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1.) (a) To prove:

$$L^{-1}\left(\frac{1}{s} \sin \frac{1}{s}\right) = t - \frac{t^3}{(3!)^2} + \frac{t^5}{(5!)^2} - \frac{t^7}{(7!)^2} + \dots$$

By the MacLaurin series expansion of  $\sin(s)$ ,

$$\sin(s) = s - \frac{s^3}{3!} + \frac{s^5}{5!} + \frac{s^7}{7!} + \dots$$

$$\therefore \sin\left(\frac{1}{s}\right) = \frac{1}{s} - \frac{1}{3!} \cdot \frac{1}{s^3} + \frac{1}{5!} \cdot \frac{1}{s^5} - \frac{1}{7!} \cdot \frac{1}{s^7} + \dots$$

$$\Rightarrow \frac{1}{s} \cdot \sin\left(\frac{1}{s}\right) = \frac{1}{s^2} - \frac{1}{3!} \cdot \frac{1}{s^4} + \frac{1}{5!} \cdot \frac{1}{s^6} - \frac{1}{7!} \cdot \frac{1}{s^8} + \dots$$

Applying Inverse Laplace Transform on both sides

$$\begin{aligned} L^{-1}\left\{\frac{1}{s} \sin \frac{1}{s}\right\} &= L^{-1}\left\{\frac{1}{s^2} - \frac{1}{3!} \cdot \frac{1}{s^4} + \frac{1}{5!} \cdot \frac{1}{s^6} - \frac{1}{7!} \cdot \frac{1}{s^8} + \dots\right\} \\ &= L^{-1}\left\{\frac{1}{s^2}\right\} - L^{-1}\left\{\frac{1}{3!} \cdot \frac{1}{s^4}\right\} + L^{-1}\left\{\frac{1}{5!} \cdot \frac{1}{s^6}\right\} - L^{-1}\left\{\frac{1}{7!} \cdot \frac{1}{s^8}\right\} + \dots \end{aligned}$$

(By Linearity Property of Inverse Laplace transform)

$$= \frac{t}{\Gamma(2)} - \frac{t^3}{\Gamma(4) \cdot 3!} + \frac{t^5}{\Gamma(6) \cdot 5!} - \frac{t^7}{\Gamma(8) \cdot 7!} + \dots$$

$$\left[ \because L^{-1}\left\{\frac{1}{s^{n+1}}\right\} = \frac{t^n}{n(n+1)} \right]$$

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$$= t - \frac{t^3}{(3!)^2} + \frac{t^5}{(5!)^2} - \frac{t^7}{(7!)^2} + \dots$$

$$(\because \Gamma(n+1) = n!)$$

$$\therefore \mathcal{L}^{-1} \left\{ \frac{1}{s} \cdot \sin \frac{1}{s} \right\} = + - \frac{t^3}{(3!)^2} + \frac{t^5}{(5!)^2} - \frac{t^7}{(7!)^2} + \dots$$

Hence Proved

1.) (b)

$$\int_0^\infty f(x) \cos \alpha x \cdot dx = e^{-\alpha}$$

$$\sqrt{\frac{2}{\pi}} \int_0^\infty f(x) \cdot \cos \alpha x \cdot dx = \sqrt{\frac{2}{\pi}} \cdot e^{-\alpha}$$

$$F_C \{f(x)\}(\alpha) = \sqrt{\frac{2}{\pi}} \cdot e^{-\alpha}$$

$$(\therefore F_C \{f(x)\}(s) = \sqrt{\frac{2}{\pi}} \cdot \int_0^\infty f(x) \cdot \cos sx \cdot dx)$$

Taking the Inverse Fourier cosine transform on both sides

$$F_C^{-1} \{ F_C \{f(x)\} \} = F_C^{-1} \left\{ \sqrt{\frac{2}{\pi}} \cdot e^{-\alpha} \right\}$$

$$f(x) = \sqrt{\frac{2}{\pi}} \int_0^\infty \sqrt{\frac{2}{\pi}} \cdot e^{-\alpha} \cdot \cos \alpha x \cdot dx$$

$$= \frac{2}{\pi} \int_0^\infty e^{-\alpha} \cdot \cos \alpha x \cdot dx = \frac{2}{\pi} \left[ e^{-\alpha} \cdot \frac{(x \cdot \sin \alpha x - \cos \alpha x)}{x^2 + 1} \right]_0^\infty$$

$$= \frac{2}{\pi} \left[ 0 - \frac{1 \cdot (0-1)}{x^2 + 1} \right]$$

Ans  $\Rightarrow$

$$f(x) = \frac{2}{\pi} \cdot \frac{1}{x^2 + 1}$$

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2.)

Given the second order linear differential equation :

$$2x^2y'' - xy' + (x-5)y = 0 \quad (1)$$

$$\Rightarrow y'' - \frac{y'}{2x} + \frac{(x-5)}{2x^2}y = 0.$$

Comparing with the standard form,

$$y'' + P(x)y' + Q(x)y = 0$$

so,

$$P(x) = \frac{-1}{2x} \quad \text{and} \quad Q(x) = \frac{(x-5)}{2x^2}$$

As we can see that both  $P(x)$  and  $Q(x)$  are not analytic at  $x=0$ .

$\therefore x=0$  is a Singular Point.

And,

$x \cdot P(x) = -1/2$  and  $x^2 Q(x) = (x-5)/2$  are analytic at  $x=0$

$\therefore x=0$  is a Regular singular Point.

$\therefore$  By Frobenius Method, let the trial solution be

$$y = \sum_{m=0}^{\infty} C_m \cdot x^{m+\alpha} \quad \alpha \in \mathbb{R}, C_0 \neq 0$$

$\therefore$  On differentiating w.r.t  $x$

$$y' = \sum_{m=0}^{\infty} C_m (m+\alpha) \cdot x^{m+\alpha-1}$$

$$y'' = \sum_{m=0}^{\infty} C_m (m+\alpha)(m+\alpha-1) \cdot x^{m+\alpha-2}$$

Substituting values of  $y$ ,  $y''$  and  $y'$  in (1)

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$$2x^2 \sum_{m=0}^{\infty} C_m(m+\alpha)(m+\alpha-1)x^{m+\alpha-2} - x \sum_{m=0}^{\infty} C_m(m+\alpha)x^{m+\alpha-1} + (x-5) \sum_{m=0}^{\infty} C_m x^{m+\alpha} = 0$$

$$2 \sum_{m=0}^{\infty} C_m(m+\alpha)(m+\alpha-1)x^{m+\alpha} - \sum_{m=0}^{\infty} C_m(m+\alpha)x^{m+\alpha} + (x-5) \sum_{m=0}^{\infty} C_m x^{m+\alpha} = 0$$

$$\sum_{m=0}^{\infty} C_m (2(m+\alpha)(m+\alpha-1) - (m+\alpha) - 5) x^{m+\alpha} + \sum_{m=0}^{\infty} C_m x^{m+\alpha+1} = 0 \quad (2)$$

Equating to zero the coefficient of the smallest power of  $x$  in the identity above i.e  $x^0$  term.

$$C_0 (2r^2 - 3r - 5) = 0, \quad C_0 \neq 0$$

This is the individual equation which gives the values of  $r$  -

$$2r^2 - 3r - 5 = 0$$

$$(2r-5)(r+1) = 0$$

$\Rightarrow r = -1, 5/2$  (real and unequal roots, not differing by an integer.)

Now in series equation

$$\sum_{m=1}^{\infty} C_m (2(m+\alpha)^2 - 3(m+\alpha) - 5) x^{m+\alpha} + \sum_{m=1}^{\infty} C_{m-1} x^{m+\alpha} = 0$$

Equating to zero the coefficients of  $x^{m+\alpha}$ ,

$$C_m (2(m+\alpha)^2 - 3(m+\alpha) - 5) + C_{m-1} = 0.$$

$$\Rightarrow C_m = \frac{-C_{m-1}}{2(m+\alpha)^2 - 3(m+\alpha) - 5}, \quad \text{or},$$

$$C_m = \frac{C_{m-1}}{5 - (m+\alpha)(2m+2\alpha-3)} \quad (\text{Recurrence Relation}).$$

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For  $\alpha = -1$ , we get recurrence relation of first solution  $y_1(x)$ -

$$C_m = \frac{C_{m-1}}{5 - (m-1)(2m-5)}, \quad m \geq 1.$$

$$C_1 = \frac{C_0}{5}, \quad C_2 = \frac{C_1}{6} = \frac{C_0}{30}, \quad C_3 = \frac{C_2}{3} = \frac{C_0}{90},$$

$$y_1(x) = \sum_{m=0}^{\infty} C_m x^{m-1} \\ = x^{-1} \left( C_0 + \frac{C_0}{5}x + \frac{C_0}{30}x^2 + \frac{C_0}{90}x^3 + \dots \right)$$

$$y_1(x) = C_0 x^{-1} \left( 1 + \frac{x}{5} + \frac{x^2}{30} + \frac{x^3}{90} + \dots \right)$$

For  $\alpha = 5/2$ , we get recurrence relation of second solution  $y_2(x)$ .

$$C_m = \frac{C_{m-1}}{5 - (2m+5)(m+1)}, \quad m \geq 1.$$

$$C_1 = \frac{C_0}{-9}, \quad C_2 = \frac{C_1}{-22} = \frac{C_0}{198}, \quad C_3 = \frac{C_2}{-39} = \frac{C_0}{7722}, \dots$$

$$y_2(x) = \sum_{m=0}^{\infty} C_m x^{m+5/2} \\ = x^{5/2} \left( C_0 - \frac{C_0 x}{9} + \frac{C_0 x^2}{198} - \frac{C_0 x^3}{7722} + \dots \right) \\ = C_0 x^{5/2} \left( 1 - \frac{x}{9} + \frac{x^2}{198} - \frac{x^3}{7722} + \dots \right)$$

Since,  $y_1(x)$  and  $y_2(x)$  are linearly independent, the general solution of the given equation is

$$y(x) = C_1 y_1(x) + C_2 y_2(x), \quad C_1, C_2 \in \mathbb{R}$$

$$y_1(x) = C_0 x^{-1} \left( 1 + \frac{x}{5} + \frac{x^2}{30} + \frac{x^3}{90} + \dots \right)$$

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$$y_2(x) = C_0 x^{5/2} \left( 1 - \frac{x}{9} + \frac{x^2}{198} - \frac{x^3}{7722} + \dots \right)$$

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$$y_1(x) = C_0 x^{\frac{5}{2}} \left[ 1 + \frac{x}{5} + \frac{x^2}{30} + \dots \right] = x^{\frac{5}{2}} \sum_{m=0}^{\infty} C_m x^m \quad \text{& } C_m = \frac{C_{m+1}}{5(m+1)(2m+5)}$$

$$y_2(x) = C_0 x^{\frac{5}{2}} \left[ 1 - \frac{x}{9} + \frac{x^2}{128} - \dots \right] = x^{\frac{5}{2}} \sum_{m=0}^{\infty} C_m x^m \quad \text{& } C_m = \frac{C_{m+1}}{5+(m+1)(2m+5)}$$

5.)

(a) Given an orthonormal set of functions  $\{\phi_0, \phi_1, \phi_2, \dots\}$  in  $L^2[-1, 1]$  such that  $\phi_n = c_n P_n$ , where  $c_n$  is a constant and  $P_n$  is Legendre's Polynomial of degree  $n$ ,  $n \in \mathbb{N} \cup \{0\}$ . So,

the set of functions must be orthogonal in  $L^2[-1, 1]$  and each of them has norm 1.

Since Legendre's polynomials are orthogonal,

$$\int_{-1}^1 P_m(x) P_n(x) dx = 0, \quad \text{if } m \neq n.$$

$$\text{So, } \int_{-1}^1 \phi_m(x) \phi_n(x) dx = \int_{-1}^1 c_m P_m(x) c_n P_n(x) dx$$

$$= \int_{-1}^1 (c_m c_n) (P_m(x) P_n(x)) dx$$

$$\Rightarrow \int_{-1}^1 \phi_m(x) \phi_n(x) dx = 0, \quad m \neq n.$$

The norm of  $f(x)$  is  $\|f(x)\|$  defined by  $\|f(x)\| = \left( \int_a^b f^2(x) dx \right)^{1/2}$  on

$a \leq x \leq b$ .

$$\text{So, } \left( \int_{-1}^1 \phi_n^2(x) dx \right)^{1/2} = \|\phi_n(x)\|$$

$$\text{Since } \|\phi_n(x)\| = 1 \text{ and } \int_{-1}^1 P_m(x) P_n(x) dx = \frac{2}{2n+1}, \text{ for } m=n,$$

$$\left( \int_{-1}^1 c_n^2 P_n^2(x) dx \right)^{1/2} = 1$$

$$\Rightarrow c_n \sqrt{\frac{2}{2n+1}} = 1 \Rightarrow c_n = \sqrt{\frac{2n+1}{2}}$$

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$$\text{So, } \int \phi_m(x) \phi_n(x) dx = \begin{cases} 0 & , \text{ if } m \neq n \\ 1 & , \text{ if } m = n \end{cases}$$

Now, by generating function of legendre's polynomials,

$$(1 - 2xz + z^2)^{-1/2} = \sum_{n=0}^{\infty} z^n P_n(x), |x| \leq 1, |z| \leq 1.$$

Putting  $x=1$

$$(1 - 2z + z^2)^{-1/2} = \sum_{n=0}^{\infty} z^n P_n(1)$$

$$\Rightarrow (1-z)^{-1/2} = \sum_{n=0}^{\infty} z^n P_n(1)$$

Since  $|z| < 1$ , binomial theorem can be used for expansion of  $(1-z)^{-1/2}$ ,

$$1 + z + z^2 + \dots + z^n = \sum_{n=0}^{\infty} P_n(1) z^n.$$

Equating the coefficient of  $z^n$  from both sides, we get

$$P_n(1) = 1.$$

Similarly for  $x=-1$ , we get  $P_n(-1) = (-1)^n$

So,

$$\begin{aligned} \text{(i)} \quad \phi_6(1) &= C_6 P_6(1) \\ &= \sqrt{\frac{2 \times 6 + 1}{2}} \cdot 1 = \pm \sqrt{\frac{13}{2}} \end{aligned}$$

$$\begin{aligned} \text{(ii)} \quad \phi_7(-1) &= C_7 P_7(-1) \\ &= \sqrt{\frac{2 \times 7 + 1}{2}} (-\frac{1}{2})^7 = \pm \sqrt{\frac{15}{2}} \end{aligned}$$

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3)(b)

To check whether  $\frac{d}{dx} (J_5^2(x)) = \frac{x}{10} [J_6^2 - J_4^2]$  is True or False.

$J_n(x)$  is the Bessel function of the first kind of order  $n$ .

We know the recurrence relations of Bessel's function:

$$\frac{d}{dx} (x^n J_n(x)) = x^n J_{n-1}(x) \quad (1) \quad \frac{d}{dx} (x^n J_n(x)) = -x^n J_{n+1}(x) \quad (2)$$

Using the recurrence relations  $2J_n'(x) = J_{n-1}(x) - J_{n+1}(x)$

$$J_n'(x) = \frac{x}{2n} [J_{n-1}(x) + J_{n+1}(x)]$$

$$\frac{d}{dx} (J_5^2(x)) = 2 \cdot J_5(x) \cdot \frac{d}{dx} (J_5(x))$$

$$= J_5(x) \cdot 2 \cdot J_5'(x)$$

$$= \frac{x}{2.5} [J_4 + J_6] \cdot [J_4 - J_6]$$

$$= \frac{x}{10} [J_4^2 - J_6^2]$$

$$\therefore \frac{d}{dx} (J_5^2(x)) = \frac{x}{10} [J_4^2 - J_6^2]$$

$\Rightarrow$  The given statement is false.

Shmili

4.) (a.)

$$H_n \left\{ \frac{f(x)}{x} \right\} = \int_0^\infty \frac{f(x)}{x} \cdot x \cdot J_n(sx) \cdot dx$$

$$= \int_0^\infty f(x) \cdot J_n(sx) \cdot dx$$

{ Using Recurrence Relation of Bessel function }

$$J_n(x) = \frac{x}{2^n} [J_{n-1}(x) + J_{n+1}(x)]$$

$$= \int_0^\infty f(x) \cdot \frac{sx}{2^n} [J_{n-1}(x) + J_{n+1}(x)] dx$$

$$= \frac{s}{2^n} \left[ \int_0^\infty f(x) \cdot x \cdot J_{n-1}(x) dx + \int_0^\infty f(x) \cdot x \cdot J_{n+1}(x) dx \right]$$

$$= \frac{s}{2^n} [H_{n-1} \{ f(x) \} + H_{n+1} \{ f(x) \}]$$

$$\therefore H_n \left\{ \frac{f(x)}{x} \right\} = \frac{s}{2^n} [H_{n-1} \{ f(x) \} + H_{n+1} \{ f(x) \}]$$

Hence, Proved.

$$H_n \left\{ \frac{\delta(x-a)}{x^{k+1}} \right\} = \int_0^\infty \frac{\delta(x-a)}{x^{k+1}} \cdot J_n(sx) \cdot dx$$

$$= \frac{s}{2^n} \left[ \int_0^\infty \frac{\delta(x-a)}{x^{k+2}} \cdot J_{n-1}(sx) dx + \int_0^\infty \frac{\delta(x-a)}{x^{k+2}} \cdot J_{n+1}(sx) dx \right]$$

$$\text{As, } \int_{-\infty}^{\infty} \delta(x-a) \cdot f(x) = f(a)$$

$$\int_{-\infty}^{a-\epsilon} \delta(x-a) \cdot f(x) + \int_{a-\epsilon}^{a+\epsilon} \delta(x-a) \cdot f(x) + \int_{a+\epsilon}^{\infty} \delta(x-a) \cdot f(x) = f(a) \quad \epsilon > 0.$$

$$\int_{-\infty}^{a-\epsilon} 0 \cdot f(x) + \int_{a-\epsilon}^{a+\epsilon} \delta(x-a) \cdot f(x) + \int_{a+\epsilon}^{\infty} 0 \cdot f(x) = f(a).$$

$$\int_{a-\epsilon}^{a+\epsilon} \delta(x-a) \cdot f(x) = f(a).$$

$$\therefore \int_{-\infty}^{\infty} \delta(x-a) \cdot f(x) = \int_{a-\epsilon}^{a+\epsilon} \delta(x-a) \cdot f(x) = f(a) \quad \epsilon > 0$$

since  $a > 0$  (given in question)

$$\therefore \int_{-\infty}^{\infty} \delta(x-a) f(x) = \int_0^{\infty} \delta(x-a) \cdot f(x) = \int_{a-\epsilon}^{a+\epsilon} \delta(x-a) \cdot f(x) = f(a) \quad \epsilon > 0$$

$$\begin{aligned} H_n \left\{ \frac{\delta(x-a)}{x^k} \right\} &= \frac{1}{2n} \left[ \int_0^{\infty} \delta(x-a) \cdot \frac{J_{n+1}(sx)}{x^{k-2}} dx + \int_0^{\infty} \delta(x-a) \cdot \frac{J_{n+1}(sx)}{x^{k-2}} dx \right] \\ &= \frac{1}{2n} \left[ \frac{J_{n+1}(sa)}{a^{k-2}} + \frac{J_{n+1}(sa)}{a^{k-2}} \right] \\ &= \frac{1}{2n \cdot a^{k-2}} \left[ J_{n+1}(sa) + J_{n+1}(sa) \right] \\ &= \frac{1}{2na^{k-2}} \cdot \frac{2n}{sa} \cdot J_n(sa) \\ &= \frac{J_n(sa)}{a^{k-1}} \end{aligned}$$

$$\therefore \boxed{H_n \left\{ \frac{\delta(x-a)}{x^k} \right\} = \frac{J_n(sa)}{a^{k-1}}}$$

4.7(b) Given differential equation

$$\frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} + \frac{\partial^2 U}{\partial z^2} = 0, \quad 0 \leq r \leq \infty, \quad z \geq 0$$

Subjected to conditions -

$$U \rightarrow 0 \text{ as } z \rightarrow \infty \text{ and } r \rightarrow \infty$$

$$U(r, 0) = f(r)$$

$U(r, z)$  is bounded.

Taking Hankel Transform of zero order on both sides of the given equation,

$$\int_0^\infty \left( \frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} + \frac{\partial^2 U}{\partial z^2} \right) r J_0(sr) dr = 0$$

where  $J_0$  is the Bessel's function of first kind of order 0.

$$\Rightarrow \int_0^\infty \left( \frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} \right) r J_0(sr) dr + \int_0^\infty \frac{\partial^2 U}{\partial z^2} r J_0(sr) dr = 0$$

We first calculate the first integral.

$$\therefore \frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial U}{\partial r} \right)$$

$$\therefore \int_0^\infty \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial U}{\partial r} \right) r J_0(sr) dr = \left( r \frac{\partial U}{\partial r} \right) J_0(sr) \Big|_0^\infty - \int_0^\infty s J_0'(sr) r \frac{\partial U}{\partial r} dr$$

Assuming  $r \frac{\partial U}{\partial r} \rightarrow 0$  as  $r \rightarrow \infty$ ,

$$0 - s \int_0^\infty J_0'(sr) r \frac{\partial U}{\partial r} dr = -s \left[ r J_0'(sr) U \right]_0^\infty - \int_0^\infty \frac{d}{dr} \left( r J_0'(sr) \right) U dr$$

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$$= -s \left[ 0 - \int_0^\infty \frac{d}{ds} (\gamma J_0'(s\gamma)) U d\gamma \right]$$

since  $U \rightarrow 0$  as  $\gamma \rightarrow \infty$  and  $U(\gamma, z)$  is bounded.

$$\therefore \int_0^\infty \frac{1}{\gamma} \frac{\partial}{\partial s} \left( \gamma \frac{\partial U}{\partial \gamma} \right) \gamma J_0(s\gamma) d\gamma = s \int_0^\infty \frac{d}{ds} (\gamma J_0'(s\gamma)) U d\gamma$$

Now from Bessel's equation,

$$x^2 J_0''(x) + x J_0'(x) + x^2 J_0(x) = 0$$

$$\Rightarrow x J_0''(x) + J_0'(x) + x J_0(x) = 0$$

$$\Rightarrow \frac{d}{dx} (x J_0'(x)) + x J_0(x) = 0$$

$$\Rightarrow \frac{1}{s} \frac{d}{ds} (s x J_0'(sx)) + s x J_0(sx) = 0$$

(Putting  $x = sx$ )

$$\Rightarrow \frac{d}{ds} (s x J_0'(sx)) = -s x J_0(sx)$$

$$\therefore \int_0^\infty s \frac{d}{ds} (\gamma J_0'(s\gamma)) U d\gamma = -s^2 \int_0^\infty \gamma J_0(s\gamma) U d\gamma$$

$$\therefore \int_0^\infty \left( \frac{\partial^2 U}{\partial \gamma^2} + \frac{1}{\gamma} \frac{\partial U}{\partial \gamma} \right) \gamma J_0(s\gamma) d\gamma = -s^2 \bar{U},$$

where  $\bar{U}$  is the Hankel transform of zero order of  $U(\gamma, z)$ .

Now,

$$-s^2 \bar{U}(s, z) + \frac{\partial^2 \bar{U}}{\partial z^2}(s, z) = 0$$

$$\text{or } \frac{\partial^2 \bar{U}(s, z)}{\partial z^2} - s^2 \bar{U}(s, z) = 0$$

which is a second order homogeneous linear differential equation whose general solution is given by,

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$$\bar{U}(s, z) = C_1 e^{sz} + C_2 e^{-sz} \quad \text{where } C_1 \text{ and } C_2 \text{ are constants.}$$

From the conditions given,

since  $U \rightarrow 0$  as  $z \rightarrow \infty$  and  $z \rightarrow -\infty$ ,  $\bar{U}(s, z)$  should also be finite as  $z \rightarrow \infty$ , which is possible only if  $C_1 = 0$ .

$$\therefore \bar{U}(s, z) = C_2 e^{-sz}$$

$$\text{Also, } U(s, 0) = f(s)$$

Taking Hankel transform of zero order on both sides,

$$\bar{U}(s, 0) = \int_0^\infty f(r) r J_0(sr) dr$$

$$\bar{U}(s, 0) = \bar{f}(s) \quad (H_0\{f(r)\} = \bar{f}(s))$$

$$\Rightarrow C_2 = \bar{f}(s)$$

$$\Rightarrow \bar{U}(s, z) = \bar{f}(s) e^{-sz}$$

Taking inverse Hankel transform of zero order on both sides,

$$U(s, z) = \int_0^\infty \bar{f}(s) e^{-sz} s J_0(sr) ds$$

is the required solution.

Shruti

Ans 5.)

$$xy'' + y' = 0 \quad ; \quad \begin{aligned} \text{i)} \quad & y(x) \text{ is bounded as } x \rightarrow 0 \\ \text{ii)} \quad & y(1) = \alpha y'(1), \alpha \neq 0 \end{aligned}$$

The above equation can be written as -

$$Ly(x) = f(x)$$

$$\text{where } L = \left[ \frac{d}{dx} \left( p(x) \cdot \frac{dy}{dx} \right) \right] \text{ and } f(x) = 0.$$

$$p(x) = x.$$

$$xy'' + y' = 0$$

$$\frac{d}{dx}(xy') = 0, \quad xy' = c \quad ; \quad c - \text{constant} = (c \in \mathbb{R})$$

$$y' = c/x.$$

$$y(x) = c \ln x + k; \quad k \in \mathbb{R}$$

For  $y(x)$  to be bounded as  $x \rightarrow 0$ ,  $c$  has to be zero.

$$\therefore c=0$$

$$\text{also, } y(1) = \alpha \cdot y'(1)$$

$$k = \alpha \cdot c/1$$

$$k=0$$

$\therefore y(x)=0$  Trivial solution is the only solution of  $Ly(x)=0$   
Let the Green's function be

$$G(x,t) = \begin{cases} G_1(x,t) = A \ln x + B, & 0 \leq x < t \\ G_2(x,t) = C \ln x + D, & t < x \leq 1 \end{cases}$$

where  $A, B, C, D$  are real constants.

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By properties of Green's function,

1)  $G_1(x, t)$  is bounded at  $x \rightarrow 0$

$G_1(x, t) = A \ln x + B$ , for  $G_1(x, t)$  to be bounded as  $x \rightarrow 0$ ,

A should be 0

$$\therefore \underline{A=0} \quad - (1)$$

$$G_2(1, t) = \alpha \cdot G_2'(1, t)$$

$$0 + D = \alpha \left[ \frac{C}{x} \right]_{x=1} \Rightarrow \underline{D=\alpha C} \quad - (2)$$

$$2) \lim_{x \rightarrow t} G_1(x, t) = \lim_{x \rightarrow t} G_2(x, t)$$

$$\lim_{x \rightarrow t} [0 + B] = \lim_{x \rightarrow t} [C \ln x + D] \quad (G_1(x, t) = A \ln x + B = 0 + B = B) \\ \text{from (1)}$$

$$B = C \ln t + D$$

$$B = C \ln t + \alpha C \quad \text{from (2)}$$

$$\underline{B = C(\ln t + \alpha)}.$$

$$3) \left| \frac{\partial G_2(x, t)}{\partial x} \right|_{x=t} - \left| \frac{\partial G_1(x, t)}{\partial x} \right|_{x=t} = \frac{-1}{P(t)}$$

$$\left| \frac{\partial}{\partial x} [C \ln x + D] \right|_{x=t} - \left| \frac{\partial}{\partial x} [B] \right|_{x=t} = \frac{-1}{P(t)}$$

$$\left| \frac{C}{x} \right|_{x=t} - D = \frac{-1}{P(t)}$$

$$\frac{C}{t} = -\frac{1}{t} \quad (\because P(t)=x)$$

$$(C = -1)$$

Shmushi

Now from (1), (2), (3), (4)

$$A = 0$$

$$c = -1$$

$$B = -(\alpha + \ln t)$$

$$D = -\alpha$$

$$\therefore g_1(x, t) = A \ln x + B = -(\alpha + \ln t)$$

$$g_2(x, t) = C \ln x + D = -\ln x - \alpha$$

$$g(x, t) = \begin{cases} -t + \ln t, & 0 \leq x < t \\ -(\alpha + \ln x), & t < x \leq 1 \end{cases}$$