1 Sterile and Massive Neutrinos

As we saw in Eq. 3 , the neutrino fields ν_{α} only couple to the associated lepton fields ℓ_{α} , conserving the lepton number L_{α} . We will now introduce two separate extensions to this part of the Standard Model.

We introduce a right-handed neutrino field, ν_R . It has the usual properties of the conventional left-handed neutrino such as hypercharge and color zero. Moreover, since the electroweak gauge group $SU(2)_L$ only couple to left-handed particles and right-handed antiparticles, it transforms as a singlet under the SM symmetry group $SU(3)_C \times SU(2)_L \times U(1)_Y$. This neutrino is *sterile* since it doesn't participate any of the SM interactions.

We extend the SM by adding a right-handed component to the Higgs-lepton Yukawa Lagrangian from Eq. 1 with neutrino Yukawa couplings $Y^{\prime\nu}_{\alpha\beta}$,

$$\mathcal{L}_{H} = -\left(\frac{v+H}{\sqrt{2}}\right) \left[\ell'_{\alpha L} Y'^{\ell}_{\alpha \beta} \ell'_{\beta R} + \nu'_{\alpha L} Y'^{\nu}_{\alpha \beta} \nu'_{\beta R}\right] \quad (1)$$

Similar to how we diagonalized the lepton Yukawa couplings $Y_{\alpha\beta}^{\prime\ell}$ in Eq. 2, we diagonalize $Y_{\alpha\beta}^{\prime\nu}$ as

$$V_{\alpha kL}^{\nu \dagger} Y_{\alpha \beta}^{\prime \nu} V_{\beta jR}^{\nu} = Y_{kj}^{\nu} \,. \tag{2}$$

Now, let the neutrino field with chriality X be denoted n_X , with components Latin numerals to distinguish them from the flavour components, i.e

$$\nu_{kX} = V_{kjX}^{\nu\dagger} \nu_{jX}' \,. \tag{3}$$

The diagonalized Lagrangian now takes the form

$$\mathcal{L}_{H} = -\left(\frac{v+H}{\sqrt{2}}\right) \left[\ell_{\alpha L}^{\prime} Y_{\alpha \beta}^{\prime \ell} \ell_{\beta R}^{\prime} + \nu_{\alpha L}^{\prime} Y_{\alpha \beta}^{\prime \nu} \nu_{\beta R}^{\prime}\right]
= -\left(\frac{v+H}{\sqrt{2}}\right) \left[\ell_{\alpha L}^{\prime} V_{\alpha \beta L}^{\ell} Y_{\alpha \beta}^{\ell} V_{\alpha \beta R}^{\ell \dagger} \ell_{\beta R}^{\prime} \right]
+ \nu_{\alpha L}^{\prime} V_{\alpha k L}^{\nu} Y_{k j}^{\nu} V_{\beta j R}^{\nu \dagger} \nu_{\beta R}^{\prime}
= -\left(\frac{v+H}{\sqrt{2}}\right) \left[\ell_{\alpha L}^{\dagger} Y_{\alpha \beta}^{\ell} \ell_{\beta R} + \nu_{k L}^{\dagger} Y_{k j}^{\nu} \nu_{j R}\right]
= -\left(\frac{v+H}{\sqrt{2}}\right) \left[\bar{\ell}_{\alpha L} Y_{\alpha \beta}^{\ell} \ell_{\beta R} + \bar{\nu}_{k L} Y_{k j}^{\nu} \nu_{j R}\right]$$
(5)

Now using the fact that Y_{kj}^{ν} is diagonal, we write it as $Y_{kj}^{\nu} = y_k^{\nu} \delta_{kj}$, leaving the Lagrangian as

$$\mathcal{L}_{H} = -\left(\frac{v+H}{\sqrt{2}}\right) \left[\bar{\ell}_{\alpha L} y_{\alpha}^{\ell} \delta_{\alpha \beta} \ell_{\beta R} + \bar{\nu}_{k L} y_{k}^{\nu} \delta_{k j} \nu_{j R}\right]
= -\left(\frac{v+H}{\sqrt{2}}\right) \left[\bar{\ell}_{\alpha L} y_{\alpha}^{\ell} \ell_{\alpha R} + \bar{\nu}_{k L} y_{k}^{\nu} \nu_{k R}\right]
= -\left(\frac{v+H}{\sqrt{2}}\right) \left[y_{\alpha}^{\ell} \bar{\ell}_{\alpha L} \ell_{\alpha R} + y_{k}^{\nu} \bar{\nu}_{k L} \nu_{k R}\right]$$
(6)

Now, the Dirac neutrino field is

$$\nu_k = \nu_{kL} + \nu_{kR} \,. \tag{7}$$

Multiplying ν_k with its conjugate $\bar{\nu}_k$, we get

$$\bar{\nu}_k \nu_k = \bar{\nu}_{kL} \nu_{kL} + \bar{\nu}_{kR} \nu_{kL} + \bar{\nu}_{kL} \nu_{kR} + \bar{\nu}_{kR} \nu_{kR}$$

$$= \bar{\nu}_{kL} \nu_{kR} + \bar{\nu}_{kR} \nu_{kL}$$

$$= \bar{\nu}_{kL} \nu_{kR} + \text{h.c.}$$
(8)

The same calculation for the charged lepton field yields the same result for ℓ_k . Substituting this result and expanding the Higgs VEV into the fields gives us

$$\mathcal{L}_{H} = -\left(\frac{v+H}{\sqrt{2}}\right) \left[y_{\alpha}^{\ell} \bar{\ell}_{\alpha} \ell_{\alpha} + y_{k}^{\nu} \bar{\nu}_{k} \nu_{k}\right]$$

$$= -\frac{y_{\alpha}^{\ell} v}{\sqrt{2}} \bar{\ell}_{\alpha} \ell_{\alpha} - \frac{y_{k}^{\nu} v}{\sqrt{2}} \bar{\nu}_{k} \nu_{k} - \frac{y_{\alpha}^{\ell}}{\sqrt{2}} \bar{\ell}_{\alpha} \ell_{\alpha} H - \frac{y_{k}^{\nu}}{\sqrt{2}} \bar{\nu}_{k} \nu_{k} H.$$
(9)

Thus, this extension to the SM generates neutrino masses by the Higgs mechanism, in the same fashion as with the charged leptons and the quarks:

$$m_k = \frac{y_k^{\nu} v}{\sqrt{2}} \tag{10}$$