

# **1 Definition of classes**

## 2 Field-theoretic treatment of interactions

Our treatment of interactions uses a field-theoretic treatment of the densities to determine the interactions between polymer segments. Following work by de Pick, *et al.* [1], we define

The simulation has a fixed volume with sides lengths  $L_x$ ,  $L_y$ , and  $L_z$ . These lengths are discretized into  $M_x$ ,  $M_y$ , and  $M_z$  bins of length  $\Delta_x = L_x/M_x$ ,  $\Delta_y = L_y/M_y$ , and  $\Delta_z = L_z/M_z$ . The bins are defined by the three indices  $i_x$ ,  $i_y$ , and  $i_z$  that run from zero to  $M_x - 1$ ,  $M_y - 1$ , and  $M_z - 1$ , respectively.

We consider the  $n$ th bead located at position  $\vec{r}^{(n)}$ . We define a weight function  $w_I(\vec{r}^{(n)})$  within the  $I$ th bin. The  $I$ th index is defined to be a superindex that combines  $i_x$ ,  $i_y$ , and  $i_z$  into a single unique index  $I = i_x + M_x i_y + M_x M_z i_z$  that runs from zero to  $M_x M_y M_z - 1$  (total of  $M_x M_y M_z$  unique indices). The total weight on the  $I$ th bin is given by the contributions from the three cartesian directions, *i.e.*  $w_I(\vec{r}^{(n)}) = w_{i_x}^{(x)}(x^{(n)})w_{i_y}^{(y)}(y^{(n)})w_{i_z}^{(z)}(z^{(n)})$ . Figure 1 shows a schematic of the  $x$ -direction weight function (same method for  $y$  and  $z$ ). This shows a linear interpolation weighting method, consistent with Ref. [1].

The number of epigenetic proteins (*e.g.* HP1) to the  $n$ th site is given by  $N_I^{(\alpha)}$ , where  $\alpha$  determines the type of epigenetic mark. The  $\alpha$ -protein density within the  $I$ th bin is given by

$$\rho_I^{(\alpha)} = \frac{1}{v_{\text{bin}}} \sum_{n=0}^{n_b-1} w_I(\vec{r}^{(n)}) N_I^{(\alpha)} \quad (1)$$

where  $v_{\text{bin}} = \Delta_x \Delta_y \Delta_z$  is the volume of a bin. The maximum number of epigenetic proteins bound  $N_{\text{max}}^{(\alpha)}$  gives an upper bound on the number of proteins that can bind to a bead, accounting for coarse graining of a bead to represent multiple nucleosomes. For discretization of one nucleosome per bead, the maximum  $N_{\text{max}}^{(\alpha)} = 2$  implies binding of a protein to the two histone tail proteins for the  $\alpha$  epigenetic mark. We define the number of  $\alpha$  marks on the  $I$ th bead as  $M_I^{(\alpha)}$ , which can take values from zero to  $N_{\text{max}}^{(\alpha)}$ .

Protein binding to a marked tail results in energy  $-\beta \epsilon_m$  [non-dimensionalized by  $\beta = 1/(k_B T)$ ], and protein binding to an unmarked tail is associated with energy  $-\beta \epsilon_u$ . The chemical potential of the  $\alpha$  protein is defined as  $\beta \mu^{(\alpha)}$ . The binding of  $N_I^{(\alpha)}$  proteins to a bead with  $M_I^{(\alpha)}$  marks results in a free energy that accounts for all of the combinatoric ways of binding.

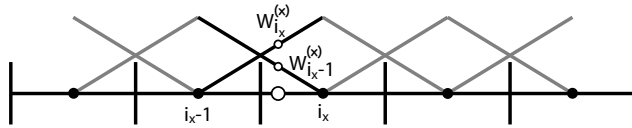


Figure 1: Schematic of the weight function  $w_{i_x}^{(x)}$  that gives the weighting of the particle in the  $i_x$  site in the  $x$ -direction based on a linear interpolation method [1].

### 3 Energetic contributions to the Monte Carlo simulation

#### 3.1 Polymer chain model: shearable, stretchable wormlike chain

We consider a polymer with  $n_b$  number of beads. We consider the shearable, stretchable wormlike chain potential, given by

$$\beta E_{\text{elas}} = \sum_{n=0}^{n_b-2} \left[ \frac{\epsilon_b}{2\Delta} \left| \vec{r}_3^{(n+1)} - \vec{r}_3^{(n)} - \eta \Delta \vec{r}_{\perp}^{(n)} \right|^2 + \frac{\epsilon_{\parallel}}{2\Delta} \left( \Delta \vec{r}^{(n)} \cdot \vec{t}_3^{(n)} - \Delta \gamma \right)^2 + \frac{\epsilon_{\perp}}{2\Delta} \left| \Delta \vec{r}_{\perp}^{(n)} \right|^2 \right], \quad (2)$$

where  $\Delta \vec{r}^{(n)} = \vec{r}^{(n+1)} - \vec{r}^{(n)}$  is the bond vector,  $\Delta \vec{r}_{\perp}^{(n)} = \Delta \vec{r}^{(n)} - (\Delta \vec{r}^{(n)} \cdot \vec{t}_3^{(n)}) \vec{t}_3^{(n)}$  is the perpendicular component of the bond vector to the tangent vector.

### References

- [1] Darin Q. Pike, François A. Detcheverry, Marcus Müller, and Juan J. de Pablo. Theoretically informed coarse grain simulations of polymeric systems. *The Journal of Chemical Physics*, 131(8):084903, August 2009.