

# **Motion-Insensitive Rapid Configuration Relaxometry**

Damien Nguyen<sup>1,2</sup>\* and Oliver Bieri<sup>1,2</sup>

**Purpose:** Triple echo steady state (TESS) uses the lowest steady state configuration modes for rapid relaxometry. Due to its unbalanced gradient scheme, however, TESS is inherently motion-sensitive. The purpose of this work is to merge TESS with a balanced acquisition scheme for motion-insensitive rapid configuration relaxometry, termed MIRACLE.

**Methods:** The lowest order steady state free precession (SSFP) configurations are retrieved by Fourier transformation of the frequency response of N frequency-shifted balanced SSFP (bSSFP) scans and subsequently processed for relaxometry, as proposed with TESS. Accuracy of MIRACLE is evaluated from simulations, phantom studies as well as in vivo brain and cartilage imaging at 3T.

**Results:** Simulations and phantom results revealed no conceptual flaw, and artifact-free configuration imaging was achieved in vivo. Overall, relaxometry results were accurate in phantoms and in good agreement for cartilage and for  $T_2$  in the brain, but apparent low  $T_1$  values were observed for brain white matter; reflecting asymmetries in the bSSFP profile.

**Conclusion:** Rapid  $T_1$  and  $T_2$  mapping with MIRACLE offers analogous properties as TESS while successfully mitigating its motionsensitivity. As a result of the Fourier transformation, relaxometry becomes sensitive to the voxel frequency distribution, which may contain useful physiologic information, such as structural brain integrity. © 2016 International Society for Magnetic Resonance in Medicine. **Magn Reson Med 78:518–526, 2017.** © **2016 International Society for Magnetic Resonance in Medicine** 

**Key words:** T<sub>1</sub> mapping; T<sub>2</sub> mapping; balanced Steady State Free Precession (bSSFP); relaxometry

#### INTRODUCTION

Quantification is thought to represent an important step toward significant improvements of the diagnostic potential of MRI, such as for the early detection of subtle or diffuse pathological changes with high specificity and sensitivity, for an unbiased assessment of treatment or drug effects, as well as for clinical trials in drug research across different sites and machines. Quantitative imaging, however, is rather time consuming and typically becomes an issue in the clinical environment, where the overall success and applicability of quantitative MRI strongly depends on the overall acquisition speed. In

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this context, SSFP-based imaging techniques (1) have shown compelling results thanks to their short scan times and high signal-to-noise ratios, e.g., for relaxation time mapping (2–8), for measuring molecular proton diffusion (9–13), for the assessment of magnetization transfer effects (14–16), or for the characterization of flow or motion (17–19).

Generally, relaxation occupies a central role within the context of NMR: it not only defines contrast in conventional MRI but also reflects the interaction of water on a molecular level. Historically, longitudinal relaxation  $(T_1)$ has been estimated by sampling the inversion-recovery curve of the longitudinal magnetization using spin-echo sequences, while the transverse relaxation  $(T_2)$  time has been estimated from the decay curve of the transverse magnetization using single-echo or, more frequently, multiecho spin-echo (SE) methods. Quantification based on the functional dependencies of the steady state, however, is much faster. One common attribute of SSFP methods is their mixed  $T_2/T_1$  imaging contrast (20); being a natural consequence of a pulse repetition time (TR) that is much shorter than  $T_2$ . Consequently, accurate quantification of relaxation times using SSFP-based imaging techniques is usually hampered by a  $T_2$ -related bias in  $T_1$  estimates for radio-frequency (RF) spoiled SSFP (21,22), or by a  $T_1$ related bias in  $T_2$ , as observed with balanced SSFP (bSSFP) (23), partially spoiled SSFP (8), and double echo steady SSFP (7). Moreover, all previously mentioned methods are sensitive to transmit field  $(B_1)$  inhomogeneities, whereas some of them show, in addition, some sensitivity to off-resonances (24) or motion (7,8).

Recently, Heule et al (25), proposed to tackle the mutual interference of  $T_1$  and  $T_2$  of coherent SSFP methods by using a triple echo steady state (TESS) imaging approach. To this end, TESS acquires the three lowest SSFP configuration modes within a single acquisition (or TR) yielding two independent ratios for simultaneous rapid quantification of both  $T_1$  and  $T_2$  using a golden section search. Quite remarkably, TESS achieves an almost completely  $B_1$ -unbiased estimation of  $T_2$ , and showed good prospects for rapid threedimensional (3D)  $T_2$  mapping of articular cartilage imaging in the clinical setting (26). Generally, TESS is also insensitive to  $B_0$  inhomogeneities but the nonbalanced gradient scheme introduces some motion sensitivity that can be softened with a single-slice version of TESS providing high quality  $T_2$  maps in the human brain even at ultra-high fields (27). More recently, a simultaneous multislice TESS sequence (28) was proposed to decrease scan time without loss of image quality.

In this work, we aim to replace the unbalanced gradient scheme of TESS by a balanced one leading to "motion-insensitive rapid configuration relaxometry", termed MIR-ACLE (29), that indirectly retrieves the basic SSFP modes from a series of bSSFP scans. The accuracy of MIRACLE-based relaxometry is evaluated from simulations and

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phantom experiments. The feasibility of high-resolution volumetric  $T_1$ - and  $T_2$ -mapping is demonstrated in vivo for the human brain and for articular cartilage at 3T.

#### **METHODS**

#### **MIRACLE**

In the following analysis, we consider an equidistant train of RF pulses with constant flip angles  $\alpha$  and constant RF phase increment φ in combination with balanced gradient moments. Finite RF pulse and diffusion effects are considered to be negligible. Immediately after the RF pulse (counterclockwise rotation around x-axis), the complex steady state magnetization  $M_{+}(t=0)$  is given by (e.g., following Ganter) (30)

$$egin{align*} M_{+}(t=0) &\equiv M_{X}(t=0) + i M_{Y}(t=0) \ &= -rac{i}{D} \cdot (1-E_{1}) \mathrm{sin} lpha (1-E_{2}e^{-i artheta}) \ D := (1-E_{1} \mathrm{cos} lpha) (1-E_{2} \mathrm{cos} artheta) - (E_{1} - \mathrm{cos} lpha) (E_{2} - \mathrm{cos} artheta) E_{2} \ \end{array}$$

where  $E_i := \exp(-TR/T_i)$  and  $\vartheta := \Phi - \varphi$  denotes the phase difference between the off-resonance related phase Φ accumulated during each repetition time (TR) interval (and is assumed to be constant in time), and the RF pulse phase increment  $\varphi$ . At a time t after the RF pulse, the magnetization is given by

$$M_{+}(t) = M_{+}(t=0)e^{-t/T_2}e^{i\phi(t)}$$
 [3]

where  $\phi(t) := t/TR \cdot \Phi$ . Generally, the steady state as given in Equations [1,2] is periodic in  $\varphi$  and can be expressed as a sum over all configuration orders  $F_n$ (30,31):

$$M_+(arphi,t) = e^{-t/T_2} e^{i\Phi t/{
m TR}} \sum_{n=-\infty}^{+\infty} F_n e^{in(\Phi-arphi)} \; .$$
 [4]

As shown previously in a seminal work by Zur et al (31), it is possible to retrieve the basic SSFP configurations or modes  $F_n$  (see Eq. [4]) from a discrete Fourier transformation of the complex bSSFP frequency response. To this end, we proceed as follows:

(1) We perform N scans with an RF phase increment of

$$\varphi_j := -\frac{2\pi}{N}j, \ j = 0, 1, \dots N - 1$$
 [5]

where i enumerates the scan. As a result, the magnetization in Equation [4] is modified to take the form

$$egin{aligned} M_+(arphi_j,t) &= \sum_{n=-\infty}^{+\infty} e^{-t/T_2} \ e^{i\Phi t/\mathrm{TR}} \ F_n \ e^{in\Phi} e^{-inarphi_j} \ &\equiv \sum_{n=-\infty}^{+\infty} M_n(\Phi,t) e^{i\left(rac{2\pi}{N}
ight)nj} \ . \end{aligned} \qquad [6]$$

(2) We calculate the N-point Fourier transform of the  $M_{+}(\varphi_{i},t)$  magnetizations

$$G(p,t) = \frac{1}{N} \sum_{j=0}^{N-1} M_{+}(\varphi_{j}, t) e^{-i\left(\frac{2\pi}{N}\right)jp}$$

$$= \frac{1}{N} \sum_{j=0}^{N-1} \sum_{n=-\infty}^{+\infty} M_{n}(\Phi, t) e^{i\left(\frac{2\pi}{N}\right)j(n-p)}$$
[7]

and because

$$\sum_{j=0}^{N-1} e^{i\left(\frac{2\pi}{N}\right)j(n-p)} = \begin{cases} N, & \text{if } \frac{n-p}{N} \text{is an integer} \\ 0, & \text{otherwise} \end{cases}$$
[8]

this yields

$$G(p,t) = e^{-t/T_2} e^{i\phi(t)} \Big( F_p e^{ip\Phi} + F_{p\pm N} e^{i(p\pm N)\Phi} + F_{p\pm 2N} e^{i(p\pm 2N)\Phi} + \dots \Big).$$
 [9]

As a result of the finite number of scans (N), aliasing occurs (see Eq. [9]). Generally, however, the mode amplitudes decrease rapidly with increasing mode order |p|and thus for large enough N,

$$G(p,t) \approx e^{-\frac{t}{T_2}} e^{i\phi(t)} F_p e^{ip\Phi}$$
. [10]

And, therefore,

$$|G(p,t)| \approx e^{-\frac{t}{T_2}} |F_p|. \tag{11}$$

Following the approach of Heule et al (25), estimation of the relaxation times  $T_1$  and  $T_2$  was then performed based on a golden-section search algorithm (32) using the two signal ratios

$$\frac{|G(1,t)|}{|G(0,t)|} \approx \frac{|F_1|}{|F_0|}$$
 [12]

$$\frac{|G(1,t)|}{|G(0,t)|} \approx \frac{|F_1|}{|F_0|}$$

$$\frac{|G(-1,t)|}{|G(0,t)| - |G(1,t)|} \approx \frac{|F_{-1}|}{|F_0| - |F_1|}$$
[13]

in an iterative approach relying on the fact that both ratios show different dependencies on  $T_1$  and  $T_2$  as was shown in Heule et al (25). From Equation [10] and in the limit where the amplitudes of higher-order  $F_n$  modes are negligible (i.e., for large enough N), only the phase of the SSFP modes depends on the local off-resonance  $\Phi$ . As a result, the signal ratios used for estimation of  $T_1$  (cf. Eq. [12]) and  $T_2$  (cf. Eq. [13]) are off-resonance insensitive. A global initial estimate of  $T_1 = 1000$  ms was used for all simulations and experiments and a tolerance of 0.1 ms was assumed as convergence criteria for the iterative search. Similarly to TESS relaxometry, we expect  $T_1$  estimates to have some remaining  $B_1$ -bias, while  $T_2$  is anticipated to be largely free of any  $B_1$  bias.

## Simulations and Imaging

All numerical simulations, data analysis and visualizations were done using MATLAB 8.5 (The MathWorks

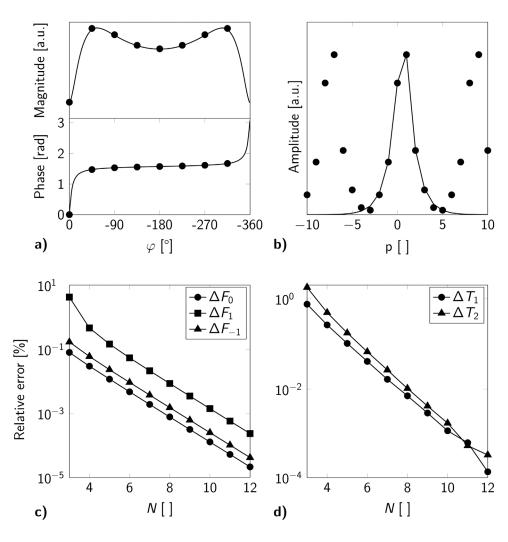


FIG. 1. a: Simulation of bSSFP steady state signal, as a function of the RF phase increment  $\varphi$  (simulation parameters:  $T_1$ /  $T_2/TR = 860/70/5.76$  ms,  $\alpha = 15^{\circ}$ , proposed optimal for TESS imaging (25)). **b**: Derived modes  $F_p$ from N-point Fourier transform using Equation [10] (dots: N=8RF phases  $\varphi_i$ , see Equation [5]; solid line: continuous RF φ; a continuous line is shown in 1b that connects the derived mode amplitudes for improved visualization). Relative mode (c) and relaxation estimation error (d) as a function of the number N of bSSFP scans performed using the same parameter set as above.

Inc., Natick, MA). Measurements and calibrations were performed on a clinical 3 Tesla (T) whole body system (Siemens MAGNETOM Prisma, Erlangen, Germany) with

actively shielded magnetic field gradient coils. Acquisitions were performed using the standard 20-channel head coil and a 15-channel Tx/Rx knee coil. Gibb's

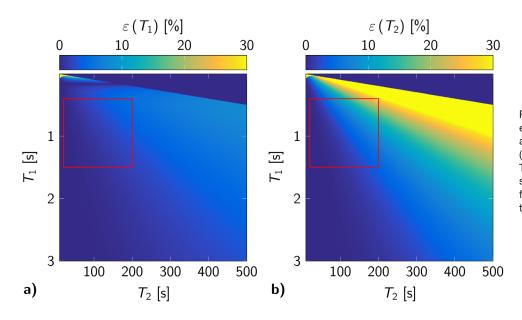


FIG. 2. MIRACLE estimation error  $\epsilon(T_i) = |T_i|_{\text{MIRACLE}} - T_i|_{\text{as}}$  a function of the simulated  $T_1$  (a) and  $T_2$  (b) for a fixed N = 8. The delimited region (red) shows typical parameter ranges for human tissues (for simulation parameters: see Figure 1).

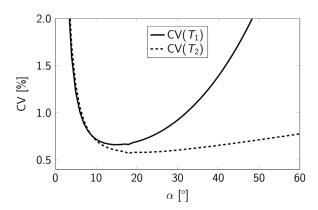


FIG. 3. Coefficient of variation (CV) for  $T_1$  and  $T_2$  as a function of the flip angle  $\alpha$ , estimated by a Monte-Carlo simulation for a N = 8 MIRACLE acquisition with 1% noise added to the bSSFP signal (for simulation parameters: see Figure 1).

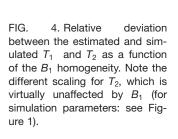
ringing was removed from base data as recently proposed by Kellner et al (33).

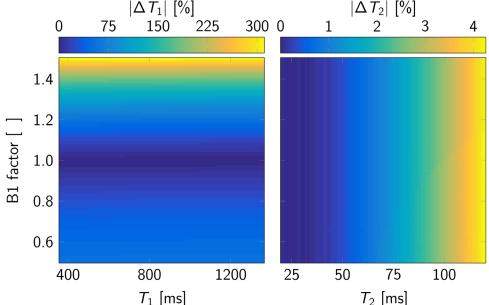
Simulations of the bSSFP signal were performed for a  $T_1/T_2$ -ratio  $\sim$  12 using sets of 4 to 12 scans with a TR/ echo time (TE) of 5.76/2.88 ms and RF phase increments  $\varphi$ , as given by Equation [5]. Evaluation of the accuracy of the proposed method for a range of  $T_1$  and  $T_2$  with an eight-point sampling scheme was performed using the same set of parameters. The influence of the flip angle on the estimation of the relaxation parameters was evaluated by performing a Monte-Carlo simulation with 100,000 independent runs for a N=8 phase cycling scheme with flip angles between  $1^{\circ}$  and  $60^{\circ}$  and the addition of 1% Gaussian white noise to simulate experimental conditions. For each flip angle, the coefficient of variation (CV) for both relaxation parameters was calculated (defined as the ratio of the standard deviation and the mean over all runs;  $CV(T_i) := \Delta T_i / \overline{T_i}, i = 1, 2$ ).

Phantom experiments were performed on a manganese-doped spherical phantom composed of  $0.125\,\mathrm{mM}$  MnCl<sub>2</sub>

dissolved in water (with nominal  $T_1/T_2$  values of 860/70 ms (25) and approximately 14 cm in diameter using an eight-point 3D MIRACLE scheme with the following protocol parameters: TR/TE = 5.10/2.55 ms,  $\alpha = 15^{\circ}$ , a resolution of  $1 \times 1 \times 2 \text{ mm}^3$  (image matrix:  $208 \times 162 \times 80$ ), a bandwidth of 401 Hz/px, elliptical scanning and RF phase increments  $\varphi$  as given by Equation [5]. With this setup, an overall scan time of approximately 7 min was obtained. The 3D TESS relaxometry was realized using the same parameters, except for TR/TE 6.19/3.17 ms, 2 averages and a bandwidth of 800 Hz/px for a total scan time of 8 min. Reference  $T_1$  relaxometry was realized using a single-slice inversion-recovery turbo spin-echo (IR-TSE) sequence with a TR/TE of 5000/13 ms, and inversion times (TI) of 50, 100, 200, 400, 800, 1600, 3200 ms. Estimation of  $T_1$  was obtained from a nonlinear fit of the recovery curve (34). Further scanning parameters include turbo factor (TF) 7,  $\alpha = 180^{\circ}$ , an in-plane resolution of 1 mm<sup>2</sup> (image matrix:  $208 \times 168$ ), slice thickness 2 mm, a bandwidth of  $130 \, \text{Hz}/$ px, GRAPPA 2 with 34 reference lines for a scan time of 1:17 min (8:59 min total). Furthermore, reference  $T_2$  relaxometry was performed using single echo spin-echo (SE) sequence with a TR of 1500 ms, and a TE of 10, 20, 40, 80, 150, 250 ms,  $\alpha = 180^{\circ}$ , an in-plane resolution of 1 mm<sup>2</sup> (image matrix: 208 × 168), slice thickness 2 mm, a bandwidth of 201 Hz/px, GRAPPA 2 with 34 reference lines resulting in 1:30 min/acquisition (9:00 min total). T2 mapping was performed using a maximum likelihood estimator approach, as described by Golub et al (35).

Exemplary in vivo human brain imaging using a 3D slab of 36 axial slices located inside the brain was performed using a 12-points bSSFP cycling scheme with a TR/TE of 5.10/2.55 ms,  $\alpha=15^{\circ}$ , a resolution of  $1\times1\times2\,\mathrm{mm}^3$  (image matrix:  $192\times150\times36$ ), a bandwidth of 400 Hz/px, elliptical scanning and RF phase increments  $\phi$  following Equation [5], resulting in a total scan time of approximately 5:30 min. Reference  $T_1$  data were acquired using a 2D single-slice IR-TSE with identical





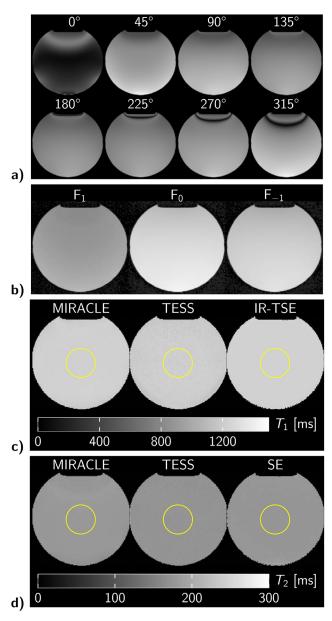


FIG. 5. Illustration of MIRACLE relaxometry calculations using an N=8 phase cycling scheme for a manganese-doped probe. **a**: Source bSSFP magnitude images. **b**: Derived three lowest SSFP mode images  $F_1$ ,  $F_0$ ,  $F_{-1}$ . **c**: Estimated  $T_1$  and  $T_2$  maps (initial guess of  $T_1=1$ s, precision enforced: 0.1 ms). **d**: Relaxation parameters were assessed for a region of interest, as indicated in the relaxation maps by the circles. (imaging parameters: N=8 of  $1 \times 1 \times 2$  mm<sup>3</sup> and TR/TE 5.00/2.50 ms).

sequence parameters as for the phantom scans, but with a resolution of  $1\,\mathrm{mm}^2$  (image matrix:  $192\times150$ ) leading to a scan time of 1:12 min (8:24 min total). Additionally, reference  $T_2$  data was obtained from a single echo 2D SE scan with  $1\,\mathrm{mm}^2$  in-plane resolution (image matrix:  $192\times150$ ) but otherwise identical to the phantom case, resulting in 1:57 min/acquisition (11:42 min total). High resolution images of the right knee were realized in axial orientations using a similar protocol with a TR/TE of 6.46/3.23 ms, a resolution of  $0.6\times0.6\times3\mathrm{mm}^3$  (image matrix:  $368\times220\times20$ ), and a bandwidth of  $300\,\mathrm{Hz/px}$  and elliptical scanning yielding a scan time of around 5:40 min.

Table 1
Estimated Relaxation Parameters from 3D MIRACLE, 3D TESS, and 2D Single Slice Reference Measurements (IR-TSE and SE) in Phantom and In Vivo

Tissue	Method	T <sub>1</sub> [ms]	T <sub>2</sub> [ms]
Phantom	MIRACLE	835 ± 16	70 ± 2
	TESS (3D)	$823 \pm 38$	$65 \pm 9$
	Reference	$850 \pm 9$	$67 \pm 1$
Brain white matter	MIRACLE	$532 \pm 56$	$44 \pm 5$
	Reference	$840 \pm 28$	$51 \pm 2$
Brain gray matter	MIRACLE	$1061 \pm 169$	$63 \pm 12$
	Reference	$1352 \pm 69$	$57 \pm 3$
Cartilage	MIRACLE	$1194 \pm 436$	$42 \pm 9$
Muscle	MIRACLE	$846 \pm 130$	$33 \pm 10$
Fat	MIRACLE	$307 \pm 23$	102 ± 13

For each series of scans, additional patient-specific  $B_1$  maps were acquired using the method proposed by Ganter et al (36). Registration of  $B_1$  data onto the on-resonant

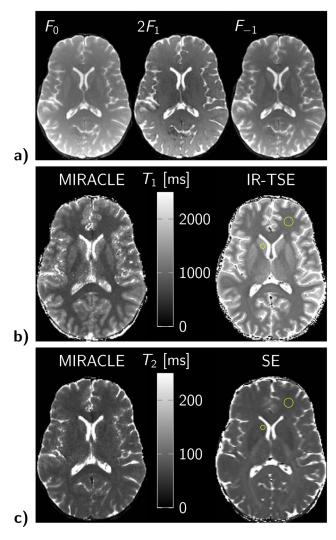


FIG. 6. Illustrative volumetric MIRACLE brain imaging of a healthy volunteer. **a**: Axial sample images from the derived three lowest SSFP mode volumes  $F_0, F_1, F_{-1}$  (note the different scaling for  $F_1$ ). **b**: Corresponding  $T_1$  (ms) map. **c**: Corresponding  $T_2$  (ms) map. (Imaging parameters: N = 12, resolution:  $1 \times 1 \times 2 \text{ mm}^3$ , TR/TE = 5.76/2.88 ms)

FIG. 7. Illustrative  $T_1$  and  $T_2$  maps (in ms), as derived from a 3D axial MIRACLE knee scan of a healthy volunteer. Note the typical decrease in both  $T_1$  and  $T_2$  for patellar cartilage from superficial to deep layers. (Imaging parameters: N=12, resolution:  $0.6 \times 0.6 \times 3 \, \text{mm}^3$  TR/TE =  $6.46/3.23 \, \text{ms}$ )

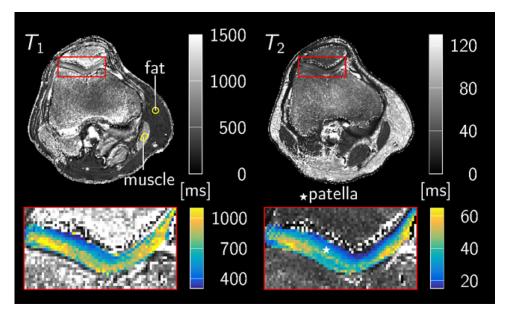


image of each MIRACLE dataset was accomplished before  $B_1$  correction using routines from the FSL libraries (37,38) for brain images or by using the elastix registration program (39,40) in other cases.

#### **RESULTS**

A noise-free simulation of the complex bSSFP frequency response (cf. Eqs. [1,2]) is presented in Figure 1a for a tissue mimicking a  $T_1/T_2$ -ratio of approximately 12 for a N=8 point phase cycling scheme and in the limit of a continuous RF phase increment. The corresponding Fourier transforms (cf. Eq. [9]) are shown in Figure 1b. Aliasing in the case of a finite N becomes evident and generally leads to a systematic deviation between the true mode amplitudes and the ones using N-point Fourier transform; especially for higher configuration orders (p) (see Figure 1b). As expected, increasing the number N of RF phases  $\varphi_i$ improves the overall accuracy of the mode amplitude estimates (cf. Figure 1c), and do not exceed 1.2% for the lowest order modes (p = -1, 0, 1) for N  $\geq$  8. The resulting  $T_1$ and  $T_2$  estimation error follows a similar trend (cf. Figure 1d), and is in the absence of noise below 0.01% for a  $T_1/T_2 \sim 12$  in combination with phase cycling acquisition schemes using  $N \ge 8$ .

Analysis of the estimation error  $\epsilon(T_i) := |T_i|_{\text{MIRACLE}}$   $-T_i|_{\text{N}}$ , as a function of  $T_1$  and  $T_2$ , is presented in Figure 2 for a fixed N=8 MIRACLE acquisition scheme. The overall error is in general much larger for  $T_2$  than for  $T_1$ , indicating the need to acquire more than 8 phase-cycles for  $T_2$  values exceeding roughly 100 ms.

Generally, the decay of the modes  $(F_p)$  with increasing mode order |p|, e.g., as observed in Figure 1b, not only depends on the relaxation, but also on the flip angle  $\alpha$  (cf. Eqs. [2,12] in Ganter) (41). As a result, different flip angle settings might be required for MIRACLE as compared to TESS. In line with the flip angle optimization results for TESS, however, Monte-Carlo simulations also revealed for MIRACLE a clear optimum for flip angles around 15  $^{\circ}$  (see Figure 3). Consequently, this flip angle

was used for all subsequent acquisitions. Moreover, in complete analogy to TESS imaging, the estimation of  $T_2$  is found to be largely insensitive to  $B_1$ , whereas the  $T_1$  estimates retain the usual  $B_1$  dependency. The bias in both  $T_1$  and  $T_2$ , as introduced by flip angle miscalibrations is shown in Figure 4. For a fixed  $T_1$ , the error in  $T_1$  scales approximately linearly with relative  $B_1$  values ranging from 0.5 to 1.5, whereas the bias in  $T_2$  from  $B_1$  (for fixed  $T_1$  values) is less than 0.1 ms over the whole investigated  $B_1$ -range.

Motion-insensitive SSFP mode imaging and subsequent relaxometry is now exemplarily illustrated in Figure 5 at 3T for a manganese-doped spherical probe. Figure 5a depicts the original N=8 phase-cycled bSSFP images, Figure 5b shows the resulting mode images for  $F_0$ ,  $F_1$ , and  $F_{-1}$ , and Figures 5c,d present the estimated  $T_1$  and  $T_2$  maps from 3D MIRACLE, 3D TESS, IR-TSE, and SE measurements. Both MIRACLE- $T_2$  and TESS- $T_2$ exhibit no spatial variations across the whole imaging volume. MIRACLE- $T_1$  and TESS- $T_1$  estimates have already been corrected for  $B_1$  inhomogeneities by using the separately acquired  $B_1$  map, which results in flat  $T_1$ profiles over the whole field-of-view. Overall, very good agreement is found between all three methods for both  $T_1$  and  $T_2$ . Evaluation of relaxation parameters within the indicated region-of-interest (ROI) are collected in Table 1 for all three methods.

High resolution in vivo 3D brain  $T_1$  and  $T_2$  mapping with MIRACLE and reference measurements is demonstrated in Figure 6 at 3T in axial slice orientation. In contrast to 3D TESS, no pulsation artifacts are noticeable in the derived base mode images (Fig. 6a and Supporting Fig. S1, which is available online, for 3D TESS brain images) and, as a result, successful  $T_1$  and  $T_2$  mapping is demonstrated in 3D (Figs. 6b,c) using motion-insensitive SSFP (31). For comparison, reference IR-TSE and SE results are also presented for the same slice. Evaluation of relaxation parameters inside highlighted ROIs can be found inside Table 1.

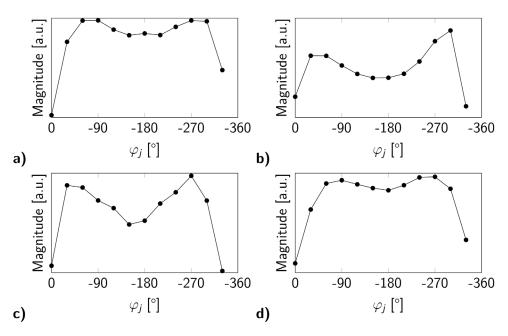


FIG. 8. Exemplary bSSFP frequency response for gray (a) and white (b) brain matter (for the definition of ROIs, see Figure 6), as well as, for patellar cartilage (c) and muscle tissues (d) (for the definition of ROIs, see Figure 7). Note the increased asymmetry in the case of white versus gray brain matter.

Finally, 3D MIRACLE imaging is demonstrated in the knee joint (Fig. 7) at 3T in axial slice orientation. See Table 1 for the evaluation of  $T_1$  and  $T_2$  within marked ROIs. Zonal variation in patellar cartilage  $T_1$  and  $T_2$  is clearly visible with a decrease in both  $T_1$  and  $T_2$  values from superficial to deep layers, as typically expected and observed for healthy cartilage (see insets, Figure 7).

#### **DISCUSSION**

Indirect estimation of the lowest order SSFP modes for rapid  $T_1$  and  $T_2$  mapping from a set of phase-cycled bSSFP scans using a Fourier transformation yields similar properties as TESS relaxometry. Furthermore, due to the balanced gradients used in the present case, MIRA-CLE relaxometry is expected to be less sensitive to diffusive effects than TESS, particularly in fluids (27). Still, it is important to remember that the accuracy of the presented method relies on the ability to correctly retrieve the SSFP mode amplitudes  $F_{-1}$ ,  $F_0$ , and  $F_1$  from the bSSFP profile. This not only depends on the number Nof RF phase increments used, but also on how fast the configurations decay with increasing order |p| (e.g., see Figure 1b). For TR  $\ll T_{1,2}$ , the decay of the mode amplitudes becomes a function of the flip angle  $\alpha$  and the relaxation time ratio  $T_1/T_2$  (41). As a result, aliasing becomes more and more an issue with decreasing  $T_1/T_2$ , e.g., as observed in Figure 4, where the error in  $T_2$ increases with increasing  $T_2$  (for a constant  $T_1$ ). Consequently, accurate relaxometry of fluids might not be granted, even for N=12. Nonetheless, the phantom data show that MIRACLE is in very good agreement with both 3D TESS and reference IR-TSE and SE methods for tissue-like  $T_1/T_2$  ratios. In addition to the aforementioned effects, patient movements between phase-cycles may affect the success of MIRACLE relaxometry, although adequate fixation and image registration could be used to mitigate such effects.

Typical  $T_1$ -values reported in the literature for white and gray brain matter are approximately 950-1000 ms (42) and 1300-1500 ms (43,44), respectively, which are similar to what was obtained during our reference measurements. Interestingly, and in contrast to the simulations as well as phantom and cartilage experiments, in vivo brain imaging with MIRACLE demonstrates a systematic underestimation of  $T_1$  even after  $B_1$  correction. This bias is likely linked to the asymmetric shape of the bSSFP frequency response, as exemplified in Figure 8 for brain tissue, but more prominently for white as compared to gray matter. Those asymmetries have already been discussed in some details by Miller et al (45,46) and are believed to be due to an inhomogeneous intra-voxel frequency content. This is not unexpected as brain tissues contain microstructural boundaries, compartments or chemical shifts that might not be properly characterized by a single pair of relaxation parameters and thus by a single-compartment signal behavior as assumed by Equations [1,2]. To further investigate these effects, noise-free simulations of the bSSFP signal in the case of a two-components model assuming Lorentzian line shapes were performed. As was initially proposed by Miller et al (46), we assume identical  $T_1$  and  $T_2$  (830/80 ms) for both components, a fixed line width of  $\Gamma_1$  = 0.1 Hz for the first component, and a volume fraction of 0.15 for the second component. Corresponding MIRACLE- $T_1$  and  $-T_2$  values as a function of the second component's width  $(\Gamma_2)$  and frequency shift  $(\Delta f)$  are shown in Figure 9a. Because this results in an overestimation rather than an underestimation of  $T_1$ , this model fails to describe the observed MIRACLE brain data.

Generally, the constraint of having the same relaxation properties for both components might be too restrictive. Deoni et al (47) suggested a two-component system with a volume fraction of roughly 0.28 for white matter, where the dominating component has a rather low  $T_1/T_2$ -ratio of 900/120 ms, in contrast to a rather high  $T_1/T_2$ -ratio of

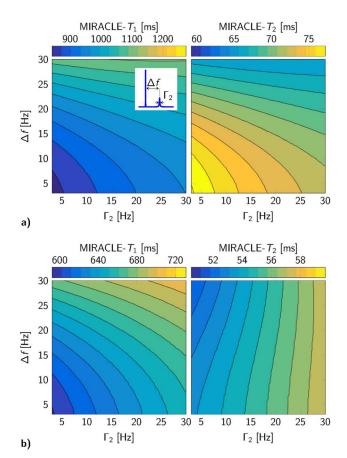


FIG. 9. Estimated MIRACLE  $T_1$  and  $T_2$  from a simulated two-component system with identical  $T_1/T_2$  values of 830/80 ms and a volume fraction of 0.15 (a), and  $T_1/T_2$  values of 900/120 ms and 380/10 ms (b), respectively, as well as a volume fraction of 0.28. The inset in (a) depicts the parameter space investigated: the first component is assumed to be on-resonant with a width  $\Gamma_1=0.1$  Hz, whereas the values for the width  $\Gamma_2$  and frequency shift  $\Delta f$  of the second component are varied as shown above. Typical values reported by Miller et al (46) for white matter are  $\Delta f=17$ –23 Hz and  $\Gamma_2=19$ –22 Hz (depending on the orientation).

380/10 ms for the smaller component, reflecting myelin. Note that contrary to the aforementioned study, we do not consider exchanges between species in the present case. We now repeat the analysis done previously and show the corresponding MIRACLE results in Figure 9b. Within this framework, we observe a shift to apparent low  $T_1$  values in combination with typical  $T_2$  values that are in good agreement with the MIRACLE- $T_1$  and  $-T_2$  values from our experiments. Consequently, the low  $T_1$  values observed in vivo are likely to originate from the presence of a myelin-like second component with different frequency distribution, in line with the observation that the resulting bias in  $T_1$  is much less pronounced for gray matter, where lower myelin contents are expected. The aforementioned suspected sensitivity of MIRACLE to tissue heterogeneity and frequency asymmetry is further corroborated by the results observed for articular cartilage. Here, the observed  $T_1$  and  $T_2$  values are in very good agreement to previously published values (25,48), which is expected because cartilage is known to be much less heterogeneous (49).

In summary,  $T_1$  and  $T_2$  mapping with MIRACLE offers analogous properties as TESS while successfully mitigating its motion-sensitivity. In contrast to the literature, however, apparent low  $T_1$  values are observed for brain white matter; reflecting the asymmetry in the bSSFP signal profile. As a result, a configuration-based relaxometry, as suggested with MIRACLE, becomes sensitive to changes in the underlying frequency spectrum and content, and might, therefore, offer improved sensitivity to diffuse pathophysiological changes in the brain. Future work will aim to explore this new frequency-sensitive relaxometry method.

#### **CONCLUSIONS**

A rapid motion-insensitive configuration-based steady state relaxometry method was presented. Compared with most SSFP methods, it offers accurate and robust  $T_2$  quantification of human tissues, even in the presence of substantial  $B_0$  and  $B_1$  field inhomogeneities, as demonstrated in the scope of this work. In contrast to contemporary single component relaxometry methods, however, our results and preliminary modelling indicate that MIR-ACLE becomes sensitive to the  $T_1$  and  $T_2$  intra-voxel frequency dispersion.

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### SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article

**Fig. S1.** Illustrative volumetric 3D-TESS and 3D-MIRACLE brain imaging of a healthy volunteer. **a:** Axial sample images of the three lowest SSFP mode volumes  $F_1$ ,  $F_0$  and  $F_{-1}$  acquired by TESS (note the different scaling for  $F_1$ ). Arrows indicate the most pronounced pulsation artifacts. **b:** Similar images derived from a MIRACLE acquisition. **c:** Corresponding  $T_1$  maps (ms) from TESS and MIRACLE d: Corresponding  $T_2$  maps (ms) from TESS and MIRACLE (MIRACLE imaging parameters identical as Figure 6; TESS imaging parameters: resolution  $1 \times 1 \times 2 \, \text{mm}^3$ , FA  $15^\circ$ , TR/TE =  $5.93/3.03 \, \text{ms}$ ,  $2 \, \text{avg.}$ ).