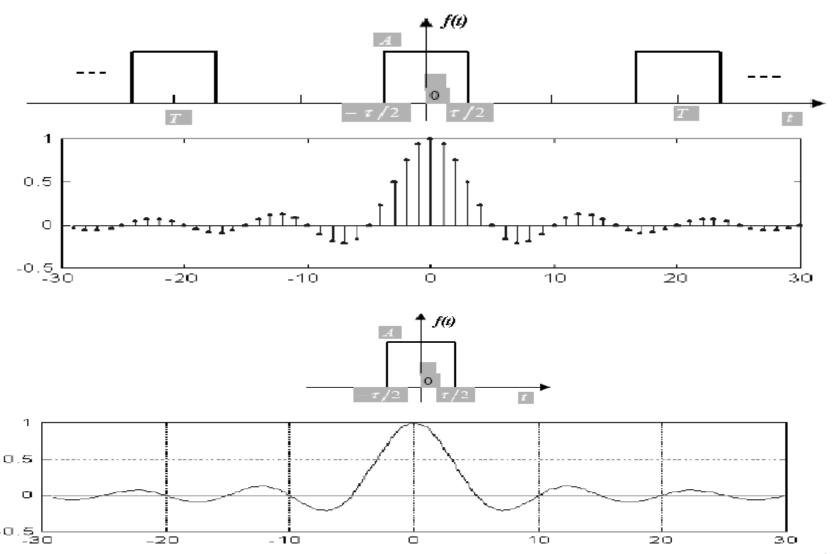


图形示意:周期——非周期信号



周期信号:
$$x(t) = \sum_{n=-\infty}^{\infty} F_n e^{in\omega_1 t}$$

其中:
$$F_n = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} \tilde{x}(t) e^{-in\omega_1 t} dt$$

$$= F(n\omega_1)$$

τ固定,B或B_f不变;

T变大时:

- ω₁减小, 谱线稠密
 - |F_n|谱线振幅减小,

收敛速率缓慢。

周期信号x(t) → 非周期信号x(t)

$$\omega_1 \to d\omega$$

离散谱 $n\omega_1 \rightarrow$ 连续谱 ω

$$\sum \rightarrow \int$$

谱线高度 $|F_n| \rightarrow 无穷小量$

Shandong Univers 日文此元穷小量间的比例关系不变

频谱密度函数

为了表明无穷小的振幅间的相对差别,有必要引入

一个新的量——称为"频谱密度函数"。

$$X(\omega) \stackrel{def}{=} \lim_{f_1 \to 0} \frac{F(n\omega_1)}{f_1} = \lim_{\omega_1 \to 0} \frac{2\pi F(n\omega_1)}{\omega_1}$$

$$= \lim_{T_1 \to \infty} T_1 F(n\omega_1) = \lim_{T_1 \to \infty} T_1 \frac{1}{T_1} \int_{-\frac{T_1}{2}}^{\frac{T_1}{2}} \tilde{x}(t) e^{-in\omega_1 t} dt$$

$$= \int_{-\infty}^{\infty} x(t) e^{-i\omega t} dt$$

$X(\omega)$ 的其它表达方式 $X(j\omega)$ 或 $X(e^{j\omega})$

$$X(\omega) \stackrel{def}{=} \lim_{f_1 \to 0} \frac{F(n\omega_1)}{f_1}$$

$$x(t) = \lim_{T \to \infty} \tilde{x}(t) = \lim_{T \to \infty} \sum_{n=\infty}^{\infty} F_n e^{in\omega_1 t}$$

$$= \lim_{f_1 \to 0} \sum_{n = -\infty}^{\infty} \frac{F(n\omega_1)}{f_1} e^{in\omega_1 t} f_1 = \int_{-\infty}^{\infty} X(2\pi f) e^{i2\pi f t} df$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{i\omega t} d\omega$$

$$F($$

$$F(\omega) = \mathscr{F}[f(t)]$$

$$f(t) = \mathscr{F}^{-1}[F(\omega)]$$

意义:任意非周期信号都可以分解为无数频率连续 可简的分布、振幅为的 $\frac{1}{2\pi}X(\omega)d\omega$ 复指数信号 $e^{i\omega t}$ 的

月刊 心力
$$f(t) \longleftrightarrow F(\omega)$$

$$X(\omega) = \int_{-\infty}^{+\infty} x(t)e^{--i\omega t}dt$$
 — — 傅立叶正变换
$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} X(\omega)e^{i\omega t}d\omega$$
 — 傅立叶逆变换

说明:

实部 虚部

1、复函数
$$F(\omega) = |F(\omega)|e^{i\phi(\omega)} = R(\omega) + iX(\omega)$$

$$|F(\omega)| = \sqrt{R^2(\omega) + X^2(\omega)}$$



相位
$$\phi(\omega) = \arctan \frac{X(\omega)}{R(\omega)}$$

Odd

Even

幅度频谱、幅度特性或幅度谱

相位频谱、相位特性或相位谱

实信号

$$f(t) = f_{\rm e}(t) + f_{\rm o}(t)$$

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt$$

$$= \int_{-\infty}^{\infty} [f_{e}(t) + f_{o}(t)] \cdot [\cos \omega t - i \sin \omega t] dt$$

$$=2\int_0^\infty f_{\rm e}(t)\cos\omega\,t\,{\rm d}\,t-i2\int_0^\infty f_{\rm o}(t)\sin\omega\,t\,{\rm d}\,t$$

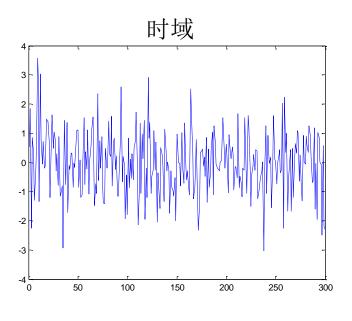


虚部,奇函数

注意: f(t)是实信号

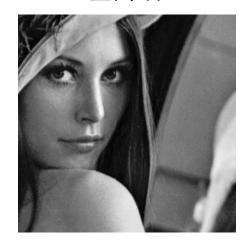


讨论理解 每一步推 导原理

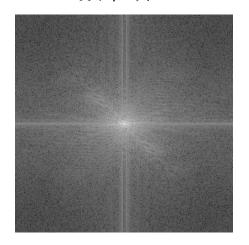


频率域
45
40
35
30
25
20
15
10
5
0
5
10
15
0
20
250
300

空间域



频率域

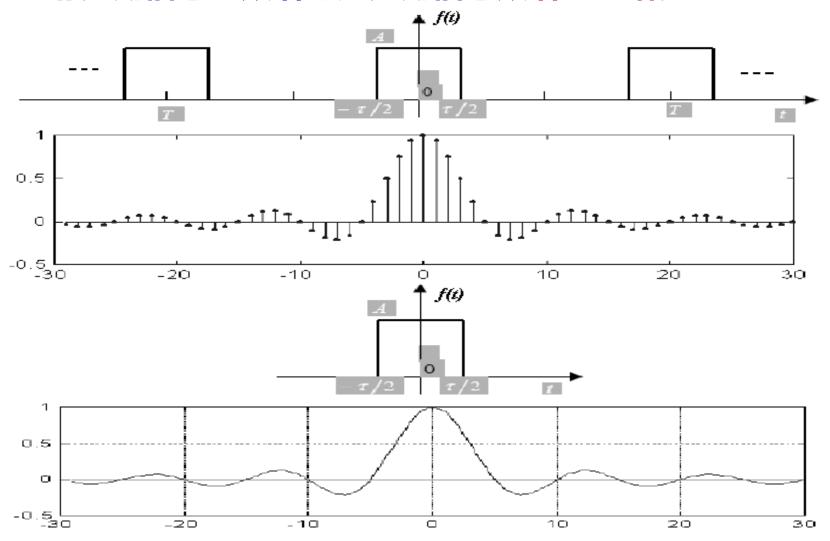


$$2, f(t) \leftarrow$$
 充分条件: Dirichlet条件 $\to F(\omega)$

3、物理意义: 非周期信号可以分解成无数连续频率分布、振幅为 $\frac{1}{2\pi}F(\omega)d\omega$ 的复指数信号之线性组合。

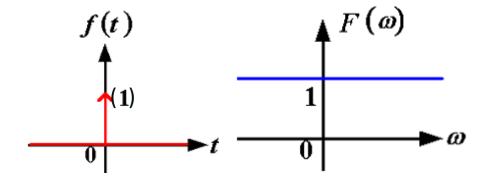
$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{i\omega t} d\omega$$

4、周期信号的频谱对应非周期信号频谱的样本; 非周期信号的频谱对应周期信号频谱的包络。



3.2.2典型非周期信号的领谱

1、单位冲激信号



$$f(t) = \delta(t) \leftrightarrow 1$$

$$F(\omega) = \int_{-\infty}^{+\infty} \delta(t) e^{-i\omega t} dt = 1$$

$$\delta(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{i\omega t} d\omega$$



阅读p10~18, 讨论 学习典型非周期信 号的频谱推导方法

尖脉冲干扰

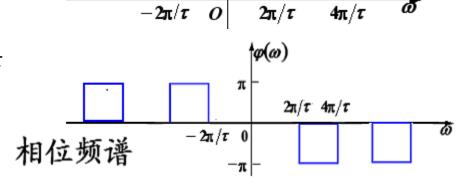
2、门函数Gτ(t)

$$F(\omega) = \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} e^{-i\omega t} dt = \frac{1}{-i\omega} e^{-i\omega t} \Big|_{-\frac{\tau}{2}}^{\frac{\tau}{2}} = \frac{\tau}{\omega \frac{\tau}{2}} \cdot \frac{e^{i\omega \frac{\tau}{2}} - e^{-i\omega \frac{\tau}{2}}}{2i}$$

$$= \tau \operatorname{Sa}\left(\frac{\omega \tau}{2}\right)$$
频谱带宽近似为 $B_{\omega} = \frac{2\pi}{\tau} (rad/s)$

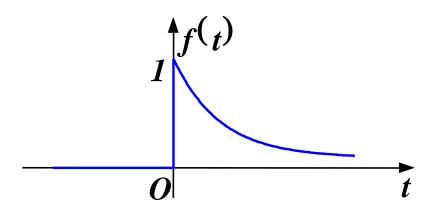
$$B_{f} = \frac{1}{\tau} (Hz)$$

对相位频谱的说明: 对于傅里叶变换为实数的相位频谱,由于相位为 π 或 $-\pi$ 均可以使幅值取反,所以在幅值为负值的频率范围内相位可以为 π 或 $-\pi$ 。但由于需要与相频特性为奇函数这一性质一致,需要在频率正半轴与负半轴分别取 $-\pi$ 和 π 。 Shandong University



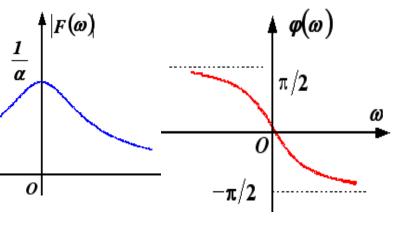
3、单边指数信号

$$f(t) = e^{-\alpha t} u(t) \quad \alpha > 0$$



$$F(\omega) = F[f(t)] = \int_0^\infty e^{-(\alpha + i\omega)t} dt = \frac{1}{\alpha + i\omega}$$

$$e^{-\alpha t}u(t)\leftrightarrow \frac{1}{\alpha+i\omega}$$



幅度频谱:

$$|F(\omega)| = \frac{1}{\sqrt{\alpha^2 + \omega^2}}$$
$$\phi(\omega) = -\arctan\frac{\omega}{\alpha}$$

相位频谱:

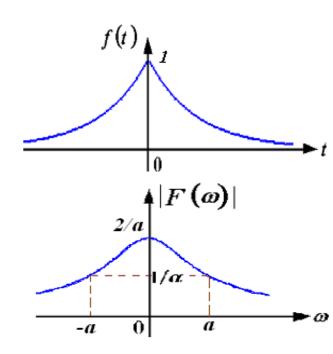
$$\phi(\omega) = -\arctan\frac{\omega}{\alpha}$$

4、双边指数信号

$$f(t) \stackrel{a>0}{=} e^{-a|t|} = e^{-at}u(t) + e^{at}u(-t)$$

$$F(\omega) = \frac{1}{a + i\omega} + \int_{-\infty}^{0} e^{\alpha t} e^{-i\omega t} dt$$

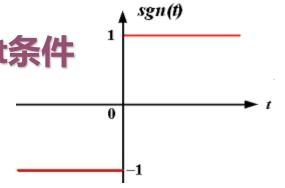
$$= \frac{1}{\alpha + i\omega} + \frac{1}{\alpha - i\omega} = \frac{2\alpha}{\alpha^2 + \omega^2}$$

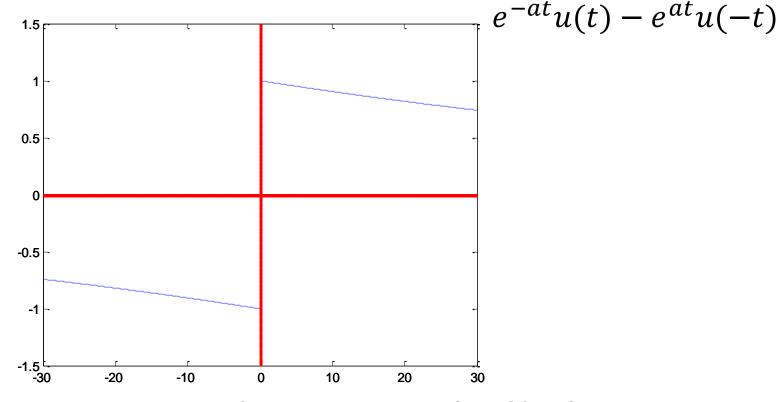


5、符号函数

不符合Dirichlet条件

$$f(t) = \operatorname{sgn}(t) = \begin{cases} +1 & (t > 0) \\ -1 & (t < 0) \end{cases}$$





5、符号函数

不符合Dirichlet条件

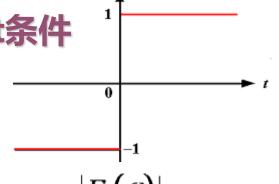
$$f(t) = \operatorname{sgn}(t) = \begin{cases} +1 & (t > 0) \\ -1 & (t < 0) \end{cases}$$

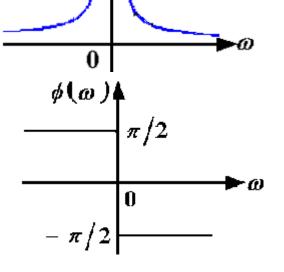
$$= \lim_{a \to 0} [e^{-at}u(t) - e^{at}u(-t)] \quad (a > 0)$$

$$=\frac{2}{i\omega}$$

$$|F(\omega)| = \frac{2}{\omega}$$
 $\phi(\omega) = \begin{cases} -\frac{\pi}{2} & (\omega > 0) \\ \frac{\pi}{2} & (\omega < 0) \end{cases}$ $-\pi/2$

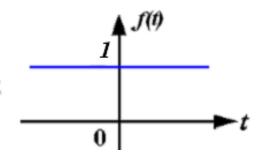
$$+\frac{\pi}{2}$$
 $(\omega < 0)$





6、直流信号

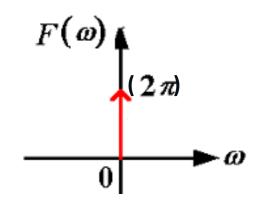
不符合Dirichlet条件



双边指数信号:
$$F(\omega) = \frac{2\alpha}{\alpha^2 + \omega^2}$$
, 当 α 趋近于零时, $\lim_{\alpha \to 0} \frac{2\alpha}{\alpha^2 + \omega^2}$ = $\begin{cases} 0, & \omega \neq 0 \\ \infty, & \omega = 0 \end{cases}$, 为冲击信号, 其强度为 $\lim_{\alpha \to 0} \int_{-\infty}^{\infty} \frac{2\alpha}{\alpha^2 + \omega^2} d\omega = 2\pi$

$$1 \leftrightarrow 2\pi\delta(\omega)$$

或
$$F^{-1}[\delta(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta(\omega) e^{i\omega t} d\omega$$
$$= \frac{1}{2\pi}$$



7、冲激函数的导数

$$\int_{-\infty}^{\infty} \delta'(t) f(t) dt = -f'(0)$$

$$F[\delta'(t)] = \int_{-\infty}^{\infty} \delta'(t) e^{-i\omega t} dt = i\omega$$

$$\delta'(t) \leftrightarrow i\omega$$

$$\delta^{(n)}(t) \leftrightarrow (i\omega)^n$$

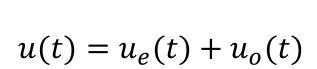
$$F^{-1}[\delta'(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} \delta'(\omega) e^{i\omega t} d\omega = \frac{1}{2\pi} (-it) = \frac{t}{i2\pi}$$

$$t \leftrightarrow i2\pi\delta'(\omega)$$

$$t^n \leftrightarrow i^n 2\pi \delta^{(n)}(\omega)$$

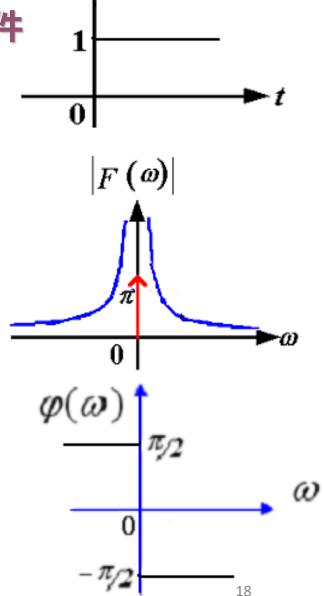
8、单位阶跃函数

不符合Dirichlet条件



$$=\frac{1}{2}+\frac{1}{2}\operatorname{sgn}(t)$$

$$F[u(t)] = \pi \delta(\omega) + \frac{1}{i\omega}$$



u(t)

表 常用傅里叶变换对

编号	f(t)	$F(j\omega)$
1	$g_{ au}(t)$	$ au$ Sa $\left(\frac{\omega au}{2}\right)$
2	$ au \operatorname{Sa}\left(\frac{ au t}{2}\right)$	$2\pi g_{\tau}(\omega)$
3	$e^{-\alpha t} \varepsilon(t), \alpha > 0$	$\frac{1}{\alpha+j\omega}$
4	$te^{-\alpha t}\varepsilon(t), \alpha > 0$	$\frac{1}{(\alpha+\mathrm{j}\omega)^2}$
5	$e^{-\alpha t }, \alpha > 0$	$\frac{2\alpha}{\alpha^2 + \omega^2}$
6	$\delta(t)$	1
7	1	$2\pi\delta(\omega)$
8	$\delta(t-t_0)$	e ^{-jat} 0
9	$\cos \omega_0 t$	$\pi\delta(\omega-\omega_{\scriptscriptstyle 0})+\pi\delta(\omega+\omega_{\scriptscriptstyle 0})$
10	$\sin \omega_0 t$	$\frac{\pi}{j} [\delta(\omega - \omega_0) - \delta(\omega + \omega_0)]$

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续表

编号	f(t)	$F(\mathrm{j}\omega)$
11	$\varepsilon(t)$	$\pi\delta(\omega) + \frac{1}{\mathrm{j}\omega}$
12	$\operatorname{Sgn}(t)$	$\frac{2}{\mathrm{j}\omega}, F(0)=0$
13	$\frac{1}{\pi t}$	-j Sgn(ω)
14	$\delta_T(t)$	$\Omega\delta_{\Omega}(\omega)$
15	$\sum_{n=\infty}^{\infty} F_n e^{jn\Omega t}$	$2\pi\sum_{n=-\infty}^{\infty}F_{n}\delta(\omega-n\Omega)$
16	$\frac{t^{n-1}}{(n-1)!}e^{-at}\varepsilon(t), a>0$	$\frac{1}{(a+\mathrm{j}\omega)^n}$

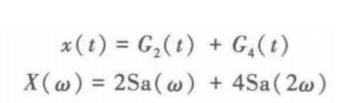
记住该表中傅立叶变换对(16可以不记)

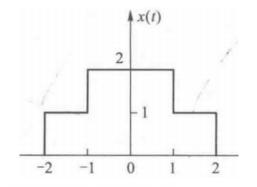
3.2.3 傅里叶变换的性质

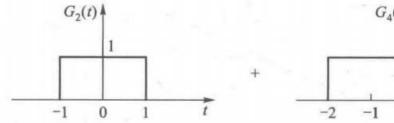
1、线性性质 设 $f(t) \leftrightarrow F(\omega)$ $af_1(t) + bf_2(t) \leftrightarrow aF_1(\omega) + bF_2(\omega)$

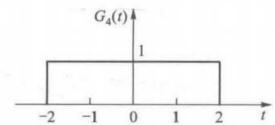
p140 例 3-8 求图 3-27 所示信号 x(t) 的频谱函数。

解:信号x(t)可以分解为两个宽度不同的门函数











2、对称性
$$F(t) \leftrightarrow 2\pi f(-\omega)$$

证明:
$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{i\omega t} d\omega$$

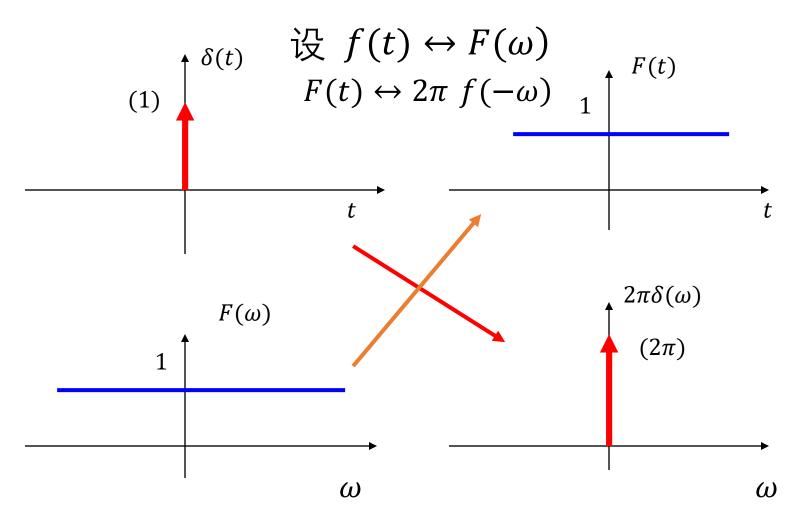
$$-t \to t$$
 $f(-t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{-i\omega t} d\omega$

$$t \leftrightarrow \omega$$
 $f(-\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(t) e^{-i\omega t} dt$

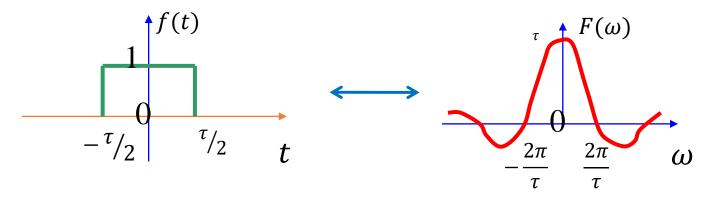
$$2\pi f(-\omega) = \int_{-\infty}^{\infty} F(t) e^{-i\omega t} dt$$

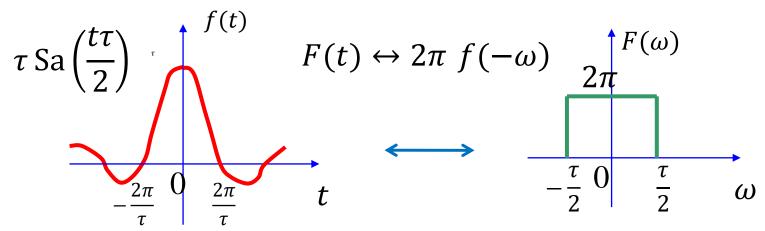
对称性举例

直流信号和冲激函数



例: 己知
$$G_{\tau}(t) \longleftrightarrow \tau \operatorname{Sa}\left(\frac{\omega \tau}{2}\right)$$





$$\tau \operatorname{Sa}\left(\frac{t\tau}{2}\right) \longleftrightarrow 2\pi G_{\tau}(-\omega) = 2\pi G_{\tau}(\omega)$$

求
$$FT$$
:

$$f_1(t) = \frac{1}{t}.$$

解:

3. 尺度变换
$$f(at) \leftrightarrow \frac{1}{|a|} F(\frac{\omega}{a})$$
 $(a \neq 0 常数)$

IIIII:
$$F[f(at)] = \int_{-\infty}^{\infty} f(at) e^{-i\omega t} dt$$

$$F[f(at)] = \frac{1}{a} \int_{-\infty}^{\infty} f(x)e^{-i\frac{\omega}{a}x} dx = \frac{1}{a}F(\frac{\omega}{a})$$

当a<0.

$$F[f(at)] = \frac{1}{a} \int_{-\infty}^{\infty} f(x)e^{-i\frac{\omega}{a}x} dx = \frac{1}{-a}F(\frac{\omega}{a})$$



理解推 导讨程

求FT:

$$\mathrm{Sa}(\omega_c t)$$

 $\omega_c \operatorname{Sa}\left(\frac{\omega_c t}{2}\right) \longleftrightarrow 2\pi G_{\omega_c}(\omega)$

$$\frac{\pi}{\omega_c} F(\omega)$$

$$-\omega_c^0 \omega_c \omega$$

$$\operatorname{Sa}\left(\frac{\omega_c t}{2}\right)$$

$$\operatorname{Sa}\left(\frac{\omega_c t}{2}\right) \longleftrightarrow \frac{2\pi}{\omega_c} G_{\omega_c}(\omega) \qquad f(t) \leftrightarrow F(\omega)$$

$$f(t) \leftrightarrow F(\omega)$$

$$Sa(\omega_c t) = Sa\left(2\frac{\omega_c t}{2}\right)$$
 根据尺度变换性质: 这里a=2

$$\longrightarrow \frac{2\pi}{|2|\omega_c} G_{\omega_c}(\frac{\omega}{2}) = \frac{\pi}{\omega_c} G_{\omega_c}(\frac{\omega}{2}) = \frac{\pi}{\omega_c} G_{2\omega_c}(\omega)$$

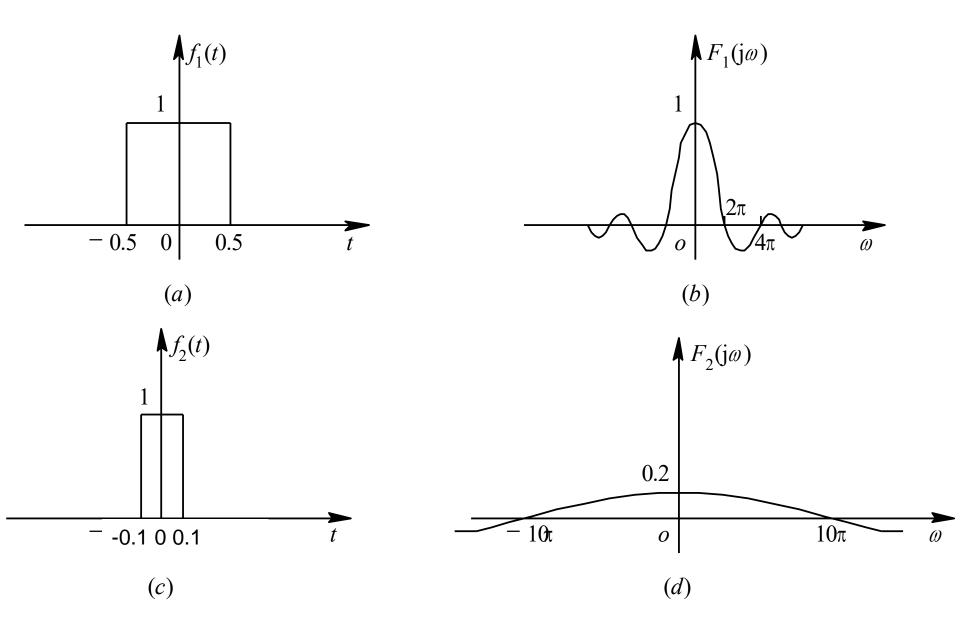
1.
$$a = -1$$
时,

1.
$$a = -1$$
时, $f(-t) \leftrightarrow F(-\omega)$

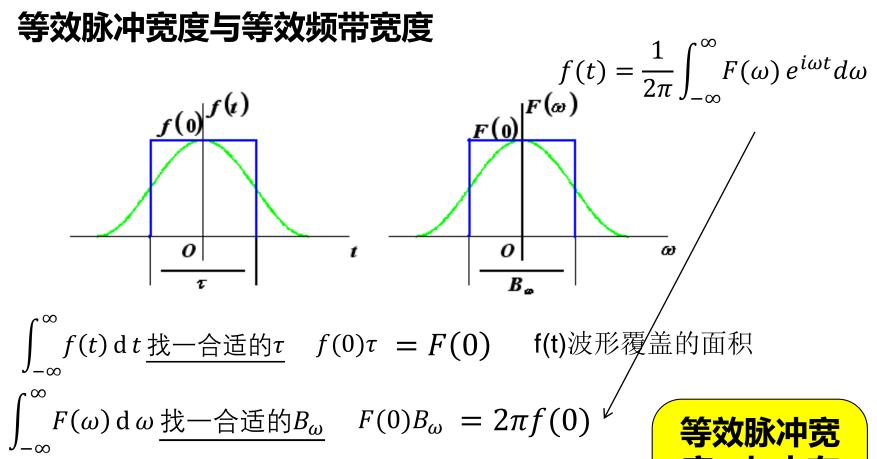
$$f(at) \leftrightarrow \frac{1}{|a|} F(\frac{\omega}{a}) \quad (a \neq 0 \text{ \text{\text{≥ 0}}})$$

compressed in time \Leftrightarrow stretched in frequency





3. 等效脉冲宽度与等效频带宽度



$$\int_{-\infty}^{\infty} F(\omega) \, \mathrm{d} \, \omega \, \underline{\mathfrak{H}} - \underline{\mathrm{Gibh}}_{\underline{\omega}} \quad F(0) B_{\omega} = 2\pi f(0)$$

$$B_{\omega} = 2\pi \frac{f(0)}{F(0)} = \frac{2\pi}{\tau}$$
 $B_f = \frac{f(0)}{F(0)} = \frac{1}{\tau}$

等效脉冲宽 的等效带宽 B_{α} 成反比。

(Pulse width in t)•(Pulse width in f)=1 Uncertainty Principle!