

16/9/19

→ Choose a measurement gate  $\Psi \rightarrow [A] \rightarrow$   
which has its own base states.

$$\Psi = a|0\rangle + b|i\rangle \quad a, b \in \mathbb{C} \quad \text{Born Rule}$$

$$|a|^2 + |b|^2 = 1 \quad |\Psi\rangle \rightarrow [A] \rightarrow P(|0\rangle) = |a|^2$$

$$P(|0\rangle) = |a|^2 \quad P(|i\rangle) = |b|^2$$

Once a qubit measured in a definite state → its state is no longer probab.  
its state becomes what is obtained in measurement

Projective Measurement or Von Neumann measurement.

Global Phase & Relative Phase:

$$|\Psi\rangle = a|0\rangle + b|i\rangle$$

$$|\Psi'\rangle = e^{i\theta} (a|0\rangle + b|i\rangle)$$

$$= e^{i\theta} |\Psi\rangle$$

all the rules studied for  $|\Psi\rangle$  also hold for  $|\Psi'\rangle$ .

- $a e^{i\theta} \xrightarrow{\text{met}} \sqrt{|a e^{i\theta}|^2} = |a e^{i\theta}| (a^* e^{-i\theta})$   
 $= |a|^2$

also  $b e^{i\theta} \rightarrow |b|^2$

$$\Rightarrow |a|^2 + |b|^2 = 1$$

- $|\Psi'\rangle \xrightarrow{\text{met}} P(|0\rangle) = |a|^2$

if  $|\Psi\rangle$  normalized

$|\Psi'\rangle$  also normalized

if  $P(|0\rangle) = |a|^2$  for  $|\Psi\rangle$

then  $P(|0\rangle) = |a|^2$  for  $|\Psi'\rangle$

for all practical purposes  $|\Psi\rangle$  &  $|\Psi'\rangle$  indistinguishable

But  $\langle \Psi | \Psi' \rangle = e^{i\theta}$

$$|\langle \Psi | \Psi' \rangle| = 1$$

$e^{i\theta}$  → Global phase  
overall multiplying factor

→ States differing by a global phase are physically indistinguishable in math  
 $\Psi$  &  $\Psi'$  are 2 diff states  
since phase is different.  
but in QC → physical sig of  $|\Psi\rangle$  &  $|\Psi'\rangle$   
but can't be measured  
∴ identical.

Global Phase: physically not imp

Relative Phase:

$$|\Psi\rangle = a|0\rangle + b|i\rangle$$

$$|X\rangle = a|0\rangle + e^{i\phi} b|i\rangle$$

→ Is  $|X\rangle$  normalized? given  $|a|^2 + |b|^2 = 1$

Yes.  
Phase factor in b component. ( $(e^{i\phi})^2 = 1$ )

→ What abt the measurement statistics?

$$P(|0\rangle) = |a|^2, \quad P(|X\rangle) = |b|^2$$

$|X\rangle$ : also both  $|X\rangle$  &  $|\Psi\rangle$  normalized  
does it imply both states are indistinguishable?

$$\langle \Psi | X \rangle = (a^* \langle 0 | + b^* \langle i |) (a|0\rangle + e^{i\phi} b|i\rangle)$$

$$= |a|^2 + e^{i\phi} |b|^2$$

Since  $|0\rangle$  &  $|i\rangle$  are orthogonal → e.o.  
depending on diff vals of  $a$  &  $b$ ,  
we'll get diff answers.

$$0 \leq |\langle \Psi | X \rangle| \leq 1$$

dep. on relative magnitude of  $a$  &  $b$   
& the val of  $\phi$ , various possibilities

$\phi = \theta \pi$

when  $a = b = \frac{1}{\sqrt{2}}$

$$|+\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$$

$$|-\rangle = \frac{|0\rangle - |1\rangle}{\sqrt{2}}$$

$$\langle +| - \rangle = 0$$

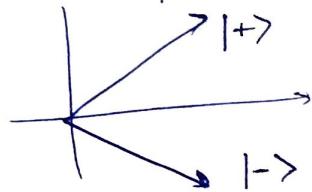
$\langle \psi | \chi \rangle$  maximally distinguishable.

when  $\langle \psi | \chi \rangle = 0$

&  $\not\equiv$  not distinguishable if

$$\langle \psi | \chi \rangle = 1$$

given  $P_\psi = P_\chi$  for both states.



states which differ by a relative phase are distinguishable physically.

Q: Are the states  $\frac{|0\rangle + |1\rangle}{\sqrt{2}}$  &  $\frac{(|0\rangle + i|1\rangle)}{\sqrt{2}}$  equivalent?

Q: For what values of  $\theta$  are the states  $|1\rangle$  &  $\frac{|+> + e^{i\theta}|-\rangle}{\sqrt{2}}$  equivalent?

Q: Consider the state  $\frac{\sqrt{3}}{2}|0\rangle - \frac{1}{2}|1\rangle$

and measurement basis  $\{|0\rangle, |1\rangle\}$ .

What are the possible measurement outcomes & what are their probab?

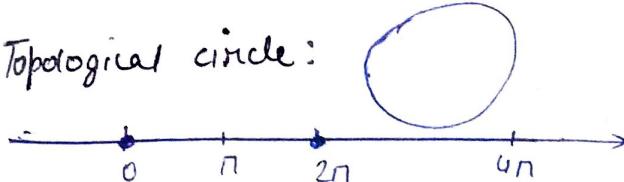
Ans: P of States  $|0\rangle$  &  $|1\rangle$  are equal.

Q: Repeat previous prob for the state

$|+\rangle$  & measurement basis

$$\left\{ \frac{1}{2}|0\rangle + \frac{\sqrt{3}}{2}|1\rangle, \quad \frac{\sqrt{3}}{2}|0\rangle - \frac{1}{2}|1\rangle \right\}$$

Topological circle:



every pt  $x$  &  $x+2\pi$  identical

\* Geometric Representa<sup>n</sup> of the states

Space of a qubit:

$$|\psi\rangle = a|0\rangle + b|1\rangle$$

represented by a vector in a 2D Complex vector space.

Why we work w  $\not\equiv$  normalized vecs?  
to reduce redundancy.

$$|\langle \psi | \psi \rangle|^2 = 1$$

$$|a|^2 + |b|^2 = 1$$

but states diff. by a global phase are also identical (measurably)

→ No one-to-one rela<sup>r</sup> b/w state & vectors  
given phy state does not uniquely  
give us vectors  $\rightarrow$  a unique vector.

Due to  $\psi$  &  $\psi'$  (global phase)  
being measurably same we have  
no redundancy.

We need to identify  $|\psi\rangle$ ,  $|\psi'\rangle$   
for all such phases (global)

in 2D complex vector space  
all lines are parallel  
No 3D & v.  
Since 13th dim can do  
∞ many many vectors  
differing by a global phase

## \* Complex Projective Space of dim-1

$[CP^1]$

not a linear space  
non-linear space  $\Rightarrow$  we don't know how to add vectors.  
 $\therefore$  difficult to work w  $CP^1$

Take  $a$  to be real & +ve.:

$$a = \frac{1}{\sqrt{1+|\alpha|^2}} \quad \& \quad b = \frac{\alpha}{\sqrt{1+|\alpha|^2}}$$

constructing inverse map:  
given  $\alpha$ , find  $a$  &  $b$ .

$$\text{if } a = 0$$

$$\alpha = \infty$$

$$b = \frac{\infty}{\infty} = 1$$

why? because  $\{\infty\}$   
included in ~~range~~ range

state  $|1\rangle$  maps to  $\infty$ .

$$\text{If } |\Psi\rangle = |1\rangle$$

$$\text{ie } a=0, b=1$$

includes pt at  $\infty$

Complex plane: 1D  $C'$  (rot  $C^2$ )  
 $\Rightarrow$  deal w only 1 complex no.

$$\therefore \text{we take } \frac{b}{a} = \alpha$$

we keep the relative phase  
(phase of  $|1\rangle$ /phase of  $|0\rangle$ )

but getting rid of global phase.

$$|a|^2 + |b|^2 = 1$$

$$a^2 + b^2 = 1$$

$$1 - a^2 = b^2$$

$$b^2 = \frac{1-a^2}{a^2}$$

$$\frac{a^2}{1-a^2} = \frac{a^2}{b^2} = \alpha^2$$

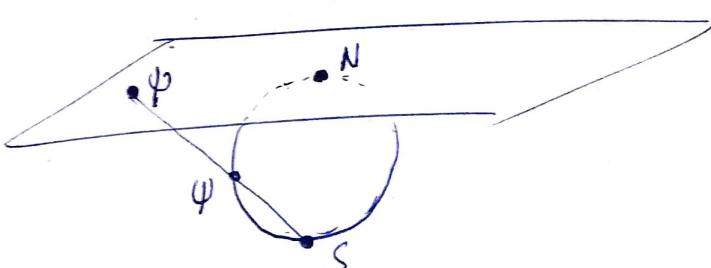
$$\frac{1-a^2}{1+a^2} = \frac{1-\alpha^2}{1+\alpha^2}$$

Block Sphere

Stereographic Projec<sup>n</sup>.

unit sphere

only big  $C'$  plane



plane tangential to unit block sphere  
at North Pole.

any pt on  $\mathbb{C}(\Psi)$  joined to

South pole cuts sphere at 1 pt.  
that is also  $\Psi$ .

Polar form of  $|\psi\rangle$

$$\begin{aligned} |\psi\rangle &= |a| e^{i\theta} |0\rangle + |b| e^{i(\theta+\phi)} |1\rangle \\ &= e^{i\theta} (|a| |0\rangle + |b| e^{i(\theta+\phi)} |1\rangle) \\ &= e^{i\theta} (|a| |0\rangle + |b| e^{i\phi} e^{i\theta} |1\rangle) \end{aligned}$$

$\Downarrow$  global phase can be discarded

$$|\psi\rangle = |a| |0\rangle + |b| e^{i\phi} e^{i\theta} |1\rangle$$

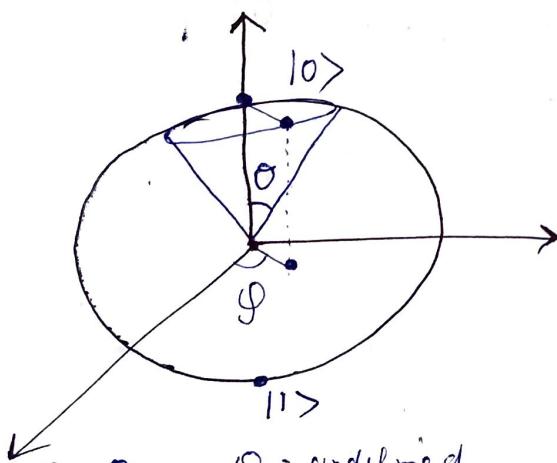
$$|a|^2 + |b|^2 = 1$$

$\therefore$  we can take  $a$  to be real & +ve.

$$|a| = \cos \theta/2 \quad |b| = \sin \theta/2$$

$$|\psi\rangle = \cos \frac{\theta}{2} |0\rangle + \sin \frac{\theta}{2} e^{i\phi} |1\rangle$$

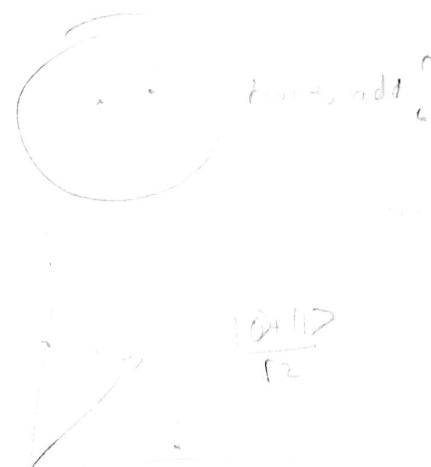
where  $\theta$  &  $\phi$  are the usual polar  
Cs on the unit sphere (Bloch sphere)



$$\theta = 0, \phi : \text{undefined}$$
$$|\psi\rangle = |0\rangle$$

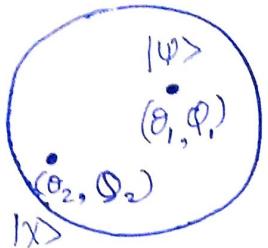
$$\theta = \pi, \phi = 0$$
$$|\psi\rangle = |1\rangle$$

ex  $|+\rangle$  &  $|-\rangle$ , identify on  
Bloch sphere.



$$\frac{|0\rangle + |1\rangle}{\sqrt{2}}$$

$$|U\rangle = |X\rangle$$



- Some rotation in 3D.
- mimicking  $3 \times 3$  real matrix.
- every transformation is a 3D rotat.

$$|\psi\rangle \xrightarrow{U} |\psi\rangle = |X\rangle$$

- changed state
- a different set of  $\theta, \phi$ .

Quantum Key Distribution: Application of extended plane.

→ Cryptography

Tx of msg in encrypted form  
lock the msg. → need a key.  
unlocking also reqs a key.

Forms of key:

Public - Private Key:

Symmetric key:

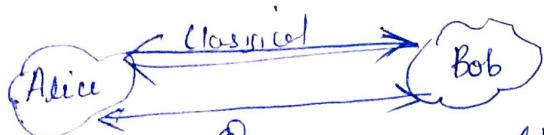
Public key:

Symmetric key:

2 parties universally called Alice & Bob.  
A & B both reqd to maintain security of key. → They generate their own key.

How to ensure security of key?

Q) encrypt works not on math formula?  
but, on properties & f's of qubits.



2 parties can commu. via:

- Alice generates a random seq of N bits.
- ~~1010001110100011~~
- (bit seq)

- A encodes this seq using qubits.
- encode bit-1  
bit 0

makes use of 2 qubits on 2 basis  
Computational & Hadamard basis

$$\begin{pmatrix} |0\rangle \\ |+\rangle \end{pmatrix}, \begin{pmatrix} |1\rangle \\ |- \rangle \end{pmatrix}$$

used to encode bit '1'. used to encode bit '0'.

$$'1' \rightarrow |0\rangle \text{ or } |+\rangle$$

which out of these 2: RANDOM

$$'0' \rightarrow |1\rangle \text{ or } |- \rangle$$

$$\begin{array}{c} \text{bit } 1010001110100011 \\ \text{qubits } 0+ - + 1 + 0 - + 1 - - + 0 \end{array}$$

Now ① transmitted to Bob  
via the Q channel.

(Bit string converted to qubits)

- Bob receives qubit seq.
- Has to make measurement

Randomly choose b/w

Computational & hadamard basis

→ to make measurement (H)  
(C)

choose random seq of basis:

~~say~~ CHCH / HCCH / HCHH / HHCC

• (if Bob choose same measur. basis  
that A used to encode : Bob will  
definitely know which qubit).

(None of A & B know the O's choice of basis)

In the meanwhile :

Over the classical channel :  $\xrightarrow{\text{No of bits}}$

→ A & B ensure that B got all bits  $\begin{cases} \text{Yes} \\ \text{No} \end{cases}$

→ share which basis used for w. bit

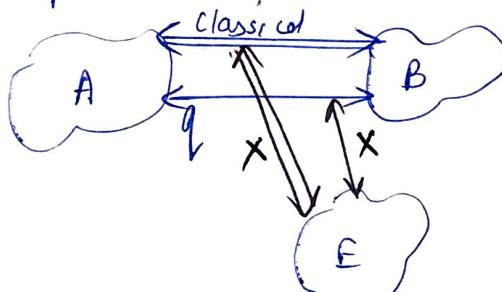
→ for those bits whose basis of encoding & measurement did not match, simply discarded those ~~basis~~ bits.

Sharing only basis info, not which exact bits sent.

On an avg 50% bits dropped.

\* 3rd party eavesdropper

Eve should not be able to tap on classical or Q channel.



Eve should not get 'key'. What is key?

'key' not generated yet.

in the process of generating "key".

⑤ M N N M | N M N N | N N M N | M M N M

M: matched, N: unmatched.

retain only those bits that correspond to 'M'.

| 0 >   | - >   | 1 >              | H >   | T >   | T > | 0 >

Now we know which qubit is representing which bit. so we get

1	0	0	1	0	0	1
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Our key.

Say eve tries to intercept:  
Classical channel :

- A & B share one of bits.  
→ doesn't help
- A & B share

To get access to key has to tap Q channel:

### ~~★~~ Gates:

- $X, Y, Z, H$ .
- $X^2 = Y^2 = Z^2 = H^2 = \text{Id}$
- $XX^\dagger = \text{Id} = YY^\dagger = \dots$

### ~~★~~ Phase Shift-gate:

$$P = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\theta} \end{pmatrix}$$

$$P|\psi\rangle = P(a|0\rangle + b|1\rangle)$$

use k. that:  $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  &  $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

$$\therefore |\psi\rangle = \begin{pmatrix} a \\ b \end{pmatrix}$$

$$\therefore P|\psi\rangle = a|0\rangle + e^{i\theta}b|1\rangle \xrightarrow{\text{relative phase}}$$

$$\bullet (U V W \dots)^+ = \dots W^+ V^+ U^+$$

can't do significant computa<sup>n</sup> task. need multiple qubits.

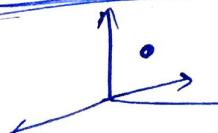
### ~~★~~ Multiple Qubits:

A qubit: 2D complex vec space or a pt on Bloch sphere.

Entanglement?

Need to know:  
How size of state space changes  
when 'N' particles in classical v/s Q.

### In Classical:

  
To represent a particle in 3D: need 3 numbers  
to specify position  
lets say state  $\langle$  position  $\rangle$  momentum  $\times$

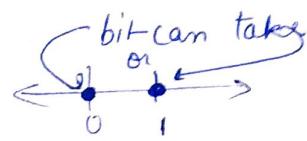
To represent 'N' particles

need  $3N$  numbers.  $\rightarrow$  linear increase in state space.

### In Classical:

to represent 2 bits

need 2 mps.



size of state space grows linearly w no. of bits/particles.

~~★~~ say: N particles

$U, V, W \dots$

$U \oplus V \oplus W \oplus \dots$

direct sum

$$\dim(U \oplus V \oplus \dots)$$

- ① first particle has state space  $U$  say  $\mathbb{R}^2$
- ② second " " " " "  $V$  say  $\mathbb{R}^3$   
for ① need to give 2 mps  
② " " " " " 3 mps

### Quantum:

2 qubits  $\uparrow \uparrow$

if  $\uparrow$  make measurement:

$$\uparrow \begin{cases} |0\rangle \\ |1\rangle \end{cases} \quad \begin{cases} |0\rangle \\ |1\rangle \end{cases} \quad |0\rangle |0\rangle$$

$$\uparrow \begin{cases} |0\rangle \\ |1\rangle \end{cases} \quad \begin{cases} |0\rangle \\ |1\rangle \end{cases} \quad |0\rangle |1\rangle$$

$$\uparrow \begin{cases} |0\rangle \\ |1\rangle \end{cases} \quad \begin{cases} |1\rangle \\ |0\rangle \end{cases} \quad |1\rangle |0\rangle$$

$$\uparrow \begin{cases} |0\rangle \\ |1\rangle \end{cases} \quad \begin{cases} |1\rangle \\ |0\rangle \end{cases} \quad |1\rangle |1\rangle$$

$$\rightarrow \alpha_{00}|00\rangle + \alpha_{01}|01\rangle + \alpha_{10}|10\rangle$$

$$+ \alpha_{11}|11\rangle$$

$$\text{where } |\alpha_{00}|^2 + \dots + |\alpha_{11}|^2 = 1$$