

Predictive Compression Dynamics: A Methodological Framework for Computable Information-Driven Modeling

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Dated: October 21, 2025

Abstract

We present *Predictive Compression Dynamics* (PCD) as a *methodological framework* for building *computable* local functionals that drive dynamics by explicit gradient flow. Rather than postulating uncomputable principles, PCD starts from an *information-motivated* surrogate codelength $\hat{\Phi}$ and specifies dynamics $\dot{x} \propto -\nabla \hat{\Phi}(x)$ with all parameters preregistered. We study two concrete classes: (i) a *fixed-graph* functional

$$\hat{\Phi}_E(x) = - \sum_{(i,j) \in E} \frac{m_i m_j}{\sqrt{\|x_i - x_j\|^2 + a^2}},$$

and (ii) a *smooth kernel* functional

$$\hat{\Phi}_K(x) = - \sum_{i < j} m_i m_j K_\sigma(\|x_i - x_j\|),$$

with K_σ compactly supported and C^1 . The first provides a transparent sandbox; the second avoids neighbor-set discontinuities and guarantees well-posedness. In the dilute two-body limit both produce an attractive $1/r^2$ *form* (after calibration), while the softening a regularizes collisions. These are *operational, information-motivated models*, not new claims about gravity; familiar “Plummer-like” roll-offs serve here as *sanity checks*, not phenomenological predictions. We give preregistered algorithms, an explicit model–data codelength decomposition, and clear falsification criteria for model instances. The framework supports reproducible exploration of *computable* compression-driven dynamics across domains.

1 Positioning and Commitments

What this paper is: a disciplined way to construct and test *computable* local functionals $\hat{\Phi}$ whose gradients define dynamics. We make code-level, preregistered choices explicit (domains, discretizations, kernels, parameters) and provide sanity checks.

What this paper is *not*: a claim of novel gravitational phenomenology, nor a derivation of GR/QM. Any resemblance of simple pair terms to familiar softened potentials is acknowledged and used only as a controlled testbed.

2 Operational Domain and Notation

We consider N point agents with positions $x_i \in \mathbb{R}^3$ and positive weights m_i (“masses”). Computations occur at finite precision: lattice spacing a_{grid} and b bits per coordinate stated *a priori*. A global calibration constant $G_{\text{eff}} > 0$ turns the dimensionless force into physical units.

3 Information-Motivated Functionals

Model–data decomposition (explicit). Following standard MDL reasoning, a description length splits as

$$L_{\text{tot}} = L(M) + L(D | M), \quad (3.1)$$

where $L(M)$ encodes the model/regularities and $L(D | M)$ the residual data given that model. A decrease $\Delta L_{\text{tot}} < 0$ corresponds to *realized compression*. In PCD we treat a computable, local $\hat{\Phi}$ as a proxy for the *negative* of an achievable ΔL_{tot} ; gradient descent in $\hat{\Phi}$ heuristically implements compression.

3.1 Fixed-graph functional

Let $E \subset \{(i, j) : 1 \leq i < j \leq N\}$ be a symmetric, degree-bounded edge set (e.g. an initial k NN graph). Define

$$\hat{\Phi}_E(x) = - \sum_{(i,j) \in E} \frac{m_i m_j}{\sqrt{\|x_i - x_j\|^2 + a^2}}, \quad a > 0. \quad (3.2)$$

Sign convention: the minus sign ensures attraction under gradient descent. Softening a regularizes collisions. The force on i is

$$\begin{aligned} F_i^{(E)}(x) &= -\nabla_{x_i} \hat{\Phi}_E(x) \\ &= - \sum_{\substack{j: \\ (i,j) \in E}} m_i m_j \frac{(x_i - x_j)}{(\|x_i - x_j\|^2 + a^2)^{3/2}}. \end{aligned} \quad (3.3)$$

In the two-body, $a \rightarrow 0$ limit with $(i, j) \in E$,

$$F_i^{(E)} \rightarrow -m_i m_j \frac{x_i - x_j}{\|x_i - x_j\|^3}, \quad (3.4)$$

i.e. an attractive $1/r^2$ *form* along the active edge. For physical units we use $F_i^{\text{phys}} = G_{\text{eff}} F_i^{(E)}$.

Remark (edge modeling). Forces vanish between pairs not in E ; this is a *modeling choice* for locality/sparsity, not a universal law.

3.2 Smooth-kernel functional

To avoid kNN neighbor-set discontinuities, choose a compactly supported, C^1 radial kernel $K_\sigma : [0, \infty) \rightarrow \mathbb{R}_{\geq 0}$ with support $[0, R\sigma]$ (e.g. a Wendland C^2 kernel). Define

$$\hat{\Phi}_K(x) = - \sum_{i < j} m_i m_j K_\sigma(\|x_i - x_j\|). \quad (3.5)$$

Then $F_i^{(K)}(x) = -\nabla_{x_i} \hat{\Phi}_K(x)$ is continuous and locally Lipschitz (away from collisions if $K'_\sigma(0)$ finite). For $K_\sigma(r) \sim (r^2 + a^2)^{-1/2}$ near zero one recovers the same regularized two-body *form* as (3.4).

3.3 A minimal codelength view

Interpreting $-\hat{\Phi}$ as a surrogate *gain in description length* when nearby pairs are represented jointly, if the expected residual codelength per pair scales like a decreasing $\ell(r_{ij})$, then

$$L_{\text{tot}} \approx \text{const} + \sum_{(i,j)} \ell(r_{ij}), \quad \Rightarrow \quad \hat{\Phi} \propto - \sum_{(i,j)} \ell(r_{ij}).$$

We refrain from claiming exact MDL optimality; $\hat{\Phi}$ is an *operational proxy* to be tested.

4 Dynamics and Preregistration

All parameters ($a_{\text{grid}}, b, a, \sigma, \Delta t, m_i, \gamma, T, \text{seeds}$) are fixed *before* experiments.

Deterministic gradient flow.

$$x_i^{(t+\Delta t)} = x_i^{(t)} + \Delta t F_i(x^{(t)}), \quad F_i \in \{F_i^{(E)}, F_i^{(K)}\}. \quad (4.1)$$

Underdamped Langevin.

$$m_i \ddot{x}_i = F_i(x) - \gamma \dot{x}_i + \xi_i(t), \quad \langle \xi_i(t) \xi_j(t') \rangle = 2\gamma k_B T \delta_{ij} \delta(t - t'). \quad (4.2)$$

5 Sanity Checks (not phenomenology)

Two-body limit. With $a \rightarrow 0$ and a single interacting pair, (3.4) holds (after calibration G_{eff}).

Softening expansion (corrected). For $r = \|x_i - x_j\| \gg a$,

$$\frac{r}{(r^2 + a^2)^{3/2}} = \frac{1}{r^2} \left(1 - \frac{3a^2}{2r^2} + O\left(\frac{a^4}{r^4}\right) \right), \quad (5.1)$$

hence

$$\|F_i^{(E)}\| = m_i m_j \frac{r}{(r^2 + a^2)^{3/2}} \stackrel{r \gg a}{\approx} m_i m_j \frac{1}{r^2} \left(1 - \frac{3a^2}{2r^2} + \dots \right). \quad (5.2)$$

We use (5.1) solely to verify numerics; *no claim* is made that nature exhibits such a roll-off at any laboratory scale.

6 Well-posedness

For $a > 0$ and bounded degree, $\hat{\Phi}_E \in C^1(\mathbb{R}^{3N} \setminus \{x_i = x_j\})$ and $F^{(E)}$ is locally Lipschitz away from collisions; similarly $F^{(K)}$ is continuous and locally Lipschitz for C^1 kernels with bounded derivative at 0. If $k\text{NN}$ is used with dynamic neighbors, forces are only piecewise smooth; then employ (i) hysteresis for neighbor swaps, (ii) temporal smoothing of weights, or (iii) prefer the smooth kernel model (3.5).

7 A Preregistered Model Card (example)

Domain. $a_{\text{grid}} = 10 \mu\text{m}$, $b = 16$.

Functional. $\hat{\Phi}_K$ with Wendland C^2 kernel of scale $\sigma = 0.5 \text{ mm}$; softening $a = 50 \mu\text{m}$.

Dynamics. Underdamped with $(m_i \equiv 1, \gamma = 0.1, T = 300 \text{ K})$, $\Delta t = 1 \times 10^{-3} \text{ s}$.

Calibration. G_{eff} fit once in a dilute two-body sandbox at $r \gg a$.

Sanity checks. Verify (3.4) form and (5.1) numerically; report seeds and residuals.

8 What to Falsify (for a given instance)

Given fixed $(\hat{\Phi}, \text{params})$, the instance is falsified if:

- (F1) Two-body trajectories disagree with the calibrated $1/r^2$ form beyond stated error under identical numerical conditions.
- (F2) Kernel vs. fixed-graph variants yield statistically inconsistent small- r behavior not explainable by topology (violates within-class robustness).
- (F3) The chosen code-length proxy (e.g. $-\hat{\Phi}$ vs. an external structural complexity metric) fails to correlate in controlled tests.

These are *model-level* falsifiers; they do not purport to test fundamental physics.

9 Discussion and Scope

PCD supplies a reproducible route from *computable* information-motivated functionals to concrete dynamics. Simple pairwise terms coincide with well-known softened interactions; this is a feature for validation, not a novelty claim. The framework can be extended to richer $\hat{\Phi}$ (e.g. graph Laplacians, learned local codes) with the same preregistration discipline.

Acknowledgments

We thank colleagues for discussions on local estimators, kernels, and N -body numerics. Any prior drafts suggesting phenomenology are superseded by this methodological formulation.

References

References

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A Schematic Field (single edge)



Figure 1: Attractive field along the active edge for (3.2); magnitude regularized near particles by a .

B Units and the corrected far-field

With x in length units, (3.2) has units of inverse length. The calibrated constant G_{eff} supplies force units so that $F^{\text{phys}} = G_{\text{eff}}F$. The large- r expansion used for numerical validation is (5.1) and (5.2).