

### Exercises 3: Adiabatic quantum computation

Lecturer: Simon Apers (*apers@irif.fr*)

**Exercise 1** (Pauli basis). Recall the unitary Pauli matrices

$$I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

- Show that the Pauli basis  $\{I, X, Y, Z\}$  forms a basis for the complex 2-by-2 matrices. I.e., any  $A \in \mathbb{C}^{2 \times 2}$  can be expanded as  $A = \alpha_1 I + \alpha_x X + \alpha_y Y + \alpha_z Z$ .
- Argue that this implies that the  $n$ -qubit Pauli basis  $\{I, X, Y, Z\}^{\otimes n}$  forms a basis for the  $2^n$ -by- $2^n$  matrices.
- Show that if  $A \in \mathbb{C}^{2^n \times 2^n}$  is Hermitian, then its coefficients in the Pauli basis are real.

**Exercise 2** (Max cut). Finding the maximum cut in a graph is a canonical problem in combinatorial optimization. For a graph  $G$  with vertex set  $[n]$  and (symmetric) edge set  $E \subseteq [n]^2$ , a maximum cut is described by a subset  $Z \subset [n]$  that cuts a maximum number of edges (edges crossing from  $Z$  to  $Z^c$ ). Equivalently, it maximizes the cut function

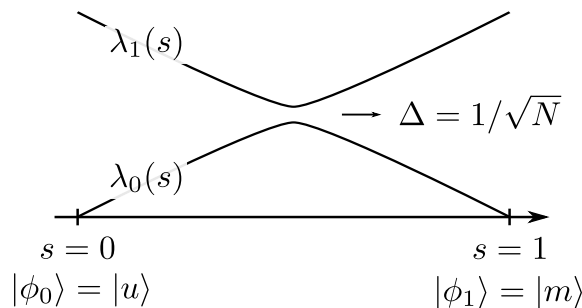
$$c(A) = \sum_{(i,j) \in E} I_{i \in A, j \notin A}.$$

- Identify a subset  $A$  with the indicator  $a \in \{0, 1\}^n$  ( $i \in A \Leftrightarrow a_i = 1$ ). Express the cut function  $c(A)$  as a degree-2 polynomial  $h$  in  $a$ .
- Rewrite the cost Hamiltonian  $H_1 = -\sum_{z \in \{0,1\}^n} c(z) |z\rangle \langle z|$  in terms of identity and Pauli- $Z$  matrices.

**Exercise 3** (Adiabatic Grover algorithm and avoided crossing). Consider the unstructured search problem over the set  $[N] = \{0, 1\}^n$ , and assume that there is a single (unknown) marked element  $m \in \{0, 1\}^n$ . We can solve this problem using the adiabatic optimization algorithm with cost function  $h(z) = 0$  if  $z = m$  and  $h(z) = 1$  otherwise. Letting  $|u\rangle = \frac{1}{\sqrt{N}} \sum_{x \in \{0,1\}^n} |x\rangle$ , we can use the Hamiltonians

$$H_0 = I - |u\rangle \langle u| \quad \text{and} \quad H_1 = I - |m\rangle \langle m|,$$

with ground states  $|\phi_0\rangle = |u\rangle$  and  $|\phi_1\rangle = |m\rangle$ , respectively. Similar to Grover, the Hamiltonian  $H(s) = (1-s)H_0 + sH_1$  only acts nontrivially in the subspace spanned by  $|u_0\rangle = \frac{1}{\sqrt{N-1}} \sum_{x \neq m} |x\rangle$  and  $|m\rangle$ . Calculate the nontrivial eigenvalues of  $H(s)$ , and show that they behave as in the following figure.<sup>1</sup>



<sup>1</sup>While this proves that the gap is inverse polynomial, and so the adiabatic algorithm solves unstructured search in  $\text{poly}(n)$ , it doesn't directly recover the quadratic Grover speedup. A more refined adiabatic algorithm does (Roland-Cerf, "Quantum search by local adiabatic evolution", Physical Review A, 2002).