



山东大学
SHANDONG UNIVERSITY

Physics I: Introduction to Wave Theory
SDU Course Number: sd01232810 (Fall 2023)

Lecture 8: Resonance

Outline

- Atomic Oscillators and Resonance
- Fabry-Perot Resonance
- Rectangular Cavity Resonator

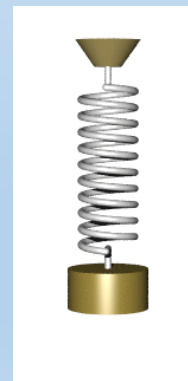
The Phenomenon of Resonance

In physics, **resonance** is a phenomenon in which a vibrating system or external force drives another system to oscillate with greater amplitude at specific frequencies.

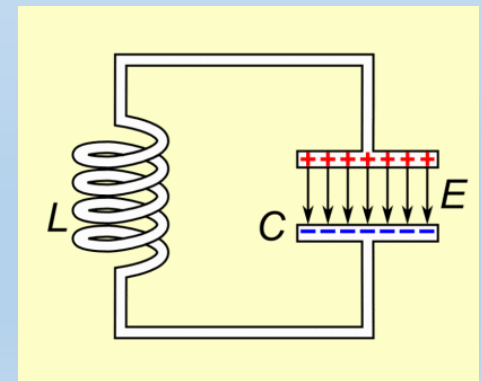
Frequencies at which the response amplitude is a relative maximum are known as the system's **resonant frequencies** or resonance frequencies.

A **resonator** is a device or system that exhibits resonance or resonant behavior, that is, it naturally oscillates at some frequencies, called its resonant frequencies, with greater amplitude than at others.

Swing

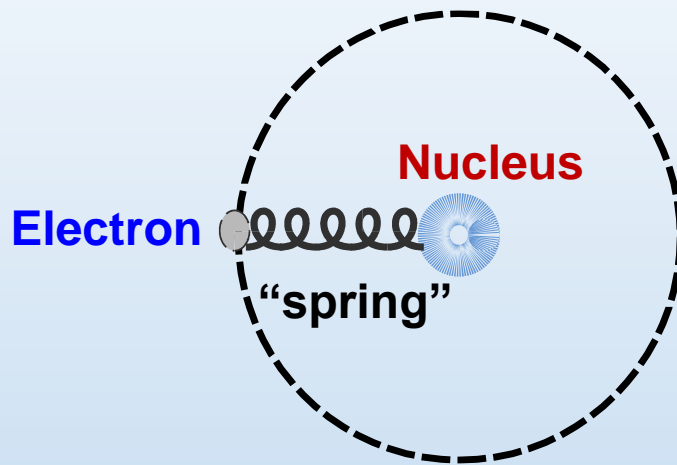


Spring



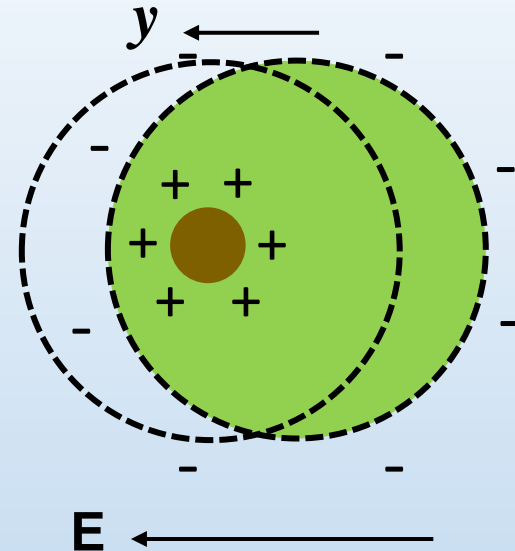
LC circuit

Atomic Oscillators and Resonance



No external E Field

$$P = -Ney$$



(Lorentz Oscillator Model)

$$\varepsilon(\omega) = \varepsilon_0 \left(1 + \frac{P}{\varepsilon_0 E} \right) = \varepsilon_0 \left(1 - \frac{\omega_p^2}{\omega^2 - \omega_0^2 - j\gamma\omega} \right)$$

$$\omega_p = \sqrt{Ne^2 / m\varepsilon_0}$$

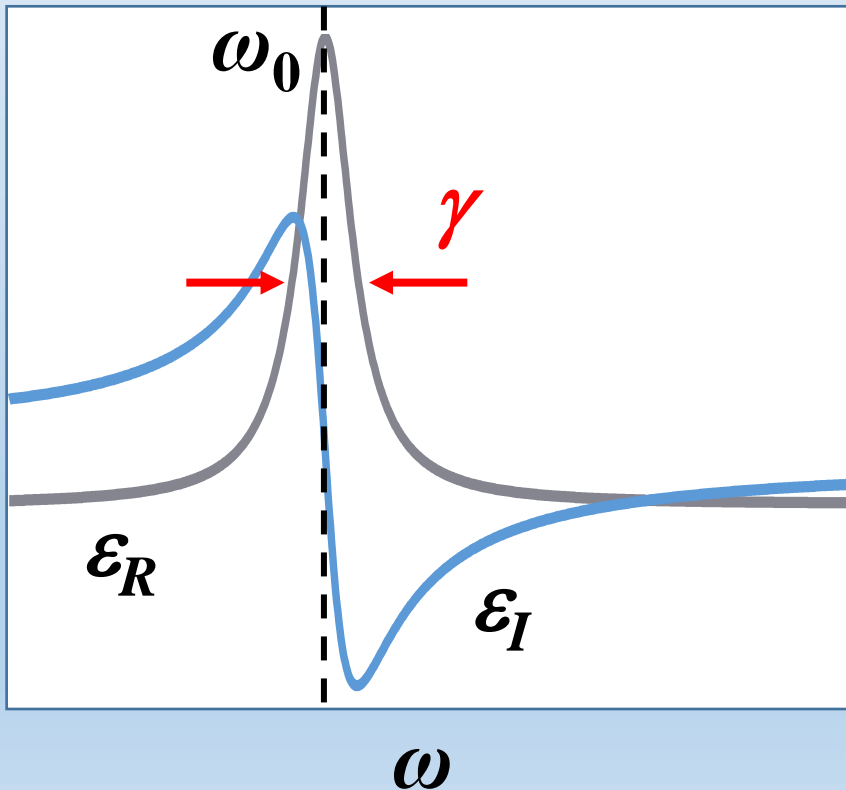
(Plasma frequency)

$$\omega_0 = \sqrt{k/m}$$

Resonant frequency
(or natural frequency)

Complex permittivity

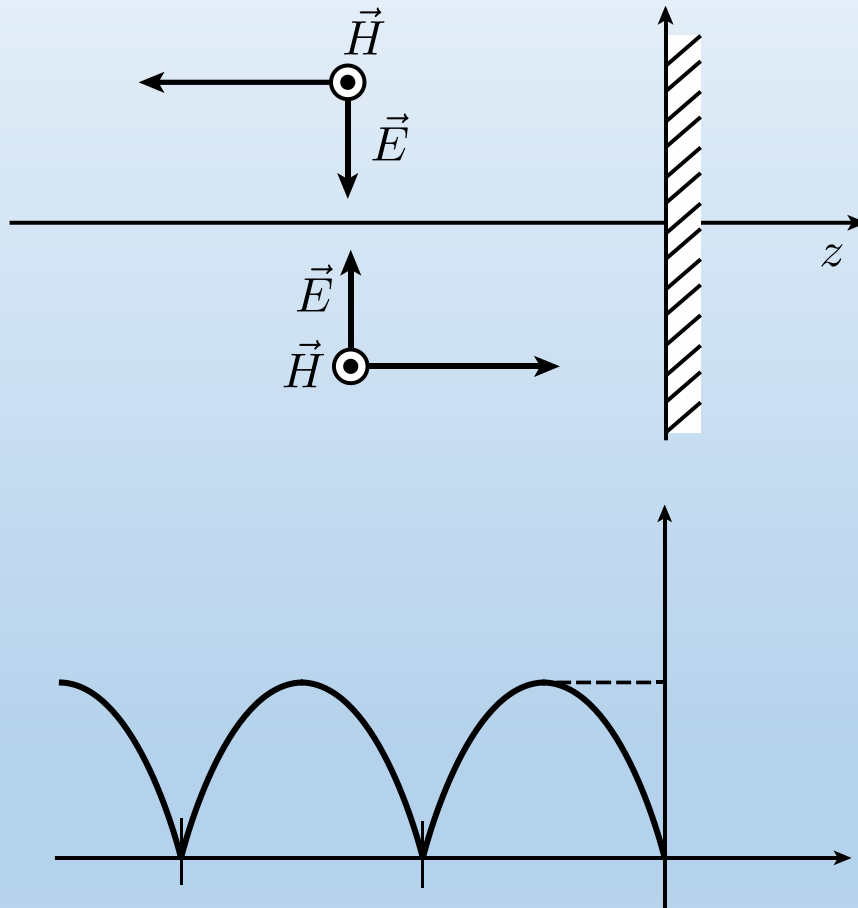
$$\varepsilon(\omega) = \varepsilon_0 \left(1 - \frac{\omega_p^2}{\omega^2 - \omega_0^2 - j\gamma\omega} \right) = \varepsilon_R - j\varepsilon_I$$



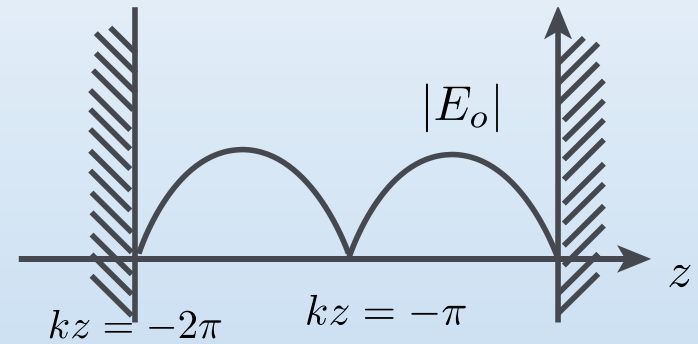
Around the resonance frequency ω_0 , the magnitude of ε_R has a drastic change and ε_I has the maximum value.

Resonators

STANDING WAVE

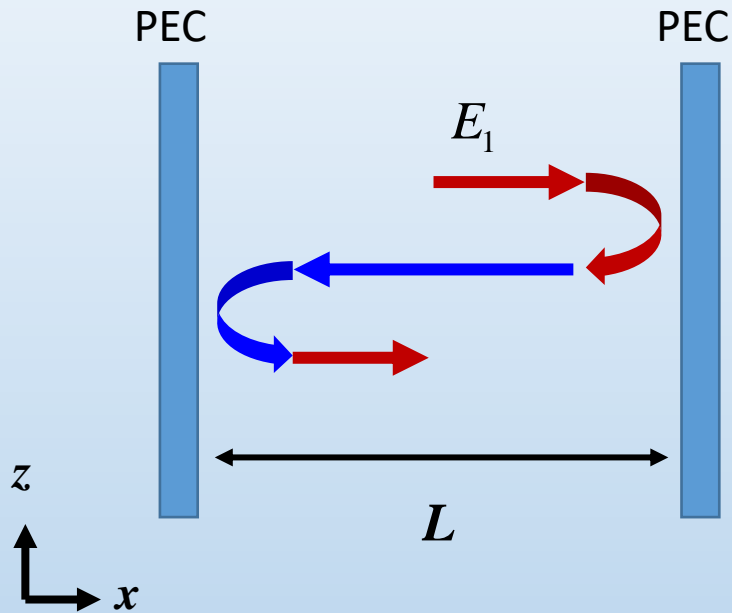


RESONATORS



Terminate the standing wave with a second wall to form a resonator

Resonator (Parallel PEC mirrors)



TE Waveguide Mode ($k_z = 0$)

$$E_y(x, z) = E_0 \sin(kx)$$

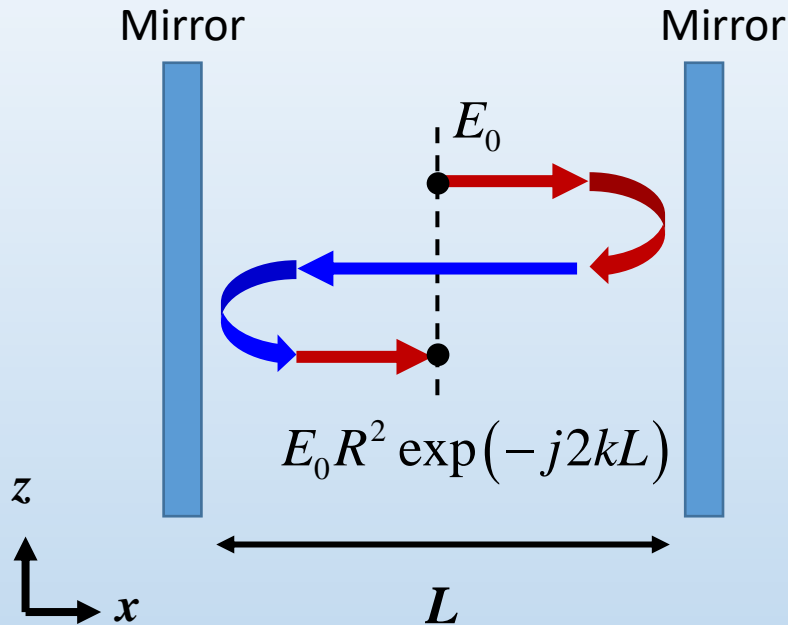
$$H_z(x, z) = -\frac{k}{j\omega\mu} E_0 \cos(kx)$$

$$kL = m\pi \quad (\text{Guidance Condition})$$

$$\Rightarrow \frac{2\pi f}{c} L = m\pi \Rightarrow \boxed{f = m \frac{c}{2L}} \quad m = 1, 2, 3, \dots$$

(Resonant frequency)

Resonator (general mirrors)



$$E_1 = E_0$$

$$E_2 = E_0 R^2 \exp(-j2kL)$$

$$E_3 = E_0 R^4 \exp(-j4kL)$$

\vdots

$$E_{cavity} = E_1 \left[1 + R^2 \exp(-j2kL) + R^4 \exp(-j4kL) + \dots \right] = \frac{E_1}{1 - R^2 \exp(-j2kL)}$$

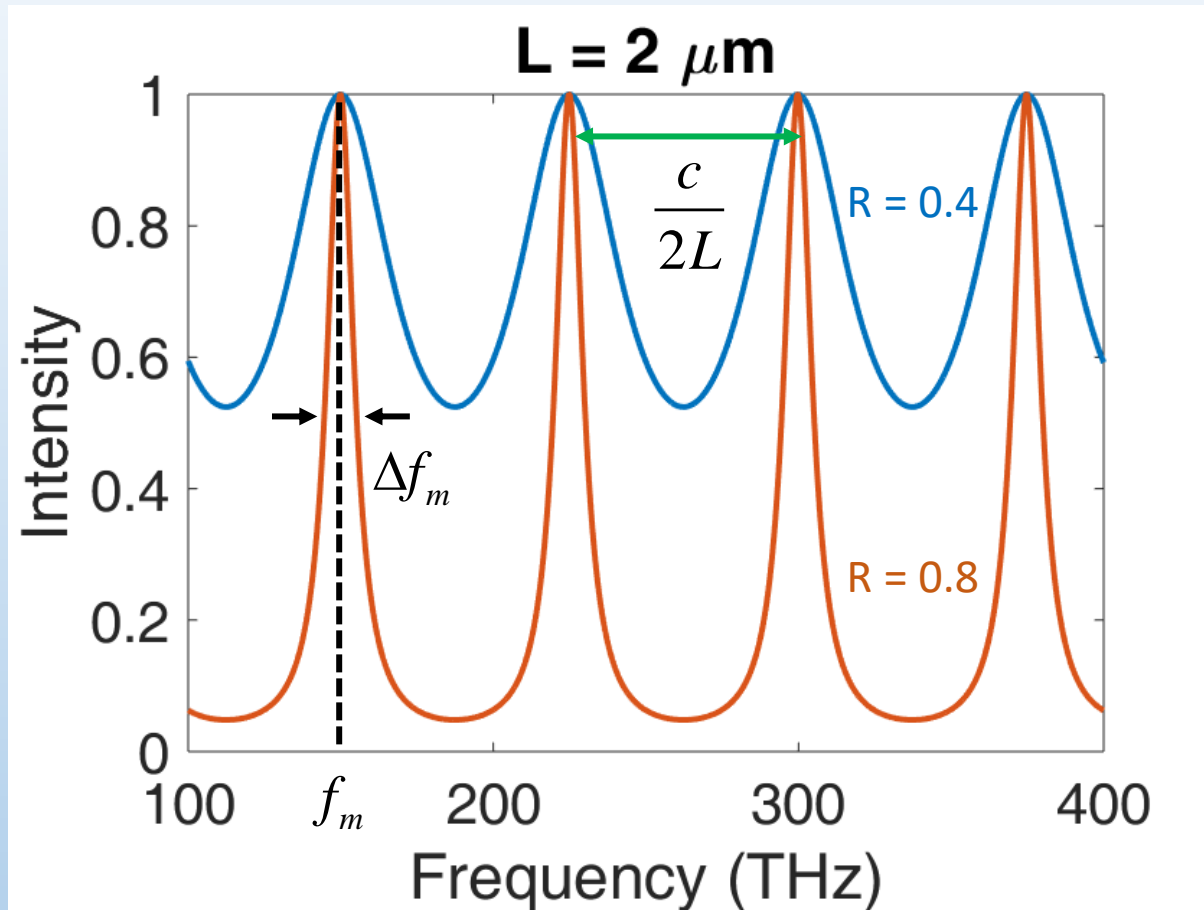
$$I_{cavity} = \frac{I_0}{\left| 1 - R^2 \exp(-j2kL) \right|^2} = \frac{I_0}{(1 - R^2)^2 + 4R^2 \sin^2(kL)}$$

Peaks of I_{cavity} :

$$kL = m\pi$$

(Fabry-Perot resonator)

Fabry-Perot resonator



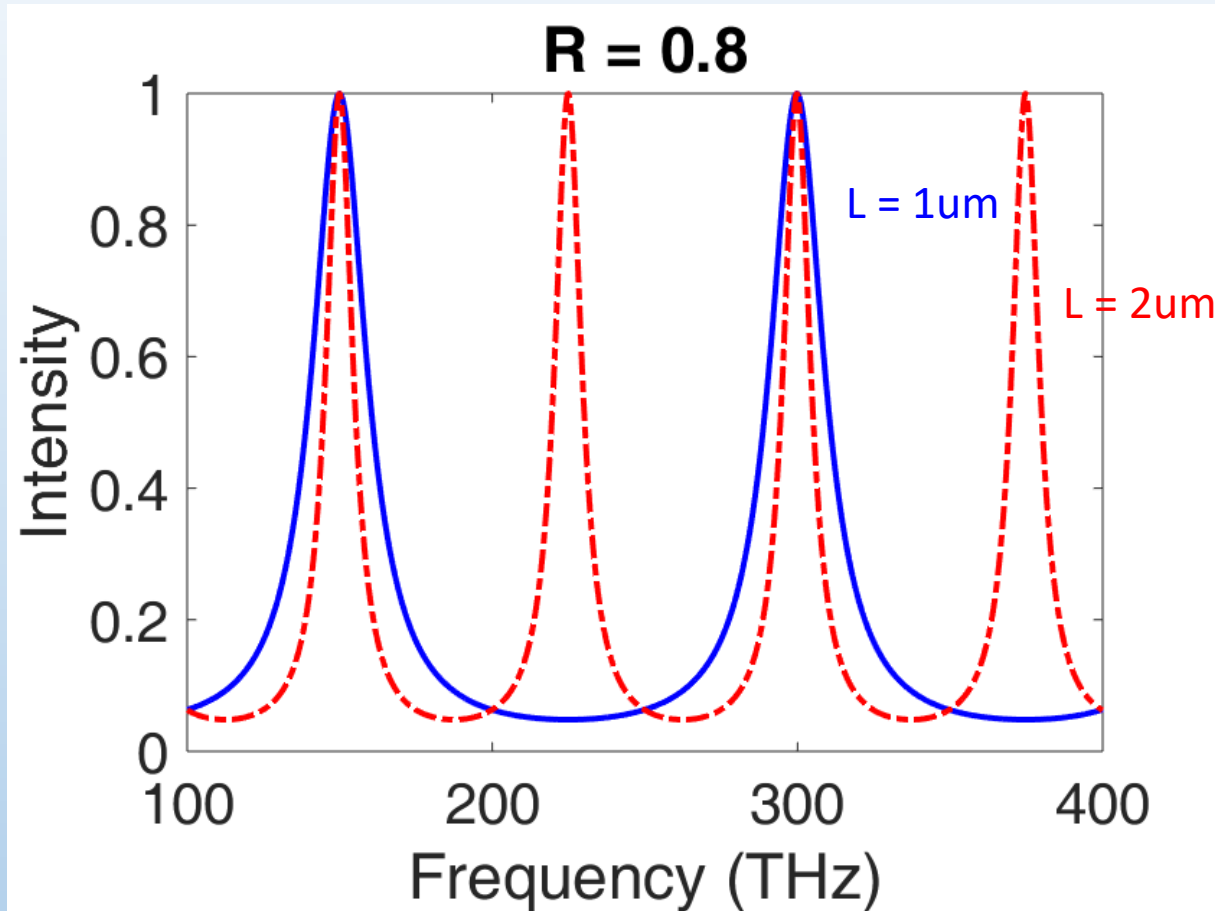
$$f_m = m \frac{c}{2L}$$

$$I_{\max} = \frac{I_0}{(1 - R^2)^2}$$

Full width at half maximum (FWHM): Δf_m

Quality Factor: $Q = f_m / \Delta f_m$

Fabry-Perot resonator



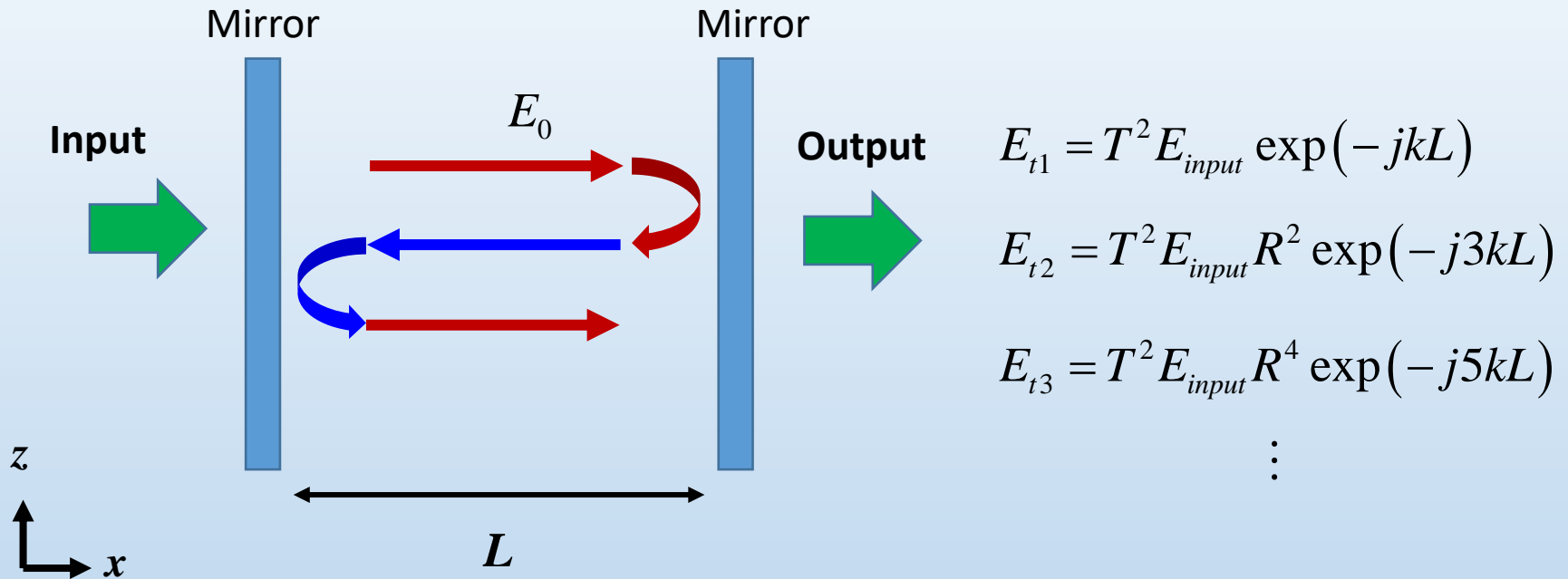
$$f_m = m \frac{c}{2L}$$

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Full width at half maximum (FWHM): Δf_m

Quality Factor: $Q = f_m / \Delta f_m$

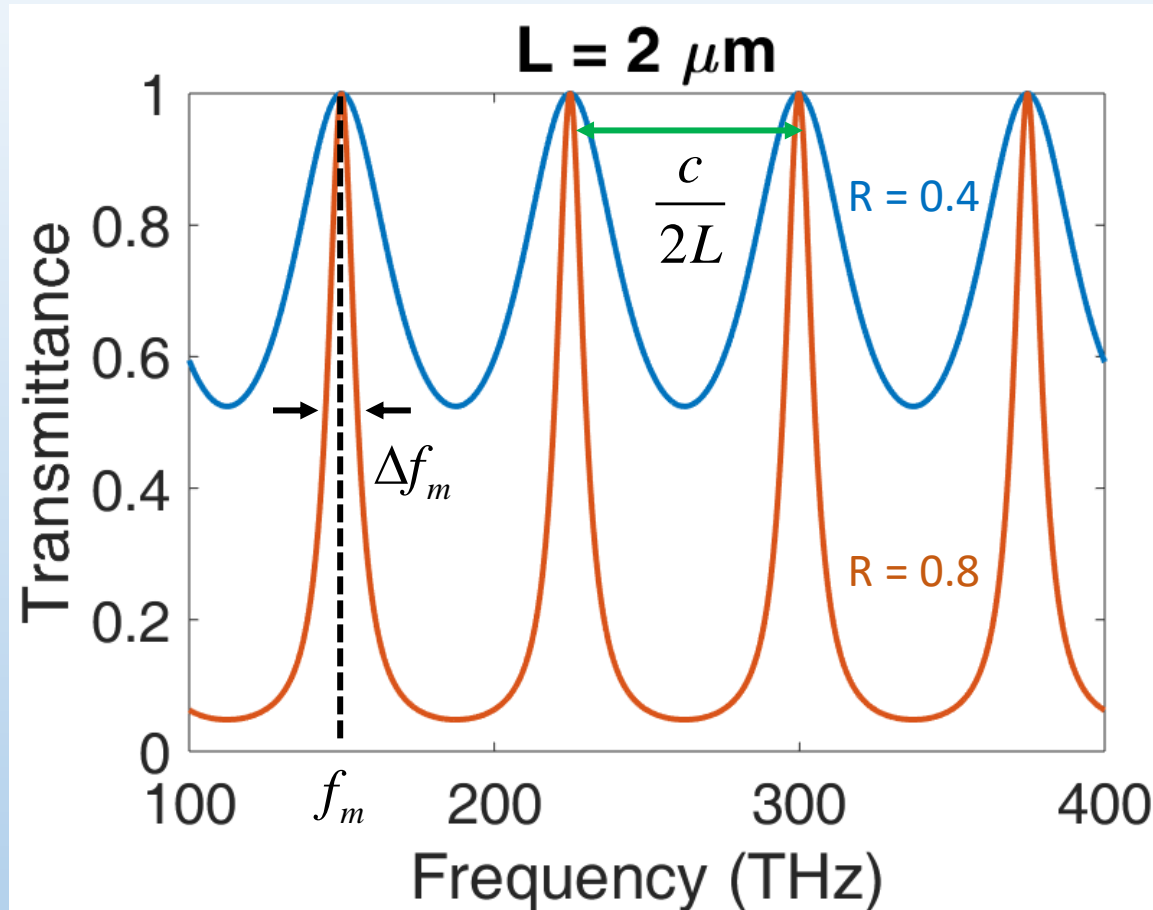
Transmitted light through a FP cavity



$$E_{transmitted} = \sum_n E_{tn} = \frac{T^2 \exp(-jkL)}{1 - R^2 \exp(-j2kL)} E_{input}$$

$$I_{transmitted} = \frac{T^4 I_{input}}{|1 - R^2 \exp(-j2kL)|^2} = \frac{(1 - R^2)^2 I_{input}}{(1 - R^2)^2 + 4R^2 \sin^2(kL)}$$

Transmitted light through a FP cavity



Maximum:

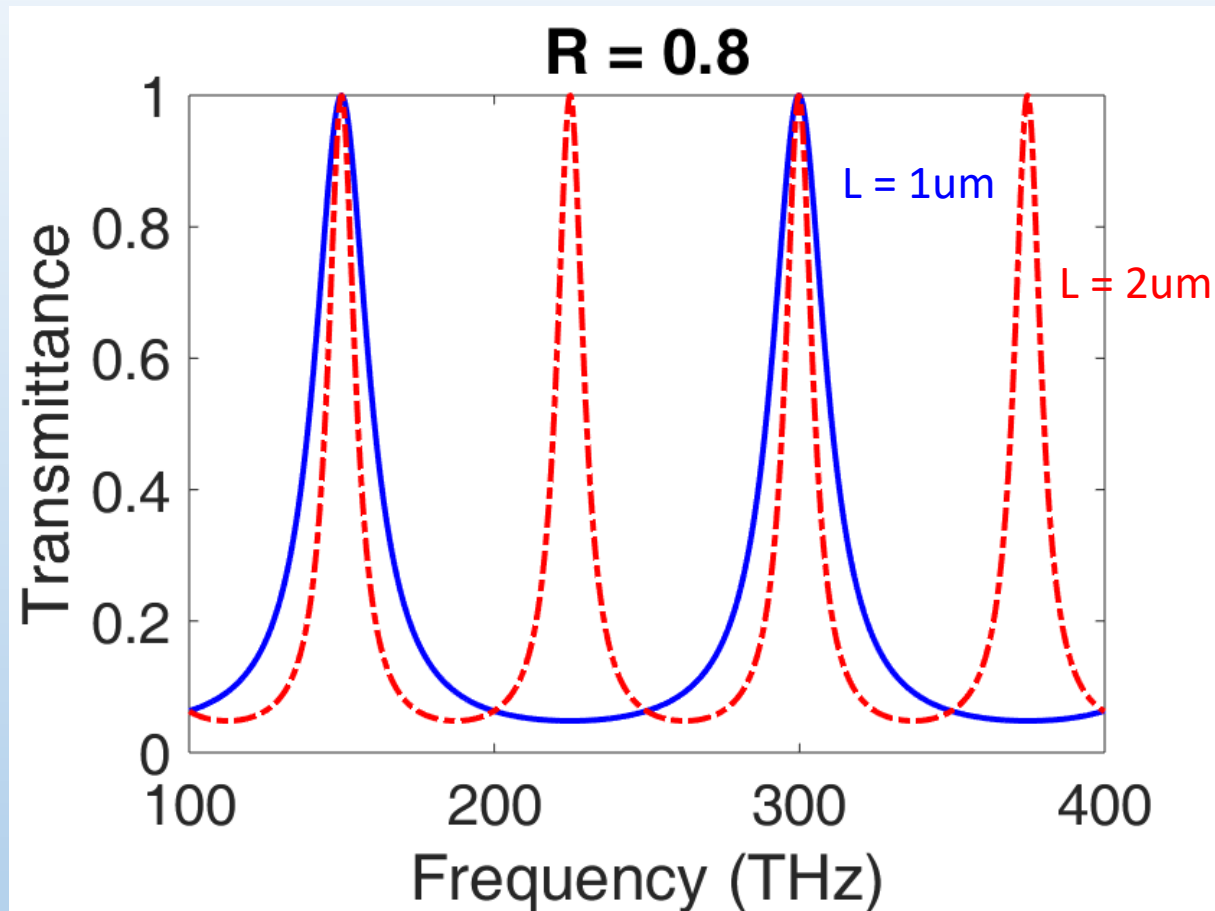
$$kL = m\pi$$

Minimum:

$$kL = \left(m + \frac{1}{2}\right)\pi$$

$$t = \frac{I_{\text{transmitted}}}{I_{\text{input}}} = \frac{(1 - R^2)^2}{(1 - R^2)^2 + 4R^2 \sin^2(kL)}$$

Fabry-Perot resonator



Maximum:

$$kL = m\pi$$

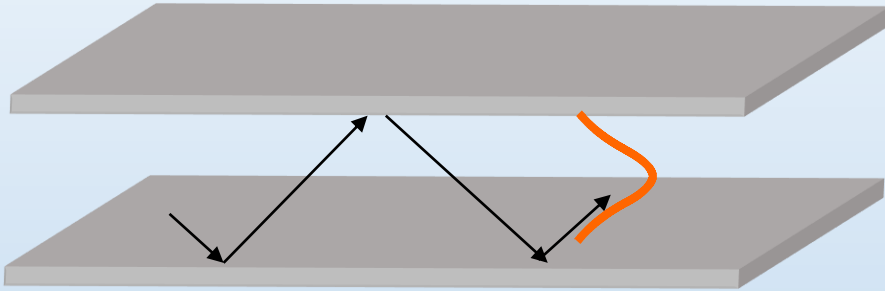
Minimum:

$$kL = \left(m + \frac{1}{2}\right)\pi$$

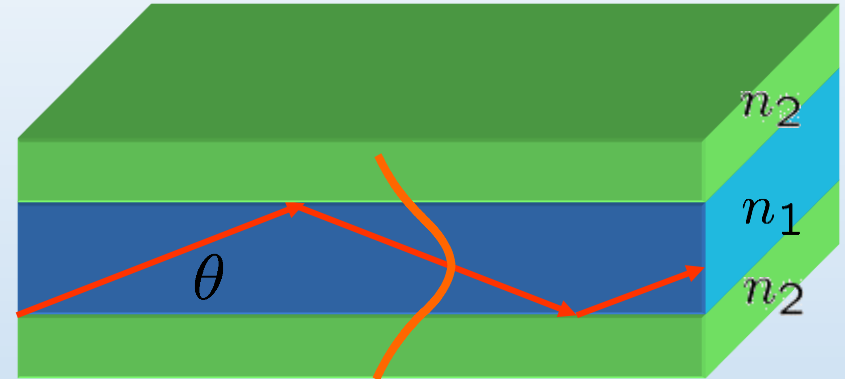
$$t = \frac{I_{\text{transmitted}}}{I_{\text{input}}} = \frac{(1 - R^2)^2}{(1 - R^2)^2 + 4R^2 \sin^2(kL)}$$

FP resonance in dielectric slab

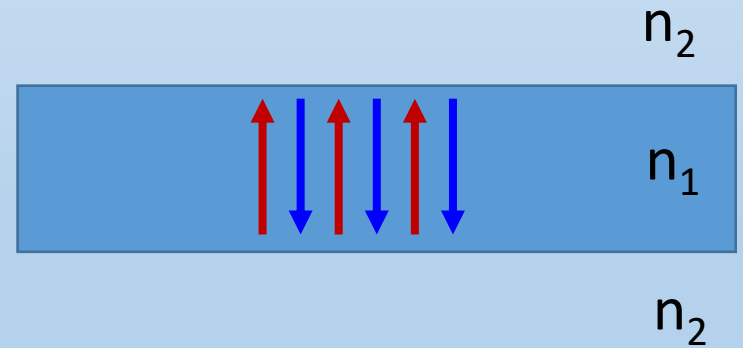
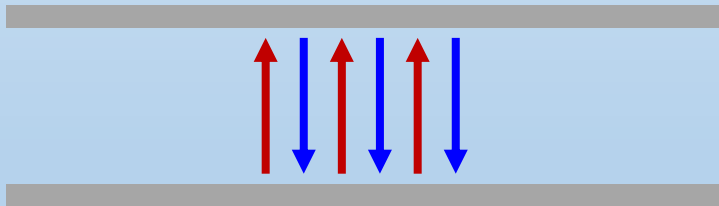
Metal waveguides



Dielectric waveguides



Resonance

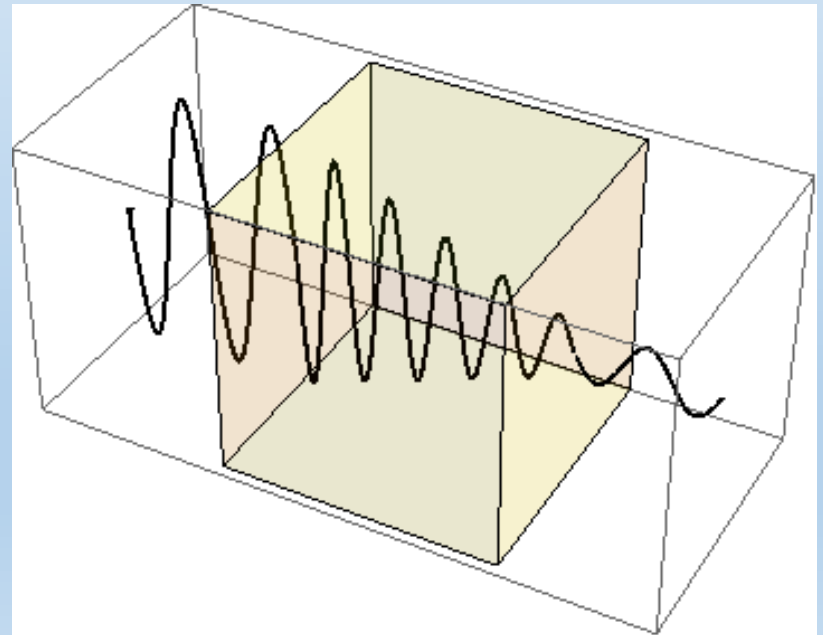
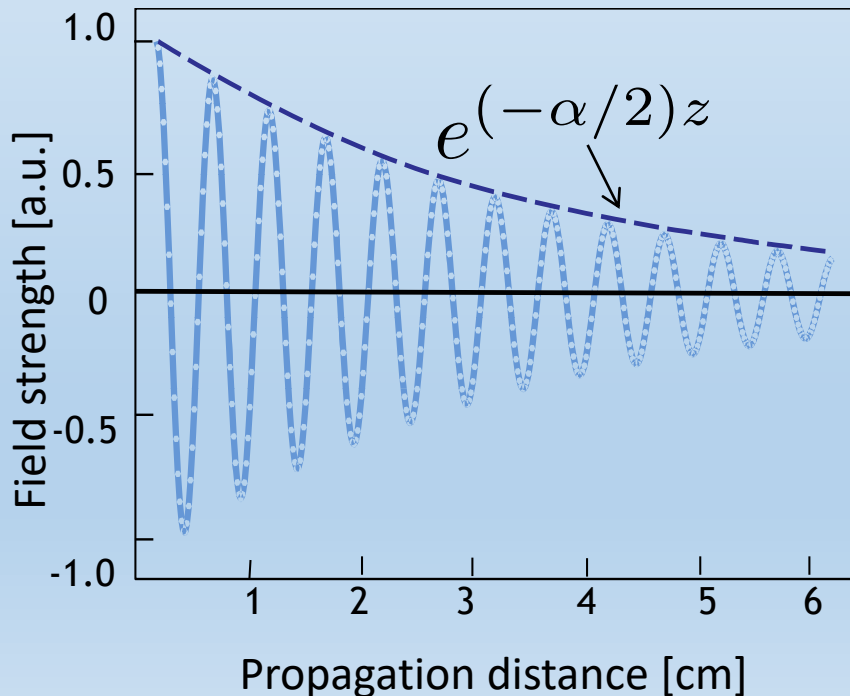


Plane Waves in Lossy Materials

Complex refractive index: $\tilde{n} = n - j\kappa$ $\kappa > 0$

Complex wave vector: $\tilde{k} = (n - j\kappa)k_0$ $\alpha = 2\kappa k_0$

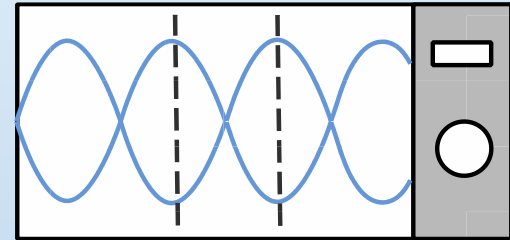
$$E_y(z, t) = \text{Re} \left\{ e^{-j(n-j\kappa)k_0 z} e^{j\omega t} \right\} = E_0 e^{-\alpha z/2} \cos(\omega t - nk_0 z)$$



Resonators with Internal Loss



Image is in the public domain



$$\frac{E_t}{E_i} = \frac{T^2 \exp(-jkL)}{1 - R^2 \exp(-j2kL)} = \frac{T^2 \exp(-\kappa k_0 L) \exp(-jnk_0 L)}{1 - R^2 \exp(-2\kappa k_0 L) \exp(-j2nk_0 L)}$$

...the EM wave loss is what heats the water inside the food

Laser Using Fabre-Perot Cavity

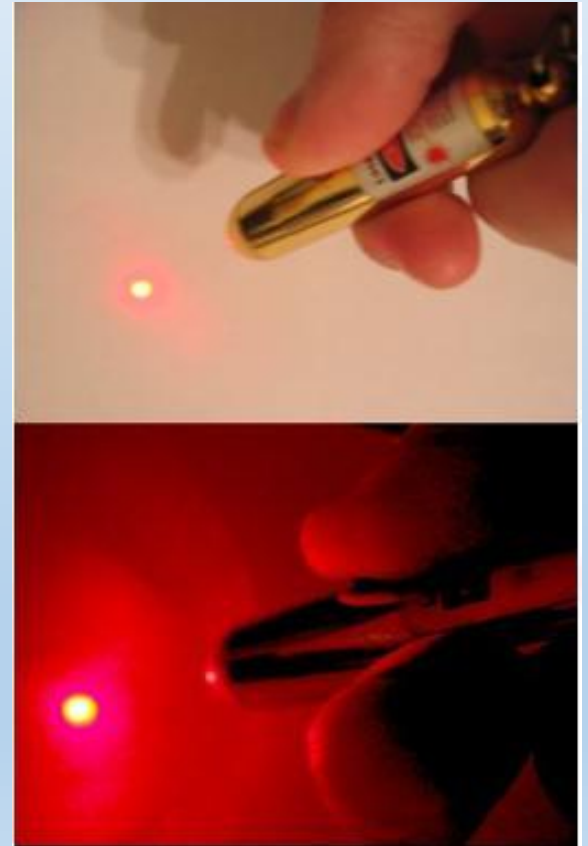
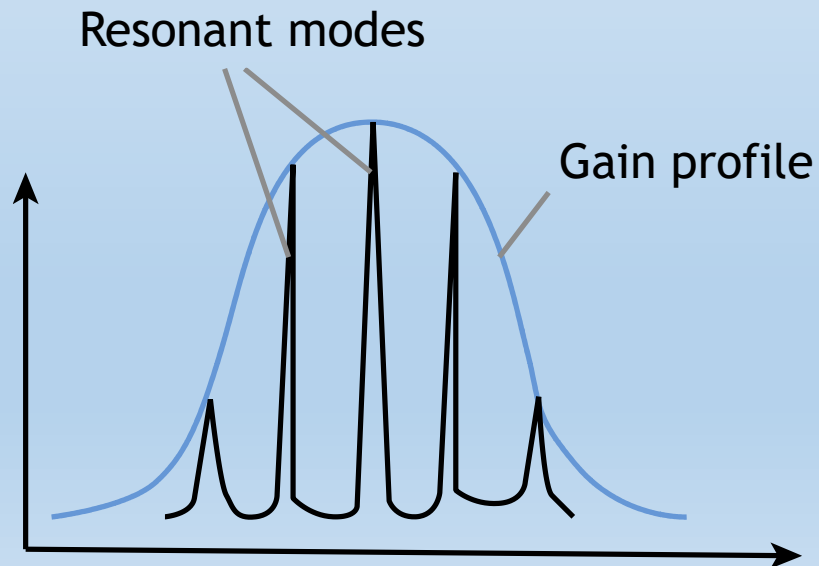
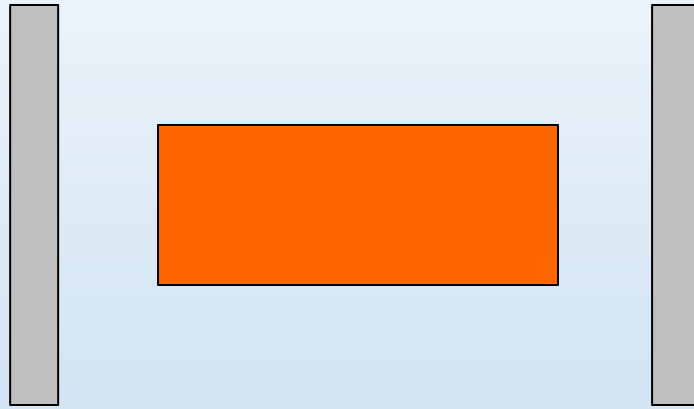


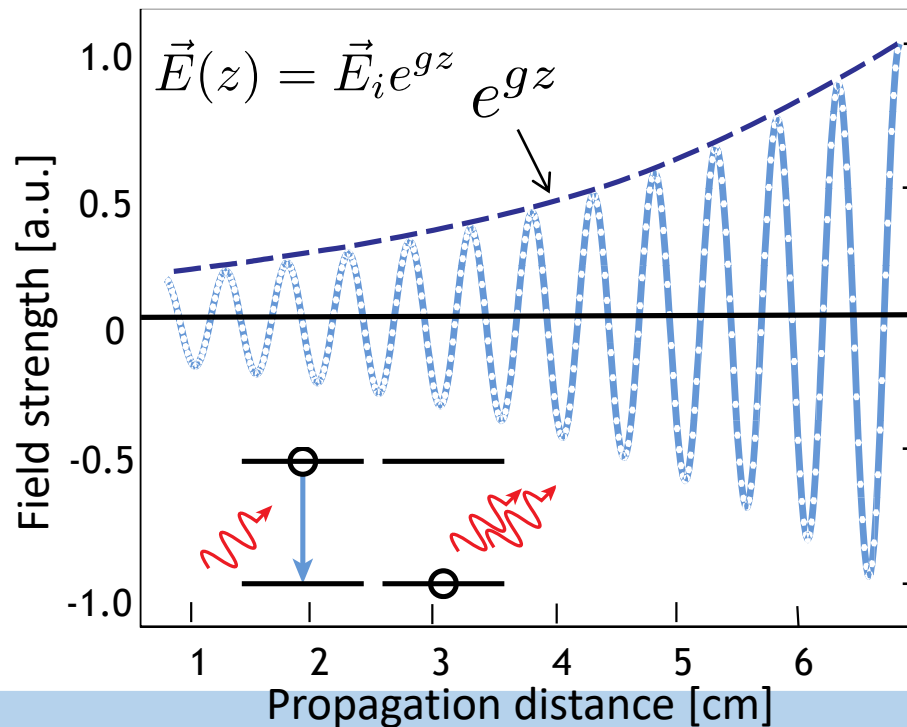
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Resonators with Internal Gain

What if it was possible to make a material with “negative absorption” so the field grew in magnitude as it passed through a material?

$$\tilde{n} = n - j\kappa$$

$$\kappa < 0$$



$$t = \frac{E_t}{E_i} = \frac{T^2 \exp(-\kappa k_0 L) \exp(-j n k_0 L)}{1 - R^2 \exp(-2\kappa k_0 L) \exp(-j 2 n k_0 L)}$$

Resonance:

$$e^{j 2 n k_0 L} = 1$$

Lasers: Something for Nothing (almost)

at resonance $e^{j2nk_0L} = 1$

$$\frac{E_t}{E_i} = \frac{T^2 \exp(-\kappa k_0 L) \exp(-jnk_0 L)}{1 - R^2 \exp(-2\kappa k_0 L) \exp(-j2nk_0 L)}$$

singularity at

$$1 = R^2 \exp(-2\kappa k_0 L) \exp(-j2nk_0 L)$$

$$\frac{E_t}{E_i} \rightarrow \infty$$

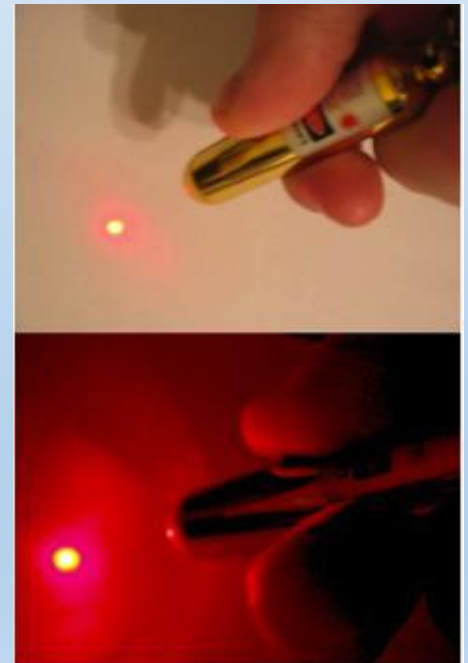
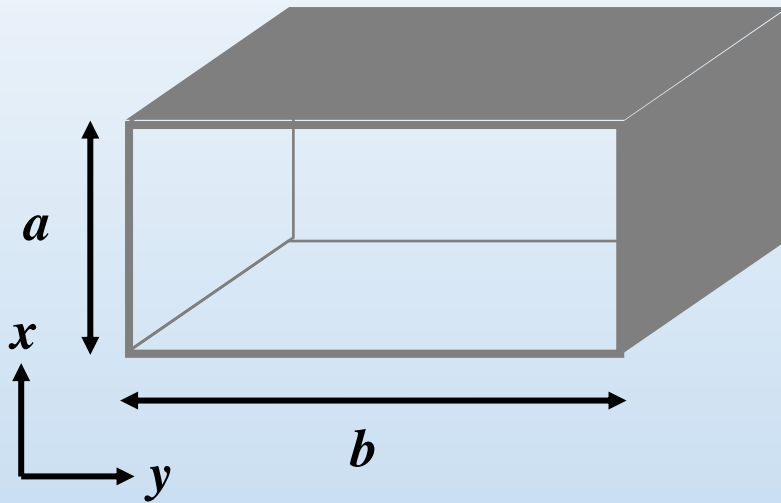


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Rectangular Waveguide - TE Modes



Boundary Conditions:

- (1) $E_x = 0$ at $y = 0$ and b
- (2) $E_y = 0$ at $x = 0$ and a

$$k_x a = m\pi$$

$$k_y b = n\pi$$

(Guidance Condition)

$$H_z = \cos(k_x x) \cos(k_y y) e^{-jk_z z}$$

$$H_x = \frac{jk_x k_z}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) e^{-jk_z z}$$

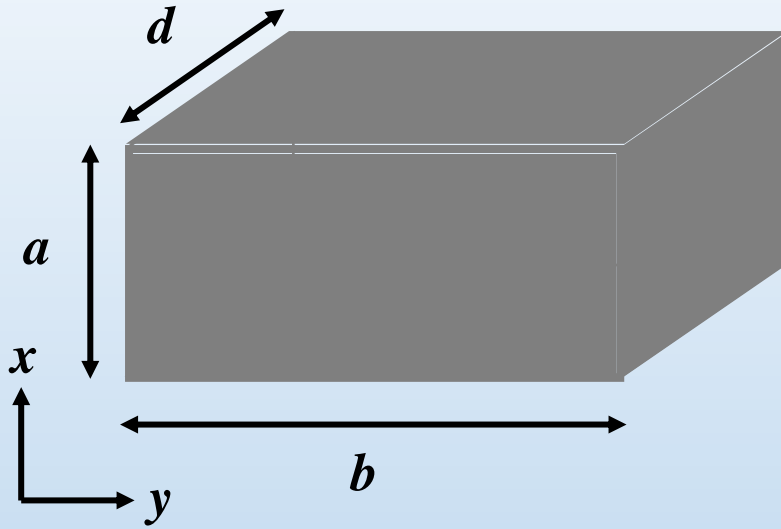
$$H_y = \frac{jk_y k_z}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) e^{-jk_z z}$$

$$E_x = \frac{j\omega \mu k_y}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) e^{-jk_z z}$$

$$E_y = \frac{-j\omega \mu k_x}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) e^{-jk_z z}$$

$$k_z = \sqrt{\left(\frac{\omega}{c}\right)^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}$$

Rectangular Cavity Resonator - TE Modes



Boundary Conditions:

- (1) $E_x = 0$ at $y = 0$ and b
- (2) $E_y = 0$ at $x = 0$ and a

$$k_x a = m\pi$$

$$k_y b = n\pi$$

$$k_z d = p\pi$$

(Resonance Condition)

$$H_z = \cos(k_x x) \cos(k_y y) \sin(k_z z)$$

$$H_x = \frac{-k_x k_z}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) \cos(k_z z)$$

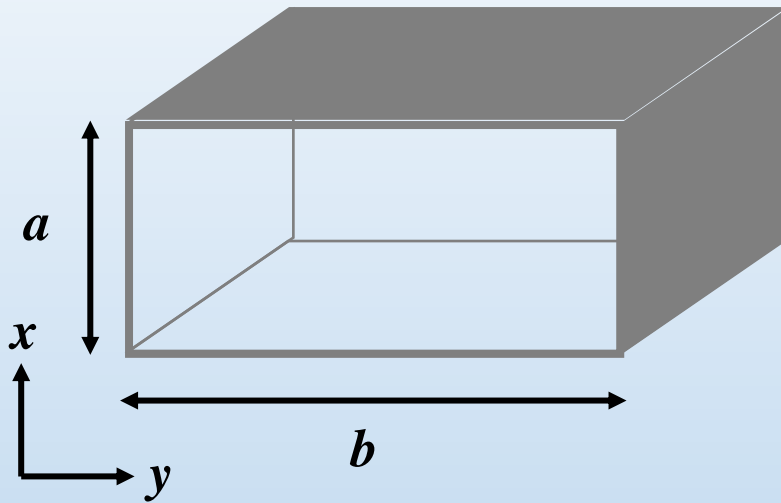
$$H_y = \frac{-k_y k_z}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) \cos(k_z z)$$

$$E_x = \frac{j\omega\mu k_y}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) \sin(k_z z)$$

$$E_y = \frac{-j\omega\mu k_x}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) \sin(k_z z)$$

$$k_r = \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 + \left(\frac{p\pi}{d}\right)^2}$$

Rectangular Waveguide - TM Modes



Boundary Conditions:

- (1) $E_x = 0$ at $y = 0$ and b
- (2) $E_y = 0$ at $x = 0$ and a

$$k_x a = m\pi$$

$$k_y b = n\pi$$

(Guidance Condition)

$$E_z(x, y, z) = \sin(k_x x) \sin(k_y y) e^{-jk_z z}$$

$$E_x = \frac{-jk_x k_z}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) e^{-jk_z z}$$

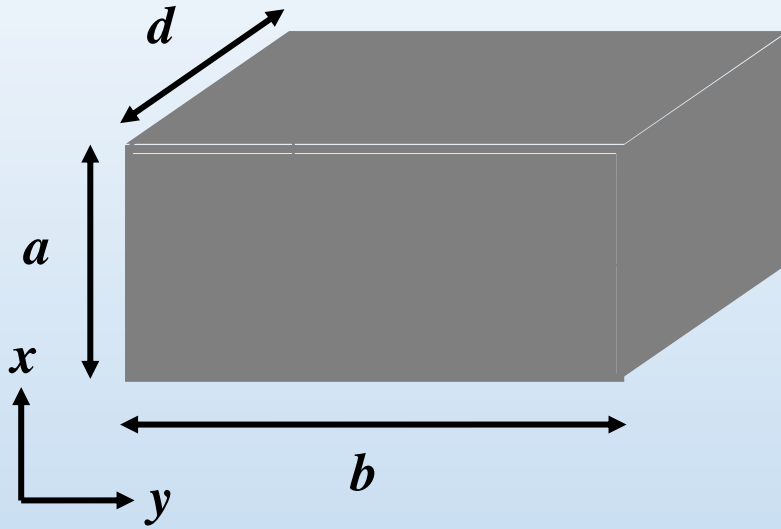
$$E_y = \frac{-jk_y k_z}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) e^{-jk_z z}$$

$$H_x = \frac{j\omega \epsilon k_y}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) e^{-jk_z z}$$

$$H_y = \frac{-j\omega \epsilon k_x}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) e^{-jk_z z}$$

$$k_z = \sqrt{\left(\frac{\omega}{c}\right)^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}$$

Rectangular Cavity Resonator - TM Modes



Boundary Conditions:

- (1) $E_x = 0$ at $y = 0$ and b
- (2) $E_y = 0$ at $x = 0$ and a

$$k_x a = m\pi$$

$$k_y b = n\pi$$

$$k_z d = p\pi$$

(Resonance Condition)

$$E_z(x, y, z) = \sin(k_x x) \sin(k_y y) \cos(k_z z)$$

$$E_x = \frac{-k_x k_z}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) \sin(k_z z)$$

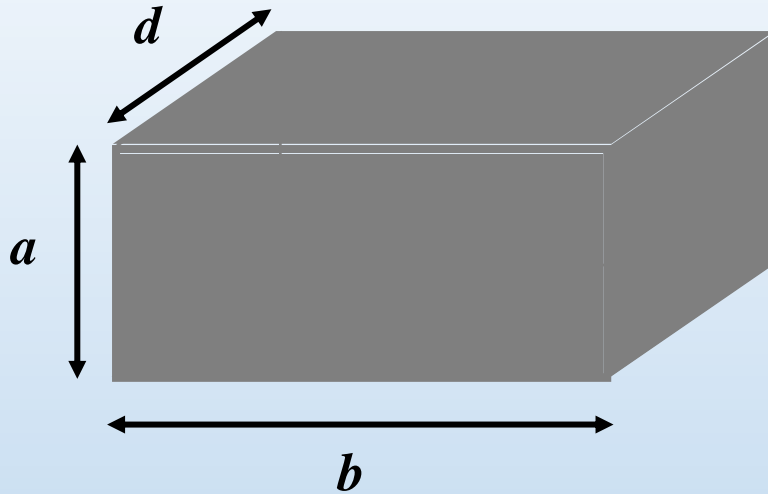
$$E_y = \frac{-k_y k_z}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) \sin(k_z z)$$

$$H_x = \frac{j\omega \epsilon k_y}{\omega^2 \mu \epsilon - k_z^2} \sin(k_x x) \cos(k_y y) \cos(k_z z)$$

$$H_y = \frac{-j\omega \epsilon k_x}{\omega^2 \mu \epsilon - k_z^2} \cos(k_x x) \sin(k_y y) \cos(k_z z)$$

$$k_r = \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 + \left(\frac{p\pi}{d}\right)^2}$$

Example



When the resonator dimensions are such that $a > b > d$, the lowest resonant spatial frequency (k) is

$$k_r = \sqrt{\left(\pi/a\right)^2 + \left(\pi/b\right)^2}$$

$$(m = n = 1, p = 0)$$

TM₁₁₀ mode:

$$E_z(x, y, z) = \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{b}\right)$$

$$H_x = \frac{j\pi k_y}{\omega\mu b} \sin\left(\frac{\pi x}{a}\right) \cos\left(\frac{\pi y}{b}\right)$$

$$E_x = E_y = 0$$

$$H_y = \frac{-j\pi}{\omega\mu a} \cos\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{b}\right)$$

TM₁₁₀ mode:

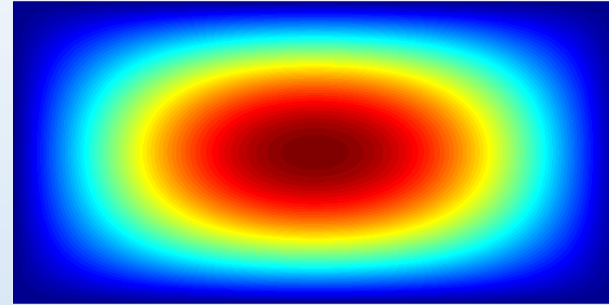
$$k_r = \sqrt{\left(\pi/a\right)^2 + \left(\pi/b\right)^2}$$

$$E_z(x, y, z) = \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{b}\right)$$

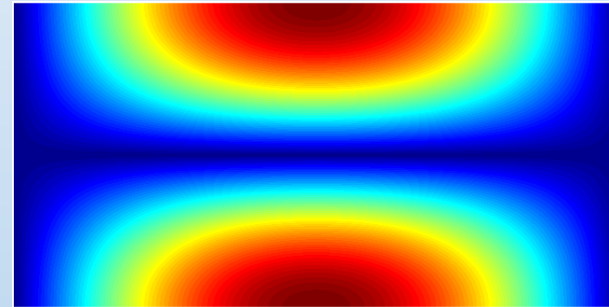
$$H_x = \frac{j\pi k_y}{\omega\mu b} \sin\left(\frac{\pi x}{a}\right) \cos\left(\frac{\pi y}{b}\right)$$

$$H_y = \frac{-j\pi}{\omega\mu a} \cos\left(\frac{\pi x}{a}\right) \sin\left(\frac{\pi y}{b}\right)$$

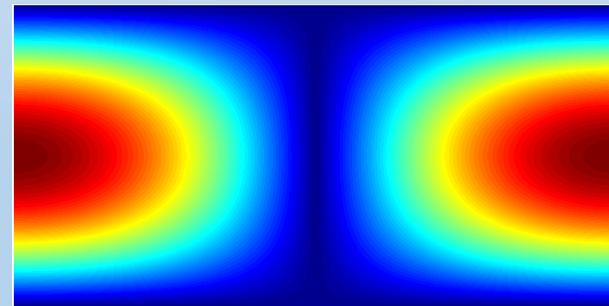
$$f_r = \frac{c_0 k_r}{2\pi n} = \frac{1}{2n} \sqrt{(c_0/a)^2 + (c_0/b)^2}$$



E_z



H_x



H_y