

考虑一维双原子链(1D diatomic chain):

- 1) 计算长波极限 $(q \to 0)$ 下声学支和光学支格波的色散关系 $\omega \sim q$ ;
- 2) 分析第一布里渊区边界处的振动特点;
- 3) 当m = M时, 画出第一布里渊区内的色散关系 $\omega \sim q$ , 并与一维单原子链 (1D monoatomic chain) 的情形进行比较.

提交时间: 3月17日之前

提交方式: 手写 (写明姓名学号) 后拍照, 通过本班课代表统一提交电子版



考虑一维双原子链(1D diatomic chain):

1) 计算长波极限 $(q \to 0)$ 下声学支和光学支格波的色散关系 $\omega \sim q$ ;

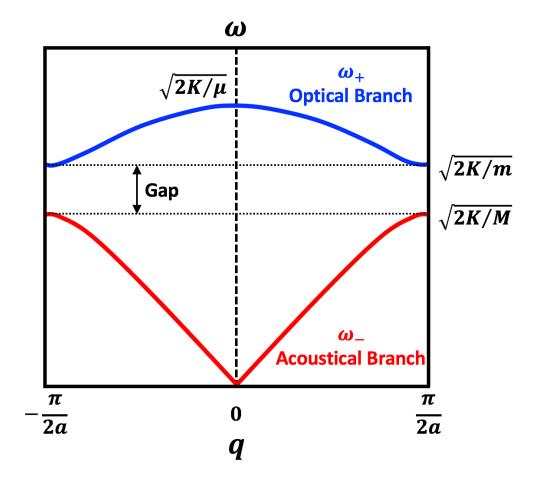
$$\omega^{2} = \begin{cases} \omega_{+}^{2} = \frac{K}{\mu} \left[ 1 + \sqrt{1 - \frac{4\mu^{2}}{mM}} \sin^{2}(aq) \right] & q \to 0 \\ \omega^{2} = \begin{cases} \omega_{+}^{2} = \frac{K}{\mu} \left[ 1 - \sqrt{1 - \frac{4\mu^{2}}{mM}} \sin^{2}(aq) \right] & q \to 0 \end{cases} & \omega_{+}^{2} \to 0 \end{cases}$$

注意对比一维单原子链: 
$$\omega = 2\sqrt{\frac{K}{m}} \left| \sin\left(\frac{1}{2}aq\right) \right| \xrightarrow{q \to 0} \omega \to a\sqrt{\frac{K}{m}} |q|$$



考虑一维双原子链(1D diatomic chain):

1) 计算长波极限 $(q \rightarrow 0)$ 下声学支和光学支格波的色散关系 $\omega \sim q$ ;





考虑一维双原子链(1D diatomic chain):

2) 分析第一布里渊区边界处的振动特点;

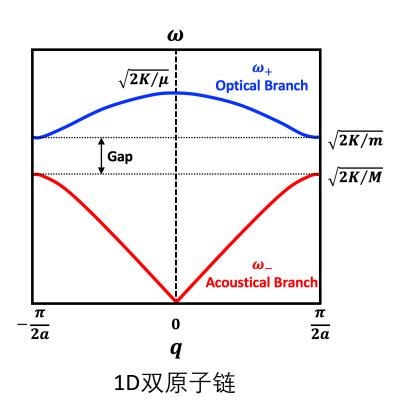
$$q = \pm \frac{\pi}{2a} \qquad \qquad \omega_+^2 \to \frac{2K}{m} \qquad \omega_-^2 \to \frac{2K}{M}$$

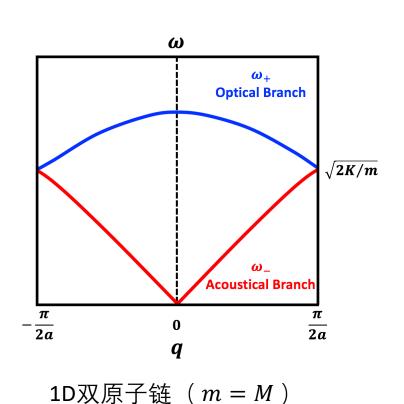


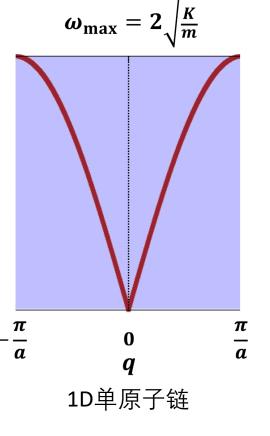
考虑一维双原子链(1D diatomic chain):

3) 当m = M时,画出第一布里渊区内的色散关系 $\omega \sim q$ ,并与一维单原子链

(1D monoatomic chain)的情形进行比较.







## Chapter 3.2: 课后作业



考虑一维量子谐振子 
$$\hat{H} = \frac{1}{2m}\hat{p}^2 + \frac{1}{2}m\omega^2\hat{x}^2$$
,

证明:

1) 
$$[\hat{a}, \hat{a}^+] = 1$$

2) 
$$[\hat{a}^{+}\hat{a}, \hat{a}^{+}] = \hat{a}^{+}$$

3) 
$$[\hat{a}^+\hat{a}, \hat{a}] = -\hat{a}$$

4) 
$$\widehat{H} = \hbar\omega \left(\widehat{a}^{+}\widehat{a} + \frac{1}{2}\right)$$

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证明要点: 利用
$$[\hat{x}, \hat{p}] = i\hbar$$

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$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{x} + \frac{i}{m\omega} \hat{p} \right) \qquad \hat{a}^{+} = \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{a}^{+} = \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

证明: 1) 
$$[\hat{a}, \hat{a}^+] = 1$$

$$[\hat{a}, \hat{a}^+] = \hat{a} \ \hat{a}^+ - \hat{a}^+ \hat{a}$$

$$=\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}+\frac{i}{m\omega}\hat{p}\Big)\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}-\frac{i}{m\omega}\hat{p}\Big)-\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}-\frac{i}{m\omega}\hat{p}\Big)\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}+\frac{i}{m\omega}\hat{p}\Big)$$

$$= -\frac{i}{\hbar}(\hat{x}\hat{p} - \hat{p}\hat{x})$$

$$= 1$$



证明要点: 利用[
$$\hat{a}$$
,  $\hat{a}^+$ ] = 1 即  $\hat{a}$   $\hat{a}^+$  —  $\hat{a}^+$   $\hat{a}$  = 1

证明: 2) 
$$[\hat{a}^+\hat{a}, \hat{a}^+] = \hat{a}^+$$

$$[\hat{a}^{+}\hat{a}, \hat{a}^{+}] = \hat{a}^{+}\hat{a}\hat{a}^{+} - \hat{a}^{+}\hat{a}^{+}\hat{a} = \hat{a}^{+}(1 + \hat{a}^{+}\hat{a}) - \hat{a}^{+}\hat{a}^{+}\hat{a} = \hat{a}^{+}$$

证明: 3) 
$$[\hat{a}^+\hat{a},\hat{a}] = -\hat{a}$$

$$[\hat{a}^{\dagger}\hat{a},\hat{a}] = \hat{a}^{\dagger}\hat{a}\hat{a} - \hat{a}\hat{a}^{\dagger}\hat{a} = (\hat{a}\hat{a}^{\dagger} - 1)\hat{a} - \hat{a}\hat{a}^{\dagger}\hat{a} = -\hat{a}$$



证明要点: 利用 
$$\hat{x} = \sqrt{\frac{\hbar}{2} \frac{1}{m\omega}} (\hat{a}^+ + \hat{a})$$
  $\hat{p} = i \sqrt{\frac{\hbar}{2}} m\omega (\hat{a}^+ - \hat{a})$   $[\hat{a}, \hat{a}^+] = 1$ 

证明: 4) 
$$\widehat{H} = \hbar\omega\left(\widehat{a}^{\dagger}\widehat{a} + \frac{1}{2}\right)$$

$$\widehat{H} = \frac{1}{2m}\widehat{p}^2 + \frac{1}{2}m\omega^2\widehat{x}^2 = -\frac{\hbar\omega}{4}(\widehat{a}^+ - \widehat{a})(\widehat{a}^+ - \widehat{a}) + \frac{\hbar\omega}{4}(\widehat{a}^+ + \widehat{a})(\widehat{a}^+ + \widehat{a})$$

$$=\frac{\hbar\omega}{2}\left(\hat{a}^{+}\hat{a}+\hat{a}\hat{a}^{+}\right)$$

$$=\frac{\hbar\omega}{2}\left(\hat{a}^{+}\hat{a}+\hat{a}^{+}\hat{a}+1\right)$$

$$=\hbar\omega\left(\hat{a}^{+}\hat{a}+\frac{1}{2}\right)$$



设晶体中每个振动模的零点振动能为 $\frac{1}{2}\hbar\omega$ ,使用德拜模型和爱因斯坦模型分别求晶体的零点振动能。

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#### 解题要点: 充分理解态密度 $g(\omega)$ 的物理意思,并利用态密度将"求和"化为"积分"

解:

$$E = \sum_{j=1}^{3N} \frac{1}{2} \hbar \omega_{j} = \int_{0}^{\omega} \frac{1}{2} \hbar \omega g(\omega) d\omega$$
 设N为原胞个数,每个原胞1个原子

$$g(\omega) = \frac{3V}{2\pi^2 \overline{c}^3} \omega^2$$

德拜模型: 
$$g(\omega) = \frac{3V}{2\pi^2 \overline{c}^3} \omega^2 \qquad \omega_{\rm D} = \overline{c} \left[ 6\pi^2 \left( \frac{N}{V} \right) \right]^{1/3}$$

$$E_{\text{Debye}} = \int_{0}^{\omega_{\text{D}}} \frac{1}{2} \hbar \omega g(\omega) d\omega = \frac{9}{8} N \hbar \omega_{\text{D}}$$



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解:

$$E = \sum_{j=1}^{3N} \frac{1}{2} \hbar \omega_{j} = \int_{0}^{\omega} \frac{1}{2} \hbar \omega g(\omega) d\omega$$

设N为原胞个数,每个原胞1个原子

爱因斯坦模型: 
$$g(\omega) = 3N\delta(\omega - \omega_{\rm E})$$

$$\int_{-\infty}^{+\infty} f(x)\delta(x - x_0) = f(x_0)$$

$$E_{\text{Einstein}} = \int_{0}^{+\infty} \frac{1}{2} \hbar \omega g(\omega) d\omega = \frac{3}{2} N \hbar \omega_{\text{E}}$$

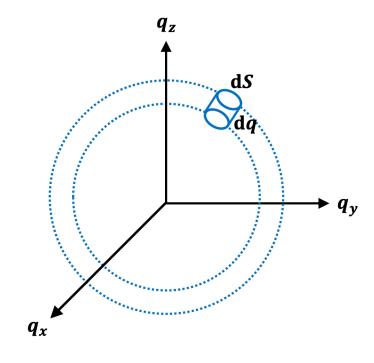


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德拜模型态密度的推导:  $g(\omega) = \frac{3V}{2\pi^2 \overline{c}^3} \omega^2$ 

德拜模型考虑线性色散:  $\omega = cq$ 

态密度可表示为:



$$g(\omega) = \frac{V}{(2\pi)^3} \oiint \frac{dS}{|\nabla_a \omega_a|} = \frac{V}{(2\pi)^3} \frac{4\pi q^2}{|\nabla_a \omega_a|} = \frac{V}{(2\pi)^3} \frac{4\pi \omega^2}{c^3} = \frac{V}{2\pi^2} \frac{\omega^2}{c^3}$$



#### 解题要点: 充分理解态密度 $g(\omega)$ 的物理意思,并利用态密度将"求和"化为"积分"

三维情况下: 1支纵波(声速 $c_L$ )和2支横波(声速 $c_T$ )

对于1支纵波,态密度: 
$$g_{\rm L}(\omega) = \frac{V}{2\pi^2} \frac{\omega^2}{c_{\rm L}^3}$$

对于2支横波,态密度: 
$$g_{\mathrm{T}}(\omega) = \frac{V}{2\pi^2} \left( \frac{\omega^2}{c_{\mathrm{T}}^3} + \frac{\omega^2}{c_{\mathrm{T}}^3} \right) = \frac{\omega^2}{c_{\mathrm{T}}^3} \frac{2\omega^2}{c_{\mathrm{T}}^3}$$

总态密度: 
$$g(\omega) = g_{L}(\omega) + g_{T}(\omega) = \frac{V}{2\pi^{2}} \left( \frac{\omega^{2}}{c_{L}^{3}} + \frac{2\omega^{2}}{c_{T}^{3}} \right) = \frac{3V}{2\pi^{2}\overline{c}^{3}} \omega^{2}$$

等效声速: 
$$\frac{1}{\overline{c}^3} = \frac{1}{3} \left( \frac{1}{c_L^3} + \frac{2}{c_T^3} \right)$$