

### Outline

- Chapter 4.1 Nearly-Free-Electron Model (近自由电子模型)
- Chapter 4.2 Tight-Binding Model (紧束缚模型)
- Chapter 4.3 Square-Potential-Well Model (方势阱模型)
- Chapter 4.4 Conductors & Nonconductors (导体与非导体)

#### **Objectives**



> To learn the **tight-binding approximation**.

> To understand the characteristics of energy band for a 1D TB model.



# Tight-Binding Approximation (紧束缚近似)



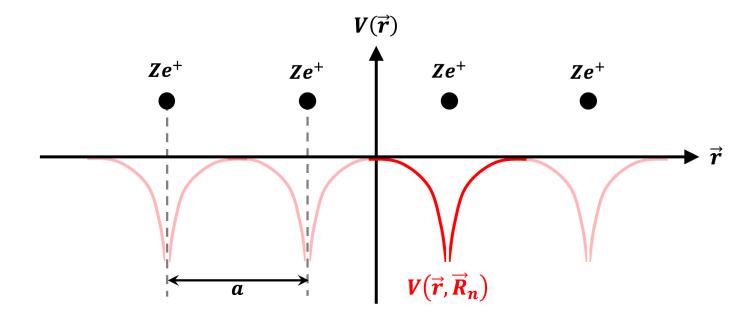
- ➤ Tight-Binding Approximation (紧束缚近似)
  - ❖ The tight-binding (TB) approximation states that:
    - The valence electrons are tightly bound to the atom to which they belong;
    - The valence electrons have limited interactions with the potentials on the surrounding atoms;
    - The potentials on the surrounding atoms can be treated as a perturbation.

The TB approximation can give a clear picture of how the energy bands of crystals are connected with the energy levels of atoms/molecules!



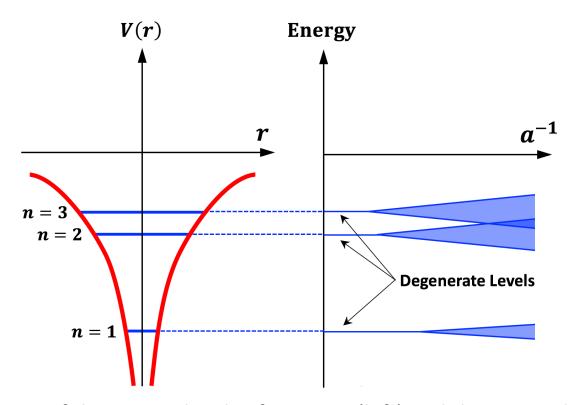
- ➤ Tight-Binding Approximation (紧束缚近似)
  - ❖ The potential of a crystal in the TB approximation:

$$V(\vec{r}) = V_0(\vec{r}) + \Delta V(\vec{r}) = V(\vec{r}, \vec{R}_n) + [V(\vec{r}) - V(\vec{r}, \vec{R}_n)]$$





- ➤ Tight-Binding Approximation (紧束缚近似)
  - ❖ The correspondence between energy levels and energy bands:



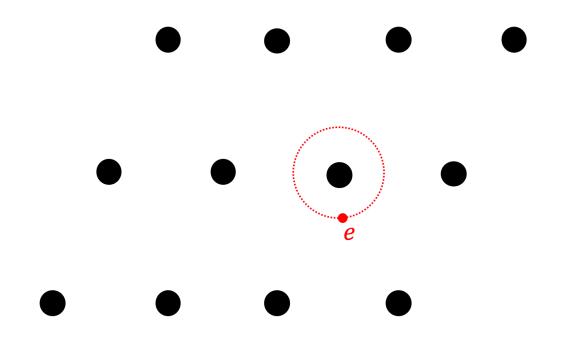
Schematic diagram of the energy levels of an atom (left) and the energy bands of a crystal consisting of the same types of atoms (right). Here, a denotes interatomic distance.



# Method of LCAO (原子轨道线性组合法)

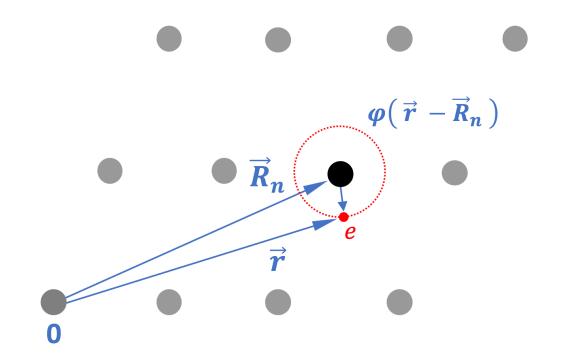


- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - \* We consider a **3D** monoatomic lattice with  $N = N_1 N_2 N_3$  atoms. Each atom has one valence electron with mass m. The electrons move in the crystal lattice with a periodic potential.





- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - \* When the interatomic interactions are fully neglected, the valence electron of an atom with position  $\vec{R}_n = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3$  moves only around the isolated atom in its localized atomic orbital  $\varphi(\vec{r} \vec{R}_n)$ .





- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - \* The Schrödinger equation for the electron in the **isolated atom** reads:

$$\left[ -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r} - \vec{R}_n) \right] \varphi(\vec{r} - \vec{R}_n) = \varepsilon \varphi(\vec{r} - \vec{R}_n)$$

Here,  $V(\vec{r} - \vec{R}_n)$  denotes the **atomic potential** of the atom at  $\vec{R}_n$  and  $\varepsilon$  the energy level of the atomic orbital.



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - \* The single-electron Schrödinger equation for the **crystal lattice** reads:

$$\left[ -\frac{\hbar^2}{2m} \nabla^2 + \mathbf{V}(\vec{r}) \right] \boldsymbol{\psi}(\vec{r}) = \mathbf{E} \boldsymbol{\psi}(\vec{r})$$

Here,  $V(\vec{r})$  denotes the **periodic potential** of the lattice and E and  $\psi(\vec{r})$  the electron energy and wavefunction in the crystal.



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

$$\widehat{H} = -\frac{\hbar^2}{2m} \nabla^2 + \mathbf{V}(\vec{r})$$

$$\widehat{\boldsymbol{H}} = \widehat{\boldsymbol{H}}_{0} + \widehat{\boldsymbol{H}}' \qquad \widehat{\boldsymbol{H}}_{0} = -\frac{\hbar^{2}}{2m} \nabla^{2} + \boldsymbol{V} (\overrightarrow{\boldsymbol{r}} - \overrightarrow{\boldsymbol{R}}_{n})$$

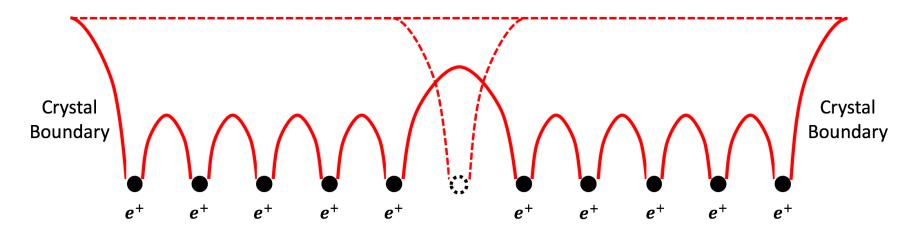
$$\widehat{H}' = V(\overrightarrow{r}) - V(\overrightarrow{r} - \overrightarrow{R}_n)$$
 (Perturbation)

Note that:  $V(\vec{r}) = \sum_{j=1}^{N} V(\vec{r} - \vec{R}_j)$ .



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

$$\widehat{H}'(\overrightarrow{r}) = V(\overrightarrow{r}) - V(\overrightarrow{r} - \overrightarrow{R}_n)$$



Schematic diagram of the perturbation potential  $\widehat{H}'$ 



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

In the zeroth-order approximation:

$$\widehat{H}_0 \varphi(\overrightarrow{r} - \overrightarrow{R}_n) = \varepsilon \varphi(\overrightarrow{r} - \overrightarrow{R}_n)$$

Since the energy level  $\varepsilon$  is **degenerate** for all the isolated atomic orbitals, we have to apply the **degenerate perturbation theory**.



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

The wave function of the crystal can be built up by a linear combination of the atomic orbitals (LCAO):

$$\psi(\vec{r}) = \sum_{n} \alpha_{n} \varphi(\vec{r} - \vec{R}_{n})$$



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

It can be expected that the coefficient  $\alpha_n$  of the linear combination satisfies:

$$\alpha_n = \frac{1}{\sqrt{N}} e^{i\vec{k}\cdot\vec{R}_n}$$

$$\psi(\vec{r}) = e^{i\vec{k}\cdot\vec{r}} \frac{1}{\sqrt{N}} \sum_{n} e^{-i\vec{k}\cdot(\vec{r}-\vec{R}_n)} \varphi(\vec{r}-\vec{R}_n) = e^{i\vec{k}\cdot\vec{r}} u(\vec{r}) \quad \text{(Bloch wave)}$$



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

By applying  $\psi(\vec{r})$  to  $\vec{H}\psi(\vec{r})=E\psi(\vec{r})$ , it is obtained:

$$\sum_{n} \alpha_{n} \left[ \varepsilon + \widehat{H}'(\vec{r}) \right] \varphi(\vec{r} - \vec{R}_{n}) = E \sum_{n} \alpha_{n} \varphi(\vec{r} - \vec{R}_{n})$$



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

Since the interactions between atoms are very weak, it can be approximated that:

$$\int \boldsymbol{\varphi}^* (\vec{r} - \vec{R}_m) \boldsymbol{\varphi} (\vec{r} - \vec{R}_n) d\vec{r} = \boldsymbol{\delta}_{mn}$$



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

By applying to the Schrödinger equation, we obtain:

$$-\sum_{n} \alpha_{n} J(\vec{R}_{m} - \vec{R}_{n}) = (E - \varepsilon)a_{m}$$

Here, 
$$-J(\vec{R}_m - \vec{R}_n) = \int \varphi^* (\vec{\xi} - (\vec{R}_m - \vec{R}_n)) \hat{H}'(\vec{\xi}) \varphi(\vec{\xi}) d\vec{\xi}$$



- Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

By applying  $\alpha_n = \frac{1}{\sqrt{N}} e^{i\vec{k}\cdot\vec{R}_n}$ , we end up with the final solutions:

$$\vec{R}_{k} = \varepsilon - \sum_{S} J(\vec{R}_{S}) e^{-i\vec{k}\cdot\vec{R}_{S}}$$

$$\vec{R}_{S} = \vec{R}_{m} - \vec{R}_{n}$$
(the separation between sites 
$$\vec{k} = \frac{h_{1}}{N_{1}} \vec{k}_{1} + \frac{h_{2}}{N_{2}} \vec{k}_{2} + \frac{h_{3}}{N_{3}} \vec{k}_{3}$$

$$(h, h_{2}h_{3} \text{ are integers})$$

$$\vec{R}_S = \vec{R}_m - \vec{R}_n$$
(the separation between sites  $m$  and  $n$ )

$$\vec{k} = \frac{h_1}{N_1} \vec{b}_1 + \frac{h_2}{N_2} \vec{b}_2 + \frac{h_3}{N_3} \vec{b}_3$$

$$(h_1 h_2 h_3 \text{ are integers})$$



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

Since the interactions between atoms are very weak, we can further obtain:

When 
$$\vec{R}_S = \vec{R}_m - \vec{R}_n = 0$$
  $J(\vec{R}_S) = -\int \boldsymbol{\varphi}^*(\vec{\xi}) \hat{\boldsymbol{H}}'(\vec{\xi}) \boldsymbol{\varphi}(\vec{\xi}) d\vec{\xi} = J_0$  (called "site energy" 在位能)

When 
$$\vec{R}_S = \vec{a}$$
 (neighbors)  $J(\vec{R}_S) = -\int \boldsymbol{\varphi}^*(\vec{\xi} - \vec{a}) \hat{H}'(\vec{\xi}) \boldsymbol{\varphi}(\vec{\xi}) d\vec{\xi} = J_1(\vec{a})$  (called "overlap integral" 交叠积分)



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Perturbation-theory** treatment:

Since the interactions between atoms are very weak, we can further obtain:

$$E_k = \varepsilon - J_0 - \sum_{j=\text{nbs}} J_1(\vec{a}_j) e^{-i\vec{k}\cdot\vec{a}_j}$$

$$\vec{k} = \frac{h_1}{N_1} \vec{b}_1 + \frac{h_2}{N_2} \vec{b}_2 + \frac{h_3}{N_3} \vec{b}_3$$
 ( $h_1 h_2 h_3$  are integers)

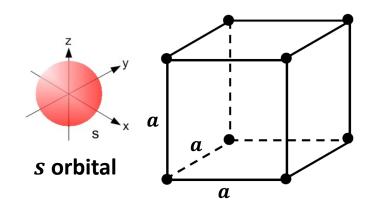


- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Example:** the energy band of *s* orbitals in a simple cubic lattice.

$$E_k = \varepsilon - J_0 - \sum_{j=\text{nbs}} J_1(\vec{a}_j) e^{-i\vec{k}\cdot\vec{a}_j}$$

For a given site (0,0,0), the coordinates of the 6 neighboring sites:

$$(\pm a, 0, 0)$$
  $(0, \pm a, 0)$   $(0, 0, \pm a)$ 



$$E_k = \varepsilon - J_0 - J_1 \left( e^{ik_x a} + e^{-ik_x a} + e^{ik_y a} + e^{-ik_y a} + e^{ik_z a} + e^{-ik_z a} \right)$$

$$E_k = \varepsilon - J_0 - 2J_1(\cos k_x a + \cos k_y a + \cos k_z a)$$



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Example:** the energy band of *s* orbitals in a simple cubic lattice.

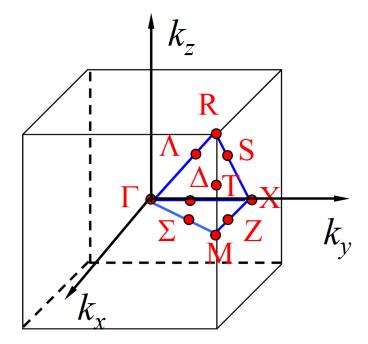
The energy at the high symmetry points in reciprocal space:

$$\Gamma(0,0,0): \quad E(\Gamma) = \varepsilon - J_0 - 6J_1$$

$$X(0,\frac{\pi}{a},0): E(X) = \varepsilon - J_0 - 2J_1$$

$$R\left(\frac{\pi}{a}, \frac{\pi}{a}, \frac{\pi}{a}\right): E(R) = \varepsilon - J_0 + 6J_1$$

$$M(\frac{\pi}{a}, \frac{\pi}{a}, 0)$$
:  $E(M) = \varepsilon - J_0 + 2J_1$ 

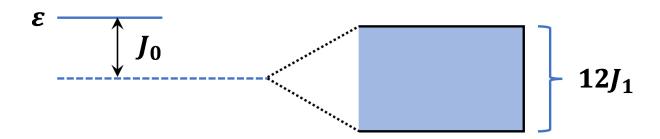


Here,  $J_0 > 0$  and  $J_1 > 0$ 



- ➤ Method of Linear Combination of Atomic Orbitals (原子轨道线性组合法)
  - **Example:** the energy band of *s* orbitals in a simple cubic lattice.

The band width (带宽):  $\Delta E = E(R) - E(\Gamma) = 12J_1$ 



The bandwidth is determined by the **magnitude of**  $J_1$ , i.e., the overlap integral between neighboring atomic orbitals!



# From Energy Levels to Energy Bands (从能级到能带)

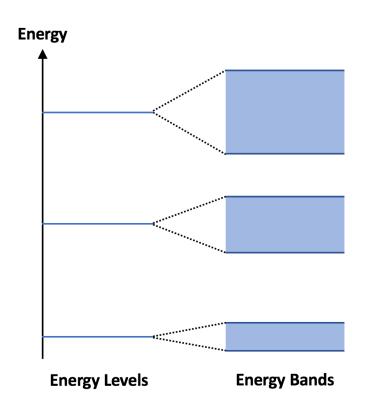


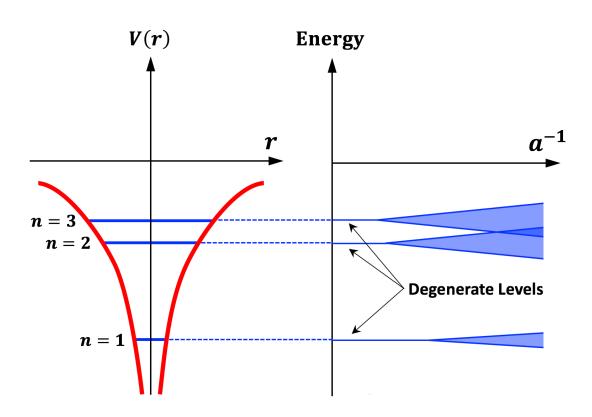
- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - **From atomic energy levels** to energy bands:

Atomic Orbitals Linear Combination Crystal "Orbitals"



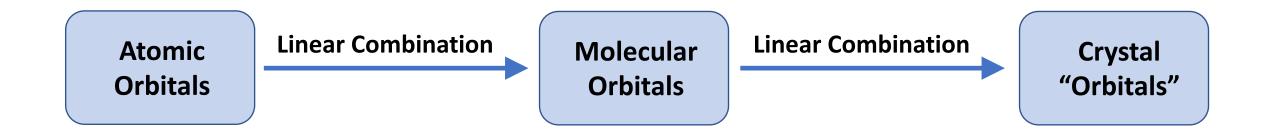
- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - **\*** From **atomic energy levels** to energy bands:





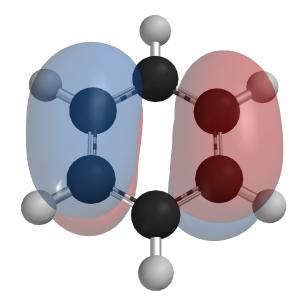


- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ From **molecular energy levels** to energy bands:





- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ From **molecular energy levels** to energy bands:
    - Molecular orbitals (分子轨道)

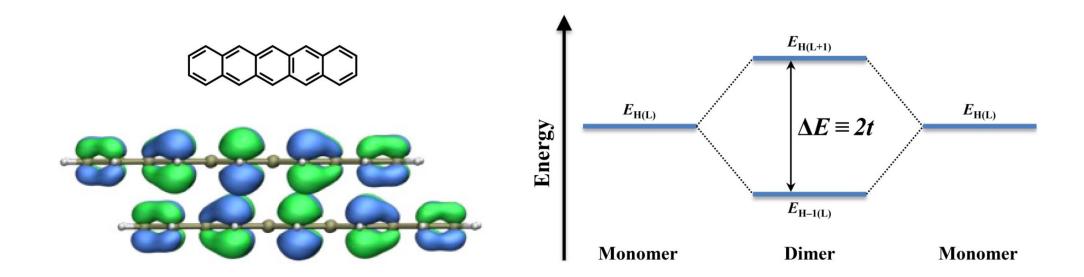


The molecular structure of **benzene (苯)**. The **higl** 

The **highest occupied molecular orbital (HOMO)** of benzene



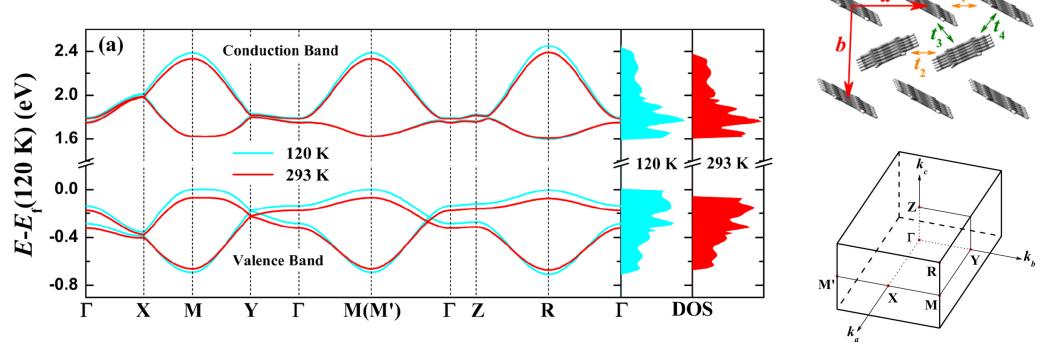
- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ From **molecular energy levels** to energy bands:
    - Interaction between molecular orbitals



The interaction between the HOMOs of two pentacene (并五苯) molecules.



- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ From **molecular energy levels** to energy bands:
    - Energy band of a molecular crystal (分子晶体)



The band structure of pentacene single crystal.



- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ Wannier function (瓦尼尔函数)

In the tight-binding model, the wavefunctions for the electronic states of the energy band are written as a linear combination of the atomic orbitals:

$$\psi_k(\vec{r}) = \frac{1}{\sqrt{N}} \sum_n e^{i\vec{k}\cdot\vec{R}_n} \varphi(\vec{r} - \vec{R}_n)$$



- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ Wannier function (瓦尼尔函数)

In general, the wavefunctions for any energy band (not only that of the TB model) can be written as a linear combination of the **Wannier functions**:

$$\psi_k(\vec{r}) = \frac{1}{\sqrt{N}} \sum_n e^{i\vec{k}\cdot\vec{R}_n} W(\vec{r} - \vec{R}_n)$$



- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ Wannier function (瓦尼尔函数)

The properties of **Wannier functions**:

$$W(\vec{r} - \vec{R}_n) = \frac{1}{\sqrt{N}} \sum_{k} e^{-i\vec{k} \cdot \vec{R}_n} \psi_k(\vec{r})$$

$$\int W^*(\vec{r} - \vec{R}_m)W(\vec{r} - \vec{R}_n)d\vec{r} = \delta_{mn}$$



- ➤ From Energy Levels to Energy Bands (从能级到能带)
  - ❖ Wannier function (瓦尼尔函数)

In the tight-binding model, the Wannier functions are just the wavefunctions of the isolated atomic orbitals:

$$W(\vec{r}-\vec{R}_n)=\varphi(\vec{r}-\vec{R}_n)$$

Wannier functions are convenient for dealing with problems associated with the localization (局域化) of wavefunctions!



# TB Model vs NFE Model (紧束缚模型与近自由电子模型对比)

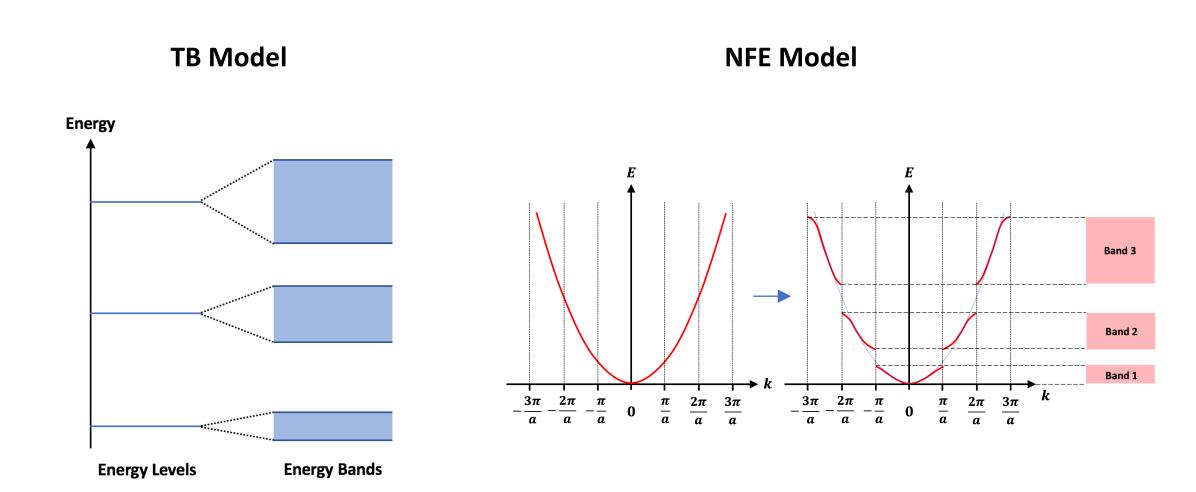


#### > TB Model vs NFE Model

Tight-Binding Model	Nearly-Free-Electron Model
Electrons are tightly bound to nuclei;	Electrons are nearly free;
Wave functions are linear combination of atomic orbitals;	Wave functions are approximated by plane waves;
Energy bands are formed as a result of interactions between atomic levels.	Energy bands are formed as a result of band gaps opened up at the boundary of Brillouin zones.



#### > TB Model vs NFE Model





## Summary (总结)

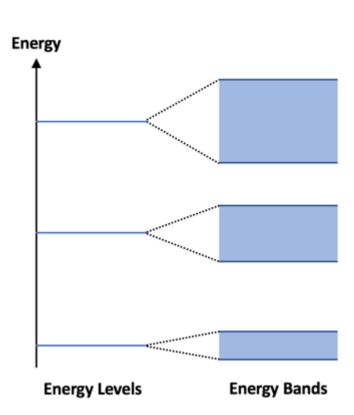


> Summary

Tight-Binding Approximation

Linear Combination of Atomic Orbitals (LCAO)
 Degenerate Perturbation Theory

From Energy Levels to Energy Bands



#### Chapter 4.2: 课后作业



考虑一维单原子链(原子间距为a,链长为Na),对于原子的s能级,利用紧束缚模型,求:

- 1. 原子链能带的色散关系E(k);
- 2. 能带的态密度g(E);
- 3. 能带的宽度.

提交时间: 4月10日之前

提交方式: 手写(写明姓名学号)后拍照,通过本班课代表统一提交电子版