

The distinctions of symmetry.

In this section we investigate the effect of a general QFT action functional. $S[\phi, \partial_\mu \phi]$.

The transformations acts on the position as well as the fields

$$\begin{aligned}x &\rightarrow x' \\ \phi(x) &\rightarrow \phi'(x')\end{aligned}$$

The post-transform field is then a function of the previous field. *Note that in many QFTs it is common to work in momentum space. This is not the case for us, and we will treat almost every situation in position space.*

This perspective on transformations is often referred to as active transformations, as opposed to passive transformations; where transformations are treated as coordinate transformations.

The action (3) transforms as the following under the general transform.

$$\begin{aligned}S &\rightarrow S' = \int d^d x \mathcal{L}(\phi'(x), \partial_\mu \phi'(x)) && \left. \begin{array}{l} \text{Change} \\ \text{of integration} \\ \text{variable.} \end{array} \right\} \\ &= \int d^d x' \mathcal{L}(\phi'(x'), \partial'_\mu \phi'(x')) \\ &= \int d^d x' \mathcal{L}(\phi(\phi(x)), \partial'_\mu \phi(\phi(x))) && \left. \begin{array}{l} \text{Jacobian} \end{array} \right\} \\ &= \int \left| \frac{\partial x'}{\partial x} \right| d^d x \mathcal{L}(\phi(\phi(x)), \frac{\partial x^\nu}{\partial x'^\mu} \frac{\partial}{\partial x^\nu} \phi(\phi(x)))\end{aligned}$$

Ex. 1 Translation; $\mathcal{Q}U$

$$x' = x + a$$

$$\phi'(x+a) = \phi(x)$$

$$S = \int \left| \frac{\partial x'}{\partial x} \right| d^d x \mathcal{L}(\phi(\phi(x)), \frac{\partial x^\nu}{\partial x'^\mu} \frac{\partial}{\partial x^\nu} \phi(\phi(x)))$$

$$\left| \frac{\partial x'}{\partial x} \right| = 1, \quad \frac{\partial x^\nu}{\partial x'^\mu} = \delta_\mu^\nu$$

$$S = \int d^d x \mathcal{L}(\phi(\phi(x)), \partial_\nu \phi(\phi(x)))$$

Remark. 1 The action is invariant if not explicitly dependent on pos.

Ex. 2 Lorentz transformation

$$x'^{\mu} = \Lambda^{\mu}_{\nu} x^{\nu}, \quad \bar{\Phi}'(\Lambda x) = \bar{\Gamma}_{\Lambda} \Phi(x)$$

$$S = \int \left| \frac{\partial x'}{\partial x} \right| d^d x \mathcal{L}(\phi(\phi(x)), \frac{\partial x^{\nu}}{\partial x'^{\mu}} \frac{\partial}{\partial x^{\nu}} \phi(\phi(x)))$$

$$\left| \frac{\partial x'}{\partial x} \right| = \det \left(\frac{\partial x'^{\mu}}{\partial x^{\nu}} \right) = \det (\Lambda^{\mu}_{\rho} \delta^{\rho}_{\nu}) \\ = \det (\Lambda^{\mu}_{\nu}) ?$$

By unitarity

Make note of confusion
Lorentz group, boost +
rotations. Rotations
are obviously unitary
but are boosts?

Ex. 3 Scale transform

We consider the transformation caused by a set of infinitesimal parameters $\{\omega_a\}$, specifically, the transformations we are referring to are

$$x^\mu \rightarrow x'^\mu = x^\mu + \omega_a \frac{\delta x^\mu}{\delta \omega_a}$$

$$\phi'(x') = \phi(x) + \omega_a \frac{\delta F}{\delta \omega_a}(x) \quad *$$

It is common to define generators (G_a) caused by a symmetry transform as follows

$$\frac{\delta \phi(x)}{\delta \omega} = \delta \omega \phi(x) \stackrel{\text{Note definition}}{=} \phi'(x) - \phi(x) \stackrel{\text{Note definition}}{=} -i \omega_a G_a \phi(x)$$

We then note that we can expand the field $\phi(x)$ around the point x' in terms of $\{\omega_a\}$

$$\phi(x) = \phi(x') - \omega_a \left(\frac{\delta x}{\delta \omega_a} \right) \left(\frac{\partial \phi(x')}{\partial x'} \right)$$

combined with * yields

$$\phi'(x') = \phi(x') - \omega_a \frac{\delta x}{\delta \omega_a} \frac{\partial \phi(x')}{\partial x'} + \omega_a \frac{\delta F}{\delta \omega_a}(x)$$

$$\phi'(x') = \phi(x') - \omega_a \frac{\delta x}{\delta \omega_a} \frac{\partial \phi(x')}{\partial x'} + \omega_a \frac{\delta F}{\delta \omega_a}(x' - \omega_a \frac{\delta x}{\delta \omega_a})$$

$$\phi'(x') = \phi(x') - \omega_a \frac{\delta x}{\delta \omega_a} \frac{\partial \phi(x')}{\partial x'} + \omega_a \frac{\delta}{\delta \omega_a} (F(x') - \omega_a \dots)$$

$$\phi'(x') = \phi(x') - \omega_a \frac{\delta x}{\delta \omega_a} \frac{\partial \phi(x')}{\partial x'} + \omega_a \frac{\delta}{\delta \omega_a} F(x')$$

$$\phi'(x') - \phi(x') = -\omega_a \left(\frac{\delta x}{\delta \omega_a} \frac{\partial \phi(x')}{\partial x'} + \frac{\delta F(x')}{\delta \omega_a} \right)$$

$$\text{Eq. b1a} \\ = -i \omega_a G_a \phi(x)$$

$$\Rightarrow i G_a \phi(x) = \frac{\delta x}{\delta \omega_a} \frac{\partial \phi(x')}{\partial x'} + \frac{\delta F(x')}{\delta \omega_a}$$

Using eq. b1a for said transform yields

$$\delta' = \int d^d x \left(1 + \partial_\mu \left(\omega_a \frac{\delta x^\mu}{\delta \omega_a} \right) \right) \times \mathcal{L} \left(\phi + \omega_a \frac{\delta F}{\delta \omega_a}, \left[\bar{\psi}^\nu - \partial_\mu \left(\omega_a \frac{\delta x^\nu}{\delta \omega_a} \right) \right] (\partial_\nu \phi + \partial_\nu \left[\omega_a \left(\frac{\delta F}{\delta \omega_a} \right) \right]) \right)$$

Where, in order to calculate the Jacobian we have used the following approximation

$$\det(1+E) \approx 1 + \text{Tr } E \quad (\text{small } E)$$

To calculate $\delta \delta$ we will need to perform an ^{first order} expansion δ' around x w.r.t. ω_a .

For this reason, it is useful to remind ourselves about the formula for multivariate Taylor expansion

$$f(a+h, b+k) = f(a, b) + \frac{\partial}{\partial x} f(a, b) h + \frac{\partial}{\partial y} f(a, b) k + \mathcal{O}(h, k) \mathcal{B}(h, k)$$

where $\mathcal{B}(h, k)$ is a compact function close to the origin.

$$\delta \delta = \int d^d x \partial_\mu \gamma^\mu_a \omega_a$$

$$\gamma^\mu_a = \left\{ \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \partial_\nu \phi - \bar{\psi}^\mu \mathcal{L} \right\} \frac{\delta x^\nu}{\delta \omega_a} - \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \frac{\delta F}{\delta \omega_a}$$

The energy-momentum tensor is the conserved current that corresponds with translational invariance.

$$\frac{\delta x^\mu}{\delta \epsilon^\nu} =$$

