

# Quantum Mechanics

## **Examples of Exactly Solvable Systems**

- 1. Free Particle**
- 2. Particle in a Square-Well Potential**
- 3. Hydrogen Atom**

## Free Particle

Time-independent Schrodinger equation

$$\hat{H}\psi = E\psi$$

$$\left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \psi(x) = E \cdot \psi(x)$$

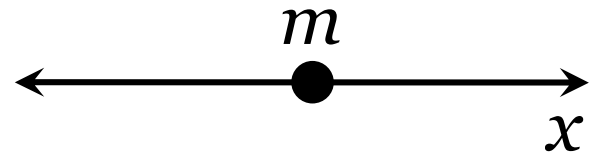
For a free particle  $V(x)=0$

There are no external forces acting

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) = E \cdot \psi(x)$$

## Free Particle

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) = E \cdot \psi(x)$$



Second-order linear differential equation

Let us assume

$$\psi(x) = A \sin kx + B \cos kx$$

Trial Solution

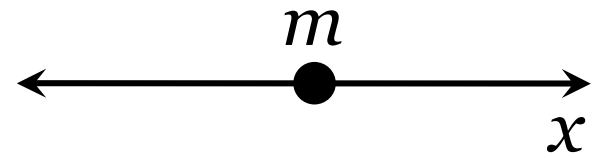
$$\psi(x) = A \sin kx + B \cos kx$$

$$\frac{\partial}{\partial x} \psi(x) = \frac{\partial}{\partial x} (A \sin kx + B \cos kx) = k (A \cos kx - B \sin kx)$$

$$\frac{\partial^2}{\partial x^2} \psi(x) = -k^2 (A \sin kx + B \cos kx) = -k^2 \psi(x)$$

## Free Particle

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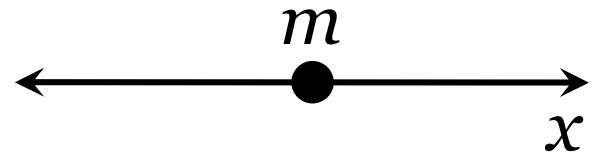
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## Free Particle

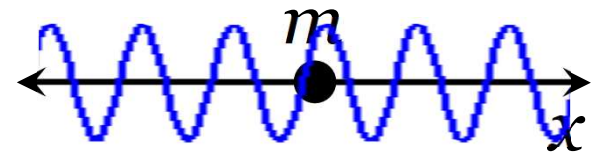
$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) = E \cdot \psi(x)$$



$$\frac{\hbar^2}{2m} k^2 \psi(x) = E \cdot \psi(x) \quad \Rightarrow \quad E = \frac{\hbar^2 k^2}{2m} \quad \Rightarrow \quad k = \pm \frac{\sqrt{2mE}}{\hbar}$$

## Free Particle

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) = E \cdot \psi(x)$$



de Broglie wave

$$\frac{\hbar^2}{2m} k^2 \psi(x) = E \cdot \psi(x) \quad \Rightarrow \quad E = \frac{\hbar^2 k^2}{2m} \quad \Rightarrow \quad k = \pm \frac{\sqrt{2mE}}{\hbar}$$

$$E = \frac{\hbar^2 k^2}{2m}$$

There are no restrictions on  $k$

$E$  can have any value

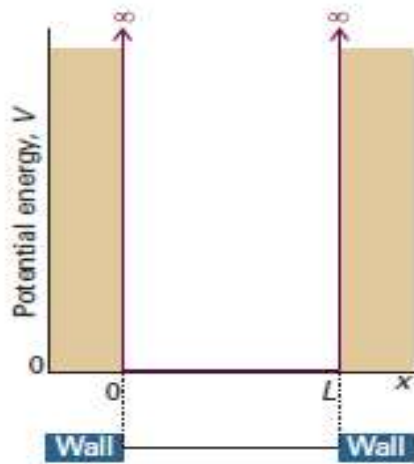
Energies of free particles are continuous

$$\psi(x) = A \sin \frac{\sqrt{2mE}}{\hbar} x + B \cos \frac{\sqrt{2mE}}{\hbar} x$$

**No Quantization**

**All energies are allowed**

# Particle in 1-D Square-Well Potential



$$V(x) = \begin{cases} \infty & x < 0 \\ 0 & 0 \leq x \leq L \\ \infty & x > L \end{cases}$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + V(x)\psi(x) = E \cdot \psi(x)$$

For regions in the space  $x < 0$  and  $x > L \Rightarrow V = \infty$

$$\frac{\partial^2}{\partial x^2} \psi(x) = \frac{2m}{\hbar^2} (V - E) \cdot \psi(x) = \infty \cdot \psi(x)$$

Normalization condition not satisfied  $\Rightarrow$

$$\psi(x < 0) = 0 \quad \text{and} \quad \psi(x > L) = 0$$

## Particle in 1-D Square-Well Potential

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + V(x)\psi(x) = E \cdot \psi(x)$$

For regions in the space  $0 \leq x \leq L \Rightarrow V = 0$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) = E \cdot \psi(x)$$

This equation is similar to free particle Schrodinger  
However, boundary conditions are present

Let is assume

$$\psi(x) = A \sin kx + B \cos kx$$

Trial Solution

$$E = \frac{\hbar^2 k^2}{2m}$$

Energy



## Particle in 1-D Square-Well Potential

$$\psi(x) = A \sin kx + B \cos kx$$

Boundary Condition  $x = 0 \Rightarrow \psi(x) = 0$

$$\psi(x) = A \sin kx \quad \because \cos 0 = 1$$

Boundary Condition  $x = L \Rightarrow \psi(L) = 0$

$$\psi(L) = 0 \Rightarrow A \sin kL = 0 \Rightarrow A = 0 \text{ or } \sin kL = 0$$

But the wavefunction  $\psi(x)$  CANNOT be ZERO everywhere

$$\sin kL = 0 \Rightarrow kL = n\pi \quad n=1,2,3,4\dots$$

Wavefunction is  $\psi(x) = A \sin kx$

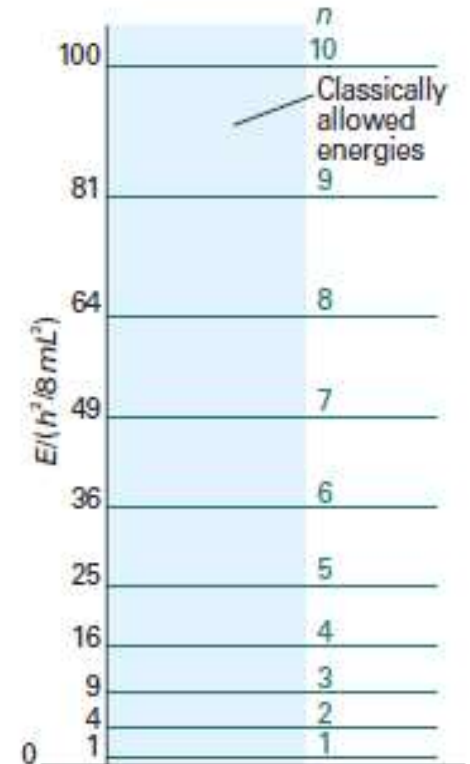
# Particle in 1-D Square-Well Potential

$$E = \frac{\hbar^2 k^2}{2m} \quad \text{and} \quad k = \frac{n\pi}{L}$$

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2} = \frac{n^2 h^2}{8mL^2} \quad n=1,2,3,4\ldots$$

Energy is no longer continuous but has discrete values; Quantization of energy

Energy separation increases with increasing values of  $n$



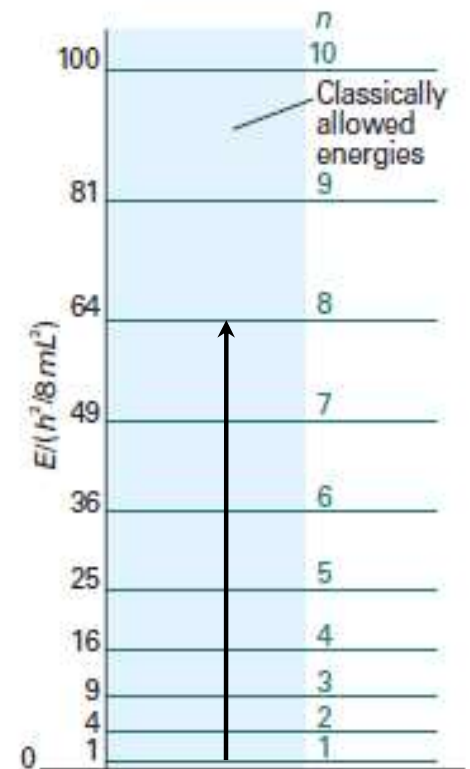
The lowest allowed energy level is for  $n=1$

$$E_1 = \frac{\pi^2 \hbar^2}{2mL^2} \quad \text{has a non zero value} \Rightarrow \text{Zero Point Energy}$$

# Particle in 1-D Square-Well Potential: Spectroscopy

$$h\nu = \Delta E = E_f - E_i = \frac{n_f^2 h^2}{8mL^2} - \frac{n_i^2 h^2}{8mL^2} = (n_f^2 - n_i^2) \frac{h^2}{8mL^2}$$

Larger the box, smaller the energy of  $h\nu$



## Particle in 1-D Square-Well Potential

Wavefunction  $\psi(x) = A \sin kx = A \sin \frac{n\pi}{L} x$

Normalization  $\int_0^L \psi^*(x) \cdot \psi(x) \cdot dx = A^2 \int_0^L \sin^2 \frac{n\pi}{L} x \cdot dx = 1$

$$A = \sqrt{\frac{2}{L}} \quad \psi(x) = \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x$$

**Homework**  
**Evaluate the above integral**

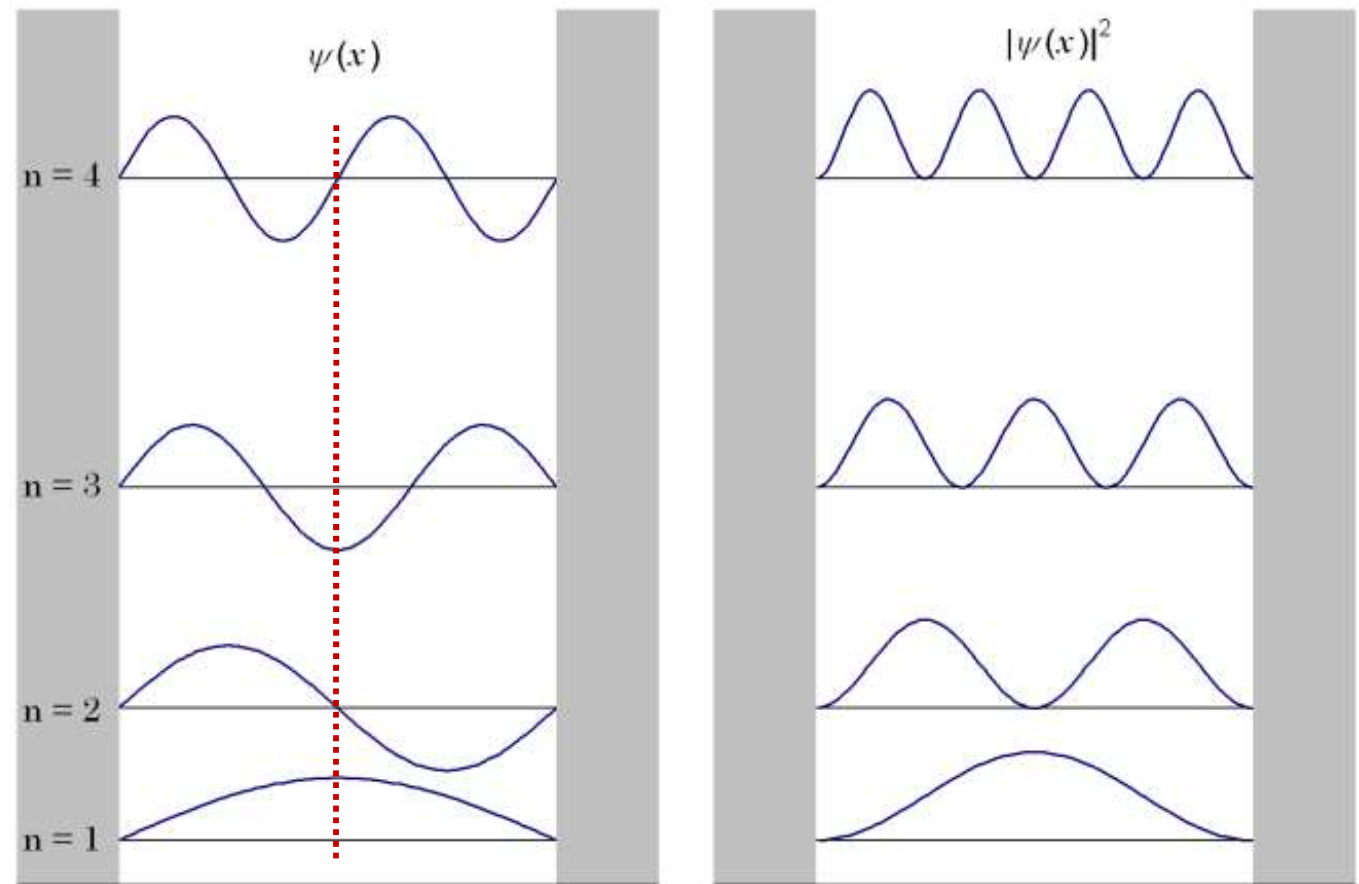
# Particle in 1-D Square-Well Potential: Spectroscopy

Wavefunction  $\psi(x) = \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x$

$n=1,3,\dots$  (odd)  
Symmetric  
(even function)

$n=2,4,\dots$  (even)  
Anti-Symmetric  
(odd function)

Number of Nodes  
(zero crossings) =  $n-1$



## Expectation values

$$\langle x \rangle = \int \psi^* \cdot x \cdot \psi \cdot dx$$

$$= \int_0^L \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x \cdot x \cdot \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x \cdot dx$$

$$= \frac{2}{L} \int_0^L x \cdot \sin^2 \frac{n\pi}{L} x \cdot dx$$

$$= \frac{L}{2}$$

**Homework  
Verify!**

## Expectation values

$$\langle p_x \rangle = \int \psi^* \cdot \left( -i\hbar \frac{\partial}{\partial x} \right) \cdot \psi \cdot dx$$

$$= -i\hbar \int_0^L \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x \cdot \frac{\partial}{\partial x} \cdot \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x \cdot dx$$

$$= \frac{-2i\hbar n\pi}{L^2} \int_0^L \sin \frac{n\pi}{L} x \cdot \cos \frac{n\pi}{L} x \cdot dx$$

$$= 0$$

**Homework  
Verify!**

## Particle in 2-D Square-Well Potential

Hamiltonian

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial y^2} = \hat{H}_x + \hat{H}_y$$

$$\hat{H} \cdot \psi(x, y) = E_n \cdot \psi(x, y)$$

Let us assume that

$$\psi(x, y) = \psi(x) \cdot \psi(y)$$



## Particle in 2-D Square-Well Potential

$$\begin{aligned}\hat{H} \cdot \psi(x, y) &= \hat{H} \cdot (\psi(x) \cdot \psi(y)) \\&= \left[ \hat{H}_x + \hat{H}_y \right] (\psi(x) \cdot \psi(y)) \\&= \psi(y) \cdot \hat{H}_x \cdot \psi(x) + \psi(x) \cdot \hat{H}_y \cdot \psi(y) \\&= \psi(y) \cdot E_x \cdot \psi(x) + \psi(x) \cdot E_y \cdot \psi(y) \\&= E_x \cdot \psi(x) \cdot \psi(y) + E_y \cdot \psi(x) \cdot \psi(y) \\&= (E_x + E_y) \cdot (\psi(x) \cdot \psi(y)) \\&= (E_x + E_y) \cdot (\psi(x, y))\end{aligned}$$

## Particle in 2-D Square-Well Potential

Hamiltonian

$$\hat{H} = -\frac{\hbar}{2m} \frac{\partial^2}{\partial x^2} - \frac{\hbar}{2m} \frac{\partial^2}{\partial y^2} = \hat{H}_x + \hat{H}_y$$

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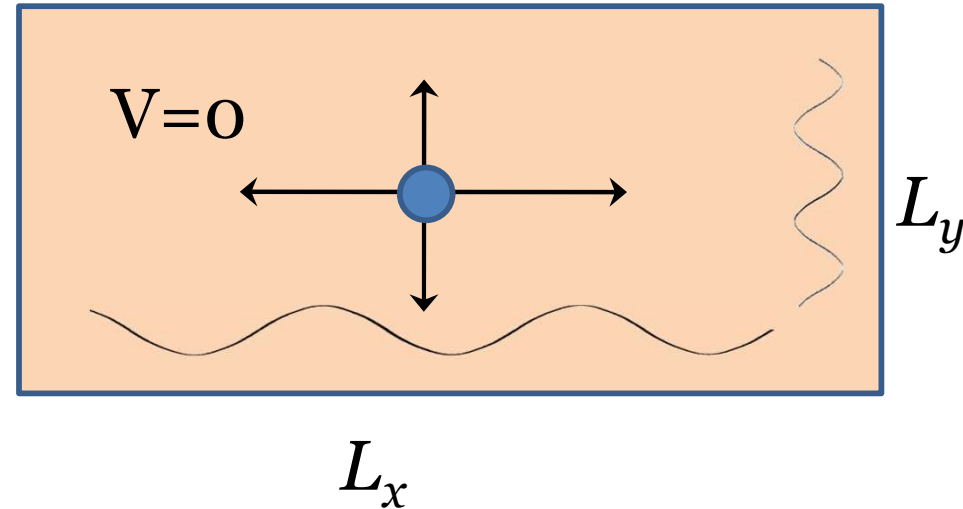
$$E_n = E_{n_x, n_y} = E_{n_x} + E_{n_y}$$

**$\psi$**  is a product of the eigenfunctions of the parts of  **$\hat{H}$**

**$E$**  is sum of the eigenvalues of the parts of  **$\hat{H}$**

# Particle in 2-D Square-Well Potential

$$\begin{aligned}\psi(x,y) &= \psi(x) \cdot \psi(y) \\ &= \sqrt{\frac{2}{L_x}} \sin \frac{n\pi}{L_x} x \cdot \sqrt{\frac{2}{L_y}} \sin \frac{n\pi}{L_y} y \\ &= \frac{2}{\sqrt{L_x L_y}} \sin \frac{n\pi}{L_x} x \cdot \sin \frac{n\pi}{L_y} y\end{aligned}$$



$$\begin{aligned}E_{n_x, n_y} &= E_{n_x} + E_{n_y} \\ &= \frac{n_x^2 h^2}{8mL_x^2} + \frac{n_y^2 h^2}{8mL_y^2} \\ &= \frac{h^2}{8m} \left( \frac{n_x^2}{L_x^2} + \frac{n_y^2}{L_y^2} \right) \quad n_x, n_y = 1, 2, 3, 4, \dots\end{aligned}$$