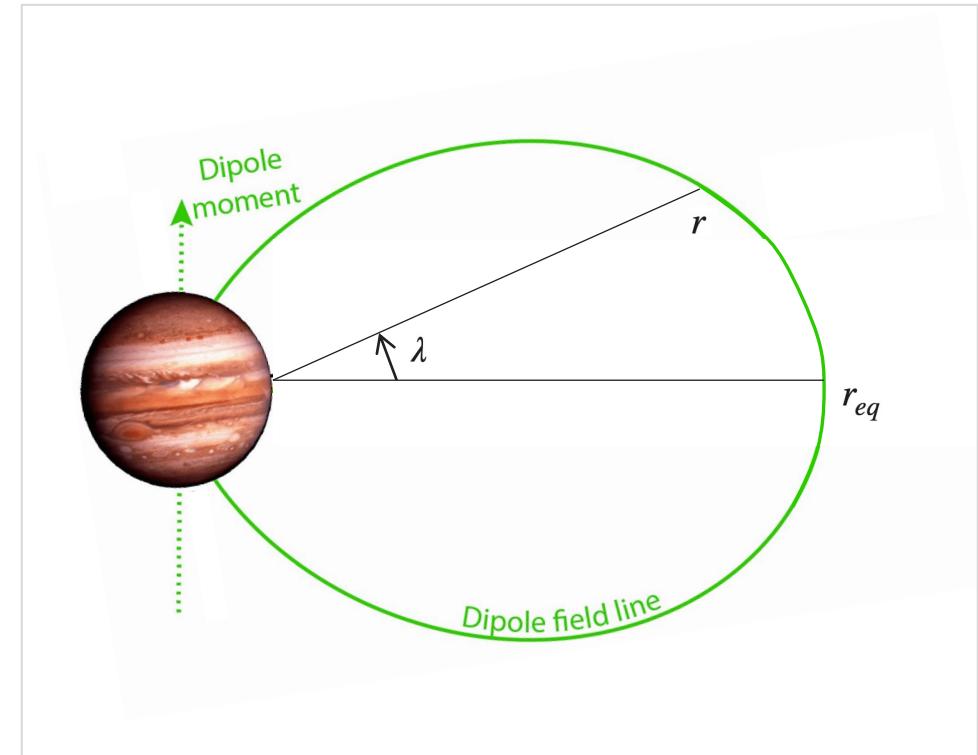
- The field

$$B = \frac{\mu_0 M}{4\pi r^3} (1 + 3 \sin^2 \lambda)^{\frac{1}{2}} \quad \lambda: \text{magnetic latitude}$$

$$\vec{B}(r, \theta) = \frac{\mu_0 M}{4\pi} \left( \frac{3Mz\vec{r}}{r^5} - \frac{M}{r^3} \vec{e}_z \right)$$

- Field line equation

$$r = r_{eq} \cos^2 \lambda$$



### Analytical Model of Static Induced Magnetic Field

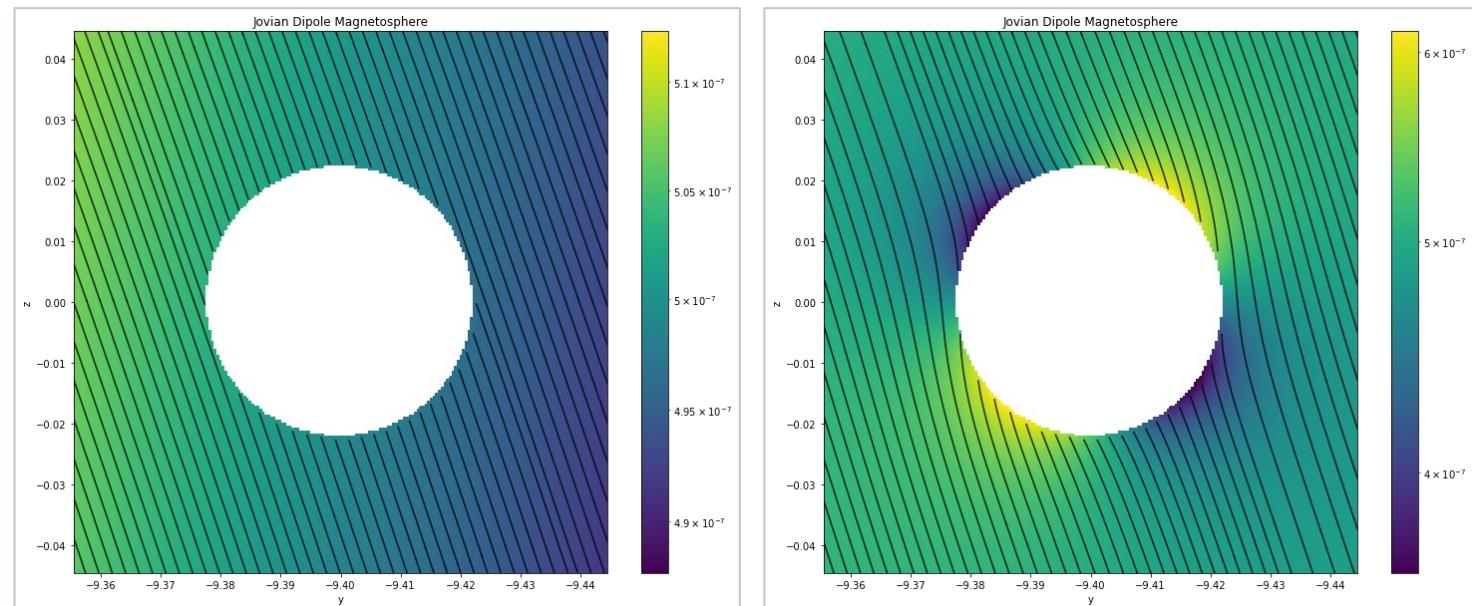
- Introduced by Neubauer+1998, Liuzzo+2016, etc.
- Treating the subsurface ocean as a highly conducting medium  
→ cancels the primary field at the "magnetic poles"

$$\mathbf{M}_{\text{ind}} = -\frac{2\pi R_C^3}{\mu_0} (B_{x,0}\hat{\mathbf{x}} + B_{y,0}\hat{\mathbf{y}})$$

(Fig)

Left: Magnetic field lines without induced magnetic field.

Right: Induced magnetic field included.



## Electrons in the Jovian Dipole Field

- Gyro-motion property

- Period

$$T_{gyro} = \frac{2\pi m}{eB} = 7.1 \times 10^{-5} \text{ sec}$$

- Radius

$$R_{gyro} = \frac{mv}{eB}$$

eV	1	10	100	1K	10K	100K	250K
ratio to light speed	1.9x10 <sup>-3</sup>	6.2x10 <sup>-3</sup>	1.9x10 <sup>-2</sup>	6.2x10 <sup>-2</sup>	1.9x10 <sup>-1</sup>	6.2x10 <sup>-1</sup>	9.8x10 <sup>-1</sup>
R_gyro [m]	6.7	21	67	213	674	2133	3373

## 4th Runge-Kutta method

- Iterative method for approximate solutions of ordinary differential equations, such as equation of motion
- Total accumulated errors on order of  $O(h^4)$  ( $h$  is the step-size.)

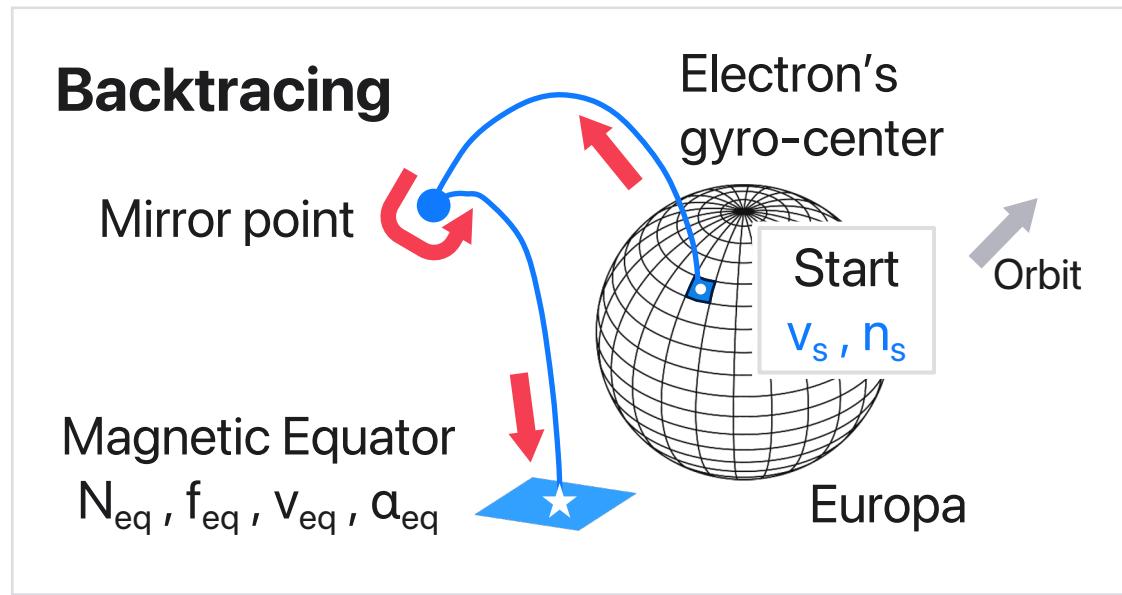
Let an initial value problem be specified as follows:

$$\frac{dy}{dt} = f(t, y), \quad y(t_0) = y_0.$$

$y$  is an unknown function of time  $t$ . Pick a step-size  $h > 0$  and for  $n = 0, 1, 2, \dots$ , define

$$\begin{aligned} k_1 &= f(t_n, y_n), \\ y_{n+1} &= y_n + \frac{1}{6}h(k_1 + 2k_2 + 2k_3 + k_4) & k_2 &= f\left(t_n + \frac{h}{2}, y_n + h\frac{k_1}{2}\right), \\ t_{n+1} &= t_n + h & k_3 &= f\left(t_n + \frac{h}{2}, y_n + h\frac{k_2}{2}\right), \\ & & k_4 &= f(t_n + h, y_n + hk_3). \end{aligned}$$

## Calculation of precipitation flux (method by Cassidy+2013)



Electron number flux into a grid labeled as "s"

$$F_s = \int_{v_{eq}} \int_{\alpha_{eq}} (\vec{v}_s \cdot \vec{n}_s) n_{eq} f_{eq}(v_{eq}, \alpha_{eq}) dv_{eq} d\alpha_{eq}$$

$v_s$ : Electron's velocity vector at the starting point on the surface

$n_s$ : Normal unit vector at the Starting point

$n_{eq}$ : Electron density at the magnetic equator

$f_{eq}$ : Electron velocity distribution at the magnetic equator

$v_{eq}$ : An electron velocity at the magnetic equator

$\alpha_{eq}$ : An electron's pitch angle at the magnetic equator

# Methods

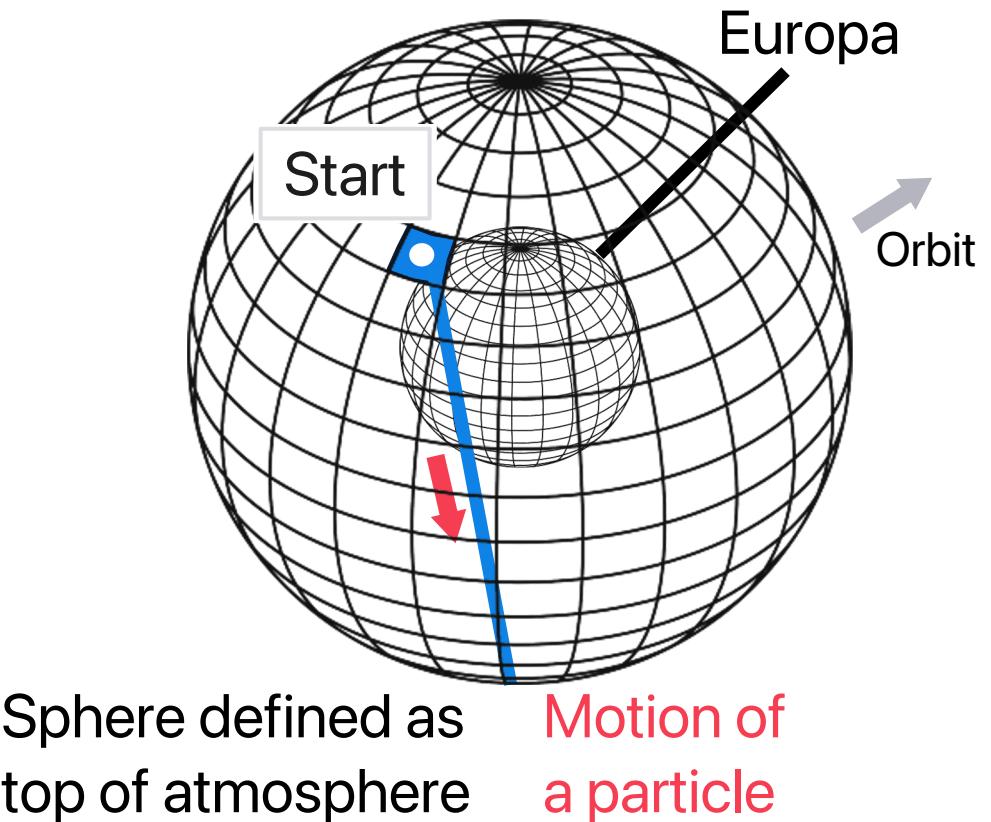
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## Calculation of Volume Emission Rate

- Same electric-magnetic field
  - Cross section for dissociative excitation of O<sub>2</sub> provided by Kanik+2003
  - No other chemical reactions included
- 
- Calculate volume emission rate each time-step
  - Integrate VER over line of sight  
→ derive brightness in Rayleigh

## Forward-tracing



# Methods

## Calculation of Volume Emission Rate

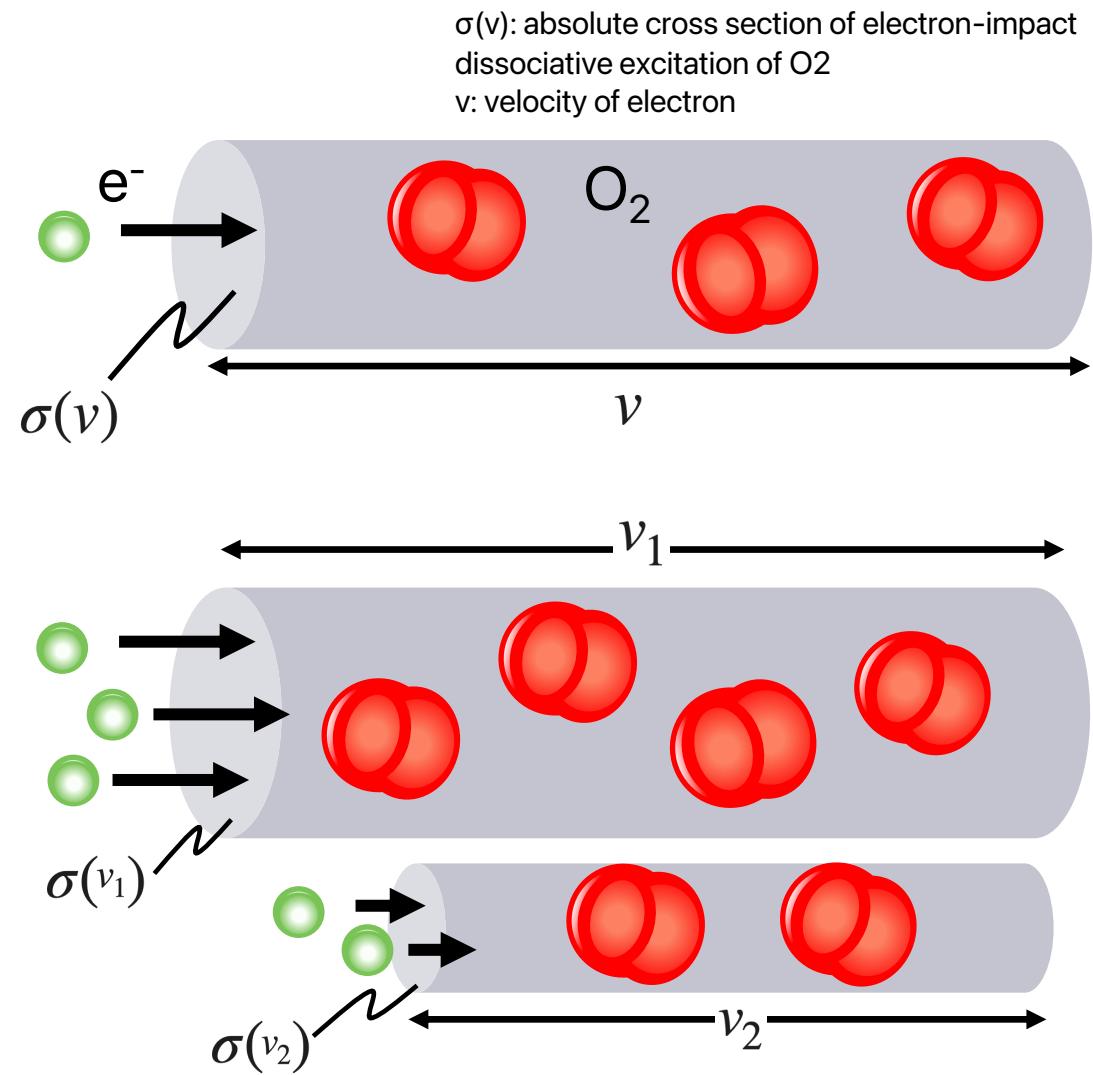
- One single electron impacting a column of  $O_2$  molecules

$$R(\vec{r}, v) = \sigma(v) v n_{O_2}(\vec{r}) \quad [\text{photons s}^{-1}]$$

- Electrons (density  $n_{e0}$ ) impacting a column of  $O_2$  molecules

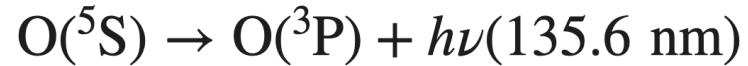
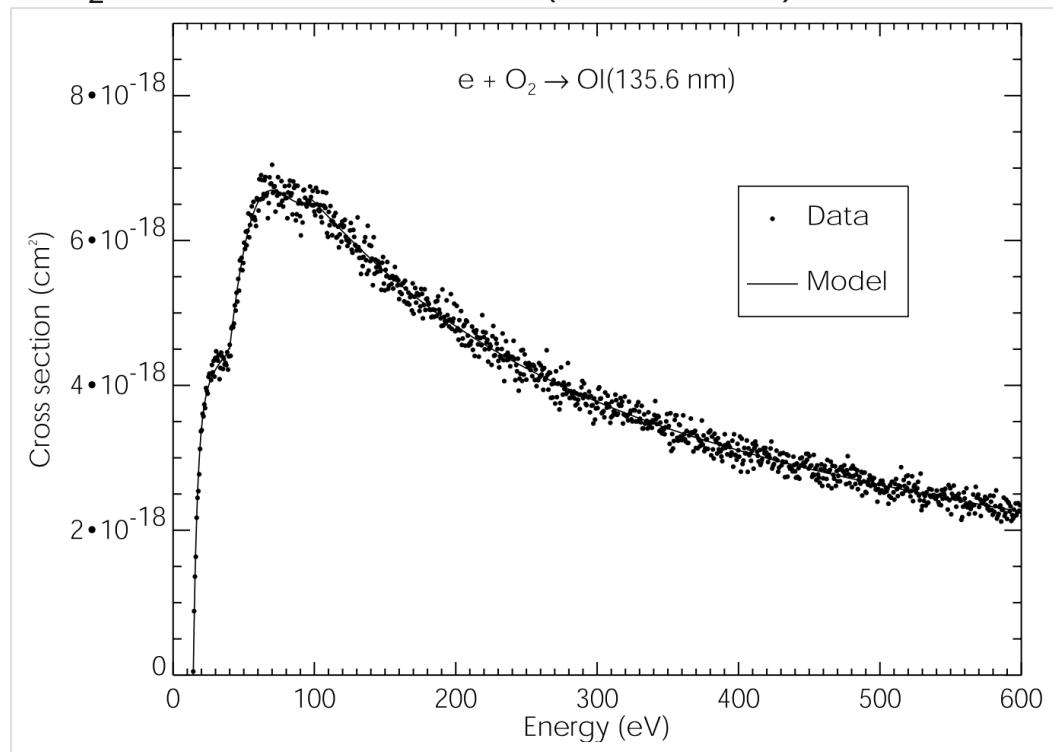
$$R'(\vec{r}) = \int_{v=\chi}^{\infty} \sigma(v) v n_{O_2}(\vec{r}) n_{e0} f(v) dv \quad [\text{photons cm}^{-3} \text{s}^{-1}]$$

Volume emission rate



## Cross Section of O<sub>2</sub> Dissociative Excitation

Laboratory data of cross section for dissociative excitation of O<sub>2</sub> with 135.6 nm emission (Kanik+2003)



- Kanik+2003 obtained the absolute value of  $\sigma(v)$  from 14.3 eV (threshold) to 600 eV.
- We adopt a linear extrapolation on the higher electron energy (> 600 eV).

## Processing of Model Spectral Images – 1 of 4

- The volume emission rate of 135.6 nm emission due to the dissociative excitation of O<sub>2</sub> is calculated by the following expression:

$$R_{135.6}(\vec{r}) = \int_{v=\chi}^{\infty} \sigma_{O_2}(v) v n_{O_2}(\vec{r}) n_e(\vec{r}) f(v) dv$$

$R_{135.6}$ : 135.6 nm volume emission rate

$\sigma_{O_2}$ : cross section for the dissociative excitation of O<sub>2</sub>

$n_e$ : local electron density

$n_{O_2}$ : local O<sub>2</sub> density

$f$ : local electron velocity distribution

$v$ : electron's velocity

- Integrating it along a line sight, we obtain the absolute brightness  $I_{135.6}$  on a pixel of the detector labeled as  $(p, q)$  in the unit of Rayleigh ( $1 R = 10^6 / 4\pi$  photons cm<sup>-2</sup> sr<sup>-1</sup> s<sup>-1</sup>).

$$I_{135.6}(p, q) = \frac{10^{-6}}{4\pi} \int_{l_{(p,q)}} R_{135.6}(\vec{r}) dl_{(p,q)}$$

$I_{135.6}(p, q)$ : 135.6 nm brightness at the  $(p, q)$  pixel

$l_{(p, q)}$ : line of sight at the  $(p, q)$  pixel

## Processing of Model Spectral Images – 2 of 4

- The obtained “raw” brightness is converted into the photon counts  $N_{\text{count}}$  by the following expression:

$$N_{\text{count}}(p, q) = \frac{10^6}{4\pi} I_{135.6}(p, q) \Omega_{\text{px}} A_{\text{eff}} \tau_e$$

$$A_{\text{eff}} = A_{\text{HST}} T_{\text{G140L}}$$

$I_{135.6}(p, q)$ : 135.6 nm brightness at the  $(p, q)$  pixel  
 $I_{(p, q)}$ : line of sight at the  $(p, q)$  pixel  
 $\Omega_{\text{px}}$ : solid angle of a single pixel of the STIS/FUV-MAMA  
calculated from the spatial plate scale of 0.0246 arcsec/pixel  
 $A_{\text{eff}}$ : the effective area of the HST  
 $A_{\text{HST}}$ : the area of the HST primary mirror,  $45239 \text{ cm}^2$   
 $\tau_e$ : exposure time, here assumed to be 45 minutes  
 $T_{\text{G140L}}$ : throughput of G140L grating, 4% for 135.6 nm

- Assuming that photon noise of the detector follows the normal distribution, the photon noise of the detector  $N_{\text{noise}}$  is calculated for the  $(p, q)$  pixel:

$$N_{\text{noise}}(p, q) = \sqrt{N_{\text{count}}(p, q)} \times (\text{random value from the normal distribution})$$

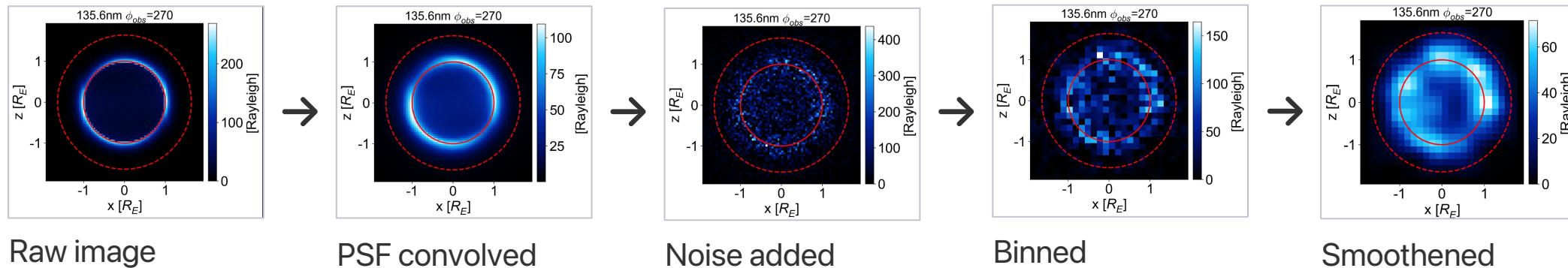
## Processing of Model Spectral Images – 3 of 4

- We use the NumPy random.randn() function to generate random values from the normal distribution.
- $N_{\text{noise}}$  is added to  $N_{\text{count}}$ , and then the summation  $N_{\text{count}} + N_{\text{noise}}$  is converted to the unit of Rayleigh → model spectral image with photon noise for STIS/FUV-MAMA
- If  $N_{\text{count}} + N_{\text{noise}}$  is negative on a pixel, we suppose that the pixel does not detect any photons at the exposure.

## Processing of Model Spectral Images – 4 of 4

- The model spectral images with photon noise will be binned in  $3 \times 3$  pixels.
- Finally, it will be smoothed for six times (Roth+2016):

$$f(i_x, i_y) = \frac{1}{8}(4f(i_x, i_y) + f(i_{x-1}, i_y) + f(i_{x+1}, i_y) + f(i_x, i_{y-1}) + f(i_x, i_{y+1})),$$

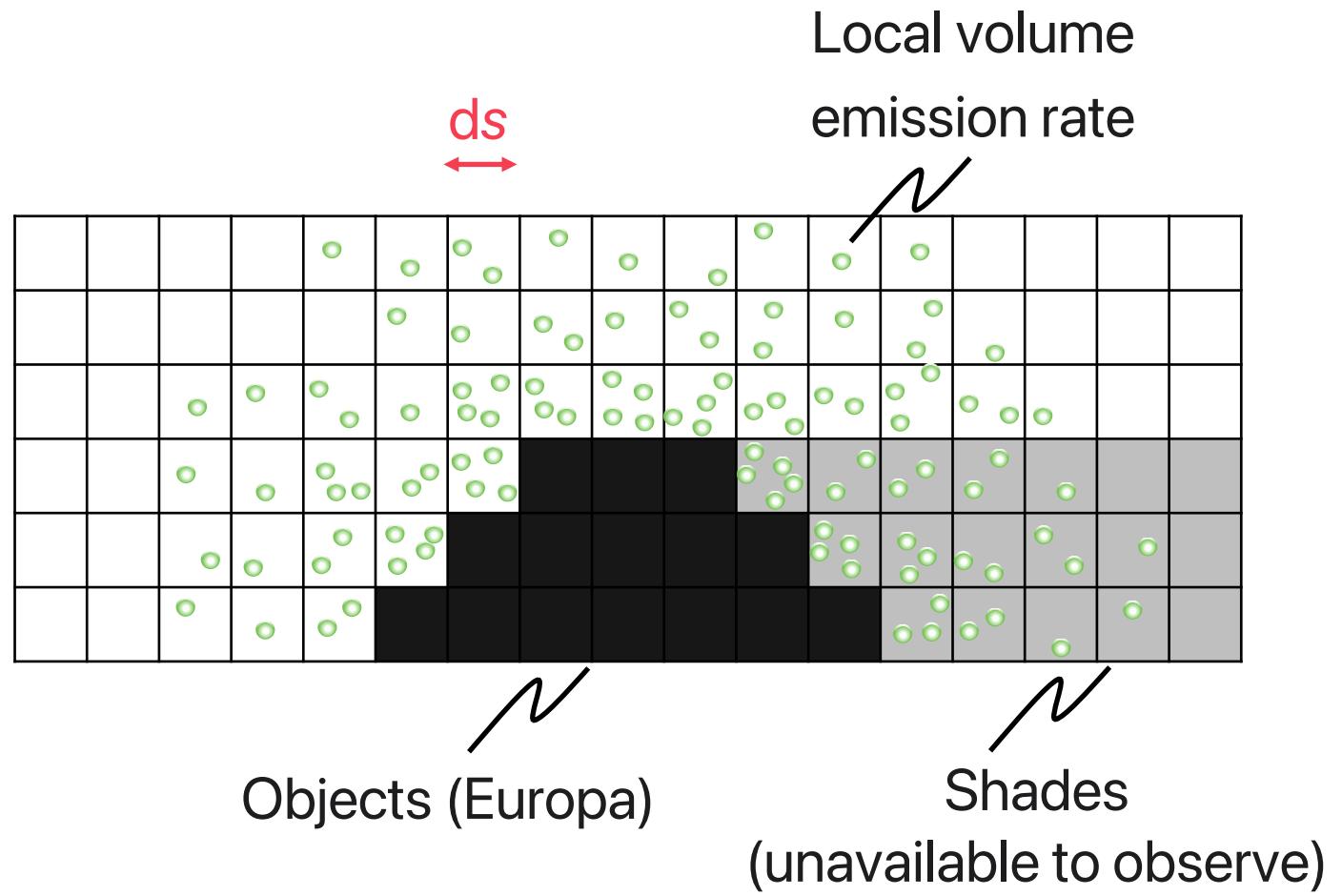


## Integration on Line of Sight



Line of sight

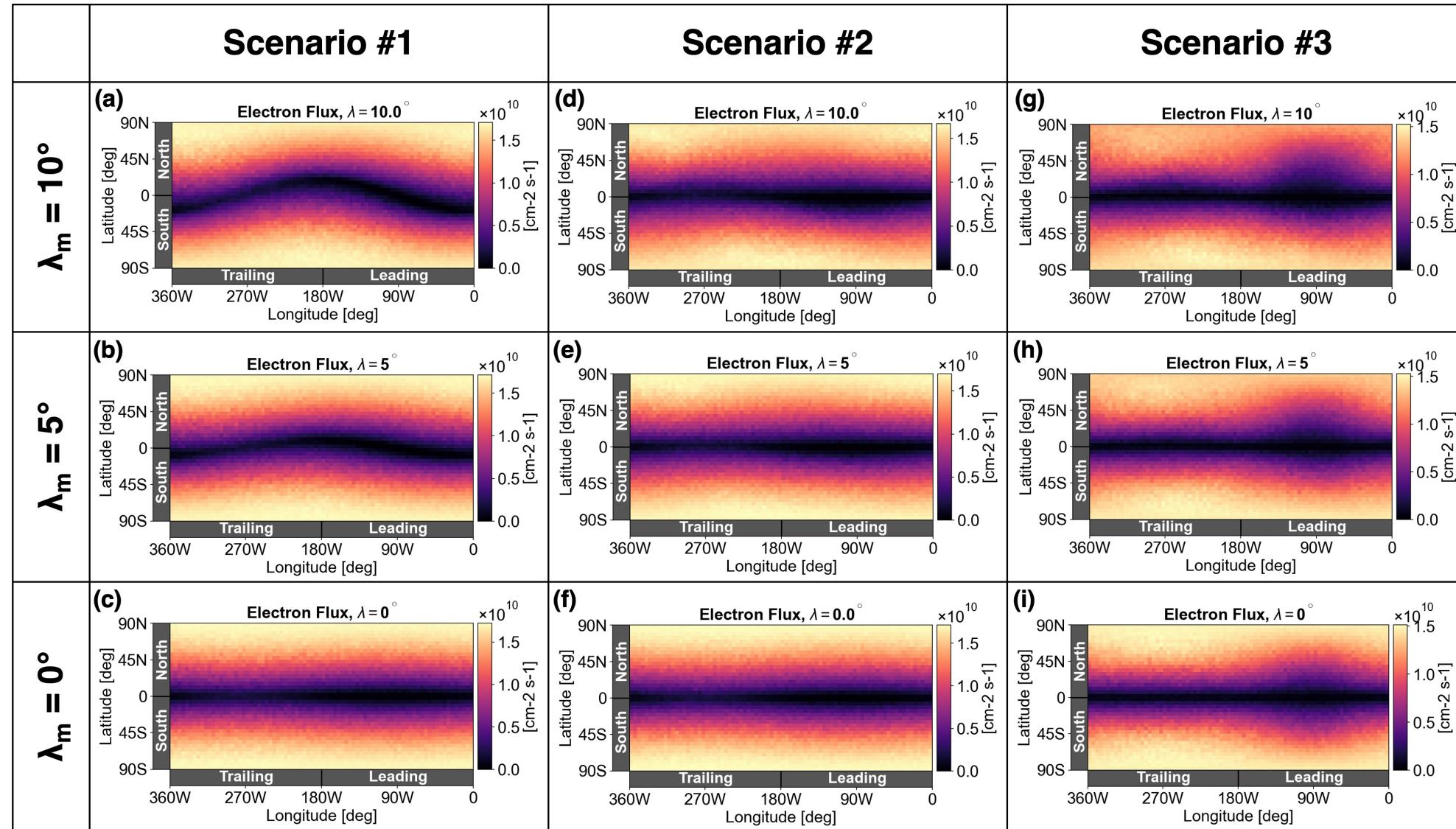
$$I = \frac{10^{-6}}{4\pi} \int_0^\infty \gamma n_e n_{O_2} ds \text{ [Rayleigh]}$$



# Results

## Appendix

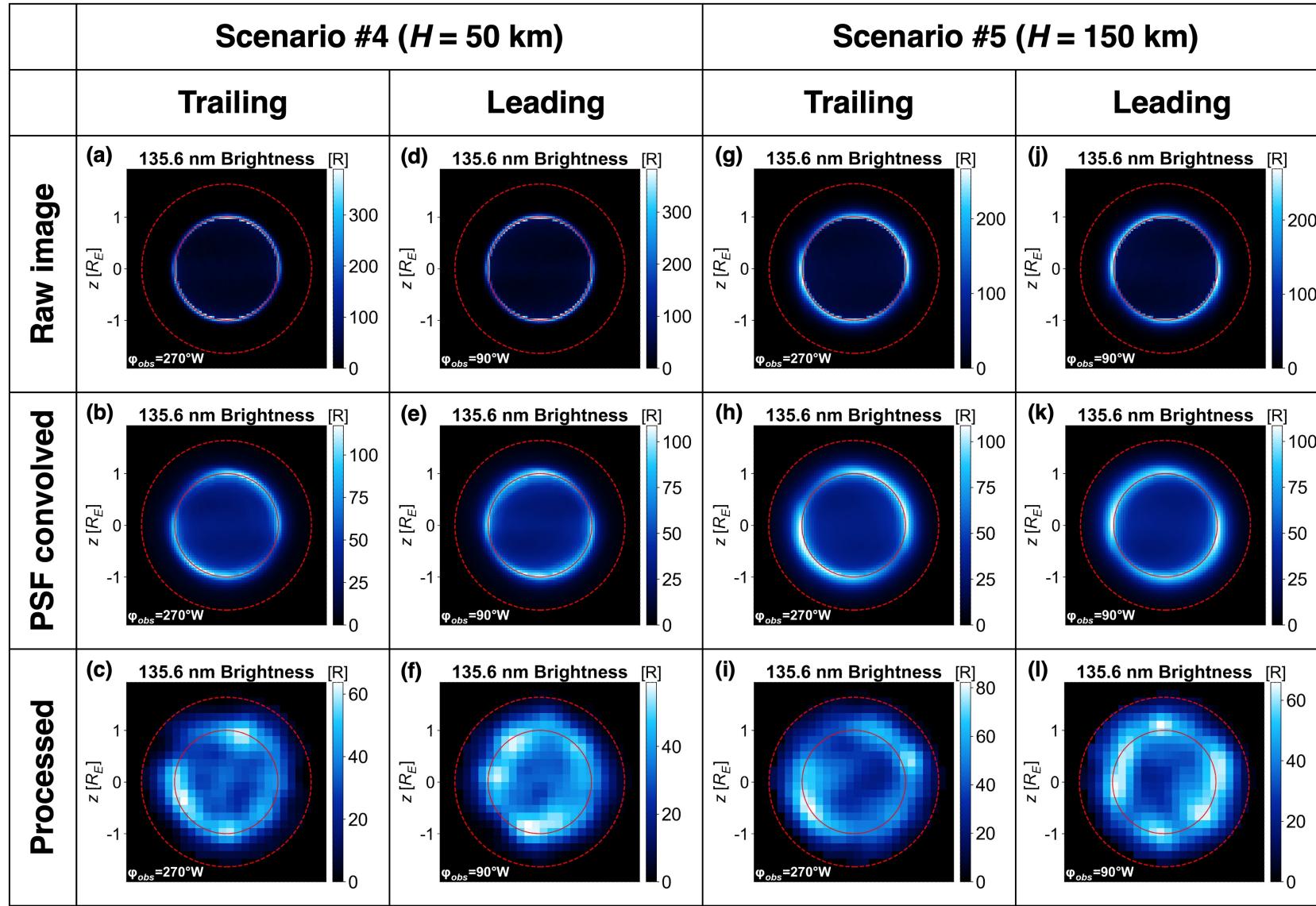
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# Results

## Appendix

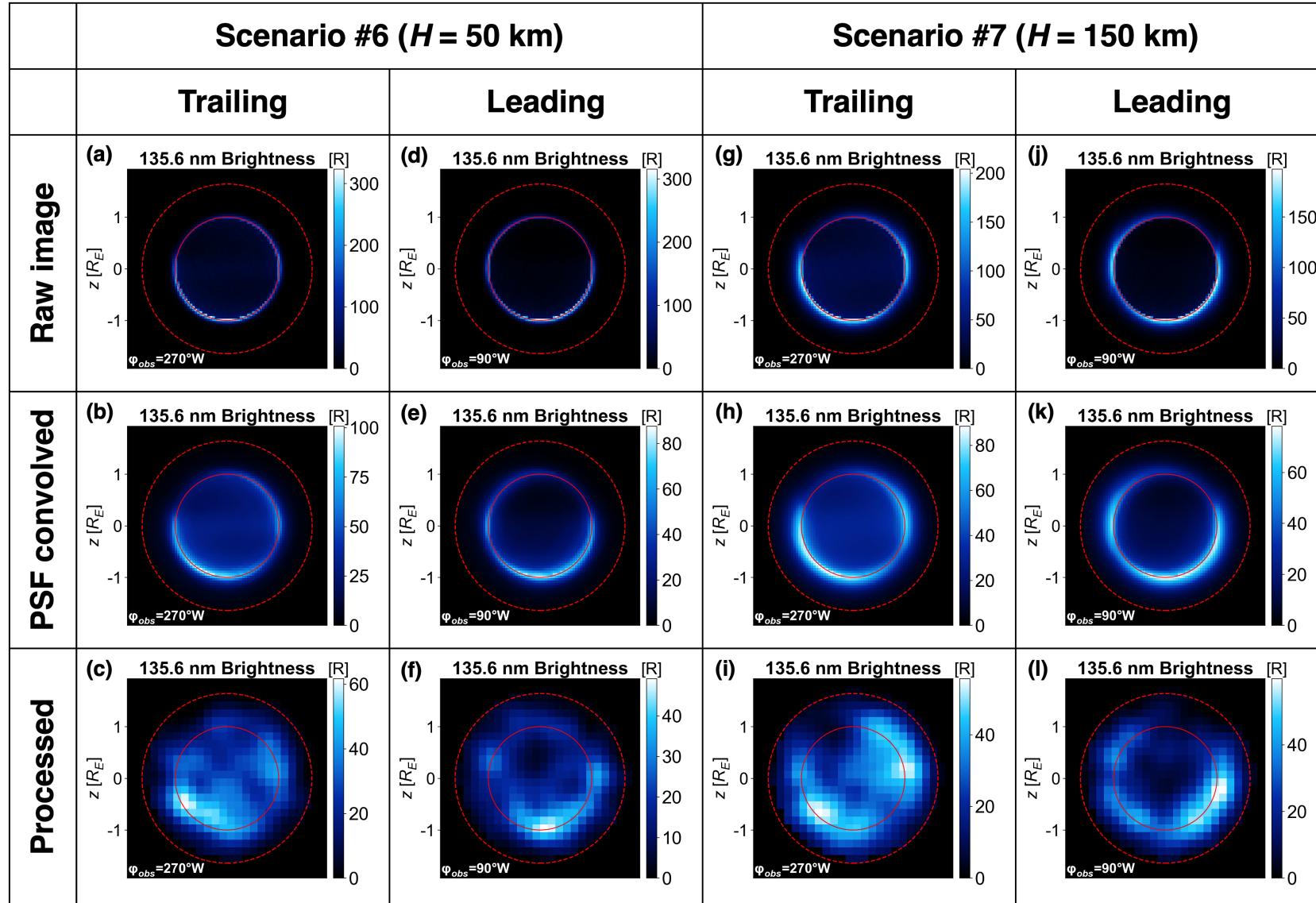
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# Results

## Appendix

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# Results

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- The limb peak brightness for  $H = 50 \text{ km}$  is smoothed more heavily and becomes lower than for the case with  $H = 150 \text{ km}$ .
- This is because of the full width at half maximum (FWHM) of the PSF, which is  $\sim 1.5$  raw pixels (Branton & Riley, 2021) and equivalent to  $\sim 113 \text{ km}$  on Europa's disk.

