

# PHSX 461: HW07

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## 3.7

- a) Suppose that  $f(x)$  and  $g(x)$  are two eigenfunctions of an operator  $\hat{Q}$ , with the same eigenvalue  $q$ . Show that any linear combination of  $f$  and  $g$  is itself an eigenfunction of  $\hat{Q}$ , with eigenvalue  $q$ .

$$\begin{aligned}\hat{Q}|f\rangle &= q|f\rangle \quad , \quad \hat{Q}|g\rangle = q|g\rangle \\ |\alpha\rangle &= a|f\rangle + b|g\rangle\end{aligned}$$

$$\begin{aligned}\hat{Q}|\alpha\rangle &= \hat{Q}[a|f\rangle + b|g\rangle] \\ &= \hat{Q}(a|f\rangle) + \hat{Q}(b|g\rangle) \\ &= a\hat{Q}|f\rangle + b\hat{Q}|g\rangle \\ &= aq|f\rangle + bq|g\rangle \\ &= q[a|f\rangle + b|g\rangle] = q|\alpha\rangle\end{aligned}$$

- b) Check that  $f(x) = \exp(x)$  and  $g(x) = \exp(-x)$  are eigenfunctions of the operator  $d^2/dx^2$ , with the same eigenvalue. Construct two linear combinations of  $f$  and  $g$  that are orthogonal eigenfunctions on the interval  $(-1, 1)$ .

$$\begin{aligned}\hat{Q}|f\rangle &\Rightarrow \frac{\partial^2}{\partial x^2}e^x = e^x \rightarrow q = 1 \\ \hat{Q}|g\rangle &\Rightarrow \frac{\partial^2}{\partial x^2}e^{-x} = (-1)^2e^{-x} = e^{-x} \rightarrow q = 1\end{aligned}$$

$$|\alpha\rangle = a|f\rangle + b|g\rangle \quad , \quad |\beta\rangle = c|f\rangle + d|g\rangle$$

For orthogonality, we need  $\langle\alpha|\beta\rangle = \int \alpha(x)^*\beta(x) dx = 0$

$$0 = \int_{-1}^1 (a^* f^*(x) + b^* g^*(x))^* (cf(x) + dg(x)) dx$$

$$0 = \int_{-1}^1 (ae^x + be^{-x})(ce^x + de^{-x}) dx$$

$$0 = \int_{-1}^1 [ace^{2x} + ad + bc + bde^{-2x}] dx$$

$$0 = \frac{ac}{2}e^{2x} + [ad + bc]x - \frac{bd}{2}e^{-2x} \Big|_{-1}^1$$

$$0 = \frac{ac}{2}e^2 + [ad + bc] - \frac{bd}{2}e^{-2} - \left[ \frac{ac}{2}e^{-2} - [ad + bc] - \frac{bd}{2}e^2 \right]$$

$$0 = e^2 \left[ \frac{ac}{2} + \frac{bd}{2} \right] + e^{-2} \left[ \frac{-bd}{2} + \frac{-ac}{2} \right] + 2[ad + bc]$$

The easiest solution to this is if we let  $ac = -bd$ . Two solutions to this question yield the linear combination:

$$|\alpha\rangle = |f\rangle + |g\rangle \quad , \quad |\beta\rangle = |f\rangle - |g\rangle$$

$$|\alpha\rangle = -|f\rangle + |g\rangle \quad , \quad |\beta\rangle = |f\rangle + |g\rangle$$

## 3.9

- a) *Cite a Hamiltonian from Chapter 2 (other than the harmonic oscillator) that has only a discrete spectrum.*
- b) *Cite a Hamiltonian from Chapter 2 (other than the free particle) that has only a continuous spectrum.*
- c) *Cite a Hamiltonian from Chapter 2 (other than the finite square well) that has both a discrete and a continuous part to its spectrum.*

### 3.13

Show that

$$\langle x \rangle = \int \Phi^* \left( i\hbar \frac{\partial}{\partial p} \right) \Phi \, dp$$

Let us start by defining what  $\Phi$  is:

$$\Phi = \frac{1}{\sqrt{2\pi\hbar}} \int e^{ipx/\hbar} \Psi(x, t) \, dx$$

So, let's plug this in and just run the calculus:

$$\begin{aligned} & \int \Phi^* \left( i\hbar \frac{\partial}{\partial p} \right) \Phi \, dp \\ &= \frac{i\hbar}{2\pi\hbar} \int \left[ \int e^{ipx/\hbar} \Psi \, dx \right]^* \frac{\partial}{\partial p} \left[ \int e^{ipx'/\hbar} \Psi \, dx' \right] \\ &= \frac{i}{2\pi} \int \left[ \int e^{ipx/\hbar} \Psi^* \, dx \right] \left[ \int \frac{-ix'}{\hbar} e^{ipx'/\hbar} \Psi \, dx' \right] \end{aligned}$$

Since these three variables,  $p, x, x'$ , are all independent, they can be moved in and out of each other's integrals.

$$\begin{aligned} &= \frac{1}{2\pi\hbar} \int e^{-i(\frac{p}{\hbar}(x-x'))} \Psi^*(x, t) x' \Psi(x', t) \, dp \, dx \, dx' \\ &= \frac{\hbar}{\hbar} \delta \left[ \frac{p}{\hbar}(x - x') \right] \int \Psi^*(x, t) x' \Psi(x', t) \, dx \, dx' \Big|_{-\infty}^{\infty} \\ &= \int \Psi^*(x, t) x \Psi(x, t) \, dx = \langle x \rangle \end{aligned}$$

### 3.26

Consider a three-dimensional vector space spanned by an orthonormal basis  $|1\rangle, |2\rangle, |3\rangle$ . Kets  $|\alpha\rangle$  and  $|\beta\rangle$  are given by

$$|\alpha\rangle = i|1\rangle - 2|2\rangle - i|3\rangle, \quad |\beta\rangle = i|1\rangle + 2|3\rangle$$

a) Construct  $\langle\alpha|$  and  $\langle\beta|$  (in terms of the dual basis  $\langle 1|, \langle 2|, \langle 3|$ ).

We can think of  $|\alpha\rangle$  and  $|\beta\rangle$  as wavefunctions, since the mechanics here are the same (orthonormal basis, linear combinations of states, ...). So, pulling that analog:

$$|\Psi\rangle = \sum c_n |f_n\rangle \Rightarrow \langle\Psi| = \sum c_n^* \langle f_n|$$

and thus:

$$\langle\alpha| = -i\langle 1| - 2\langle 2| + i\langle 3| \quad , \quad \langle\beta| = -i\langle 1| + 2\langle 3|$$

b) Find  $\langle\alpha|\beta\rangle$  and  $\langle\beta|\alpha\rangle$ , and confirm that  $\langle\beta|\alpha\rangle = \langle\alpha|\beta\rangle^*$ .

$$\begin{aligned} \langle\alpha|\beta\rangle &= (-i\langle 1| - 2\langle 2| + i\langle 3|)(i|1\rangle + 2|3\rangle) \\ &= -i^2\langle 1|1\rangle + 2i\langle 3|3\rangle \\ &= 1 + 2i \end{aligned}$$

$$\begin{aligned} \langle\beta|\alpha\rangle &= (-i\langle 1| + 2\langle 3|)(i|1\rangle - 2|2\rangle - i|3\rangle) \\ &= -i^2\langle 1|1\rangle - 2i\langle 3|3\rangle \\ &= 1 - 2i \end{aligned}$$

$$\langle\alpha|\beta\rangle^* = (1 + 2i)^* = 1 - 2i = \langle\beta|\alpha\rangle$$

c) Find all nine matrix elements for the operator  $\hat{A} = |\alpha\rangle\langle\beta|$ , in this basis, and construct the matrix  $A$ . Is it hermitian?

$$\hat{Q} \Rightarrow Q = \begin{pmatrix} Q_{11} & Q_{12} & \cdots \\ Q_{21} & Q_{22} & \cdots \\ \cdots & \cdots & \cdots \end{pmatrix}, \text{ where } Q_{mn} = \langle e_m | \hat{Q} | e_n \rangle$$

We will use this to calculate the  $A_{13}$  element, then point out the pattern and fill out the rest of the matrix.

$$\begin{aligned} A_{13} &= \langle 1 | \alpha \rangle \langle \beta | 3 \rangle \\ &= \langle 1 | (i | 1 \rangle - 2 | 2 \rangle - i | 3 \rangle) (-i \langle 1 | + 2 \langle 3 |) | 3 \rangle \\ &= (i)(2) = 2i \end{aligned}$$

So, an element  $A_{mn}$  can be found by taking the product of the  $m$ th coefficient from  $|\alpha\rangle$  and the  $n$ th coefficient from  $\langle\beta|$ . Using this method:

$$\hat{A} = \begin{pmatrix} (i)(-i) & (i)(0) & (i)(2) \\ (-2)(-i) & (-2)(0) & (-2)(2) \\ (-i)(-i) & (-i)(0) & (-i)(2) \end{pmatrix} = \begin{pmatrix} 1 & 0 & 2i \\ 2i & 0 & -4 \\ -1 & 0 & -2i \end{pmatrix}$$

To test if this is hermitian,  $\hat{A}^\dagger = \hat{A}$

$$\hat{A}^\dagger = \begin{pmatrix} 1 & 0 & 2i \\ 2i & 0 & -4 \\ -1 & 0 & -2i \end{pmatrix}^\dagger = \begin{pmatrix} 1 & 0 & 2i \\ 2i & 0 & -4 \\ -1 & 0 & -2i \end{pmatrix}^{*T} = \begin{pmatrix} 1 & -2i & -1 \\ 0 & 0 & 0 \\ -2i & -4 & -2i \end{pmatrix} \neq \hat{A}$$

So,  $A$  is **not** hermitian.

## Question 5.

*Prove that the momentum operator,  $\hat{p}$  is Hermitian.*

**Hint:** *you will need to assume that any functions you use are normalizable. You may also use the results from the previous homework assignment.*

We want to show that  $\hat{p}^\dagger = \hat{p}$ . To do this, first we can identify that a hermitian operator in momentum space is also a hermitian operator in position space; so, we will be doing this problem with  $\Phi$ , instead of  $\Psi$ .

$$\begin{aligned} & \langle \Phi | \hat{p}^\dagger | \Phi \rangle \\ &= \langle \hat{p} \Phi | \Phi \rangle \\ &= \int (\hat{p} \Phi)^* \Phi \, dp \\ &= \int (p \Phi)^* \Phi \, dp \\ &= \int \Phi^* p \Phi \, dp \\ &= \int \Phi^* \hat{p} \Phi \, dp \\ &= \langle \Phi | \hat{p} | \Phi \rangle \rightarrow \hat{p}^\dagger = \hat{p} \end{aligned}$$

### 3.33

An operator  $\hat{A}$ , representing observable  $A$ , has two (normalized) eigenstates  $\psi_1$  and  $\psi_2$ , with eigenvalues  $a_1$  and  $a_2$ , respectively. Operator  $\hat{B}$ , representing observable  $B$ , has two (normalized) eigenstates  $\phi_1$  and  $\phi_2$ , with eigenvalues  $b_1$  and  $b_2$ . The eigenstates are related by

$$\psi_1 = (3\phi_1 + 4\phi_2)/5, \quad \psi_2 = (4\phi_1 - 3\phi_2)/5$$

- a) *Observable  $A$  is measured, and the value  $a_1$  is obtained. What is the state of the system (immediately after this measurement)?*

The wavefunction has been observed and thus has been collapsed to  $\psi_1$ . So, the state is  $\phi_1$

- b) *If  $B$  is now measure, what are the possible results, and what are their probabilities?*

Since we are completely in state  $\psi_1$ , the probabilities will only be decedents of that equation. There will be a  $(\frac{3}{5})^2 = \frac{9}{25} = 35\%$  chance for  $b_1$  and a  $(\frac{4}{5})^2 = \frac{16}{25} = 65\%$  chance for  $b_2$ .

- c) *Right after the measurement of  $B$ ,  $A$  is measured again. What is the probability of getting  $a_1$ ? (Note that the answer would be quite different if I had told you the outcome of the  $B$  measurement.*

If we are asking for the probability of  $a_1$ , we are thinking about  $\psi_1$ :

$$\begin{aligned} \psi_1^2 &= \frac{9\phi_1^2 + 24\phi_1\phi_2 + 16\phi_2^2}{25} \\ &= \frac{3^4 + 23^2 4^2 + 4^4}{5^4} = 1 \end{aligned}$$

To do a quick check to make sure that we don't need to renormalize:

$$\begin{aligned} \psi_2^2 &= \frac{16\phi_1^2 - 24\phi_1\phi_2 + 9\phi_2^2}{25} \\ &= \frac{3^2 4^2 - 23^2 4^2 + 3^2 4^2}{5^4} = 0 \end{aligned}$$

So, it would seem that there is a 100% chance that we get back  $a_1$ .



But what about that hint? What would it be then? Well, if we say that  $B = b_1$ ,  $\phi_1 = 1$  and  $\phi_2 = 0$ . Thus,  $\psi_1^2 = \frac{9}{25} = 35\%$  and  $\psi_2^2 = \frac{16}{25} = \frac{16}{25} = 65\%$ . This is a very different answer from what we got, so the “100%” is consistent with the hint.