

# Chapter 12 – Rotation of a Rigid Body

- Centre of mass and moment of inertia
- Torque and cross product
- Rolling motion and rotational energy
- Angular momentum

$$\theta = s/R$$

$$\omega = v/R$$

$$\alpha = a/R$$

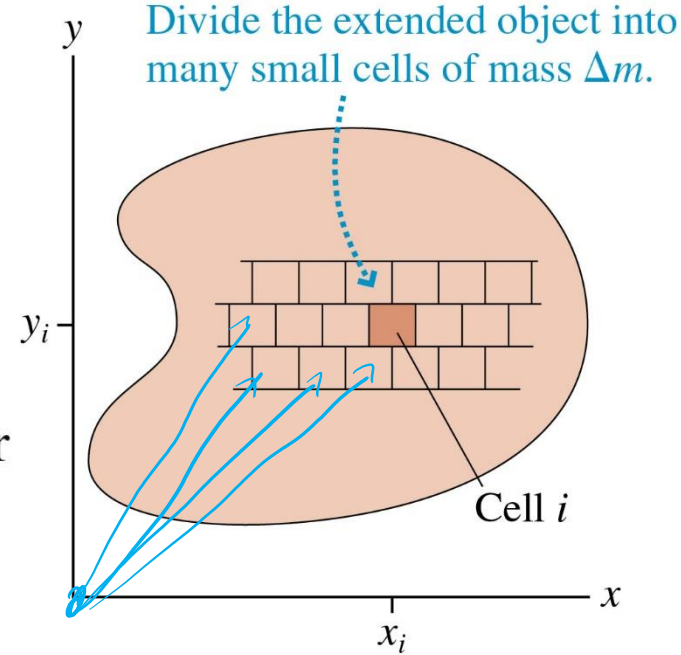


A system of particles on which there is no net force undergoes unconstrained rotation about the **center of mass**:

$$x_{\text{cm}} = \frac{1}{M} \int x \, dm \quad y_{\text{cm}} = \frac{1}{M} \int y \, dm$$

The gravitational torque on a body can be found by treating the body as a particle with all the mass  $M$  concentrated at the center of mass.

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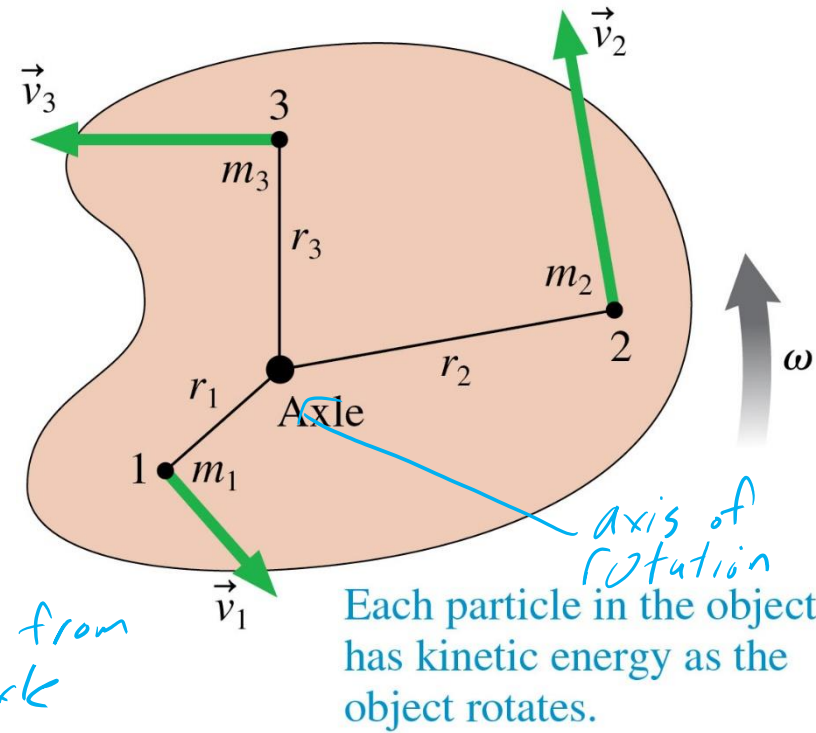
$$x_{\text{cm}} = \frac{1}{M} \sum_{i=1}^N x_i m_i$$

$$y_{\text{cm}} = \frac{1}{M} \sum_{i=1}^N y_i m_i$$

$$M = \sum_{i=1}^N m_i$$

$$K = \sum_{i=1}^N \frac{1}{2} m_i v_i^2 = \frac{1}{2} \sum m_i (r_i \omega)^2 \quad v = R\omega$$

$$= \frac{1}{2} (\sum m_i r_i^2) \omega^2 = \frac{1}{2} I \omega^2$$



**The moment of inertia**  
(rotational inertia)

$$I = \sum_i m_i r_i^2 = \int r^2 dm$$

distance from the axis

$$x_{cm} = \int x dm$$

is the rotational equivalent of mass. The moment of inertia depends on how the mass is distributed around the axis. If  $I_{cm}$  is known,  $I$  about a parallel axis distance  $d$  away is given by the **parallel-axis theorem**:

**theorem:**  $I = I_{cm} + Md^2$ .

$I_{cm}$  rotational

$Md^2 \rightarrow$  point approx far away

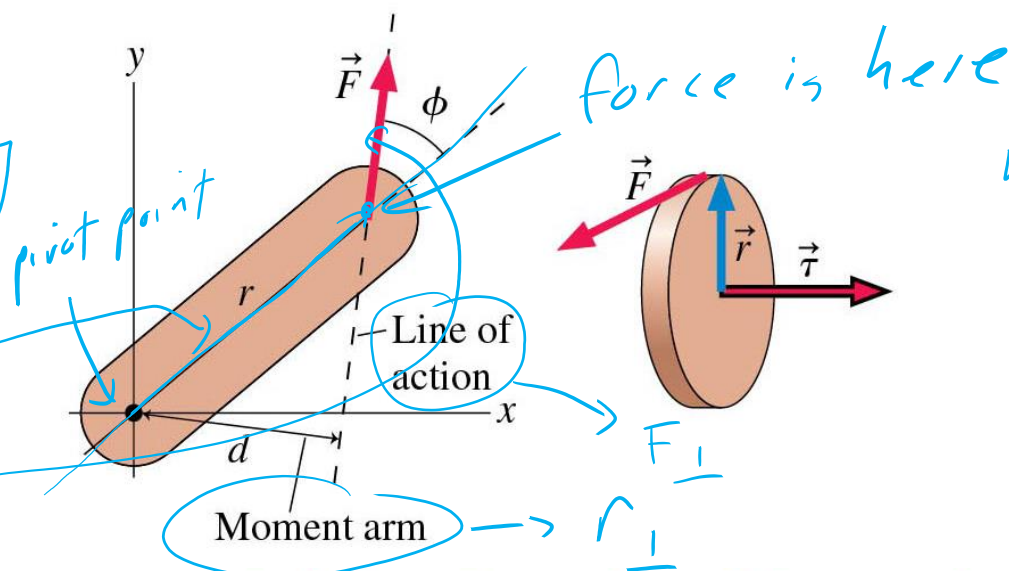
**Torque** is the rotational equivalent of force:

$$\tau = rF \sin \phi = rF_t = dF$$

The vector description of torque is

$$\vec{\tau} = \vec{r} \times \vec{F}$$

$$|\vec{r}| |\vec{F}| \sin \phi$$



## Solving Rotational Dynamics Problems

**MODEL** Model the object as a **rigid body**.

**VISUALIZE** Draw a pictorial representation.

**SOLVE** Use **Newton's second law** for rotational motion:

$$\alpha = \frac{\tau_{\text{net}}}{I} \quad \text{like} \quad a = \frac{F_{\text{net}}}{m}$$

Use rotational kinematics to find angles and angular velocities.

**REVIEW** Is the result reasonable?

$$\vec{r} \parallel \vec{F} \rightarrow \sin 0 = 0$$

$\vec{r} \cdot \vec{F}$  is how much of

$\vec{r}$  is  $\parallel \vec{F}$ ?

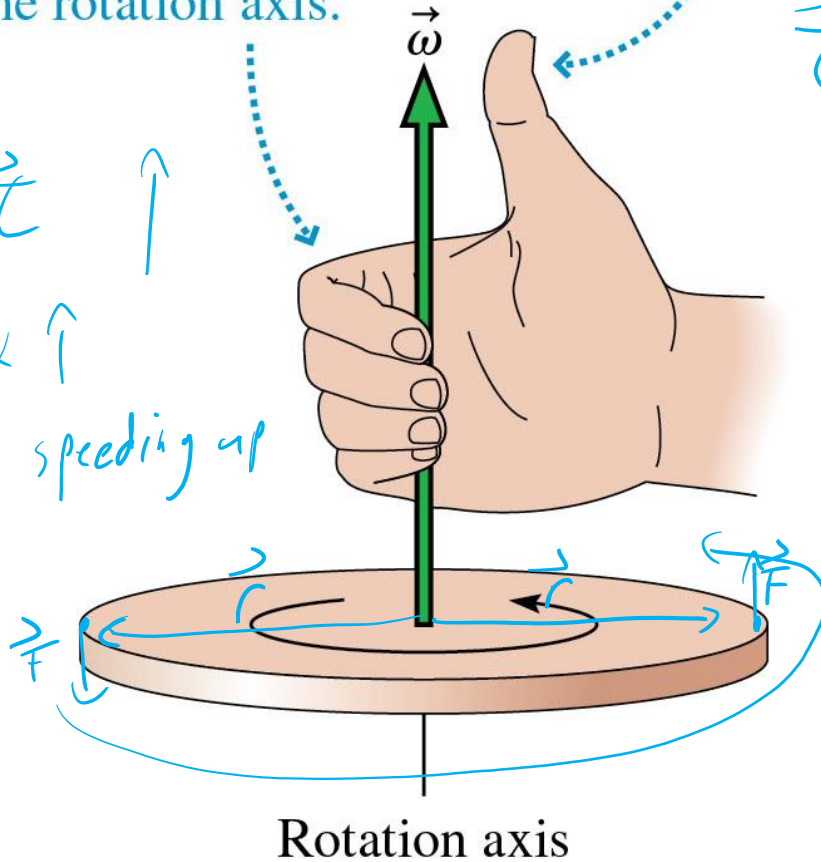
$\vec{r} \times \vec{F}$  is how much of  $\vec{r}$  is  $\perp \vec{F}$ ?

1. Using your right hand, curl your fingers in the direction of rotation with your thumb along the rotation axis.

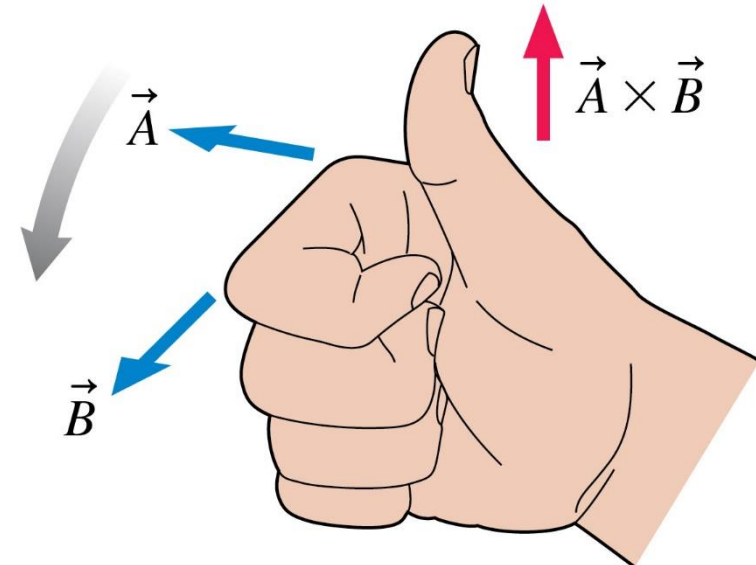
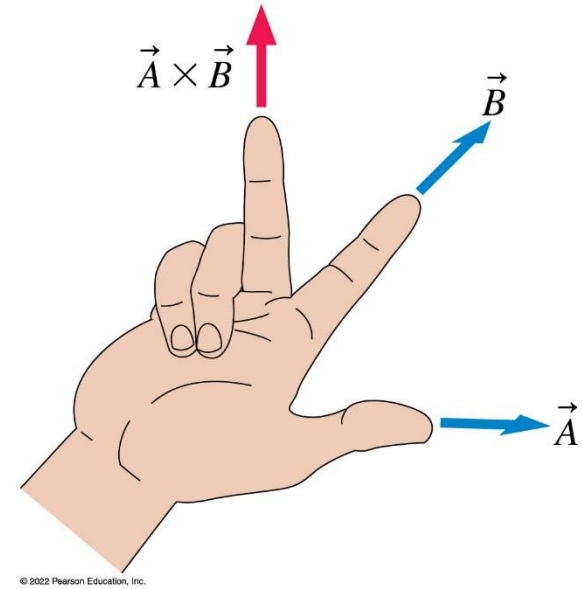
2. Your thumb is then pointing in the direction of  $\vec{\omega}$ .

$$\vec{\tau} = \vec{r} \times \vec{F}$$

$\vec{\tau} \uparrow$   
 $\vec{\omega} \uparrow$   
 $\therefore$  speeding up



Prof Wilson approved





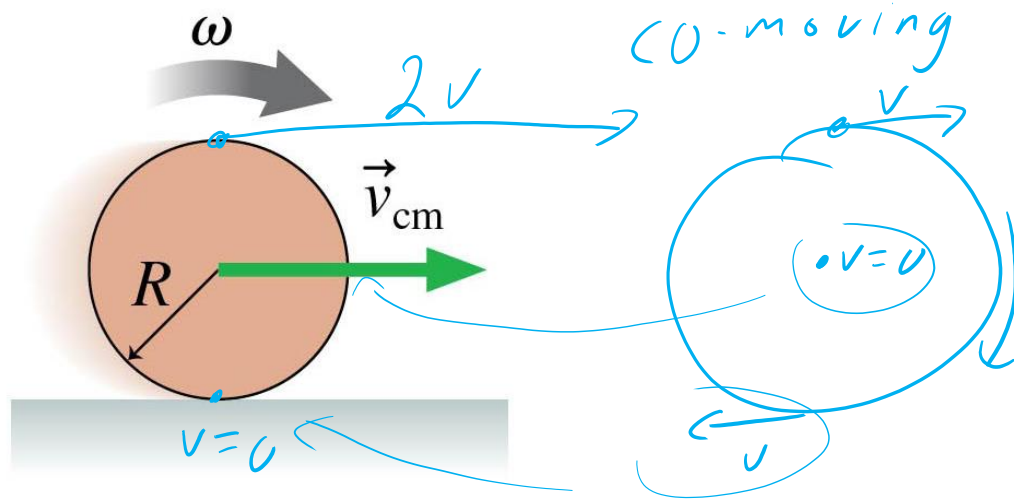
For an object that rolls without slipping

$$v_{\text{cm}} = R\omega$$

$$K = K_{\text{rot}} + K_{\text{cm}}$$

co-moving  
spin

point-like



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Energy is conserved for an isolated system.

- Pure rotation  $E_{\text{mech}} = K_{\text{rot}} + U_G = \frac{1}{2}I\omega^2 + Mgy_{\text{cm}}$  rare
- Rolling  $E_{\text{mech}} = K_{\text{rot}} + K_{\text{cm}} + U_G = \frac{1}{2}I\omega^2 + \frac{1}{2}Mv_{\text{cm}}^2 + Mgy_{\text{cm}}$

Angular velocity  $\vec{\omega}$  points along the rotation axis in the direction of the right-hand rule.

For a rigid body rotating about a fixed axle, the angular momentum is  $\vec{L} = I\vec{\omega}$ .

Newton's second law is  $\frac{d\vec{L}}{dt} = \vec{\tau}_{\text{net}}$ .

$$\frac{d\vec{p}}{dt} = \vec{F}_{\text{net}}$$

$$\begin{aligned} m &\longleftrightarrow I \\ \vec{v} &\longleftrightarrow \vec{\omega} \\ \vec{a} &\longleftrightarrow \vec{\alpha} \\ \vec{p} &\longleftrightarrow \vec{L} \\ \vec{F} &\longleftrightarrow \vec{\tau} \end{aligned}$$

