



Appendix 3

Quantum Harmonic Oscillator (量子谐振子)

- To learn the **quantum description** of harmonic oscillator.
- To understand the properties of **quantum harmonic oscillator**.

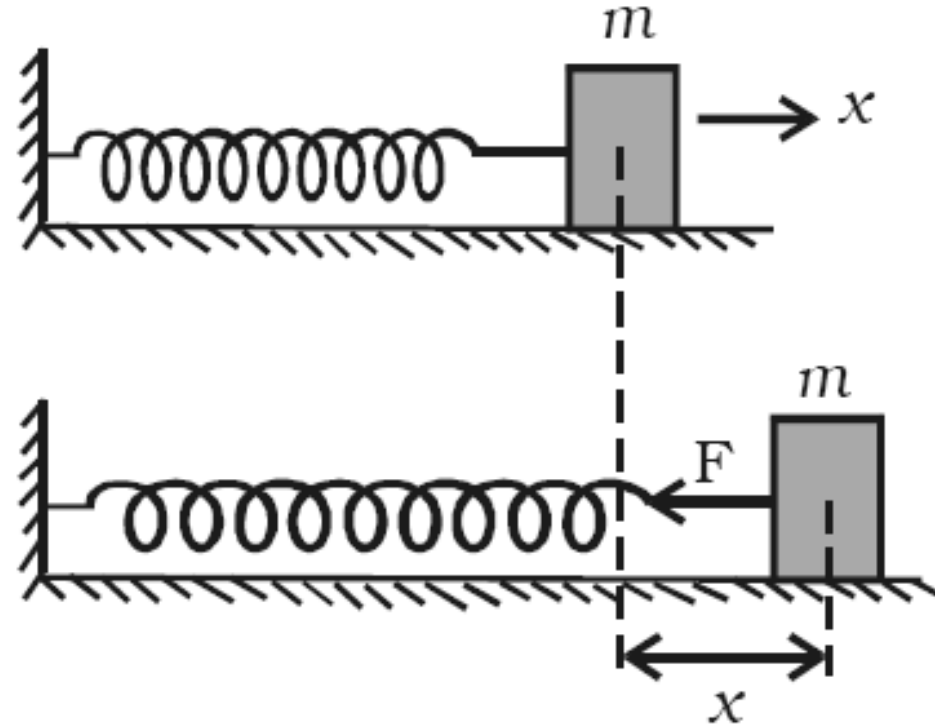
Appendix 3: Quantum Harmonic Oscillator (量子谐振子)



➤ Simple Harmonic Oscillator (简谐振子)

❖ Hook's law (胡克定律):

$$F = -Kx$$



Appendix 3: Quantum Harmonic Oscillator (量子谐振子)



➤ Description of Classical Mechanics (经典力学描述)

❖ Equation of motion:

$$m \frac{d^2 x}{dt^2} = F = -Kx$$

→ $\frac{d^2 x}{dt^2} + \omega^2 x = 0 \quad \omega \equiv \sqrt{K/m}$

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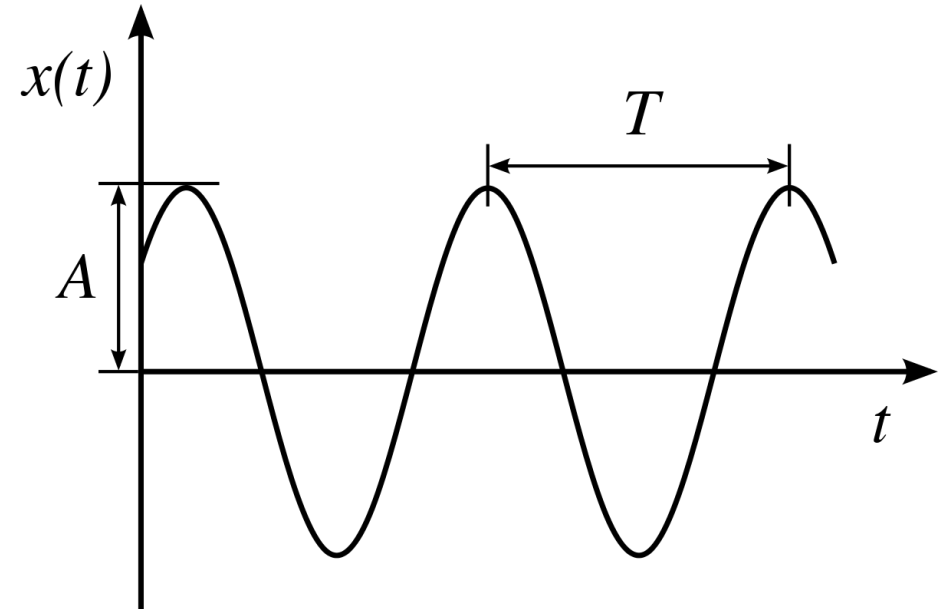
➤ Description of Classical Mechanics (经典力学描述)

❖ Solution:

$$x = Ae^{i(\omega t + \varphi)}$$



$$\text{Re}(x) = A \cos(\omega t + \varphi)$$



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➤ Description of Quantum Mechanics (量子力学描述)

❖ Kinetic energy:

$$T = \frac{1}{2m} p^2 \longrightarrow \hat{T} = \frac{1}{2m} \hat{p}^2$$

❖ Potential energy:

$$V = \frac{1}{2} K x^2 \longrightarrow \hat{V} = \frac{1}{2} m \omega^2 \hat{x}^2$$

❖ Hamiltonian:

$$\hat{H} = \hat{T} + \hat{V} = \frac{1}{2m} \hat{p}^2 + \frac{1}{2} m \omega^2 \hat{x}^2$$

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➤ Description of Quantum Mechanics (量子力学描述)

❖ The stationary Schrodinger equation (定态薛定谔方程):

$$\hat{H}\psi = E\psi$$



E ? ψ ?

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Algebraic Method (代数解法)

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➤ Algebraic Method (代数解法)

❖ For the Hamiltonian $\hat{H} = \frac{1}{2m}\hat{p}^2 + \frac{1}{2}m\omega^2\hat{x}^2$, we can define a new pair of operators:

$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} + \frac{i}{m\omega} \hat{p} \right) \quad \hat{a}^+ = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

→
$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}^+ + \hat{a}) \quad \hat{p} = i \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}^+ - \hat{a})$$

Note that \hat{a}^+ and \hat{a} are not Hermitian!

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➤ Algebraic Method (代数解法)

❖ By applying the new forms of \hat{x} and \hat{p} to the Hamiltonian \hat{H} , we can obtain:

$$\hat{H} = \hbar\omega \left(\hat{a}^+ \hat{a} + \frac{1}{2} \right)$$

Note that a commuting relation $[\hat{a}, \hat{a}^+] = 1$ is used.

❖ Then, the Schrödinger equation reads:

$$\hbar\omega \left(\hat{a}^+ \hat{a} + \frac{1}{2} \right) \psi = E\psi$$

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➤ Algebraic Method (代数解法)

❖ Properties of operators \hat{a}^+ and \hat{a} :

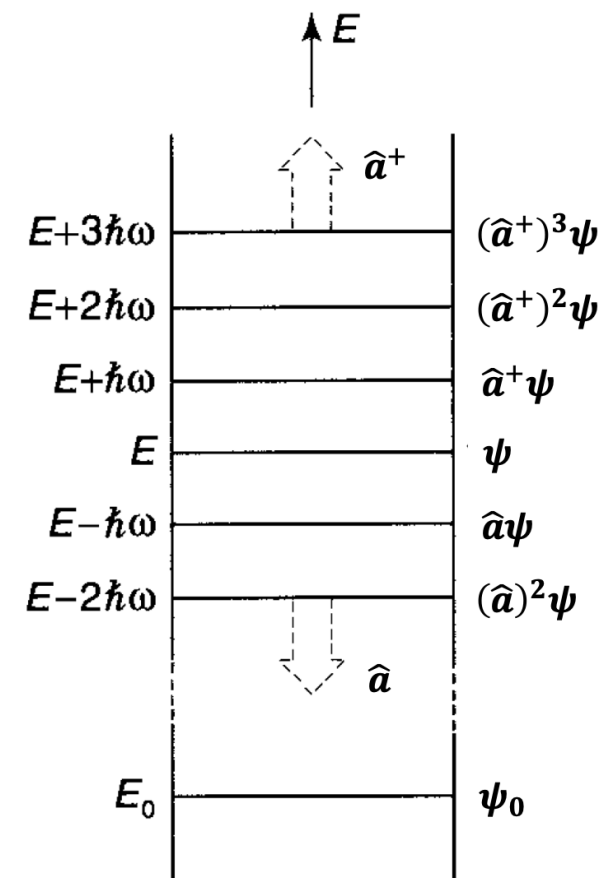
If ψ satisfies the Schrödinger equation with eigenvalue E ,

- $\hat{a}^+ \psi$ satisfies the Schrödinger equation with eigenvalue $E + \hbar\omega$;
- $\hat{a} \psi$ satisfies the Schrödinger equation with eigenvalue $E - \hbar\omega$;

As a result, \hat{a}^+ and \hat{a} are also called **ladder operators** (阶梯算符)!

\hat{a}^+ is called the **raising operator** (升算符);

\hat{a} is called the **lowering operator** (降算符).



Appendix 3: Quantum Harmonic Oscillator (量子谐振子)

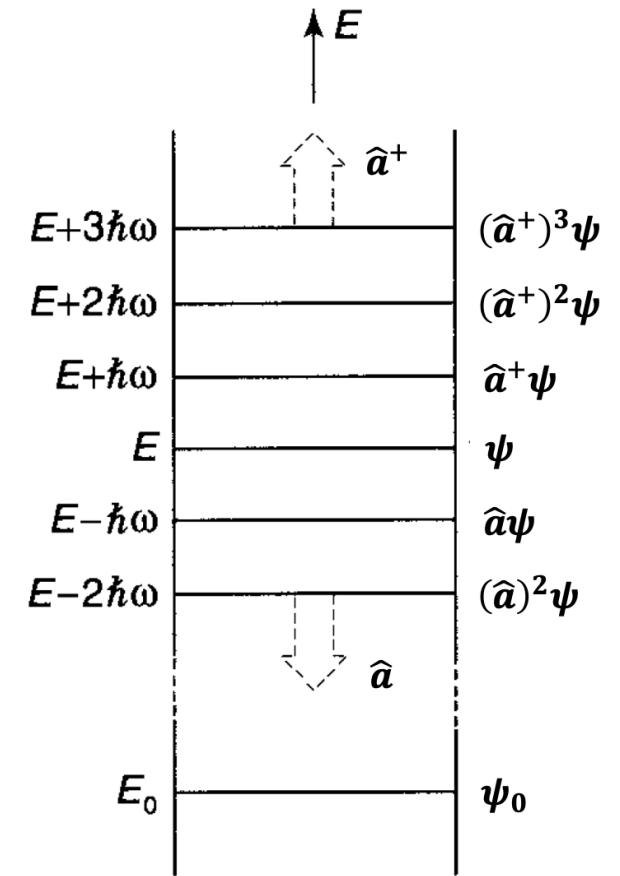


➤ Algebraic Method (代数解法)

❖ The **ground state** (基态) ψ_0 of the harmonic oscillator:

- If \hat{a} is repeatedly applied to ψ , it would end up with a state with energy less than zero, which is meaningless.
- Thus, the lowest possible state ψ_0 (i.e., the **ground state**) of the harmonic oscillator must satisfy:

$$\hat{a}\psi_0 = 0$$



Appendix 3: Quantum Harmonic Oscillator (量子谐振子)



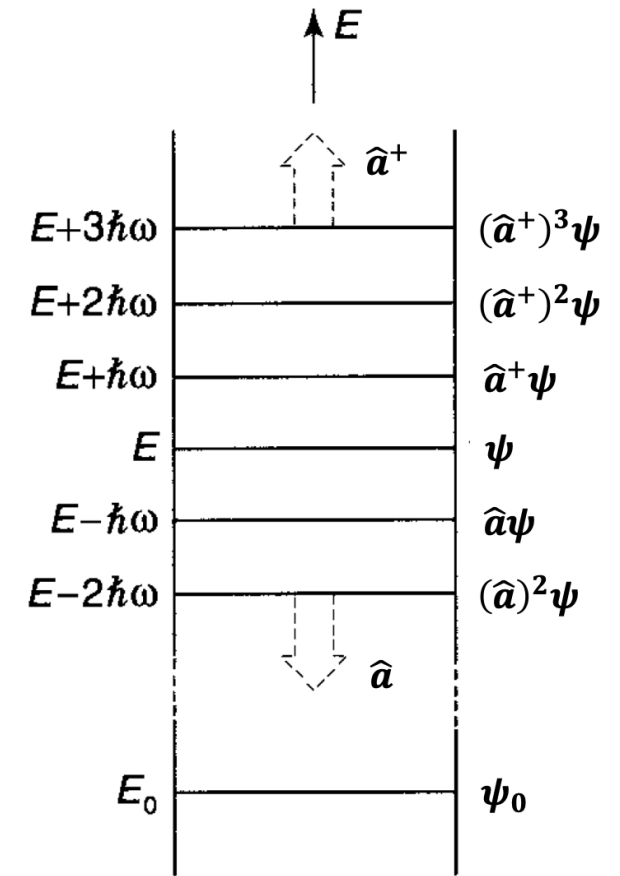
➤ Algebraic Method (代数解法)

❖ The ground state (基态) ψ_0 of the harmonic oscillator:

$$\longrightarrow \hat{a}\psi_0 = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} + \frac{i}{m\omega} \hat{p} \right) \psi_0 = 0$$

$$\hat{p} = -i\hbar \frac{\partial}{\partial x} \longrightarrow \frac{d\psi_0}{dx} = -\frac{m\omega}{\hbar} x\psi_0$$

$$\longrightarrow \psi_0(x) = A_0 e^{-\frac{m\omega}{2\hbar} x^2} \quad E_0 = \frac{1}{2} \hbar \omega$$



Appendix 3: Quantum Harmonic Oscillator (量子谐振子)



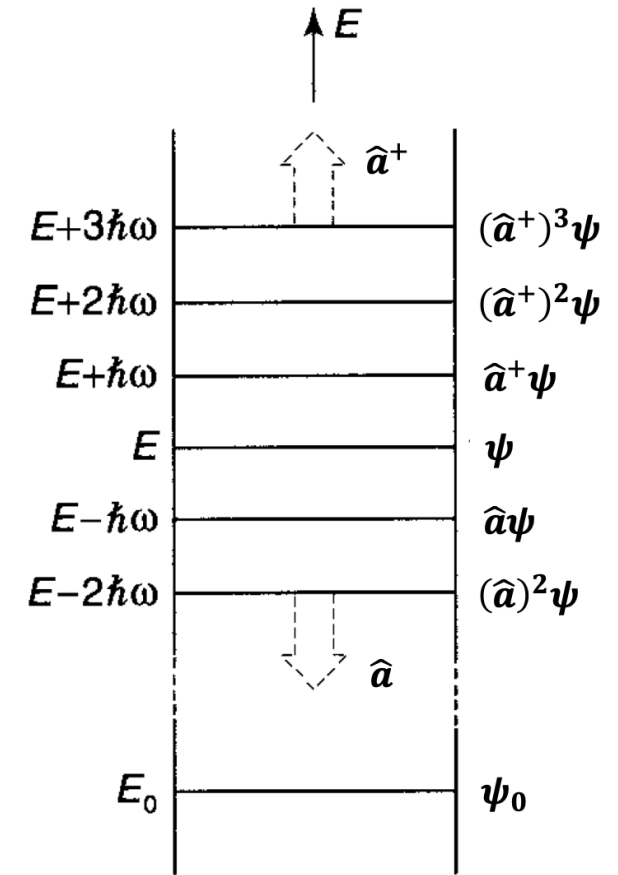
➤ Algebraic Method (代数解法)

❖ The **excited states** (激发态) ψ_n of the harmonic oscillator:

- All the states higher in energy than the ground state, i.e., the **excited states**, can be obtained by repeatedly applying \hat{a}^+ to ψ_0 :

$$\psi_n(x) = A_n (\hat{a}^+)^n e^{-\frac{m\omega}{2\hbar}x^2}$$

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega \quad n = 0, 1, 2, 3, \dots$$



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➤ Algebraic Method (代数解法)

❖ If we define a number operator $\hat{N} = \hat{a}^+ \hat{a}$ and let $\psi_n = |n\rangle$, it can be obtained that:

$$\hat{H} = \hbar\omega \left(\hat{N} + \frac{1}{2} \right)$$

$$\hat{N}|n\rangle = n|n\rangle$$

$$\hat{a}^+ |n\rangle = \sqrt{n+1} |n+1\rangle \quad \hat{a} |n\rangle = \sqrt{n} |n-1\rangle$$

$$|n\rangle = \frac{(\hat{a}^+)^n}{\sqrt{n!}} |0\rangle$$



Analytic Method (解析解法)

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➤ Analytic Method (解析解法)

❖ In the **position representation** (位置表象), the Schrödinger equation reads:

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} m\omega^2 x^2 \psi = E\psi$$

By introducing two dimensionless variables:

$$\xi = \sqrt{\frac{m\omega}{\hbar}} x \quad \varepsilon = \frac{2E}{\hbar\omega}$$



$$\frac{d^2\psi}{d\xi^2} = (\xi^2 - \varepsilon)\psi$$

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➤ Analytic Method (解析解法)

- ❖ Given the fact that $\psi \rightarrow 0$ when $x \rightarrow \pm\infty$, by applying some complicated mathematical techniques (omitted here), the **normalized stationary states** (归一化定态) and the corresponding **eigenvalues** of the harmonic oscillator are:

$$\psi_n(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(\xi) e^{-\xi^2/2}$$

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega \quad n = 0, 1, 2, 3, \dots$$

Here, $H_n(\xi)$ denotes the so-called **Hermite polynomials** (厄米多项式), e.g.,

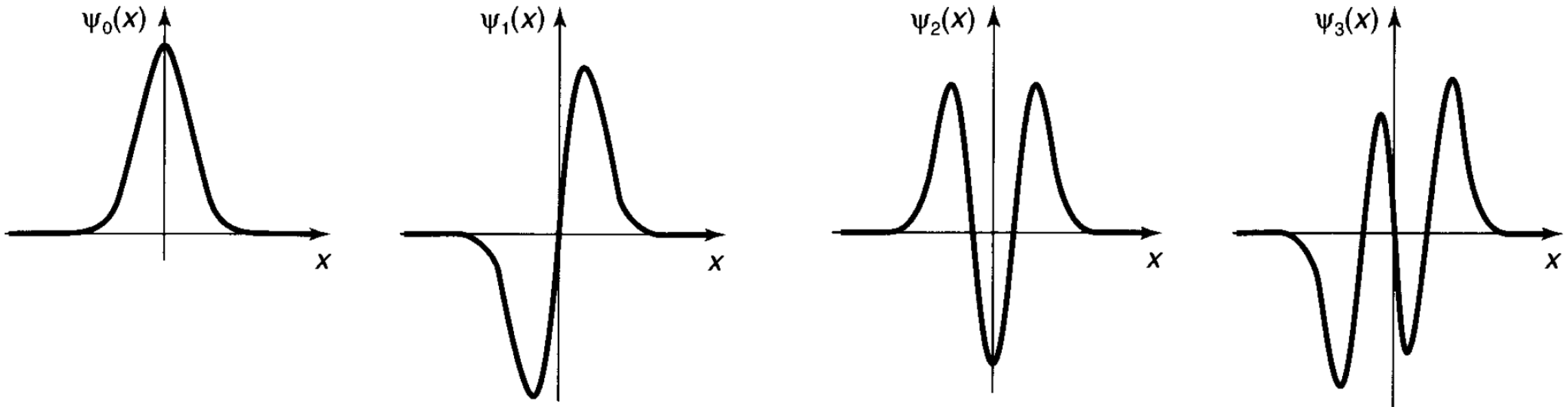
$$H_0(\xi) = 1 \quad H_1(\xi) = 2\xi \quad H_2(\xi) = 4\xi^2 - 2 \quad H_3(\xi) = 8\xi^3 - 12\xi \quad \dots$$

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➤ Analytic Method (解析解法)

❖ The spatial distribution of the wave functions:

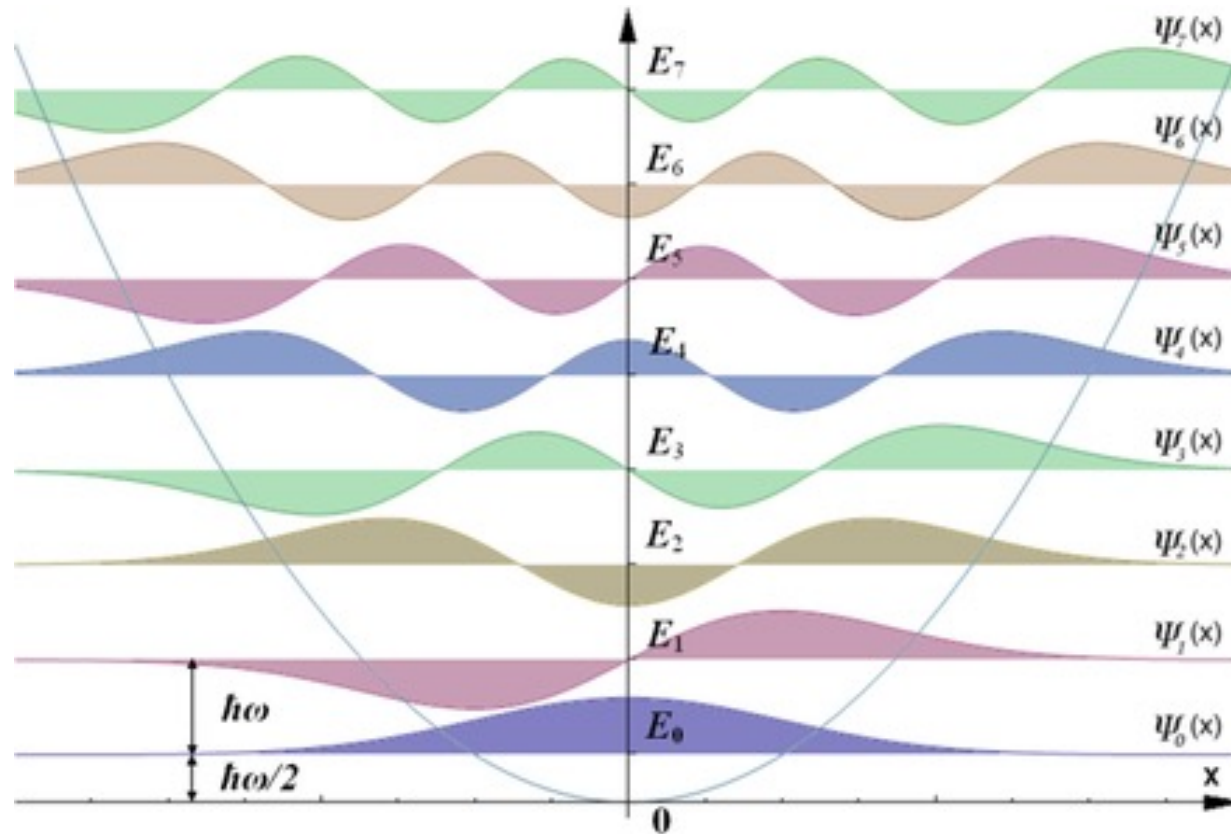


Appendix 3: Quantum Harmonic Oscillator (量子谐振子)



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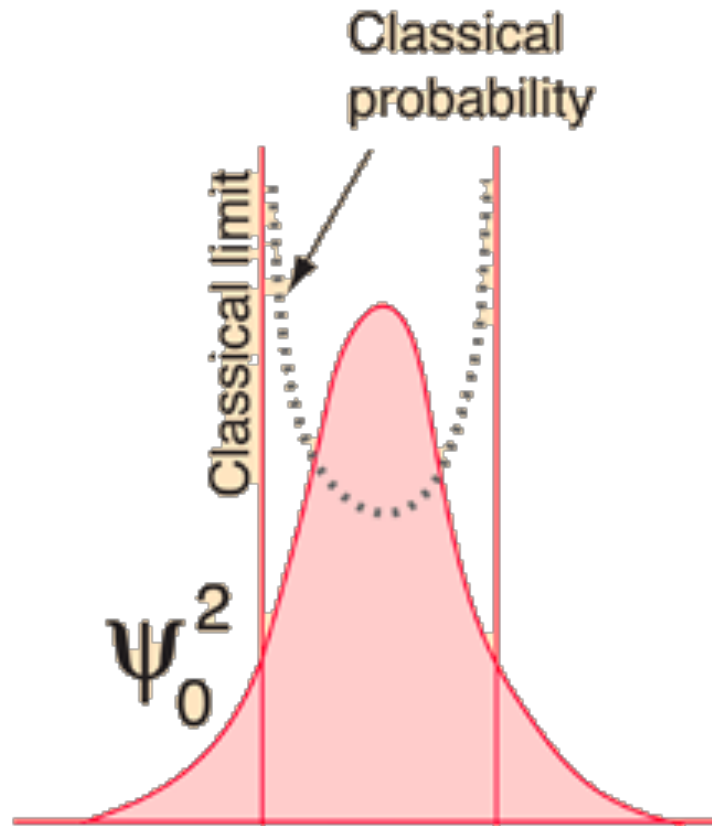


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➤ Analytic Method (解析解法)

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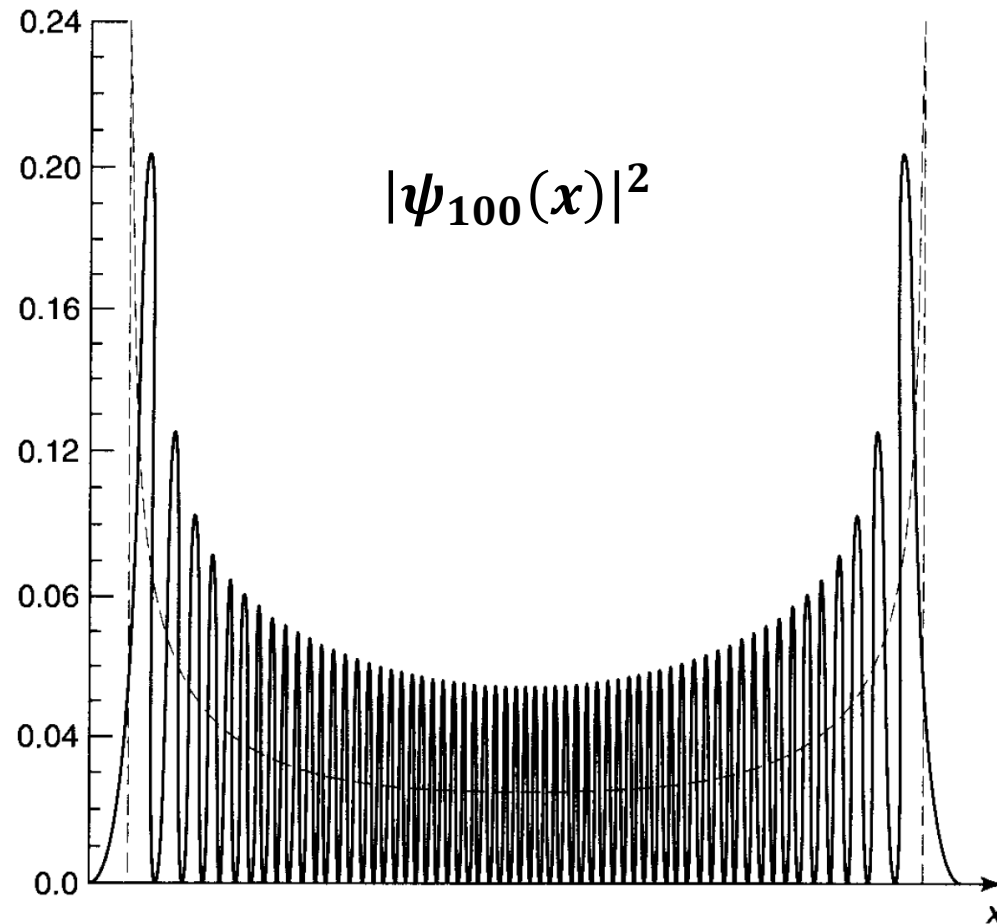


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➤ Analytic Method (解析解法)

❖ The spatial distribution of the wave functions:





Summary (总结)

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➤ Summary (总结)

❖ The quantum description of harmonic oscillator.

$$\hat{H}\psi = E\psi \quad \hat{H} = \frac{1}{2m}\hat{p}^2 + \frac{1}{2}m\omega^2\hat{x}^2$$

$$E_n = \left(n + \frac{1}{2}\right)\hbar\omega \quad n = 0, 1, 2, 3, \dots$$

❖ Algebraic method: **ladder operators**

❖ Analytic method: **wave functions**