

# Arbeid en energie

$$v \cdot w = \frac{\vec{v} \cdot \vec{w}}{\cos \theta}$$

Arbeid bij constante kracht

$$W = F d \cos \theta = \vec{F} \cdot \vec{d}$$

$\hookrightarrow$  scalar, in Nm of J

check of  
van richting  
arbeid

Arbeid bij veranderlijke kracht

$$W \approx \sum_i F_i \cos \theta_i \Delta l_i$$

$$\leadsto W = \int_{x_1}^{x_2} F \cos \theta \, dl = \int_{x_1}^{x_2} \vec{F} \cdot d\vec{l}$$

In een veer:

$$F_p = kx$$

$$F_s = -kx$$

$$W_p = \int_0^{x_b} \vec{F}_p \cdot d\vec{x} = \int_0^{x_b} F_p \, dx$$

$$= \int_0^{x_b} -kx \, dx = \frac{1}{2} k x_b^2$$

$\theta = 0^\circ$

1. Punt massa onderhevig aan  
uitwendige kracht

2de WN:

$$m \frac{d\vec{v}}{dt} = \sum_i \vec{F}_i(\vec{r}, \vec{v}, t) \\ = \vec{F}(\vec{r}, \vec{v}, t)$$

Scaliere verm met  $\vec{v}$

$$\vec{v} \cdot m \frac{d\vec{v}}{dt} = \vec{v} \cdot \vec{F}(\vec{r}, \vec{v}, t) \\ \Rightarrow \frac{d}{dt} \left( \frac{1}{2} m \vec{v}^2 \right) = \vec{F} \cdot \vec{v} \\ \int_{t_0}^{t_1} \frac{d}{dt} \left( \frac{1}{2} m \vec{v}^2 \right) dt = \int_{t_0}^{t_1} \vec{F} \cdot \vec{v} dt \\ \underbrace{\qquad\qquad\qquad}_{\text{kin energie} = T} = \int_{t_0}^{t_1} \vec{F} \cdot \vec{v} dt \\ \qquad\qquad\qquad = \frac{m \cdot v^2}{2} \qquad \vec{v} \cdot \vec{v} = v \cdot v \cdot \cos 0^\circ = v^2$$

$$T_1 - T_0 = \int_{t_0}^{t_1} \vec{F} \cdot \vec{v} dt$$

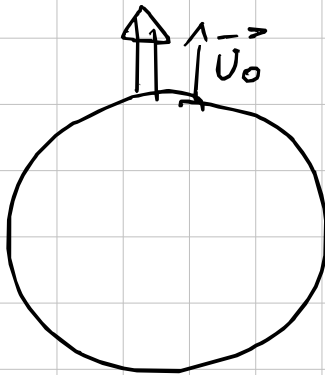
$$\begin{aligned}
 T_1 - T_0 &= \int_{t_0}^{t_1} \vec{F} \cdot \frac{d\vec{r}}{dt} \cdot dt \\
 &= \int_{r_0}^{r_1} \vec{F} \cdot d\vec{r} \\
 &= W
 \end{aligned}$$

Arbeid energie principe

arbeid =  $\Delta E_{kin}$   
 $\rightarrow$  verandering kinetische energie

Voorbeeld:

Minimale snelheid  
 om de rodel weg  
 te laten glijden



2de WN:

$$m_R \vec{a} = \vec{F}$$

$$m_R \frac{d\vec{u}}{dt} \overset{\substack{\text{1. reducing} \\ \text{du} \rightarrow \vec{u} \text{ neg}}}{=} \vec{F}_G$$

$$m_R \frac{d^2 x}{dt^2} = - \frac{G \cdot m_R \cdot M_\oplus}{(R_\oplus + x)^2}$$

$$m_R \frac{dx}{dt} \cdot \frac{d^2 x}{dt^2} = \frac{dx}{dt} \cdot "$$

$$\frac{d}{dt} \left( \frac{1}{2} m_R \left( \frac{dx}{dt} \right)^2 \right) = " \quad " \quad \frac{dx}{dt}$$

$$= \frac{d}{dt} \left( \frac{G m_R M_\oplus}{R + x} \right)$$

$$\int_{t_0}^{t_1} " dt = \int_{t_0}^{t_1} " dt$$

$$\frac{1}{2} m_R v^2 \Big|_{v_0}^{v_1} = \frac{G m_R M_\oplus}{R + x} \Big|_{x=0}^{x_E}$$

$$\frac{1}{2} m_R v_1^2 - \frac{1}{2} m_R v_0^2 = \frac{G M_R M_\oplus}{R+x} - \frac{G M_\oplus R}{R}$$

$$v_1^2 = v_0^2 + 2 \left( \frac{G M_\oplus}{R+x} - \frac{G M_\oplus}{R} \right)$$

$$v_1 \geq 0$$

$$\Rightarrow \geq 0$$

$$v_0^2 \geq \frac{G M_\oplus}{R} - \frac{G M_\oplus}{R+x}$$

$x \rightarrow \infty$

$$v_0^2 \geq 2 \frac{G M_\oplus}{R}$$

$$v_0^2 = 2 \frac{G M_\oplus}{R}$$

$$v_0 = \sqrt{\frac{2 G M_\oplus}{R}}$$

$$v_0 = \sqrt{2 g R_\oplus} \simeq 11.2 \text{ km/s}$$

Zelfde voorbeeld met Arbeids energie principe

$$\begin{aligned} T_1 - T_0 &= \int_{r_0}^{r_1} \vec{F} \cdot d\vec{r} \\ &= \int_{t_0}^{t_1} \vec{F} \cdot \frac{d\vec{r}}{dt} dt \\ &= \int_{t_0}^{t_1} \vec{F} \cdot \vec{v} \cdot dt \end{aligned}$$

$$\begin{aligned} \frac{1}{2} m_R v_1^2 - \frac{1}{2} m_R v_0^2 &= \int_{x=0}^{x_1} \frac{-G M_\oplus m_R}{(R+x)^2} dx \\ &= \frac{G M_\oplus m_R}{R+x_1} - \frac{G M_\oplus m_R}{R} \end{aligned}$$

eee >  $v_0 = \sqrt{2gR_\oplus}$

## 2. Behand van energie:

Intermezzo afgeleiden

$$f(x) = \frac{1}{2} k x^2$$
$$\frac{df}{dx} = kx$$

$$f(x, y, z) = 2x^2 y^3 z$$

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz$$
$$= 4xy^3z + 6x^2y^2z + 2x^2y^3$$

$$\nabla f = \frac{\partial f}{\partial x} \hat{i} + \frac{\partial f}{\partial y} \hat{j} + \frac{\partial f}{\partial z} \hat{k}$$
$$= 4xy^3z \hat{i} + 6x^2y^2z \hat{j} + 2x^2y^3 \hat{k}$$

$$\nabla U : d\vec{r} \quad \text{met } d\vec{r} = dx\hat{i} + dy\hat{j} + dz\hat{k}$$

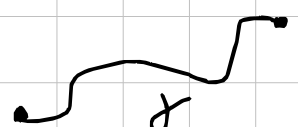
$$\hookrightarrow = \frac{\partial U}{\partial x} dx + \frac{\partial U}{\partial y} dy + \frac{\partial U}{\partial z} dz$$


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$$= dU$$

$\vec{F}(\vec{r}, t) \equiv$  conservatieve kracht (identieke energieën)

$$W = \int_{r_0}^{r_1} \vec{F} \cdot d\vec{r} = \int_{t_0}^t \vec{F} \cdot \frac{d\vec{r}}{dt} dt$$

$$= \int_{\gamma} \vec{F} \cdot d\vec{r}$$


conservatieve kracht maakt keuze van  $\gamma$  niet uit.

Bij conservatieve kracht  $\vec{F}(\vec{r}, t)$

$$\hookrightarrow \exists! U : \vec{F} = -\nabla U$$

$\hookrightarrow$  potentiële energie

$\hookrightarrow$  niet of constante na



$$W = \int_{r_1}^{r_2} \vec{F} \cdot d\vec{r}$$

$$= \int_{r_1}^{r_2} -\nabla U \cdot d\vec{r}$$

$$= - \int_{U_1}^{U_2} dU$$

$$= U_1 - U_2 = W = T_2 - T_1$$

$$\Rightarrow T_2 + U_2 = T_1 + U_1$$

$\hookrightarrow$  behoud mechanische energie

word niet behouden als er wrijving is  
bevoordeld

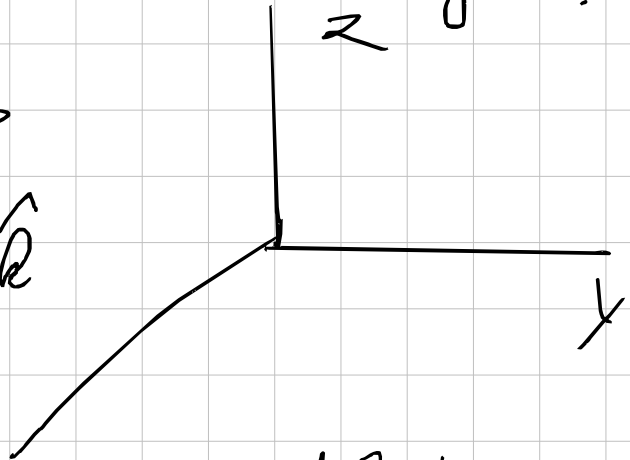
$$E^{\text{tot}} = E^{\text{mech}} + \underbrace{E^{\text{inw}}}_{\text{"warmte"}}$$

Als krachten ~~bevoordeld~~ ~~bevoordeld~~ zijn op de loopban, dan maakt het niet uit de kracht conservatief is of niet.  $E^{\text{mech}}$  zal behouden worden

Zelfde voorbeeld. Nu met wet van behoud van mechanische energie.

$$\textcircled{1} \quad \vec{F} = m \vec{g}$$

$$\vec{F} = -mg \hat{k}$$



$$U = \int \vec{F} \cdot d\vec{r} \quad \left( = dx\hat{i} + dy\hat{j} + dz\hat{k} \right)$$

$$U = - \int m g \hat{k} \cdot d\vec{r}$$

$$= - \int m g dz$$

$$= mgz + C$$

Harmonische Kraft

$$\vec{F} = -k\vec{r} = -k r \hat{r}$$

$$U = - \int \vec{F} \cdot d\vec{r}$$

$$= \int k r \hat{r} \cdot dr \hat{r}$$

$$= \int k r dr$$

$$= \frac{1}{2} k r^2 + C$$

Allgemeine Gravitationskraft

$$\vec{F}_G = - \frac{G m M}{r^2} \hat{r}$$

$$U = - \int \vec{F} \cdot d\vec{r} = \int \frac{G m M}{r^2} \hat{r} \cdot dr \hat{r}$$

$$= \int \frac{G m M}{r^2} dr = - \frac{G m M}{r} + C$$

