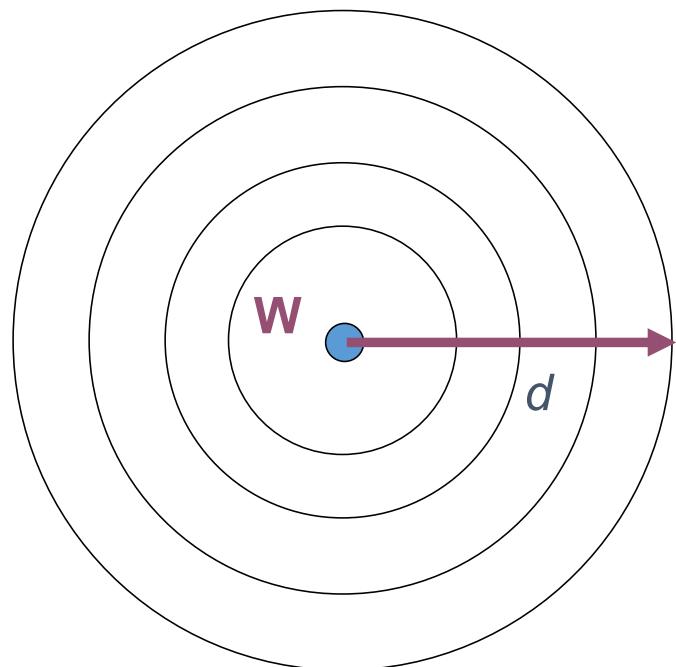


Sound propagation



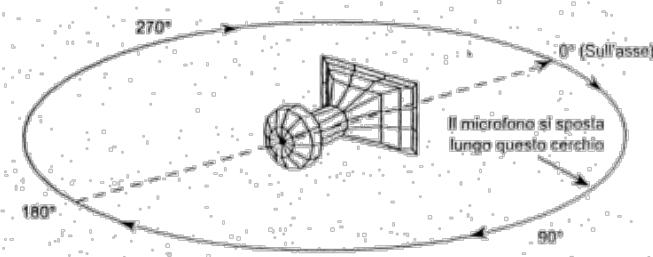
$$L_W = 10 \log \frac{W}{W_{RIF}}$$

$$I = \frac{W}{4\pi d^2} \quad (Wm^{-2})$$

$$L_I = L_W - 10 \log(4\pi) - 10 \log d^2$$

$$L_p = L_W - 20 \log d - 11$$

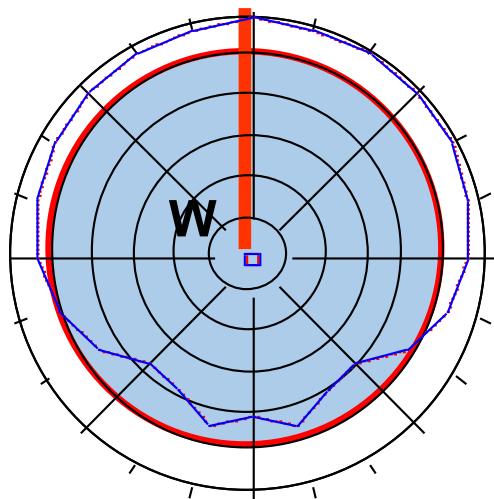
Sound propagation



$$Q = \frac{I}{I_s} - \frac{p_{eff}^2}{p_{effS}^2}$$

$$DI = 10 \log Q$$

$$L_p = L_W + DI - 20 \log d - 11$$



Sound propagation

$$L = L_W + DI - 20 \log r - 11 - \sum_i Att_i$$

- Attenuations are due to:
 - Atmospheric absorption
 - Ground interference/absorption
 - Barriers and other obstacles
 - Weather-dependent phenomena (wind, temperature)

Sound propagation

$$L = L_W + DI - 20 \log r - 11 - \sum_i Att_i$$

- Attenuations are due to:
 - Atmospheric absorption
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Sound propagation

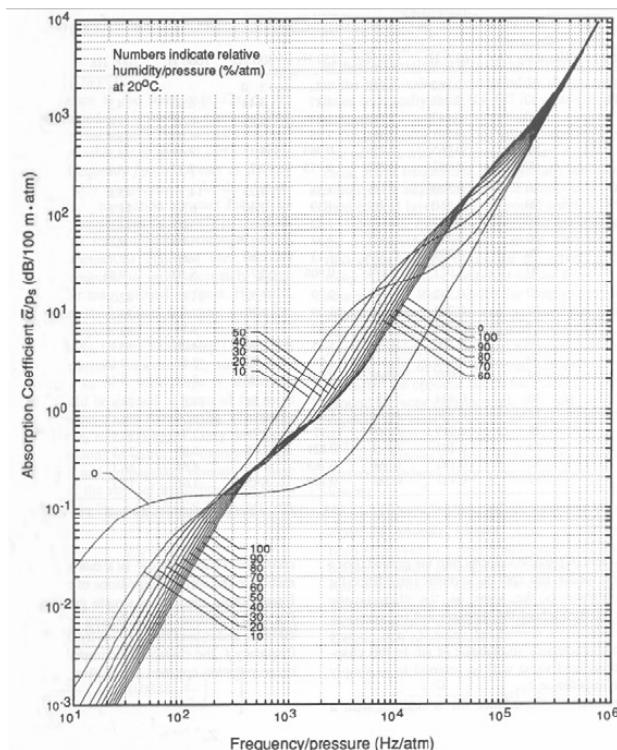
$$Att_{air} = a \cdot r \quad [dB]$$

$$I(r) = I_0 e^{-mr}$$

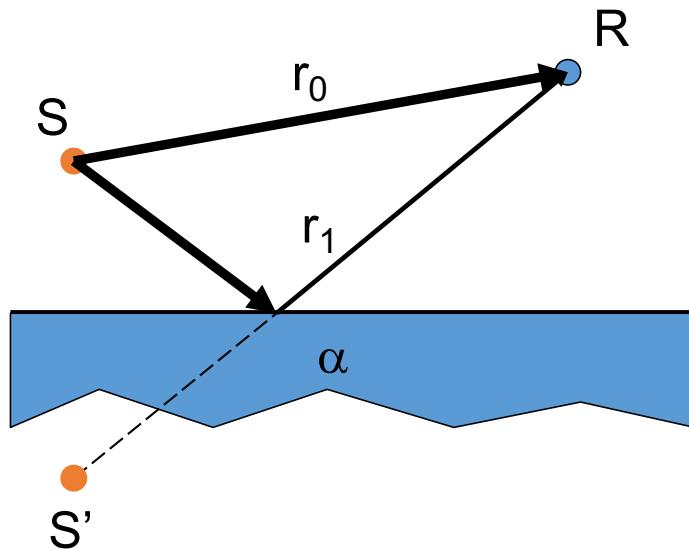
$$Att_{air} = -10 \log_{10} \frac{I(r)}{I_0} = -10 \log_{10} e^{-mr}$$

$$\log_{10} x = \ln x / \ln 10$$

$$Att_{air} = -10(\ln e^{-mr} / \ln 10) = 4.34m \cdot r$$



Sound propagation



$$I_0 = \frac{W}{4\pi r_0^2} \quad \left(\frac{W}{m^2} \right)$$

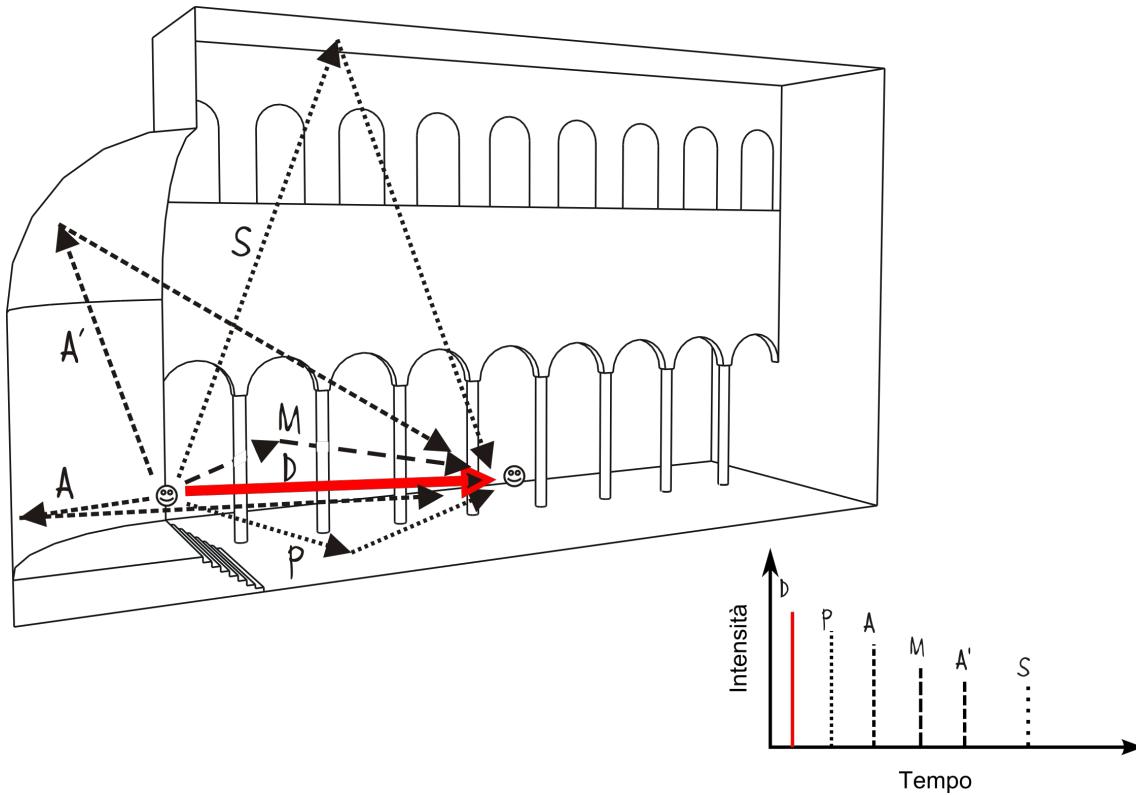
$$I_1 = \frac{W(1 - \alpha)}{4\pi r_1^2} \quad \left(\frac{W}{m^2} \right)$$

$$I_0 = \frac{W}{4\pi r_0^2} \exp(-mr_0) \quad \left(\frac{W}{m^2} \right)$$

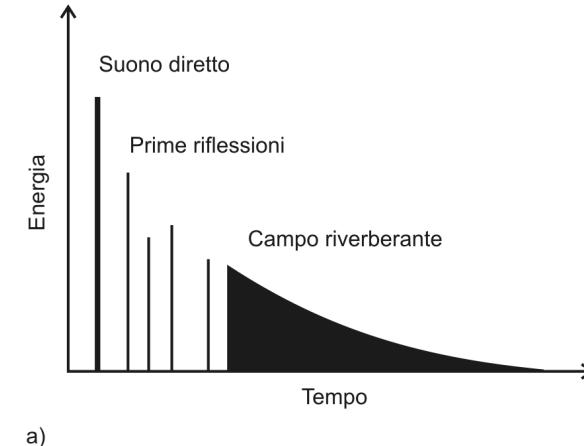
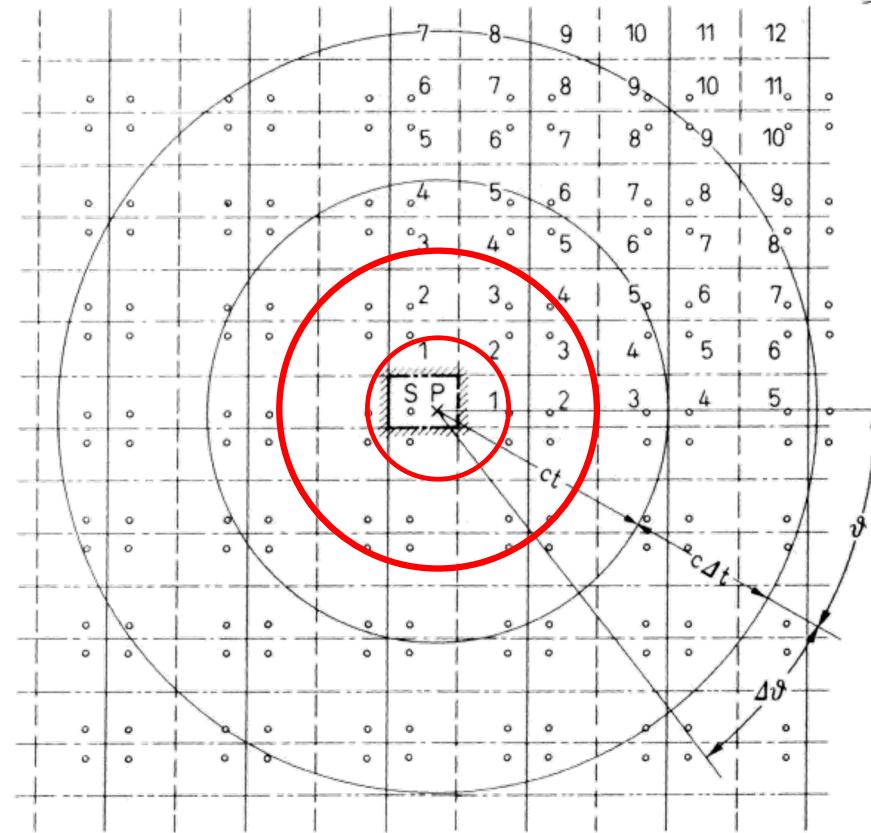
$$I_1 = \frac{W(1 - \alpha)}{4\pi r_1^2} \exp(-mr_1) \quad \left(\frac{W}{m^2} \right)$$

N.B.: Large surface compared to wavelength!

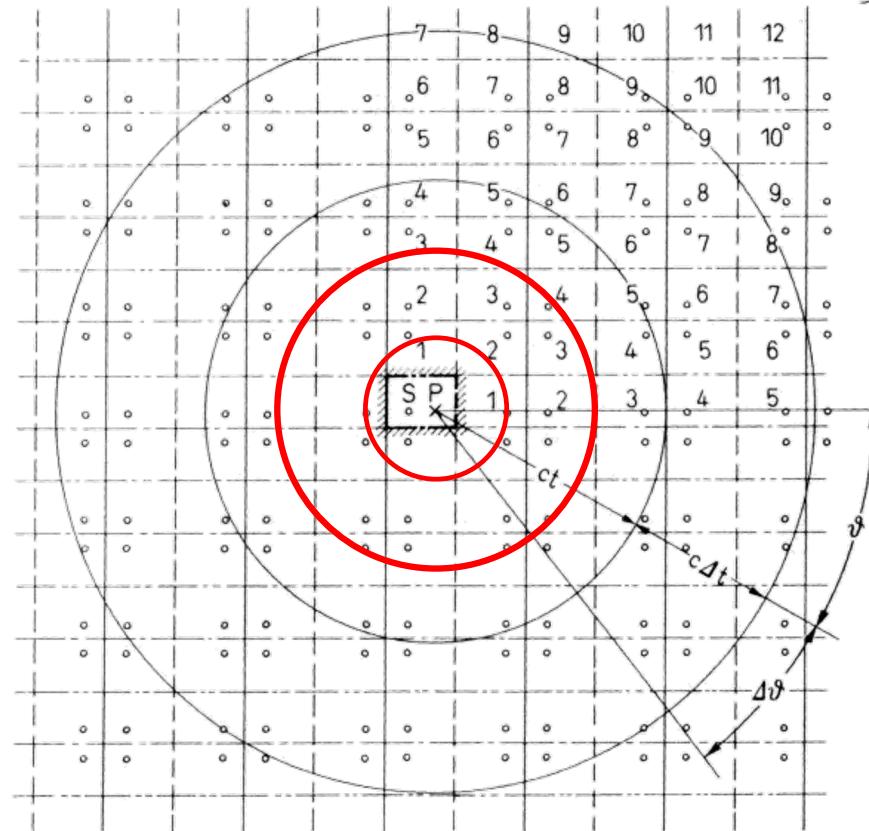
Sound propagation



Sound propagation



Sound propagation



$$I(t) = \frac{W(1-\alpha)^{nt}}{4\pi(ct)^2} \exp(-mct) dN$$

$$dN = \frac{4\pi(ct)^2 \cdot cdt}{V}$$

$$I(t) = \frac{W(1-\alpha)^{nt}}{4\pi(ct)^2} \exp(-mct) \times \frac{4\pi(ct)^2 cdt}{V}$$

$$I(t) = \frac{Wcdt}{V} (1-\alpha)^{nt} \exp(-mct)$$

$$I(t) = I_0 (1-\alpha)^{nt} \exp(-mct)$$

Sound propagation

Assumptions:

- The **Mean Free Path**, can be derived from kinetic theory of gases

$$l = \frac{4V}{S} \quad [\text{m}]$$

Mean time interval between two subsequent reflections

$$t = \frac{l}{c} = \frac{4V}{cS}$$

Mean number of surfaces hit per second

$$n = \frac{1}{t} = \frac{cS}{4V}$$

Sound propagation

$$I(t) = I_0 (1 - \alpha)^{nt} \exp(-mct)$$

$$n = \frac{1}{t} = \frac{cS}{4V}$$

$$(1 - \alpha)^{nt} = \exp[\ln(1 - \alpha)^{nt}] = \exp[nt \cdot \ln(1 - \alpha)]$$

$$I(t) = I_0 \exp[-mc + n \ln(1 - \alpha)] t$$

$$I(t) = I_0 \exp\left[-ct \frac{4mV - S \ln(1 - \alpha)}{4V}\right]$$

Sound propagation

$$L_I(t) = L_0 - 4.34 \left[ct \frac{4mV - S \ln(1-\alpha)}{4V} \right]$$

$$T_{60} = \frac{60}{4.34c} \frac{4V}{4mV - S \ln(1-\alpha)} = 0.161 \frac{V}{4mV - S \ln(1-\alpha)}$$
Eyring's formula

$$\ln(1+x) = x - \frac{1}{2}x^2 + \frac{1}{3}x^3 - \dots \Rightarrow \ln(1-\alpha) \approx -\alpha$$

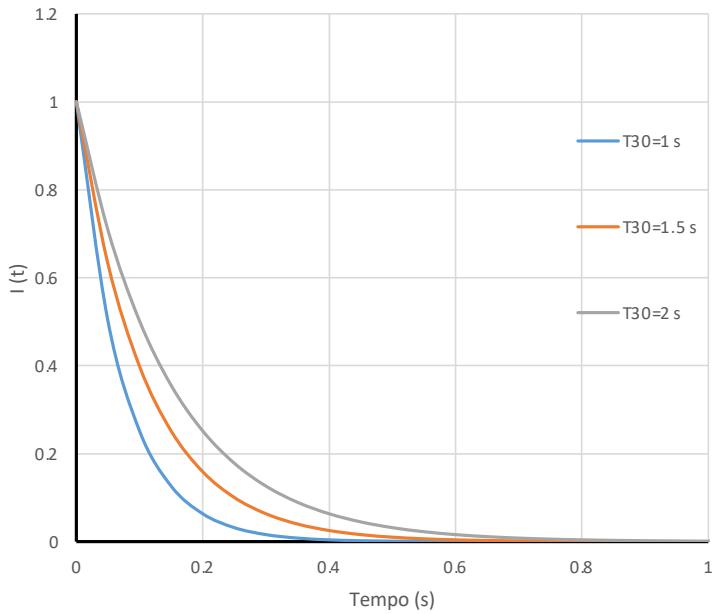
$$T_{60} = 0.161 \frac{V}{4mV + S\alpha} = 0.161 \frac{V}{A}$$
Sabine's formula

Sound propagation

$$T_{60} = 0.161 \frac{V}{4mV - S \ln(1 - \alpha)}$$

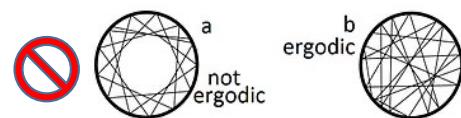
$$I(t) = I_0 \exp \left[-ct \frac{4mV - S \ln(1 - \alpha)}{4V} \right]$$

$$I(t) = I_0 \exp \left[-ct \frac{0.161}{4T_{60}} \right] = I_0 \exp[-13.8t/T_{60}]$$

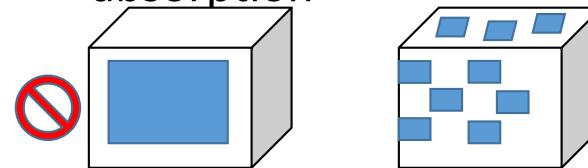


Sound propagation

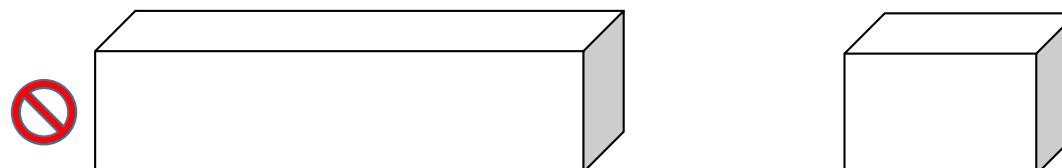
Ergodicity



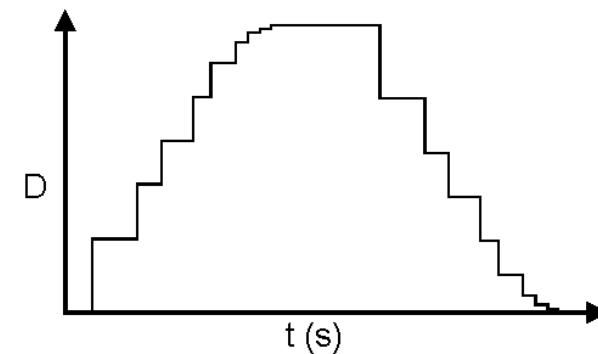
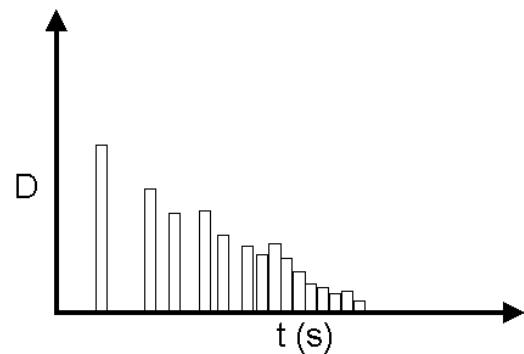
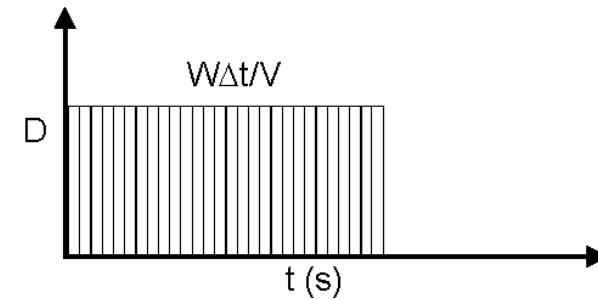
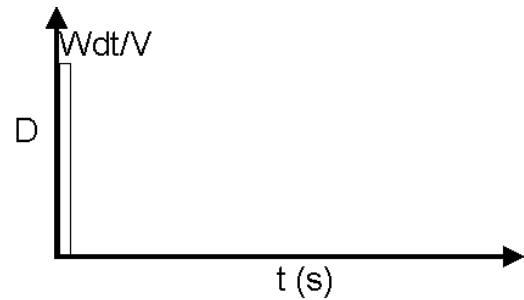
Mixing and distributed absorption



Proportionate



Sound propagation



Sound propagation

$$w_D = \frac{W}{4\pi r^2 c} \quad w_R = \frac{4W}{cR} \quad R = \frac{\bar{\alpha}S}{1 - \bar{\alpha}} \approx A$$

$$w = w_D + w_R = \frac{W}{c} \left(\frac{1}{4\pi r^2} + \frac{4}{R} \right)$$

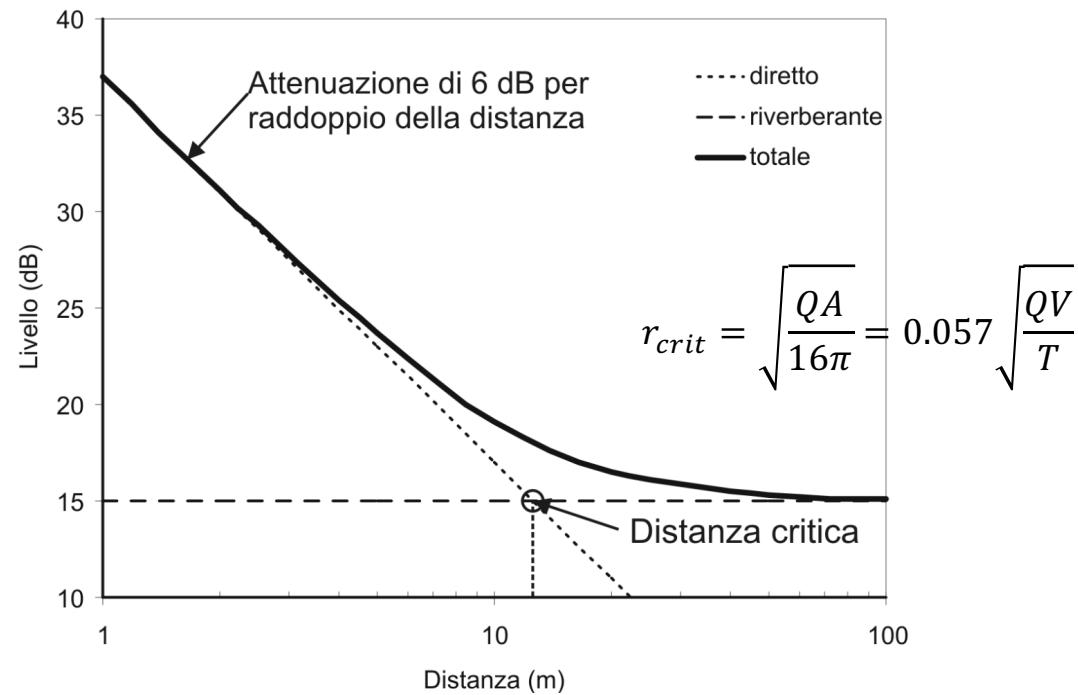


$$L_p = L_w + 10 \log \left(\frac{1}{4\pi r^2} + \frac{4}{A} \right)$$



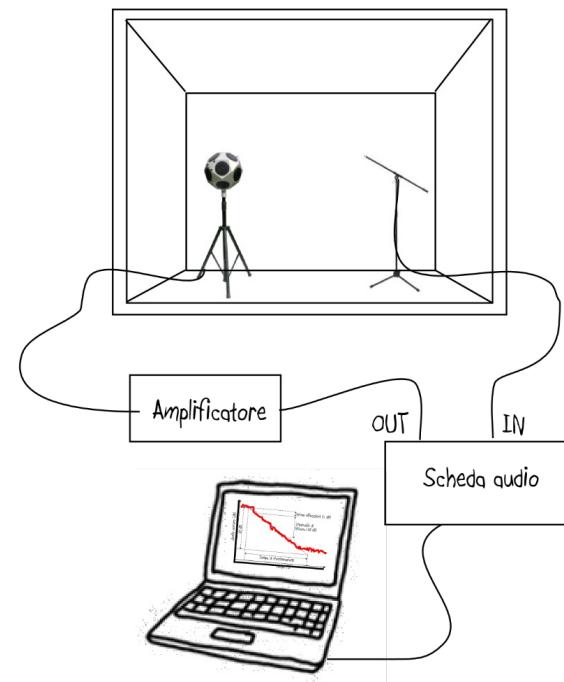
$$L_p = L_w + 10 \log \left(\frac{Q}{4\pi r^2} + \frac{4}{A} \right)$$

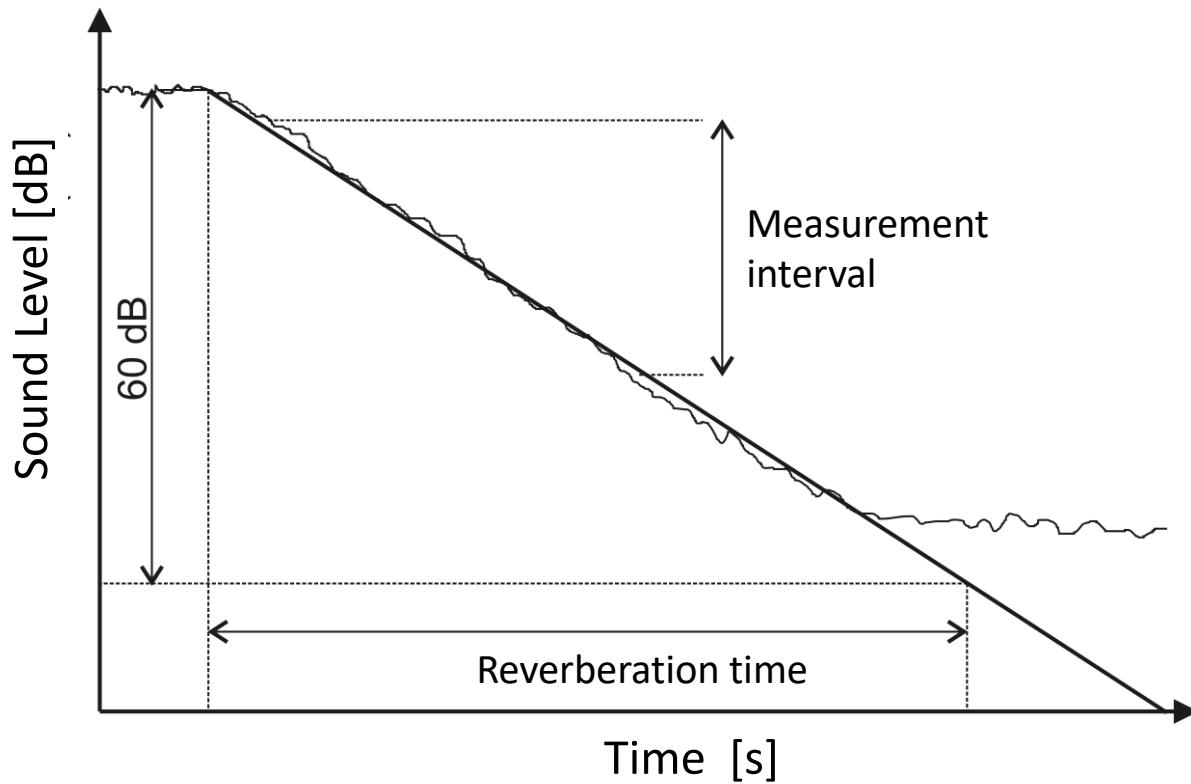
Sound propagation



How to measure reverberation time (and impulse responses)

- Measurement techniques
 - Noise emission (standing)
 - Impulsive source
 - Electro acoustic excitation (deconvolution)
- ISO 3382-1:2009

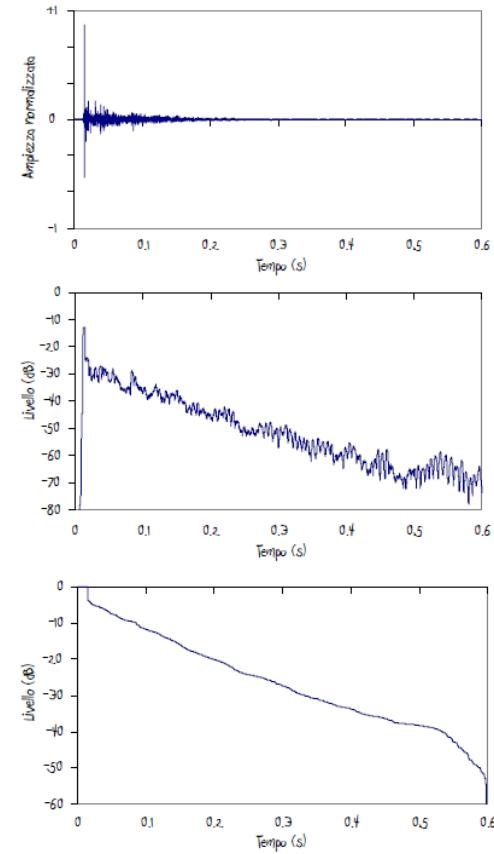




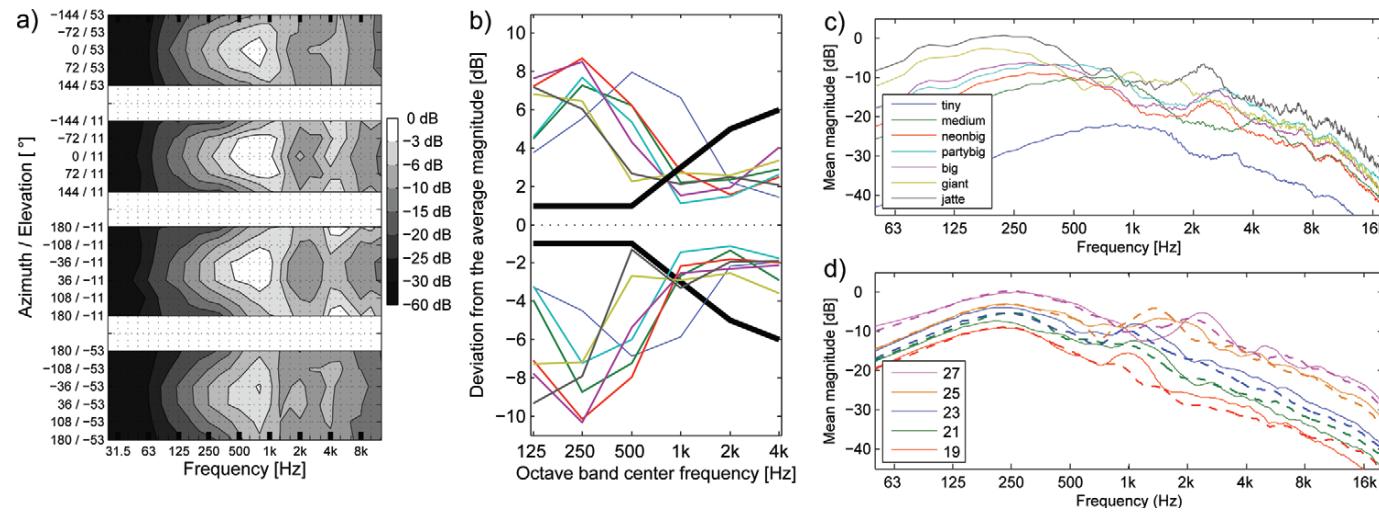
Backward integration

$$h^2(t) = Ae^{-kt} \longrightarrow \int Ae^{-kt} dt = -\frac{A}{k} e^{-kt}$$

$$\int_t^\infty Ae^{-kt} dt = -\frac{A}{k} e^{-k \cdot \infty} + \frac{A}{k} e^{-kt}$$



Balloons



The clapper

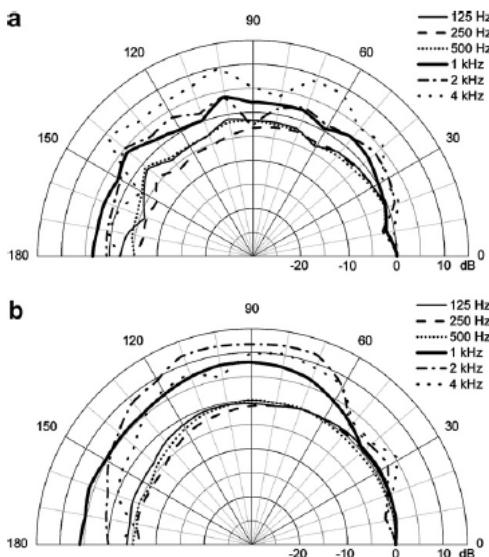
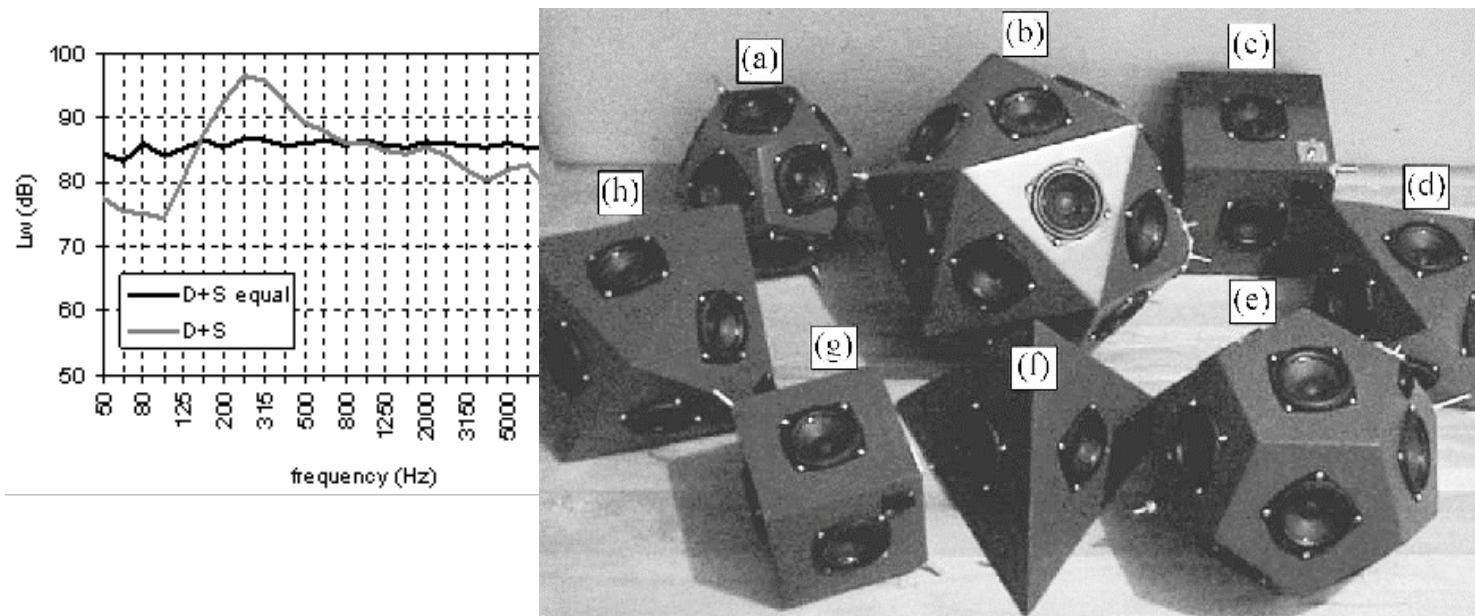


Fig. 5. Directivity of the clapper – variations in the maximum reached signal value: (a) frontal plane; and (b) medial plane.

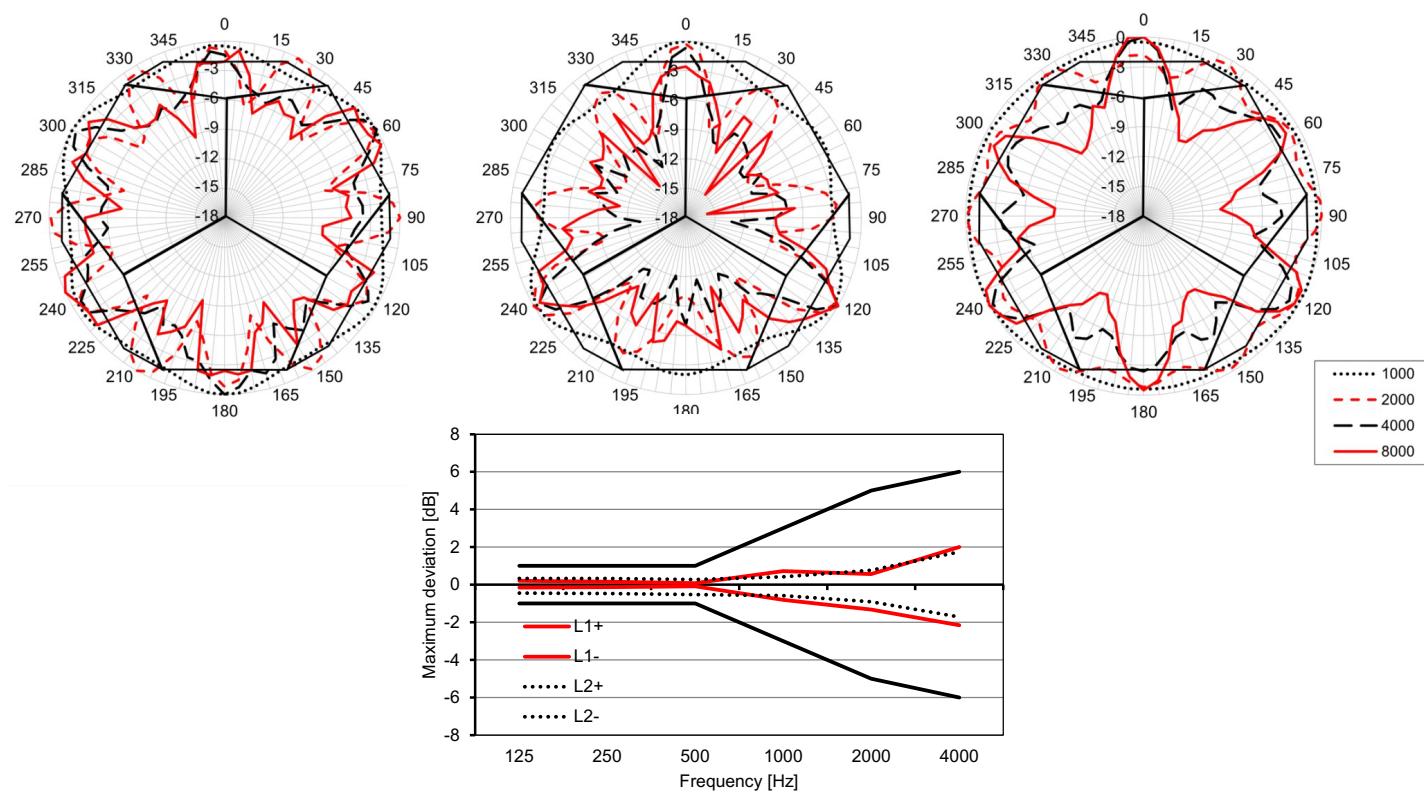
Impulsive sound source

- Pros
 - Simple procedure and lightweight sources
 - Yields a usable impulse response (with limitations)
 - Smooth decay (thanks to backward integration)
- Cons
 - Poor repeatability
 - Poor control of radiated spectrum and source directivity
 - Signal-to-noise ratio depending on the source power

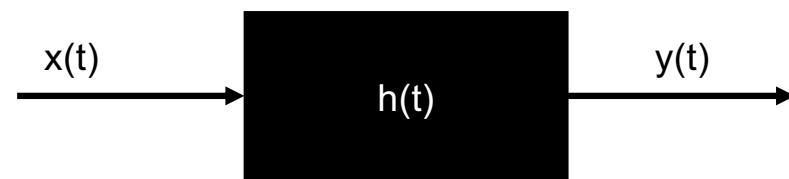
Electro acoustic source



Electro acoustic source



Convolution/Deconvolution

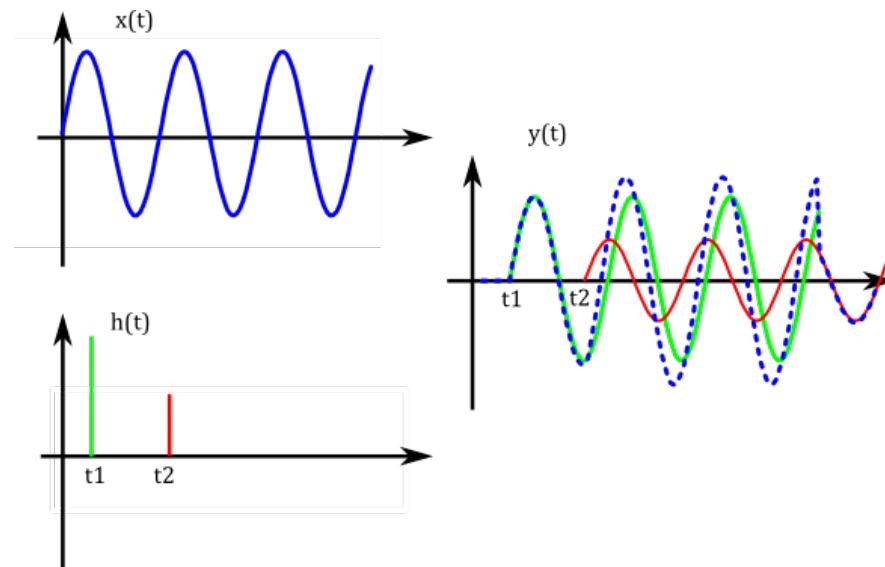


$$y(t) = x(t) \otimes h(t) = \int_0^{\infty} h(\tau) \cdot x(t - \tau) d\tau$$


A blue downward-pointing arrow labeled "FFT" is positioned between the convolution integral and the final equation.

$$Y(f) = X(f) \cdot H(f)$$

Convolution



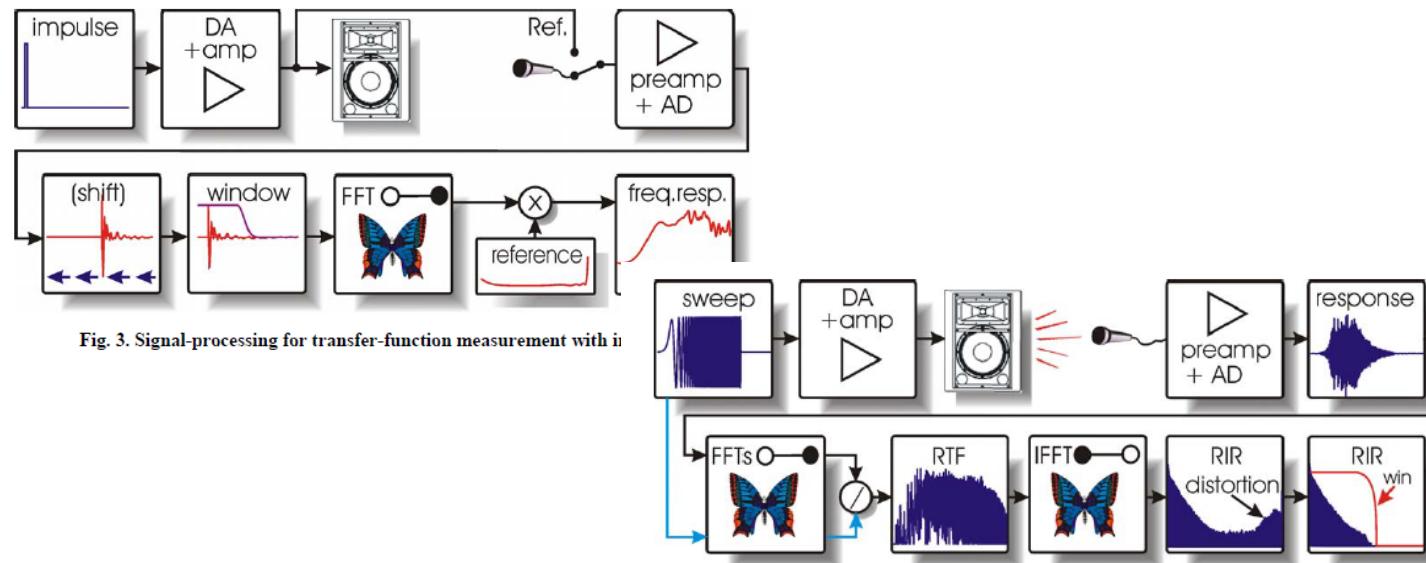
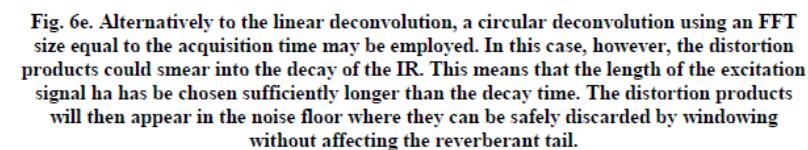


Fig. 3. Signal-processing for transfer-function measurement with i



The diagram illustrates an alternative signal processing method for transfer-function measurement using circular deconvolution. It shows a feedback loop where an 'excitation signal' (labeled 'sweep') is converted to digital by a 'DA + amp' and sent to a speaker. A microphone picks up the sound and feeds into a 'preamp + AD' block, producing a 'response' plot. The 'excitation signal' is also processed through an 'FFT O' (Fast Fourier Transform Out) block, followed by a circular convolution operation (\circlearrowright). The result is an 'RTF' (Reverberation Transfer Function) plot. This is then converted back to the time domain using an 'IFFT' (Inverse Fast Fourier Transform Out) block. Finally, the signal is processed through an 'RIR distortion' block and a windowing operation (RIR^{win}), resulting in a 'RIR' (Reverberation Impulse Response) plot.

Fig. 6e. Alternatively to the linear deconvolution, a circular deconvolution using an FFT size equal to the acquisition time may be employed. In this case, however, the distortion products could smear into the decay of the IR. This means that the length of the excitation signal ha has to be chosen sufficiently longer than the decay time. The distortion products will then appear in the noise floor where they can be safely discarded by windowing without affecting the reverberant tail.

Sine sweeps (or chirps)

- Pros
 - Reject distortion;
 - High signal-to-noise ratio;
 - Immune to noise (more or less);
 - Possibility to equalize the source;
- Cons
 - Needs a LTI (linear-time-invariant) system