

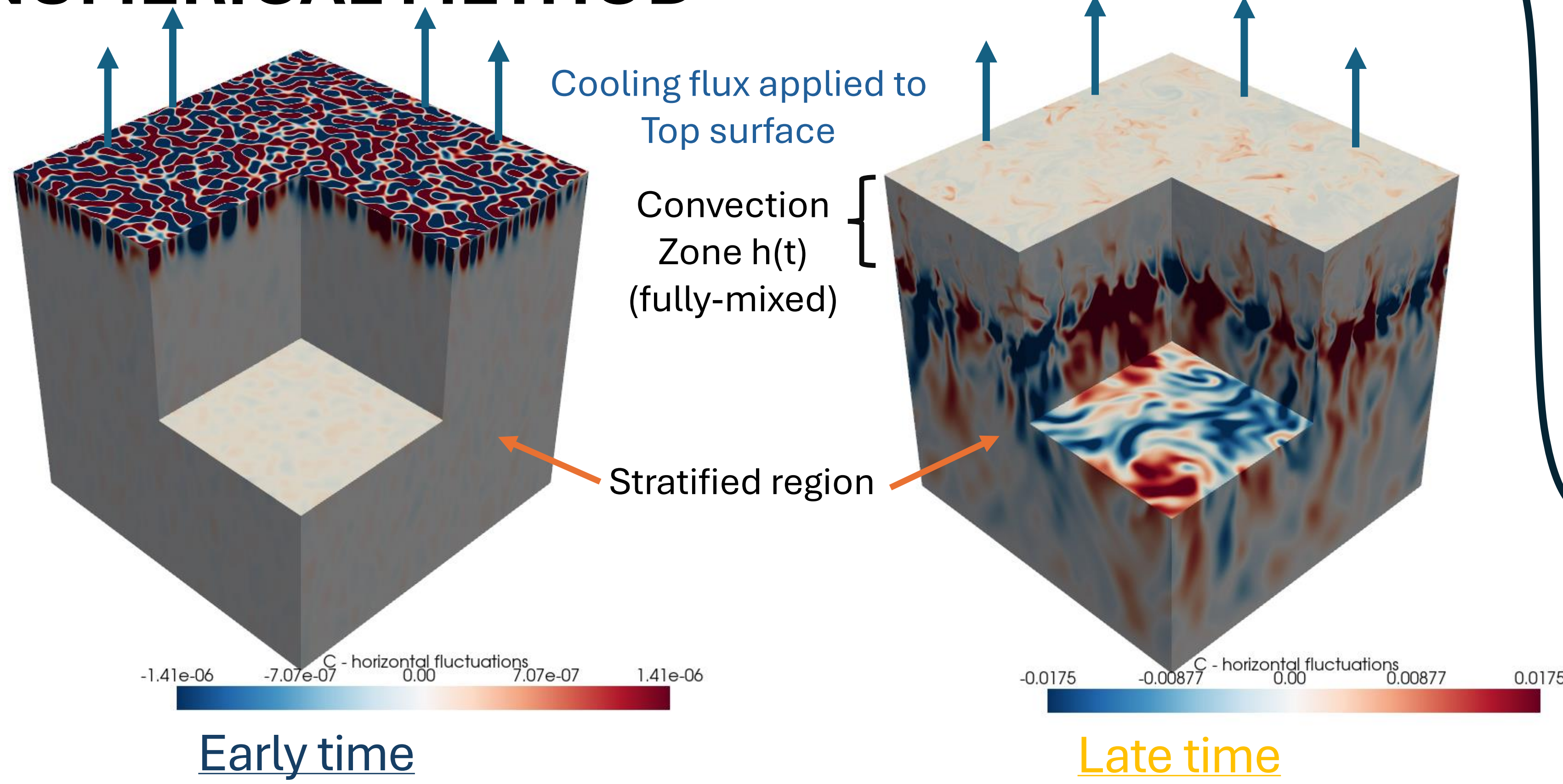
# 3D Simulations Support Stalled Convective Mixing of Jupiter's Interior

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## NUMERICAL METHOD



We simulate entrainment and the penetration of a convection zone into a stably-stratified region, using Dedalus [4]. In doing so, we adopt the Boussinesq approximation, where density variations are neglected except in the buoyancy term ( $\rho = \rho_0(\beta S - \alpha T)$ ). The domain is a 3D box with equal dimensions  $H$ , periodic boundary conditions in the horizontal directions, and impermeable, stress-free boundaries in the vertical direction. Convection is driven by imposing a cooling flux at the top boundary.

**Boussinesq equations** in non-dimensionalized form describe the flow:

$$\begin{aligned} \nabla \cdot \tilde{u} &= 0 \\ \frac{\partial \tilde{u}}{\partial \tilde{t}} + (\tilde{u} \cdot \nabla) \tilde{u} &= -\nabla P + \text{Pr} \cdot \mathcal{R}(\tilde{T} - \tilde{S})\hat{z} - \left(\frac{\text{Pr}}{\text{Ek}}\right) \tilde{u} \times \hat{z} + \text{Pr} \nabla^2 \tilde{u} \\ \frac{\partial \tilde{S}}{\partial \tilde{t}} + (\tilde{u} \cdot \nabla) \tilde{S} &= \tau \nabla^2 \tilde{S} \\ \frac{\partial \tilde{T}}{\partial \tilde{t}} + (\tilde{u} \cdot \nabla) \tilde{T} &= \tau \nabla^2 \tilde{T} \end{aligned} \quad \begin{aligned} \mathcal{R} &= \frac{g\beta H^3 S_0}{k_T \nu} = 2 \cdot 10^9 \\ \text{Pr} &= \frac{\nu}{k_T} = 0.1 \\ \tau &= \frac{k_S}{k_T} = 0.1 \\ \text{Ek} &= \frac{\nu}{2\Omega H^2} = 3 \cdot 10^{-6} \end{aligned}$$

Parameters

Hindman & Fuentes [1] estimated the convective velocity using mixing-length theory

Non-rotating	$U_{NR} \sim \left(\frac{g\alpha}{\rho_0 c_P}\right)^{1/3} (hF)^{1/3}$
Rotating	$U_R \sim \left(\frac{g\alpha}{\rho_0 c_P}\right)^{2/5} \left(\frac{hF^2}{2\Omega}\right)^{1/5}$

With the assumption that the increase in gravitational potential is equal to the flux of kinetic energy of the flow,

$$g\rho h \frac{dh}{dt} = \gamma U^3,$$

we predict the size of the convection zone  $h(t)$  for 4 different configurations:

	Non-rotating	Rotating
Constant Flux $F = F_0$	$\left[(2C) \left(\frac{F_0}{F_c}\right) t\right]^{1/2}$	$\left[(2C) \left(\frac{F_0}{F_c}\right)^{6/5} \left(\frac{\mathcal{R} \cdot \text{Ek}^3}{\text{Pr}^2}\right)^{1/5} t\right]^{5/12}$
Dwindling Flux $F = F_0 \left(\frac{t_0}{t_0 + t}\right)$	$\left[(2C) \left(\frac{F_0}{F_c}\right) t_0 \ln \left \frac{t+t_0}{t}\right \right]^{1/2}$	$\left\{ (2C) \left(\frac{F_0}{F_c}\right)^{6/5} \left(\frac{\mathcal{R} \cdot \text{Ek}^3}{\text{Pr}^2}\right)^{1/5} t_0 \left[1 - \left(\frac{t_0}{t_0 + t}\right)^{1/5}\right] \right\}^{5/12}$

The flux  $F_0$  is normalized with  $F_c = \frac{\rho_0 c_P k_T \beta S_0}{\alpha H}$ .

We adopt the efficiency coefficient  $C = 1 - \epsilon + 2\gamma$  from the 2D simulations of Fuentes et al. [5]. We take  $\epsilon = 0.4$  and  $\gamma = 0.9$  based on the values they measured in their simulations. These values give good agreement with the numerical results.

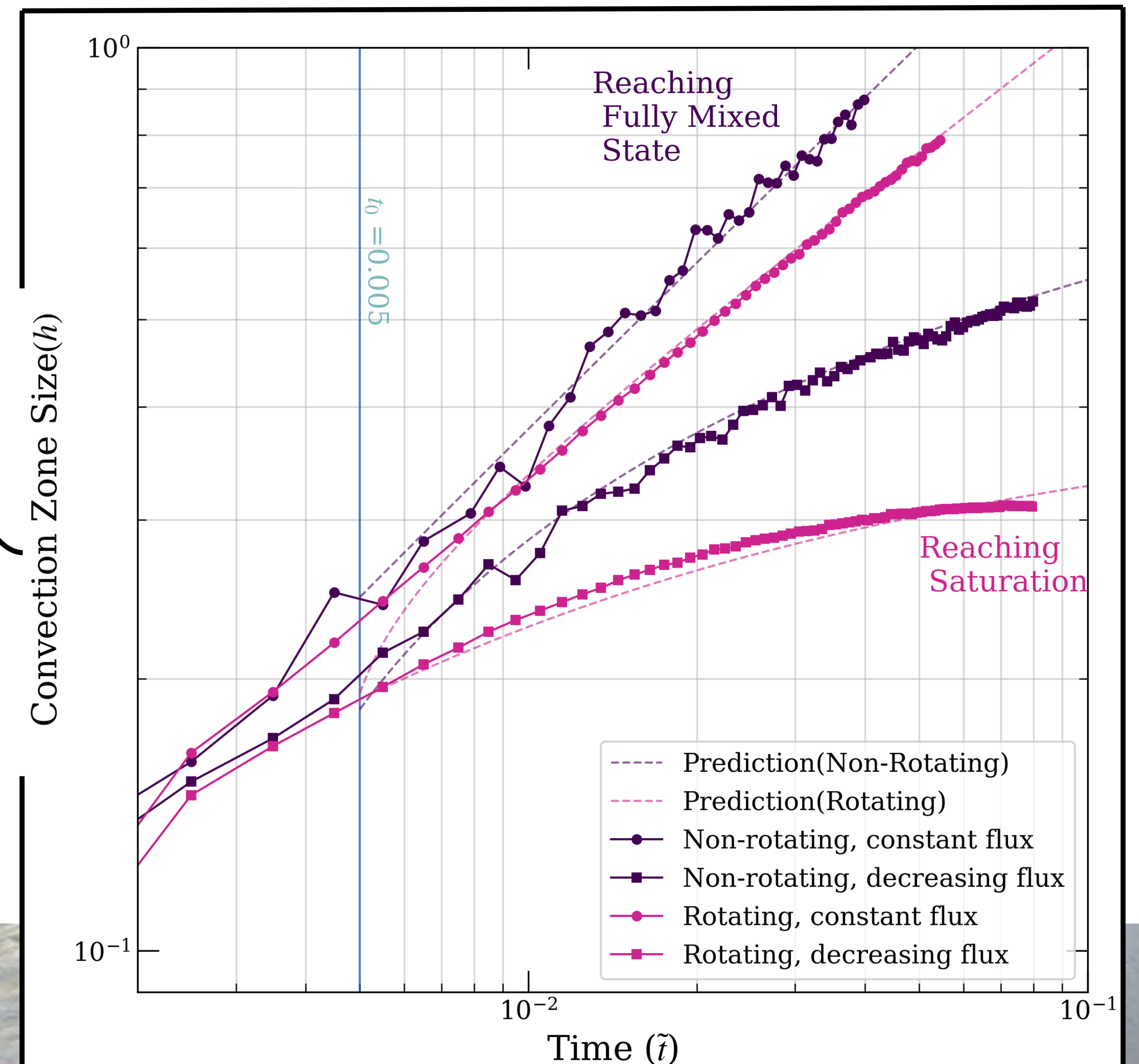
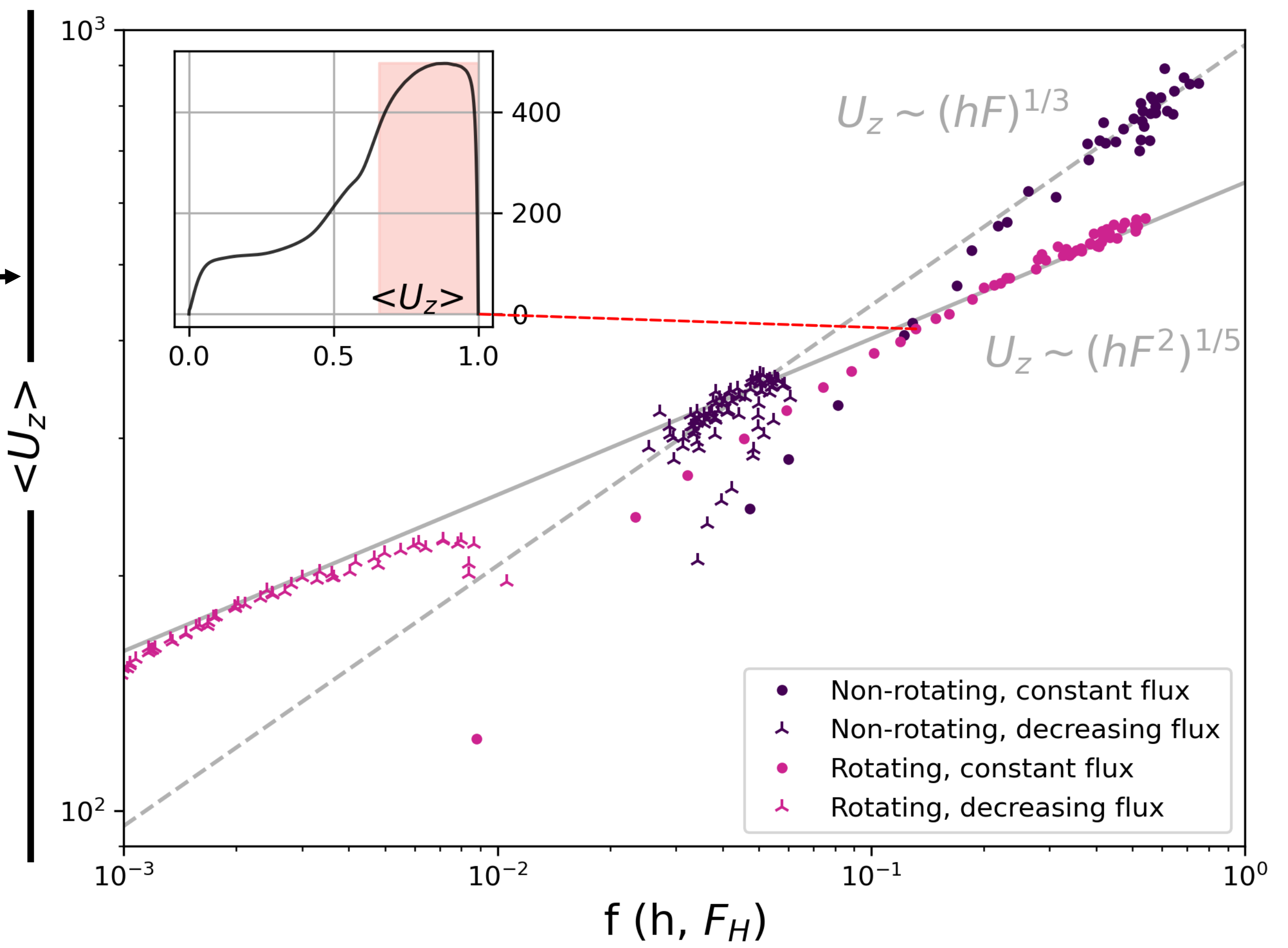
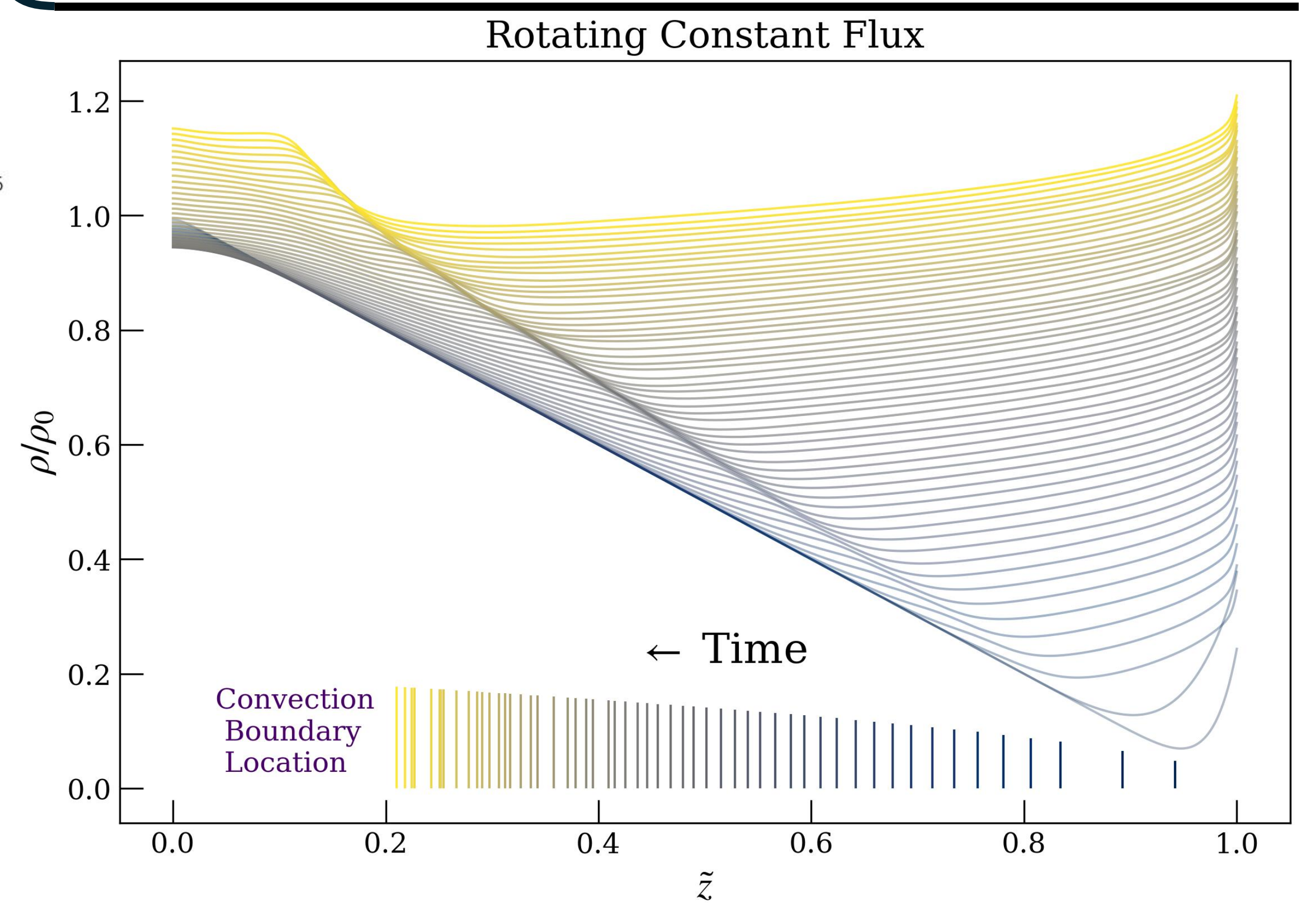
## ANALYTIC MODEL

## CONCLUSION

**Rotation**, when combined **with dwindling surface flux**, results in a finite convective mixing depth, potentially explaining Jupiter's diffuse core. **Both** rotation and decreasing flux are essential for the convection zone to stall.

## MOTIVATION

Observations of Jupiter's gravitational moments by Juno imply that, instead of a distinct dense core surrounded by fully mixed light elements, Jupiter still possesses a stably-stratified core [1]. What prevents the primordial concentration gradient from mixing by convection is not understood. Without composition gradients, the convection zone would mix the entire planet in less than a million years [2]. Hindman & Fuentes (2023) used an analytic entrainment model to show that rapid rotation combined with the dwindling surface cooling flux of Jupiter would stall the growth of the convection zone [3]. **Here we test these analytic predictions by carrying out 3D simulations with both a time-dependent cooling flux and rotation.**



- [1] Howard, S. et al. 2023, A&A, 672, A33  
 [2] Muller, S., Helled, R., & Cumming, A. 2020, A&A, 638, A121  
 [3] Hindman, B.W. & Fuentes, J.R. 2023, ApJL, 957, L23  
 [4] Burns, K.J. et al. 2020, Phys. Rev. Res., 2, 023068  
 [5] Fuentes, J.R., Cumming, A., & Anders, E.H. 2020, Phys. Rev. Fluids, 5, 124501