

# Ultrafast manipulation of Heisenberg exchange and Dzyaloshinskii Moriya interactions in antiferromagnetic insulators

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# Spintronics

Study and use the intrinsic spin of the electrons, in addition to their charge, to build devices. Example: Magnetoresistive random access memory.

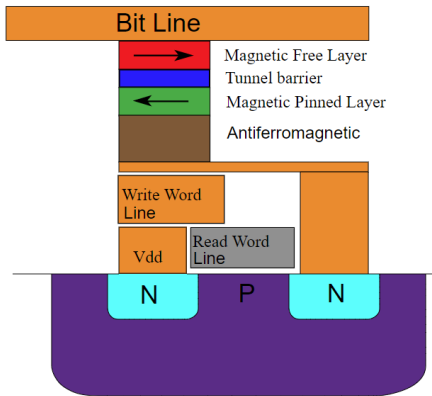


Figure:  
Source:  
Wikipedia

# Antiferromagnetic Spintronics

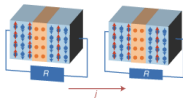
The development of information technology demands devices with high storage density, high energy efficiency and high write-read speeds. Antiferromagnets have appear as a natural alternative to ferromagnets due to some advantages:

- ▶ Robustness against magnetic disturbances.
- ▶ No hysteresis loss.
- ▶ Produce no stray fields.
- ▶ Ultrafast dynamics (femtosecond).

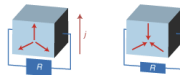
# Antiferromagnetic Spintronics

Electrical detection

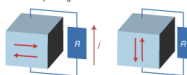
a Magnetoresistance in antiferromagnetic junctions



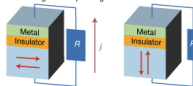
c Anomalous Hall effect



b Anisotropic magnetoresistance



d Tunneling anisotropic magnetoresistance



Electrical switching

e Current-induced torque in antiferromagnetic junction



f

Edelstein spin-orbit torque



g

SHE spin-orbit torque

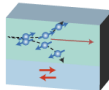


Figure:  
Source:  
Nature, Spin  
transport and  
spin torque in  
antiferromag-  
netic  
devices

# Manipulation of the spin state

## Heisenberg model

$$\hat{H} = \sum_{\langle i,j \rangle} J \mathbf{S}_i \mathbf{S}_j$$

In antiferromagnets  $J > 0$  ( $J < 0$  for ferromagnets). Mentink, J. H. et al. showed that it is possible to manipulate the exchange constant by applying a laser field.  $J = J_0 + \Delta J$ .

$$\hat{H} = \sum_{\langle i,j \rangle} J \mathbf{S}_i \mathbf{S}_j + \sum_{\langle i,j \rangle} \vec{D}_{ij} \mathbf{S}_i \times \mathbf{S}_j$$

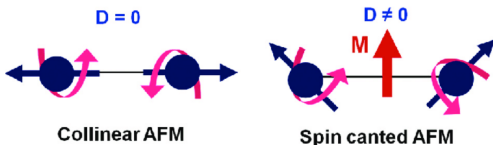


Figure:  
Source:  
Nanoscale

# Manipulation of the spin state

If we can manipulate  $J$  and  $D$  we can control the canting angle.

**Initial state**

$$E_{ex} = JS_iS_j + D[S_i \times S_j]$$



$2\theta$

$t < 0$

**Laser excitation**

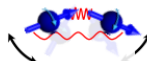
$$E_{ex} = (J + \Delta J)S_iS_j + (D + \Delta D)[S_i \times S_j]$$



$t = 0$

**Final state**

$$E_{ex} = JS_iS_j + D[S_i \times S_j]$$



$t \gg \tau$

Figure: Source: Kimmel presentation

## Manipulation of the spin state

We need  $\frac{J}{D} \neq \frac{J+\Delta J}{D+\Delta D}$ .

However,  $\frac{J}{D} = \frac{J+\Delta J}{D+\Delta D}$ .



# Manipulation of the spin state

## Kane-Mele-Hubbard model

$$\hat{H}_{\text{KMH}} = -t_1 \sum_{\langle ij \rangle \sigma} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + i\Delta \sum_{\langle\langle ij \rangle\rangle \sigma \sigma'} \hat{c}_{i\sigma}^\dagger \nu_{ij} \sigma_{\sigma\sigma'}^z \hat{c}_{j\sigma'} + U\hat{D}$$

It leads to a spin Hamiltonian:

$$\hat{H}_{\text{KMH}}^{\text{eff}} = \sum_{\langle i,j \rangle} J_{ij} \mathbf{S}_i \mathbf{S}_j + \sum_{\langle\langle i,j \rangle\rangle} \mathbf{S}_i \boldsymbol{\Gamma}_{ij} \mathbf{S}_j$$

# Manipulation of the spin state

## Modified Kane-Mele-Hubbard model

$$\hat{H} = -t_1 \sum_{\langle i,j \rangle, \sigma} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{\langle\langle i,j \rangle\rangle, \sigma} (t_2 + i\Delta\nu_{ij}\sigma_{\sigma,\sigma}^z) \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + U\hat{D}$$

It leads to a spin Hamiltonian:

$$\hat{H}_{\text{eff}}(t) = \sum_{\langle i,j \rangle} J_{1,ij} \mathbf{S}_i \mathbf{S}_j + \sum_{\langle\langle i,j \rangle\rangle} \{ J_{2,ij} \mathbf{S}_i \mathbf{S}_j + \mathbf{D}_{2,ij} \mathbf{S}_i \times \mathbf{S}_j + \mathbf{S}_i \mathbf{\Gamma}_{ij} \mathbf{S}_j \}$$

What we wanted, in this case if we apply light  $\frac{J}{D} \neq \frac{J+\Delta J}{D+\Delta D}$ .

# The End