



STATISTICS OF BURSTY EVENTS

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Abstract

The prediction of extreme events resulting from human behaviour is becoming significant in areas as far reaching as the control of disease spread, resource allocation, and emergency response. Classical methods in extreme value theory fail to recognise that most human created events do not occur uniformly, but in bursts. In this thesis, we aim to design statistical methods for the inference and prediction of Continuous Time Random Maxima (CTRM).

Contents

Chapter 1	Introduction	1
Chapter 2	Preliminaries	2
2.1	Probability and Measure Theory	2
2.2	Stochastic Processes	3
2.3	Weak Convergence	4
Chapter 3	Scaling Limit Theorems	6
3.1	Problem Formulation	6
3.2	Background	8
3.3	Stable Distributions	11
3.4	Scaling Limit of the Sum of Waiting Times	14
3.5	Scaling Limit of the Maximum of Event Magnitudes	16
3.6	Scaling Limit of CTRM	16
Chapter 4	Distribution of CTRM	17
4.1	Distribution of Exceedances	17
4.2	Distribution of Exceedance Durations	17
Chapter 5	Statistical Inference	18
5.1	Log Moment Estimator for Mittag-Leffler Random Variables	18
5.2	Maximum Likelihood Estimator for Generalised Pareto Random Variables	18
5.3	Stability of Parameter Estimates	18
5.4	Goodness of Fit	18
Chapter 6	Model Fitting	19
6.1	Model Fit of Simulated Data	19
6.2	Model Fit of Real Data	20
Chapter 7	Conclusion	21
	References	22

CHAPTER 1

Introduction

Time series displaying inhomogeneous behaviour have received strong interest in the recent statistical physics literature, [Bar05, OB05, VOD⁺06, VRLB07, OS11, MGV11, KKP⁺11, BB13], and have been observed in the context of earthquakes, sunspots, neuronal activity, human communication etc., see [KKBK12, VTK13] for a list of references. Such time series exhibit high activity in some ‘bursty’ intervals, which alternate with other, quiet intervals. Although several mechanisms are plausible explanations for bursty behaviour (most prominently self-exciting point processes [Haw71]), there seems to be one salient feature which very typically indicates the departure from temporal homogeneity: A heavy-tailed distribution of waiting times [VOD⁺06, KKBK12, VTK13]. A simple renewal process with heavy-tailed waiting times captures these dynamics. For many systems, the renewal property is appropriate, as can be checked by a simple test: the dynamics do not change significantly if the waiting times are randomly reshuffled [KKBK12].

When a magnitude can be assigned to each event in the renewal process, such as for earthquakes, sun flares, neuron voltages or the impact of an email, two natural and basic questions to ask are: What is the distribution of the largest event up to a given time t ? What is the probability that an event exceeds a given level ℓ within the next t units of time? A probabilistic extreme value model which assumes that the events form a renewal process is available in the literature. This model has been studied under the names “Continuous Time Random Maxima process” (CTRM) [BSM07, ?, HS15, ?], “Max-Renewal process” [Sil02, ST04, BŠ14], and “Shock process” [EM73, SS83, SS84, SS85, And87, Gut99]. This article aims to develop concrete statistical inference methods for this model, a problem which has seemingly received little attention by the statistical community.

CHAPTER 2

Preliminaries

In this chapter, we introduce the notation and review the background theory of probability and stochastic processes that are necessary in developing the statistical methods outlined in this thesis.

2.1 Probability and Measure Theory

Definition 2.1.1. Let Ω be a nonempty set and let \mathcal{F} be a collection of subsets of Ω . We say that \mathcal{F} is a **σ -algebra** if it satisfies the following conditions

1. The empty set $\emptyset \in \mathcal{F}$,
2. If $A \in \mathcal{F}$ then the complement $A^c \in \mathcal{F}$, and
3. If $A_1, A_2, \dots \in \mathcal{F}$ then their union $\bigcup_{n=1}^{\infty} A_n \in \mathcal{F}$.

It follows that a σ -algebra is also closed under intersection.

Definition 2.1.2. Let Ω be a nonempty set and let \mathcal{F} be a σ -algebra on Ω . Then the pair (Ω, \mathcal{F}) is called a **measurable space**.

Definition 2.1.3. Let (Ω, \mathcal{F}) be a measurable space. A **measure** μ is a real-valued set function on \mathcal{F} satisfying the following conditions

1. For every set $A \in \mathcal{F}$, $\mu(A) \geq 0$, and
2. If A_1, A_2, \dots is a sequence of disjoint sets in \mathcal{F} then

$$\mu \left(\bigcup_{n=1}^{\infty} A_n \right) = \sum_{n=1}^{\infty} \mu(A_n).$$

Now that the concept of a measure has been defined we can interpret a σ -algebra as the set of all events that can be measured from the underlying set Ω .

Definition 2.1.4. Let (Ω, \mathcal{F}) be a measurable space. A **probability measure** \mathbb{P} is a measure such that $\mathbb{P}(\Omega) = 1$.

Definition 2.1.5. Let μ be a measure defined on the measurable space (Ω, \mathcal{F}) . We call the triplet $(\Omega, \mathcal{F}, \mu)$ a **measure space**. In particular, if \mathbb{P} is a probability measure then we call $(\Omega, \mathcal{F}, \mathbb{P})$ a **probability space**.

A particular example of a probability measure that we will use later on is the Dirac measure.

Definition 2.1.6. Let (Ω, \mathcal{F}) be some measurable space. A **Dirac measure** δ_x is a measure defined on a given point $x \in \Omega$ and a set $A \in \mathcal{F}$

$$\delta_x(A) = \begin{cases} 0, & x \notin A; \\ 1, & x \in A. \end{cases}$$

A useful result of Dirac measures is the identity

$$\int \delta_x(dy) f(y) = f(x).$$

The Dirac measure can be thought of as the indicator function. Another example of a measure is the Lebesgue measure, which is the standard measure on the Euclidean space \mathbb{R}^d . The rigorous definition of a Lebesgue measure is fairly analytical and is omitted. The Lebesgue measure is essentially an extension of the standard notions of length, area and volume to d -dimensional Euclidean spaces.

Definition 2.1.7. Let (Ω, \mathcal{F}) and (E, \mathcal{E}) be measurable spaces. Then a function $f : \Omega \rightarrow E$ is **$(\mathcal{F}, \mathcal{E})$ -measurable** if $f^{-1}(S) \in \mathcal{F}$ for every $S \in \mathcal{E}$.

Definition 2.1.8. Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space and (E, \mathcal{E}) a measurable space. A **random variable** is a function $X : \Omega \rightarrow E$ that is $(\mathcal{F}, \mathcal{E})$ -measurable. We call (E, \mathcal{E}) the **state space**.

Definition 2.1.9. Brownian Motion

2.2 Stochastic Processes

Generally speaking, a stochastic process $\{X_t\}_{t \geq 0}$ is a collection of random variables that represent the evolution of some object over time. As this thesis focuses heavily on stochastic processes it is important that we rigorously define what a stochastic process is.

Definition 2.2.1. A **stochastic process** is a collection of random variables $\{X_t\}_{t \in \mathcal{I}}$ that is defined on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ and is indexed by an ordered set \mathcal{I} .

If we fix a $t_0 \in \mathcal{I}$ then $X(t_0, \omega)$ is simply a random variable. If we fix a $\omega_0 \in \Omega$ then $X(t, \omega_0)$ is a function with respect to t and is the trajectory of a single realisation of the stochastic process $X(t)$. From now on we will only consider stochastic processes where our ordered set $\mathcal{I} = \mathbb{R}^+$ represents an interval in continuous time. We now write $\{X(t)\}_{t \geq 0} : \mathbb{R}^+ \times \Omega \rightarrow \mathbb{R}$.

Definition 2.2.2. A stochastic process $\{X(t)\}_{t \geq 0}$ has a **jump** at time t if the process is discontinuous at that point. More precisely $X(t) \neq X_-(t)$.

Definition 2.2.3. A stochastic process $\{X(t)\}_{t \geq 0}$ is **càdlàg** (derived from the French term "continue droite, limite gauche") if it has right continuous paths with left hand limits. i.e $X_-(t)$ exists and $X_+(t)$ exists and is equal to $X(t)$.

Definition 2.2.4. A **filtration** $\{\mathcal{F}_t\}_{t \geq 0}$ is a collection of σ -algebras such that for any $0 \leq s \leq t$, $\mathcal{F}_s \subseteq \mathcal{F}_t$.

Note that with each successive σ -algebra the amount of sets that can be measured is non-decreasing. Thus a filtration can be interpreted as the accumulation of information about our stochastic process with respect to time. We write $\sigma(X_s, s \leq t)$ as the smallest σ -algebra generated by $\{X_s : s \leq t\}$.

Definition 2.2.5. Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a probability space with the filtration $\mathcal{F} = \{\mathcal{F}_t\}_{t \geq 0}$ and let (E, \mathcal{E}) a measurable space. We say the stochastic process $\{X_t\}_{t \geq 0}$ is **\mathcal{F} -adapted** if for every $t \geq 0$, $X_t : \Omega \rightarrow E$ is $(\mathcal{F}_t, \mathcal{E})$ - measurable.

2.3 Weak Convergence

The notion of weak convergence or convergence in distribution is also required to understand results utilised in this thesis.

Definition 2.3.1. A **metric** d on a set E is a non-negative real-valued function on the product space $E \times E$ such that

- $d(x, y) = 0$ iff $x = y$,
- $d(x, y) = d(y, x)$,
- $d(x, z) \leq d(x, y) + d(y, z)$.

A metric is a function that defines the distance between two points in a set.

Definition 2.3.2. Let E be a set and let d be a metric on E . Then the ordered pair (E, d) is a **metric space**.

Definition 2.3.3. A sequence $\{x_n\}_{n \geq 1}$ in a metric space (E, d) **converges to a limit** $x \in E$ if, for all $\epsilon > 0$, there exists an integer n_0 such that $d(x_n, x) < \epsilon$ for all $n \geq n_0$.

Definition 2.3.4. A sequence of probability measures $\{\mathbb{P}_n\}_{n \geq 1}$ on the measurable space (Ω, \mathcal{F}) **converges weakly** to a probability measure \mathbb{P} on (Ω, \mathcal{F}) if

$$\lim_{n \rightarrow \infty} \int_E f d\mathbb{P}_n = \int_E f d\mathbb{P}$$

for all continuous and bounded real-valued functions f on Ω . We then write

$$\mathbb{P}_n \rightarrow \mathbb{P}.$$

Note that this definition is just 2.3.3 with E as the set of all probability measures on the measurable space (Ω, \mathcal{F}) and $d(\mathbb{P}, \mathbb{Q}) = |\int_E f d\mathbb{P} - \int_E f d\mathbb{Q}|$, where \mathbb{Q} is another probability measure on (Ω, \mathcal{F}) .

Definition 2.3.5. Let $X_n : \Omega \rightarrow E$ be a sequence of random variables with a corresponding sequence of probability measures $\{\mathbb{P}_{X_n}\}_{n \geq 1}$ on the measurable space (Ω, \mathcal{F}) . Then $\{X_n\}_{n \geq 1}$ **converges in distribution** to a random variable $X : \Omega \rightarrow E$ if

$$\mathbb{P}_{X_n} \rightarrow \mathbb{P}_X$$

where \mathbb{P}_X is the probability measure of X . We then write $X_n \xrightarrow{d} X$.

Definition 2.3.6. Let $X_1 : \Omega \rightarrow E$ be a random variable with a corresponding probability measure \mathbb{P}_{X_1} on the measurable space (Ω, \mathcal{F}) . Then X_1 is **equal in distribution** to a random variable $X_2 : \Omega \rightarrow E$ if

$$\mathbb{P}_{X_1} = \mathbb{P}_{X_2}$$

where \mathbb{P}_{X_2} is the probability measure of X_2 . We then write $X_1 \stackrel{d}{=} X_2$ or $X_1 \sim X_2$.

CHAPTER 3

Scaling Limit Theorems

In this chapter we will provide an analysis of some well known scaling limit theorems, eventually arriving at a scaling limit for CTRM. We require these scaling limits before we can develop our statistical inference methods for CTRM.

3.1 Problem Formulation

In this section we rigorously define our problem as well as establish some preliminary notation that we will use to solve the problem.

Definition 3.1.1. *Assume identically and independently distributed (i.i.d.) pairs of random variables (J_k, W_k) , $k = 1, 2, \dots$ where $W_k > 0$ represents the inter-arrival times of certain events and $J_k \in \mathbb{R}$ the corresponding event magnitudes. Now write*

$$S(n) = \sum_{i=1}^n W_i \quad (3.1.1)$$

as the sum of the first n inter-arrival times. Also write

$$N(t) = \max\{n \in \mathbb{N} : S(n) \leq t\} \quad (3.1.2)$$

for the renewal process associated with the W_k . Then the process

$$M(t) = \bigvee_{k=1}^{N(t)} J_k = \max\{J_k : k = 1, \dots, N(t)\}, \quad t \geq 0. \quad (3.1.3)$$

is called a CTRM (Continuous Time Random Maxima process), where the maximum of the empty set is set to 0.

It is clear that the sample paths of $N(t)$ and $M(t)$ are right-continuous with left-hand limits. If W_k is interpreted as the time *leading up* to the event with magnitude J_k , then $M(t)$ is the largest magnitude observed until time t . The alternative case where W_k represents the inter-arrival time *following* J_k is termed “second type”

(in the shock model literature) or OCTRM (overshooting CTRM), and the largest magnitude up to time t is then given by

$$\tilde{M}(t) = \bigvee_{k=1}^{N(t)+1} J_k, \quad t \geq 0. \quad (3.1.4)$$

Finally, the model is called *coupled* when W_k and J_k are not independent. In this thesis we focus on the uncoupled case, for which it can be shown that the processes $M(t)$ and $\tilde{M}(t)$ have the same limiting distributions at large times [?], and hence we focus on the CTRM $M(t)$.

However deriving the distribution of $M(t)$ is difficult for a number of reasons that will be explained in further detail in 3.5 **not sure if this is true**. It is for this reason that we choose to make inference on the distribution of the following quantities:

Definition 3.1.2. *Let $M(t)$ be an uncoupled CTRM whose magnitudes J_k are supported on the interval $[x_0, x_F]$. Then the exceedance time of level $\ell \in [x_0, x_F]$ is the random variable*

$$T_\ell = \inf\{t : M(t) > \ell\}$$

and the exceedance is

$$X_\ell = M(T_\ell) - \ell.$$

These quantities are chosen as it is natural to be concerned with both the size and the timing of events of large magnitude. Possessing a strong understanding of these large events will aid in managing the consequences associated with extreme natural phenomena, financial loss and human behaviour.

Lemma 3.1.3. *Given a level $\ell \in [x_0, x_F]$, exceedance X_ℓ and exceedance time T_ℓ are independent. Moreover,*

$$\begin{aligned} \mathbf{P}[X_\ell > x] &= \frac{\bar{F}_J(\ell + x)}{\bar{F}_J(\ell)}, \quad x > 0, \\ \mathbf{P}[T_\ell > t] &= \bar{F}_J(\ell) \sum_{n=1}^{\infty} \int_0^{\infty} \bar{F}_W(t - t') \mathbf{P}[L_{n-1} \leq \ell] \mathbf{P}[S_{n-1} \in dt'], \quad t > 0 \end{aligned}$$

where $L_n = \bigvee_{k=1}^n J_k$, $L_0 = x_0$, and $S_n = \sum_{k=1}^n W_k$, $S_0 = 0$.

Proof. Let $\tau_\ell = \min\{k : J_k > x\}$. Then **rigorize this**

$$\begin{aligned} T_\ell > t, X_\ell > x &\iff S_{\tau_\ell} > t, L_{\tau_\ell} > \ell + x \iff \exists n : L_{n-1} \leq \ell, J_n > \ell + x, S_n > t \\ &\iff \exists n : L_{n-1} \leq \ell, J_n > \ell + x, W_n > t - S_{n-1} \end{aligned}$$

and such n must be unique, since $x > 0$. Thus

$$\begin{aligned} \mathbf{P}[T_\ell > t, X_\ell > x] &= \sum_{n=1}^{\infty} \mathbf{P}[J_n > \ell + x, W_n > t - S_{n-1}, L_{n-1} \leq \ell] \\ &= \sum_{n=1}^{\infty} \int_0^\ell \int_0^\infty \mathbf{P}[J_n > \ell + x, W_n > t - t' | L_{n-1} = m', S_{n-1} = t'] \mathbf{P}[L_{n-1} \in dm', S_{n-1} \in dt'] \\ &= \sum_{n=1}^{\infty} \int_0^\infty \bar{F}_J(\ell + x) \bar{F}_W(t - t') \mathbf{P}[L_{n-1} \leq \ell] \mathbf{P}[S_{n-1} \in dt'] \end{aligned}$$

since the sequences J_k and W_k are i.i.d. and independent of each other. Letting $t \downarrow 0$, we see

$$\begin{aligned} \mathbf{P}[X_\ell > x] &= \sum_{n=1}^{\infty} \bar{F}_J(\ell + x) \mathbf{P}[L_{n-1} \leq \ell] = \bar{F}_J(\ell + x) \sum_{n=1}^{\infty} F_J(\ell)^{n-1} \\ &= \frac{\bar{F}_J(\ell + x)}{\bar{F}_J(\ell)}, \end{aligned}$$

and letting $x \downarrow 0$, one gets $\mathbf{P}[T_\ell > t]$. One checks that $\mathbf{P}[T_\ell > t, X_\ell > x] = \mathbf{P}[T_\ell > t] \mathbf{P}[X_\ell > x]$, implying independence. \square

The distribution of the exceedance is hence, as expected, simply the conditional distribution of a magnitude J_k given $J_k > \ell$. We will look at the distributions of exceedances and exceedance times in further detail in chapter 4.

3.2 Background

A stochastic process limit is the limit of a sequence of stochastic processes. Before we examine the stochastic process limits that we will utilise in this thesis it will be illustrative to analyse one of the most well known examples of a stochastic process limit - Brownian Motion. One of the methods used to derive Brownian motion is devised by trying to seek a continuous time representation of the discrete time random walk. Recall that a discrete time random walk is defined in the following manner.

Definition 3.2.1. Let $\{U_i\}_{1 \leq i \leq n}$ be i.i.d. $U(0,1)$ random variables. Write the partial sums at time k as

$$S_k = \sum_{i=1}^k U_i,$$

then successive partial sums form a **random walk** with S_n being the position after n steps.

The general method of achieving a continuous limit of a discrete function is to consider a continuous analogue of the sum S_n as $n \rightarrow \infty$ and then rescaling to ensure the continuous analogue remains finite. We can then appeal to the Central Limit Theorem which gives us the discrete limit

$$\frac{S_n - n\mu}{\sqrt{n\sigma^2}} \xrightarrow{d} Z \sim N(0,1) \text{ as } n \rightarrow \infty \quad (3.2.1)$$

where $\mu := \mathbb{E}[U_k] = 1/2$ and $\sigma^2 := \text{Var}(U_k) = 1/12$. The random walk only accepts integer arguments so in order to extend this discrete limit to a continuous limit we first consider two different continuous time representations of a discrete time random walk. One obvious way of forming a process from the random walk is to perform linear interpolation between each discrete point. Thus, we can consider

$$\tilde{S}(t) := (t - \lfloor t \rfloor)U_{\lfloor t \rfloor + 1} + S_{\lfloor t \rfloor} \quad \text{for all } t \geq 0, \quad (3.2.2)$$

or more simply, we can consider the step function

$$S_{\lfloor t \rfloor} = U_1 + \cdots + U_{\lfloor t \rfloor} \quad \text{for all } t \geq 0. \quad (3.2.3)$$

maybe some graphs here As it turns out for large values of t both processes are equivalent, see [Whi01] for more details. Although the interpolation function is more intuitive, the step function is easier to work with and so we only utilise the step function from now on. We now add a time scaling constant c and consider

$$S_{\lfloor ct \rfloor} = U_1 + \cdots + U_{\lfloor ct \rfloor} \quad \text{for all } t \geq 0. \quad (3.2.4)$$

Whereas in 3.2.3 a step can only occur at the end of every time period, the inclusion of c in 3.2.4 allows us to adjust how often steps occur. For example, if we multiply c by 60, then steps that were occurring every minute now occur every second. Clearly as c approaches infinity steps start to occur continuously. Now that we have a continuous analogue of the discrete random walk we wish to find the limiting

process. Following 3.2.1 we would expect that

$$\frac{S_{[ct]} - [ct]\mu}{\sqrt{c\sigma^2}} \xrightarrow{d} B_t \sim N(0, t) \text{ as } c \rightarrow \infty, \quad (3.2.5)$$

where B_t is the standard Brownian Motion. An easy way to check that the above weak convergence is correct is to consider the Fourier transforms of both sides.

Definition 3.2.2. Let $Y : \Omega \rightarrow E$ be a random variable with probability density function $f_Y(x)$. Then the **Fourier Transform** of $f_Y(x)$ is

$$\hat{f}_Y(s) = \mathbb{E}[e^{-isY}] = \int_{\Omega} e^{-isx} f_Y(x) dx.$$

When the first two moments are finite we can use the Taylor series expansion of $e^z = 1 + z + z^2/2! + \dots$ to see that

$$\hat{f}_Y(s) = \int_{\Omega} \left(1 - isx + \frac{1}{2!}(-isx)^2 + \dots \right) f_Y(x) dx \quad (3.2.6)$$

$$= \int_{\Omega} f_Y(x) dx - is \int_{\Omega} x f_Y(x) dx - \frac{s^2}{2} \int_{\Omega} x^2 f_Y(x) dx + \dots \quad (3.2.7)$$

$$= 1 - is\mathbb{E}[Y] - \frac{s^2}{2}\mathbb{E}[Y^2] + o(s^2), \quad (3.2.8)$$

where $o(s^2)$ is a function that approaches zero faster than s^2 as $s \rightarrow 0$.

Suppose $Y \sim Y_k = U_k - \mu$ then

$$\hat{f}_Y(s) = 1 - \frac{s^2}{2}\sigma^2 + o(s^2), \quad (3.2.9)$$

since $\mathbb{E}[Y] = 0$, and $\mathbb{E}[Y^2] = \text{Var}(Y) + \mathbb{E}[Y]^2 = \sigma^2$. We now have the necessary tools to calculate the Fourier transform of the rescaled $S_{[ct]}$ defined in 3.2.5. Using the fact that the Fourier transform of a sum is simply the product of the individual Fourier transforms we then have

$$\mathbb{E} \left[\exp \left(-is \frac{S_{[ct]} - [ct]\mu}{\sqrt{c\sigma^2}} \right) \right] = \mathbb{E} \left[\exp \left(-is \frac{Y_1 + \dots + Y_{[ct]}}{\sqrt{c\sigma^2}} \right) \right] \quad (3.2.10)$$

$$= \mathbb{E} \left[\exp \left(\frac{-isY}{\sqrt{c\sigma^2}} \right) \right]^{[ct]} \quad (3.2.11)$$

$$= \hat{f}_Y \left(\frac{s}{\sqrt{c\sigma^2}} \right)^{[ct]}. \quad (3.2.12)$$

Given that $(1 + (r/n) + o(n^{-1}))^n \rightarrow e^r$ as $n \rightarrow \infty$ for any $r \in \mathbb{R}$ we then have

$$\mathbb{E} \left[\exp \left(-is \frac{S_{\lfloor ct \rfloor} - \lfloor ct \rfloor \mu}{\sqrt{c\sigma^2}} \right) \right] = \hat{f}_Y \left(\frac{s}{\sqrt{c\sigma^2}} \right)^{\lfloor ct \rfloor} \quad (3.2.13)$$

$$= \left(1 - \frac{s^2}{2c} + o(c^{-1}) \right)^{\lfloor ct \rfloor} \quad (3.2.14)$$

$$= \left[\left(1 - \frac{s^2}{2c} + o(c^{-1}) \right)^c \right]^{\frac{\lfloor ct \rfloor}{c}} \quad (3.2.15)$$

$$\rightarrow e^{-ts^2/2} \quad \text{as } c \rightarrow \infty, \quad (3.2.16)$$

since $\lfloor ct \rfloor / c \rightarrow t$ as $c \rightarrow \infty$.

We now refer to the Lévy Continuity Theorem to prove the weak convergence outlined in 3.2.5.

Definition 3.2.3. [MS01] **Lévy Continuity Theorem.** *If X_n, X are random variables on \mathbb{R} , then $X_n \xrightarrow{d} X$ implies that $\hat{f}_{X_n}(s) \rightarrow \hat{f}_X(s)$ for each $s \in \mathbb{R}$, uniformly on compact subsets. Conversely, if X_n is a sequence of random variables such that $\hat{f}_{X_n}(s) \rightarrow \hat{f}_X(s)$ for each $s \in \mathbb{R}$, and the limit is continuous at $s = 0$, then $\hat{f}_X(s)$ is the Fourier transform of some random variable X , and $X_n \xrightarrow{d} X$.*

This theorem essentially states that $X_n \xrightarrow{d} X$ if and only if $\hat{f}_{X_n}(s) \rightarrow \hat{f}_X(s)$. We previously showed that the limit of the Fourier transform of the rescaled $S_{\lfloor ct \rfloor}$ is $e^{-ts^2/2}$ which is also the Fourier transform of a $N(0, t)$ random variable. Thus we must have the result outlined in 3.2.5.

Now that we have shown how one can derive Brownian Motion as a stochastic process limit of the rescaled $S_{\lfloor ct \rfloor}$, the stochastic process limits presented in later sections will be more accessible.

3.3 Stable Distributions

In order for our inter-arrival waiting times W_k to be modelled correctly they should be heavy tailed and completely right skewed so as not to violate our initial bursty assumption. Stable distributions are an extensive class of probability distributions that allow for both skew and heavy tails and are thus a perfect distribution to describe our waiting times with. As we will find out later in this section they also have many useful properties. The class was first defined by Paul Lévy in the early 20th century but was under utilised due to the lack of closed formulas for densities and distribution functions for all but a few stable distributions (Gaussian, Cauchy and Lévy). However its use has increased since advances in computing methods expand now allow for the computation of stable densities, distribution functions

and quantiles.

Definition 3.3.1. [Nol15] A random variable X is **sum-stable** if for X_1 and X_2 independent copies of X and any positive constants a and b we have

$$aX_1 + bX_2 \stackrel{d}{=} cX + d$$

for some positive c and some $d \in \mathbb{R}$. We then say that X follows a **stable distribution**. Note that the symbol $\stackrel{d}{=}$ was defined in Definition 2.3.6.

Remark 3.3.2. Note that other texts refer to what we have just defined as **sum-stable** as just **stable**. We make this modification in order to easily differentiate it from the term **max-stable** which we will define in section 3.5.

As a closed form for the density of a stable random variable cannot be found for all cases we instead choose to specify a stable random variable by its characteristic function.

Definition 3.3.3. Let $X \sim S(\beta, \alpha, \gamma, \delta)$ be a stable random variable. Then the characteristic function $\phi_X(s) = \mathbb{E}[e^{isX}]$ is given by

$$\phi_X(s) = \begin{cases} \exp(-\gamma^\beta |s|^\beta [1 - i\alpha \operatorname{sign}(s) \tan(\frac{\pi\beta}{2})] + i\delta s) & \beta \neq 1 \\ \exp(-\gamma |s| [1 + i\alpha \operatorname{sign}(s) \frac{2}{\pi} \log(s)] + i\delta s) & \beta = 1. \end{cases} \quad (3.3.1)$$

There are over half a dozen parameterisations of a stable random variable but the well informed reader will note that the parameterisation we use is consistent with [MS08] & [ST94]. Definition 3.3.3 then implies that a stable distribution possesses four different parameters:

- An index of stability $\beta \in (0, 2]$ which determines heavy tailedness
- A skewness parameter $\alpha \in [-1, 1]$
- A scale parameter $\gamma \geq 0$
- A location parameter $\delta \in \mathbb{R}$

As we are using the stable distribution to model bursty waiting times we need a completely right skewed distribution with a heavy tail (infinite mean). Thus we consider waiting times drawn from a stable distribution with:

- Index of stability $\beta \in (0, 1)$ to ensure heavy tailedness

- Skewness parameter $\alpha = 1$ to ensure a right skewed distribution
- Standard scale parameter $\gamma = 1$ and standard location parameter $\delta = 0$

To check that the mean is infinite when $\beta \in (0, 1)$ we first see from [JW94] [check parameterisation](#) that if a Central Limit Theorem type argument is used we have

$$\lim_{x \rightarrow \infty} x^\beta \mathbb{P}(X > x) = C_\beta \frac{1 + \alpha}{2} \gamma^\beta, \quad (3.3.2)$$

$$\lim_{x \rightarrow \infty} x^\beta \mathbb{P}(X < -x) = C_\beta \frac{1 - \alpha}{2} \gamma^\beta, \quad (3.3.3)$$

where

$$C_\beta = \left(\int_0^\infty x^{-\beta} \sin(x) dx \right)^{-1} = \frac{2}{\pi} \Gamma(\beta) \sin \frac{\pi\beta}{2}.$$

When $\alpha = 1$ we have $\lim_{x \rightarrow \infty} \mathbb{P}(X < -x) = 0$, as one would expect. Since a Pareto random variable T with parameter β has a tail function $\mathbb{P}(T > t) = ct^{-\beta}$ for some $c \in \mathbb{R}$ and $t > c^{1/\beta}$, we can then observe that the tails of a stable distribution are asymptotically equivalent to a Pareto distribution and thus exhibit power-law behaviour. [Thus if we show that a Pareto distribution with \$\beta \in \(0, 1\)\$ has an infinite mean then the corresponding stable distribution also has an infinite mean.](#) Now if $p(t)$ is the density of T we have that

$$\mathbb{E}[T] = \int_0^\infty tp(t)dt = \int_0^\infty t \frac{c\beta}{t^{\beta+1}} dt = c\beta \int_0^\infty t^{-\beta} dt = c\beta \left[\frac{t^{1-\beta}}{1-\beta} \right]_0^\infty.$$

The above expectation is clearly infinite if $0 < \beta \leq 1$, but if $\beta = 1$ our initial stable distribution reduces to a Cauchy distribution and thus we exclude that case. [why ignore Cauchy?? also maybe talk about variance here](#)

One of the reasons the stable distribution is so useful is that we can now generalise the classical Central Limit Theorem. Recall that the Central Limit Theorem says that the normalised sum of i.i.d. random variables with a finite variance converges to a normal distribution. To be more precise, let $\{X_i\}_{1 \leq i \leq n}$ be i.i.d. random variables with mean μ and variance σ^2 . Then the classical Central Limit Theorem states that

$$\frac{\bar{X}_n - \mu}{\sigma/\sqrt{n}} \xrightarrow{d} Z \sim N(0, 1) \text{ as } n \rightarrow \infty$$

where $\bar{X}_n = (\sum_{i=1}^n X_i)/n$. Rewriting the above to match the notation we will use from now on we have

$$\frac{\sum_{i=1}^n X_i - a_n}{b_n} \xrightarrow{d} Z \sim N(0, 1) \text{ as } n \rightarrow \infty$$

where $a_n = \sqrt{n}\sigma/\mu$ and $b_n = \sigma\sqrt{n}$.

The Generalised Central Limit Theorem generalises the Central Limit theorem by dropping the assumption that the variance is finite.

Theorem 3.3.4. Generalised Central Limit Theorem *Let $\{X_i\}_{1 \leq i \leq n}$ be an i.i.d. sequence of random variables. There exists constants $a_n > 0, b_n \in \mathbb{R}$ and a nondegenerate random variable Z with*

$$\frac{\sum_{i=1}^n X_i - a_n}{b_n} \xrightarrow{d} Z$$

if and only if Z is a stable random variable.

The Generalised Central Limit theorem is so useful as it allows us to find the limiting distribution of rescaled random variables with undefined moments. This means we can apply our theorem to our inter-arrival waiting times W_k with infinite means. As discussed in Section 3.2 finding the limiting distribution is a prerequisite for finding any stochastic process limit. The Generalised Central Limit Theorem also naturally leads to the following definition.

Definition 3.3.5. [Nol15] *A random variable X is in the **domain of attraction** of Z if there exists constants $a_n > 0, b_n \in \mathbb{R}$ with*

$$\frac{\sum_{i=1}^n X_i - a_n}{b_n} \xrightarrow{d} Z,$$

where X_1, X_2, X_3, \dots are independent identically distributed copies of X . We then say $X \in DOA(Z)$.

The above definition will be useful in the next section where we present a stochastic process limit for the the sum of our inter-arrival waiting times.

3.4 Scaling Limit of the Sum of Waiting Times

As mentioned in the previous section, we assume a heavy-tailed distribution for the waiting times, more precisely that the tail function $\bar{F}_W := 1 - F_W$ of the CDF of W is regularly varying:

$$t \mapsto \bar{F}_W(t) \in RV(-\beta), \quad \beta \in (0, 1) \tag{3.4.1}$$

meaning that [Sen76, MS01]

$$\lim_{t \rightarrow \infty} \frac{\overline{F}_W(\lambda t)}{\overline{F}_W(t)} = \lambda^{-\beta}.$$

By [MS01, Cor. 8.2.19] (also compare [Whi01, Th. 4.5.1]), this is equivalent to W_1 being in the (sum-) domain of attraction of a positively skewed stable law D with stability parameter β , which means the weak convergence

$$b(n)(W_1 + \dots + W_n) \Rightarrow D, \quad n \rightarrow \infty. \quad (3.4.2)$$

The i.i.d. nature of the W_k allows for a more general, functional limit theorem [MS01, Ex. 11.2.18], [Whi01, Th. 4.5.3]:

$$b(c)(W_1 + \dots + W_{\lfloor cr \rfloor}) \Rightarrow D(cr), \quad c \rightarrow \infty, \quad (3.4.3)$$

where $\{D(r)\}_{r \geq 0}$ is a stable subordinator, defined via

$$\mathbf{E}[e^{-\lambda D(r)}] = e^{-r\lambda^\beta}, \quad (3.4.4)$$

where $\lfloor r \rfloor$ denotes the largest integer not larger than r , and where convergence holds on the stochastic process level, i.e. in the sense of weak convergence in Skorokhod space endowed with the J_1 topology [Whi01, Sec. 3.3].

The renewal process paths $N(t)$ result, up to a constant, from inverting the non-decreasing sample paths of the sum $r \mapsto W_1 + \dots + W_{\lfloor r \rfloor}$. Accordingly, the scaling limit of the renewal process is [MS04]

$$N(ct)/\tilde{b}(c) \Rightarrow E(t), \quad c \rightarrow \infty \quad (3.4.5)$$

where $E(t)$ denotes the inverse stable subordinator [MS13]

$$E(t) = \inf\{r : D(r) > t\}, \quad t \geq 0, \quad (3.4.6)$$

and where $\tilde{b}(c)$ is asymptotically inverse to $1/b(c)$, in the sense of [Sen76, p.20]:

$$1/b(\tilde{b}(c)) \sim c \sim \tilde{b}(1/b(c)) \quad (3.4.7)$$

where a \sim symbol indicates that the quotient of both sides converges to 1 as $c \rightarrow \infty$. Note that $\tilde{b}(c) \in RV(\beta)$.

The inverse stable subordinator $E(t)$ governs the temporal dynamics of the scaling limit of the CTRM process $M(t)$, see Theorem ???. It is self-similar with exponent β [MS04], non-decreasing, and the (regenerative, random) set \mathcal{R} of its

points of increase is a fractal with dimension β [Ber99]. $E(t)$ is a model for time series with intermittent, ‘bursty’ behaviour, for the following reasons:

- i) Conditional on $t \geq 0$ being a point of increase, any interval $(t, t + \epsilon)$ almost surely contains uncountably many other points of \mathcal{R} ;
- ii) \mathcal{R} has Lebesgue measure 0, and hence $E(t)$ is constant at any “randomly” chosen time t .

Having only two parameters ($\beta \in (0, 1)$ and a scale parameter) the inverse stable subordinator hence models scaling limits of heavy-tailed waiting times parsimoniously.

3.5 Scaling Limit of the Maximum of Event Magnitudes

3.6 Scaling Limit of CTRM

The Continuous Time Random *Walk* (CTRW) has been a highly successful model for anomalous diffusion in the past two decades [MK00, HTSL10], likely due to its tractable and flexible scaling properties. The stochastic process we study in this article is conceptually very close to the CTRW, since essentially the jumps J_k are reinterpreted as magnitudes, and instead of the cumulative sum, one tracks the cumulative maximum. Similarly tractable scaling properties apply to the CTRM, see below. For the above reasons, we use the name CTRM in this article.

CHAPTER 4

Distribution of CTRM

4.1 Distribution of Exceedances

4.2 Distribution of Exceedance Durations

An Overview of Generalized Gamma MittagLeffler Model and Its Applications

CHAPTER 5

Statistical Inference

5.1 Log Moment Estimator for Mittag-Leffler Random Variables

5.2 Maximum Likelihood Estimator for Generalised Pareto
Random Variables

5.3 Stability of Parameter Estimates

5.4 Goodness of Fit

We will test Goodness of Fit as Mainardi et al with the reduced chi squared test

CHAPTER 6

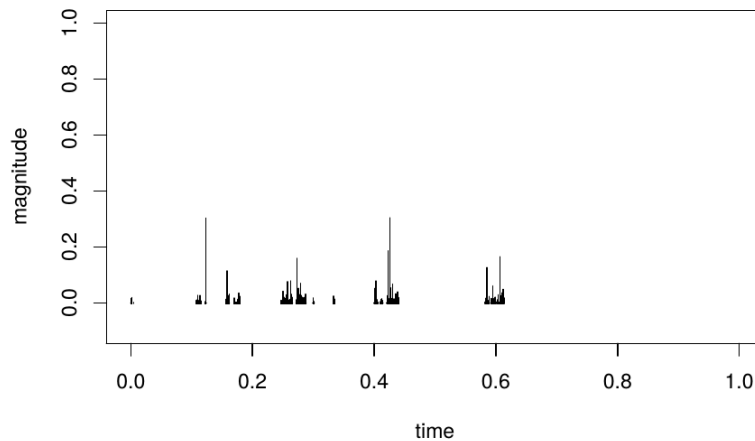
Model Fitting

The accuracy of the Mittag-Leffler log moment estimator and the Generalised Pareto MLE has been established in [Cah13] and [?] respectively. However it is yet to be confirmed that these estimation methods work for exceedances and exceedance durations. Thus before we apply our inference methods to actual data, we will check that we can correctly estimate the theoretical Mittag-Leffler and Generalised Pareto parameters derived from a simulated bursty process. If we are able to accurately estimate these parameters we can then proceed to apply the estimation methods to actual data, confident in knowing that the actual parameters will be estimated.

6.1 Model Fit of Simulated Data

We assume n i.i.d. waiting times are drawn from the positively skewed stable distribution with stability parameter $\beta = 0.8$, scaled with $n^{-1/\beta}$ to ensure convergence to a stable subordinator. This defines a renewal process, at whose renewal times we assume n i.i.d. magnitudes are drawn from a Generalized Extreme Value Distribution with shape parameter $\xi = 0.7$.

Figure 6.1: A simulated bursty process



Recall that we have defined the exceedance duration of level $\ell \in [x_0, x_F]$ as the random variable

$$T_\ell = \inf\{t : M(t) > \ell\}$$

and the exceedance as

$$X_\ell = M(T_\ell) - \ell.$$

As these random variables are the quantities of interest it is natural for us to then consider the (simulated) sample exceedances and their corresponding durations over the threshold level ℓ .

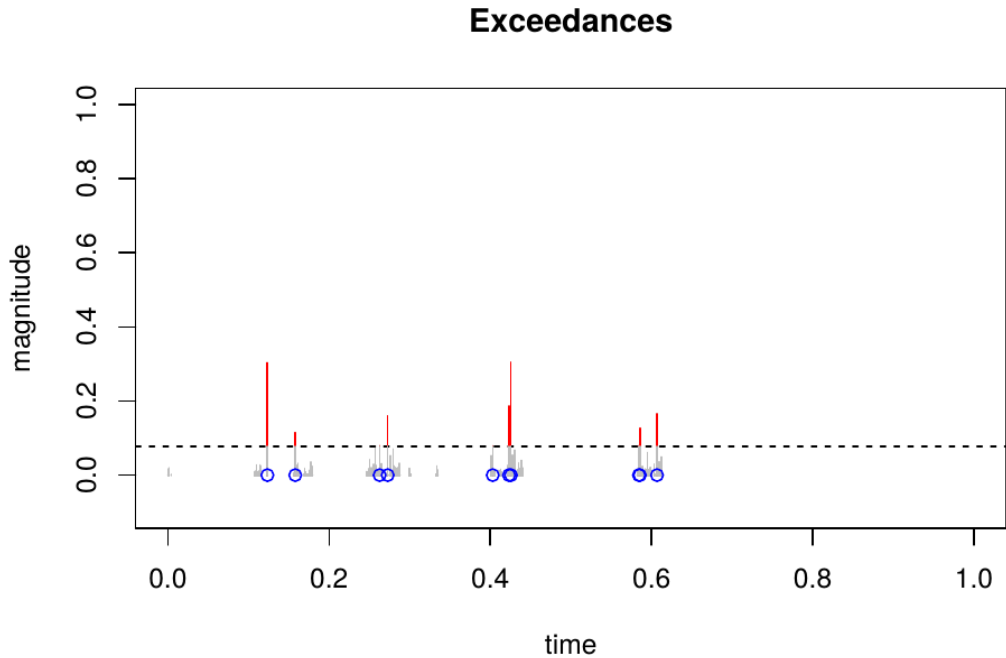


Figure 6.2: Exceedance durations (blue circles) and Exceedance sizes (red lines).

Assume now a time series of magnitudes, and that interest lies in the estimation of the timings of the large magnitudes. Consider a minimum threshold ℓ_0 , e.g. at the 95% quantile. Vary the threshold ℓ on the interval $[\ell_0, x_F]$, and consider the resulting sequences of exceedance sizes and exceedance times $\{(X_{\ell,i}, T_{\ell,i})\}$. Due to the renewal property each sequence is i.i.d.. Now $T_{\ell,1}, T_{\ell,2}, \dots$ can be modelled by a Mittag-Leffler distribution and $J_{\ell,1}, J_{\ell,2}, \dots$ can be modelled by a Generalised Pareto distribution as shown in chapter 4.

6.2 Model Fit of Real Data

CHAPTER 7

Conclusion

References

- [And87] Kevin K Anderson. Limit Theorems for General Shock Models with Infinite Mean Intershock Times. *J. Appl. Probab.*, 24(2):449–456, 1987.
- [Bar05] Albert-László Barabási. The origin of bursts and heavy tails in human dynamics. *Nature*, 435(May):207–211, 2005.
- [BB13] James P. Bagrow and Dirk Brockmann. Natural emergence of clusters and bursts in network evolution. *Phys. Rev. X*, 3(2):1–6, 2013.
- [Ber99] Jean Bertoin. *Subordinators: examples and applications*, volume 1717 of *Lecture Notes in Mathematics*. Springer Berlin Heidelberg, Berlin, Heidelberg, 1999.
- [BŠ14] Bojan Basrak and Drago Špoljarić. Extremal behaviour of random variables observed in renewal times. jun 2014.
- [BSM07] David A Benson, Rina Schumer, and Mark M Meerschaert. Recurrence of extreme events with power-law interarrival times. *Geophys. Res. Lett.*, 34(116404):DOI:10.1029/2007GL030767, aug 2007.
- [Cah13] Dexter O. Cahoy. Estimation of Mittag-Leffler Parameters. *Commun. Stat. - Simul. Comput.*, 42(2):303–315, feb 2013.
- [EM73] J. D. Esary and A. W. Marshall. Shock Models and Wear Processes, 1973.
- [Gut99] Allan Gut. Extreme Shock Models. *Extremes*, (1983):295–307, 1999.
- [Haw71] Alan G Hawkes. Point spectra of some mutually exciting point processes. *J. R. Stat. Soc. Ser. B*, pages 438–443, 1971.
- [HS15] Katharina Hees and H.P. Scheffler. Coupled continuous time random maxima. pages 1–24, 2015.
- [HTSL10] B.I. Henry, T. A.M. Langlands, Peter Straka, and T.A.M. Langlands. An introduction to fractional diffusion. In R L. Dewar and F Detering, editors, *Complex Phys. Biophys. Econophysical Syst. World Sci. Lect. Notes Complex Syst.*, volume 9 of *World Scientific Lecture Notes in Complex Systems*, pages 37–90, Singapore, 2010. World Scientific.
- [JW94] Aleksander Janicki and Aleksander Weron. Can one see α -stable variables and processes? *Statist. Sci.*, 9(1):109–126, 1994.

- [KKBK12] Márton Karsai, Kimmo Kaski, Albert László Barabási, and János Kertész. Universal features of correlated bursty behaviour. *Sci. Rep.*, 2, 2012.
- [KKP⁺11] M. Karsai, M. Kivelä, R. K. Pan, K. Kaski, J. Kertész, Albert László Barabási, and J. Saramäki. Small but slow world: How network topology and burstiness slow down spreading. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 83:1–4, 2011.
- [MGV11] Byungjoon Min, K. I. Goh, and Alexei Vazquez. Spreading dynamics following bursty human activity patterns. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 83(3):2–5, 2011.
- [MK00] Ralf Metzler and Joseph Klafter. The random walk’s guide to anomalous diffusion: a fractional dynamics approach. *Phys. Rep.*, 339(1):1–77, dec 2000.
- [MS01] Mark M Meerschaert and H.P. Scheffler. *Limit Distributions for Sums of Independent Random Vectors: Heavy Tails in Theory and Practice*. Wiley-Interscience, New York, first edition, jul 2001.
- [MS04] Mark M Meerschaert and H.P. Scheffler. Limit Theorems for Continuous-Time Random Walks with Infinite Mean Waiting Times. *J. Appl. Probab.*, 41(3):623–638, sep 2004.
- [MS08] Mark M Meerschaert and Stilian A Stoev. Extremal limit theorems for observations separated by random power law waiting times. *J. Stat. Plan. Inference*, 139(7):2175–2188, jul 2008.
- [MS13] Mark M Meerschaert and Peter Straka. Inverse Stable Subordinators. *Math. Model. Nat. Phenom.*, 8(2):1–16, apr 2013.
- [Nol15] J. P. Nolan. *Stable Distributions - Models for Heavy Tailed Data*. Birkhauser, Boston, 2015. In progress, Chapter 1 online at academic2.american.edu/~jpnolan.
- [OB05] J Oliveira and Albert László Barabási. Darwin and Einstein correspondence patterns. *Nature*, 437(October):1251, 2005.
- [OS11] Takahiro Omi and Shigeru Shinomoto. Optimizing Time Histograms for Non-Poissonian Spike Trains. *Neural Comput.*, 23(12):3125–3144, 2011.
- [Sen76] E Seneta. *Regularly Varying Functions*, volume 508 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 1976.
- [Sil02] Dmitrii S Silvestrov. *Limit Theorems for Randomly Stopped Stochastic Processes*. Springer (Berlin, Heidelberg), 2002.
- [SS83] J. George Shanthikumar and Ushio Sumita. General shock models associated with correlated renewal sequences. *J. Appl. Probab.*, 20(3):600–614, 1983.

- [SS84] J. George Shanthikumar and Ushio Sumita. Distribution Properties of the System Failure Time in a General Shock Model. *Adv. Appl. Probab.*, 16(2):363–377, 1984.
- [SS85] J. George Shanthikumar and Ushio Sumita. A class of correlated cumulative shock models. *Adv. Appl. Probab.*, 17(2):347–366, 1985.
- [ST94] Gennady Samorodnitsky and Murad S. Taqqu. *Stable non-Gaussian random processes*. Stochastic Modeling. Chapman & Hall, New York, 1994. Stochastic models with infinite variance.
- [ST04] Dmitrii S Silvestrov and Jozef L. Teugels. Limit theorems for mixed max-sum processes with renewal stopping. *Ann. Appl. Probab.*, 14(4):1838–1868, nov 2004.
- [VOD⁺06] a Vasquez, J G Oliveira, Z Dezso, K-I Goh, I Kondor, and Albert-László Barabási. Modeling bursts and heavy tails in human dynamics. *Phys. Rev. E*, 73:361271–3612718, 2006.
- [VRLB07] Alexei Vazquez, Balázs Rácz, András Lukács, and Albert-László Barabási. Impact of non-poissonian activity patterns on spreading processes. *Phys. Rev. Lett.*, 98(APRIL):1–4, 2007.
- [VTK13] Szabolcs Vajna, Bálint Tóth, and János Kertész. Modelling bursty time series. *New J. Phys.*, 15(10):103023, oct 2013.
- [Whi01] Ward Whitt. *Stochastic-Process Limits: An Introduction to Stochastic-Process Limits and their Application to Queues*. Springer, New York, 1st edition, nov 2001.