


Simple harmonic motion



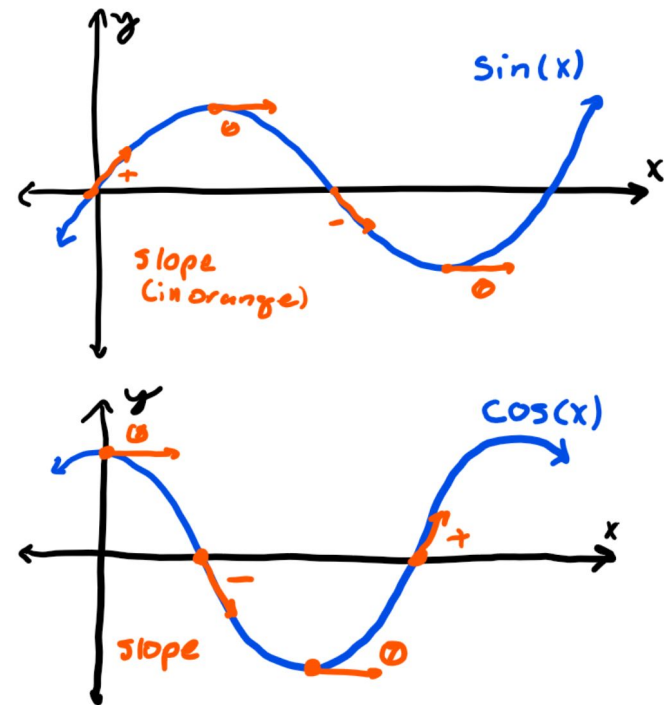
Trigonometric derivatives

Slopes of sine and cosine

The slope of the sine graph starts positive, goes toward zero, then goes negative, then rises back to zero.

The cosine function also follows the same pattern: in fact, the derivative of sine is cosine.

However, the slope of the cosine function first starts zero, then falls, then rises, which is what negative sine looks like, so the derivative of cosine is actually negative sine.



Circular motion

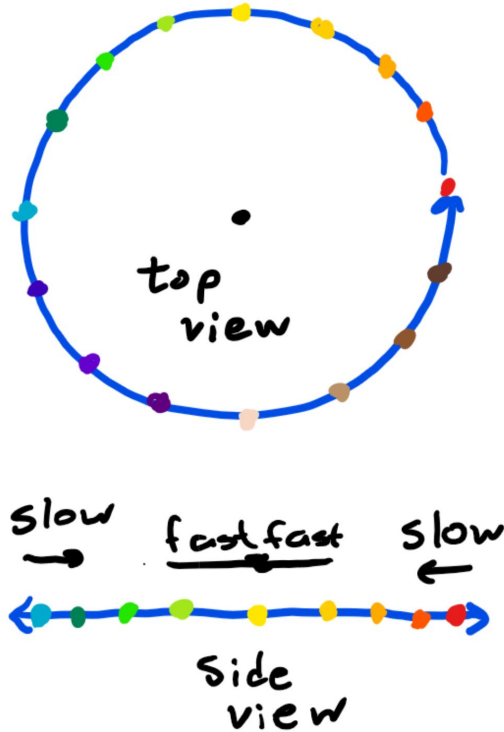
What it looks like

Circular motion is just an object travelling in a circle at a constant angular velocity.

What would it look like if we saw it from the side?

In the middle it would travel slow, but on the outside it would travel fast.

This “side view” is just the horizontal portion of the circular motion.

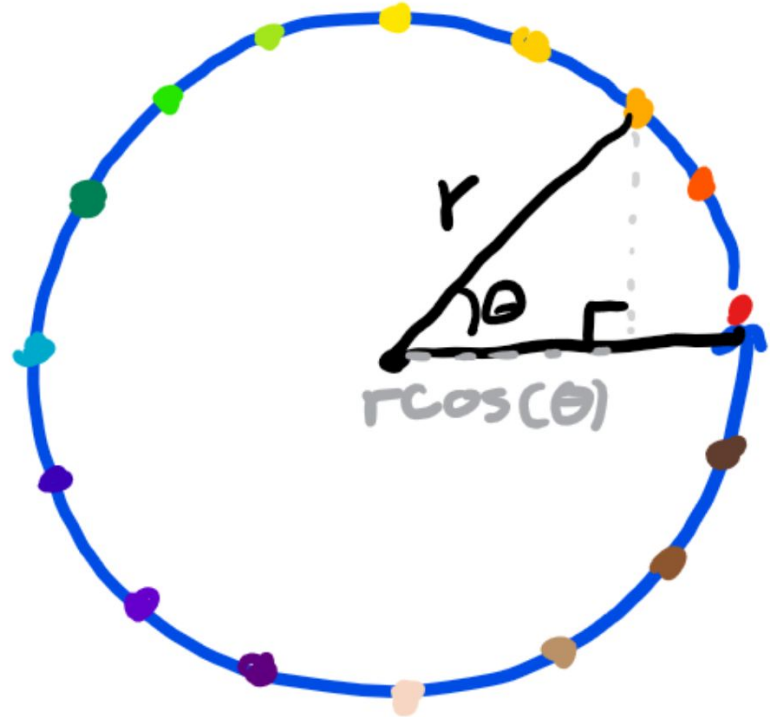


Horizontal position

We can write an equation for the horizontal position of the object.

Since the horizontal portion of the unit circle is $\cos(\theta)$, the horizontal portion of a circle of radius r is $r\cos(\theta)$.

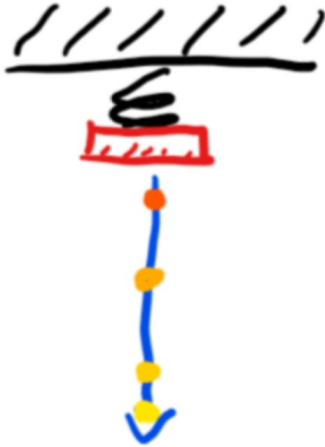
Note that $r\cos(\theta)$ is maximized at r and minimized at $-r$ since $\cos(\theta)$ ranges from -1 to 1 .



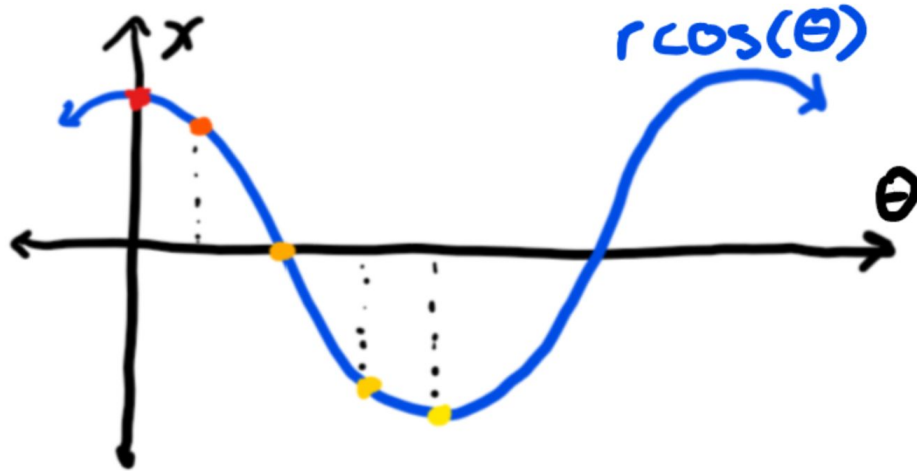
Spring SHM

Graphs

If stretch out a spring and watch it oscillate, we'll see it bounce up and down over time.



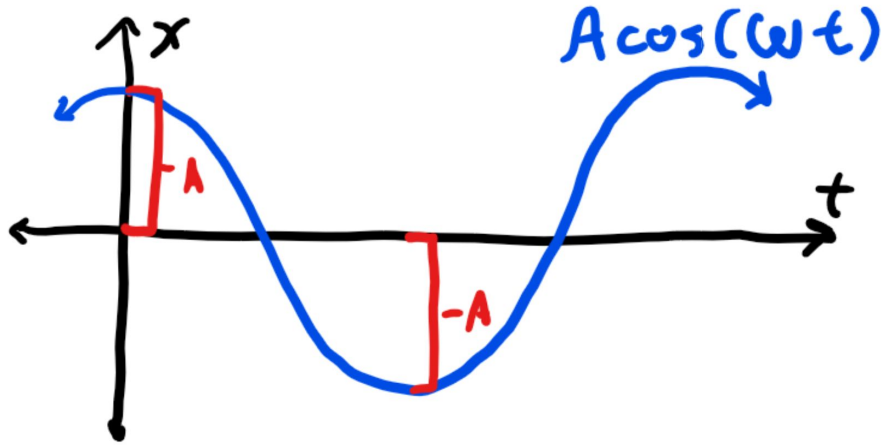
If we plot it against time we'll notice it looks exactly like a cosine graph.



Definitions

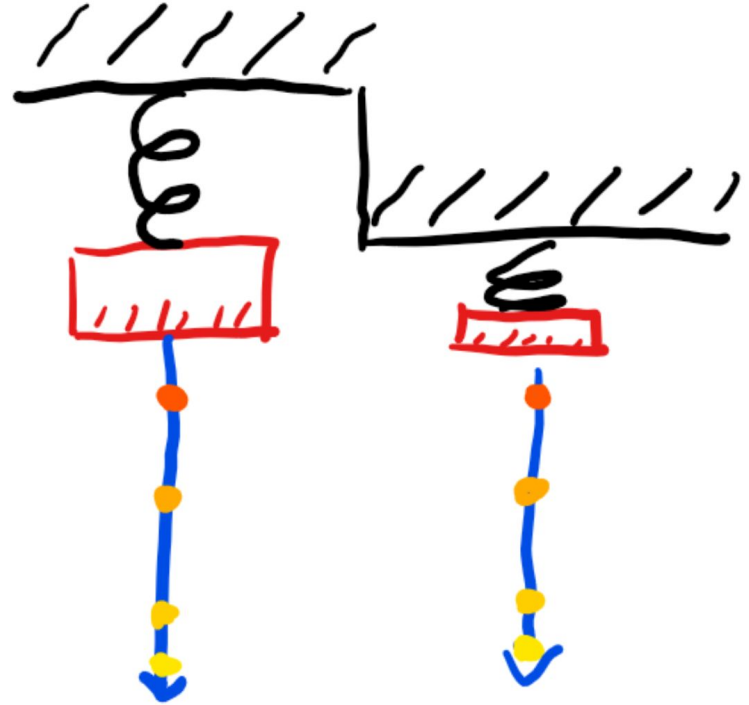
Amplitude: how far an object travels from equilibrium. On the graph, these are the highest/lowest positions and is symbolized by a letter A.

We know $r\cos(\theta)$ has highest and lowest point with magnitude r , which is what we defined amplitude, and we know $\theta = \omega t$, so $r\cos(\theta)$ can be rewritten as the equation $A\cos(\omega t)$.



Gravity

Since $F = -kx$ increases linearly with distance displaced, if a mass hanging from it causes the spring to reach a new equilibrium, as long as the spring follows Hooke's law, the spring will still follow Hooke's law, just with the new equilibrium as, well, the equilibrium.

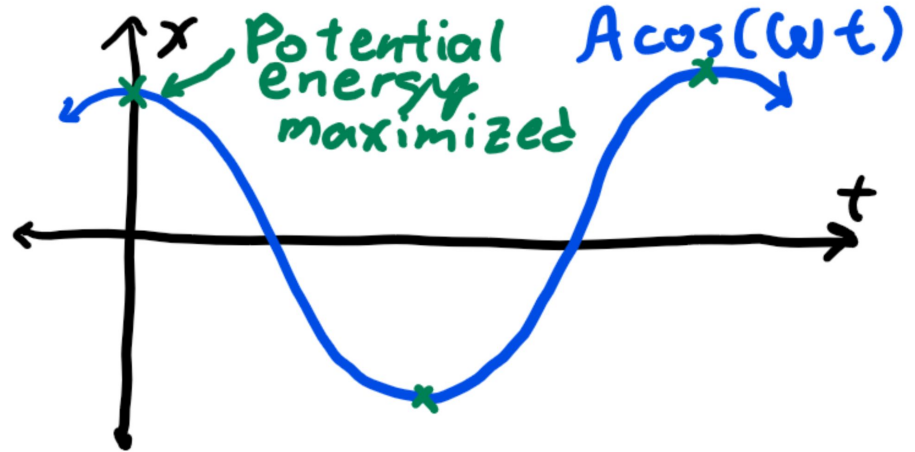


Position

The position function can be written as $x(t) = A\cos(\omega t)$.

The function $x(t)$ is maximized at A , so by Hooke's law, the maximum force on the spring is $F = -kx = -kA$.

The maximum spring potential energy is thus $\frac{1}{2} kx^2 = \frac{1}{2} kA^2$.

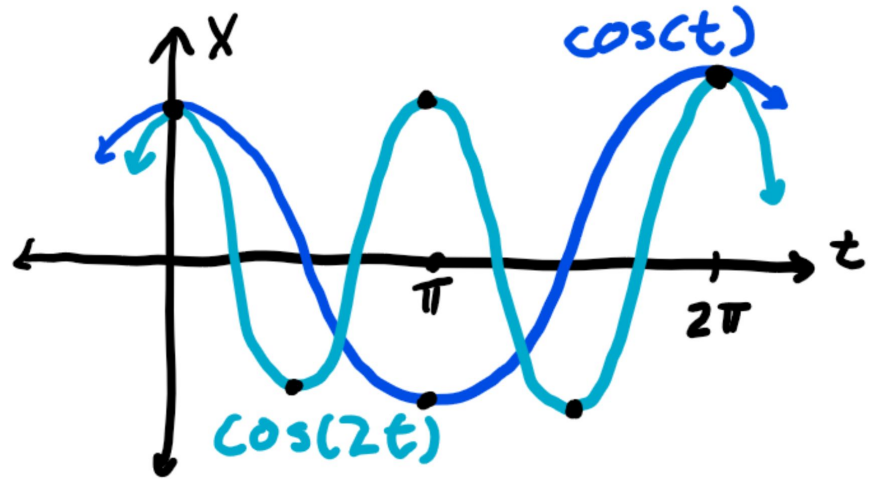


Modified cosine functions

It takes the cosine function 2π before it starts to repeat, so we say it has a period of 2π since it repeats its cycle after this much time has passed.

The normal cosine function $\cos(t)$ has a period of 2π , but $\cos(\omega t)$ goes by ω times as fast, so it has a period of $2\pi/\omega$.

For example, $\cos(2t)$ has a period of π since it takes half the time before the graph repeats.

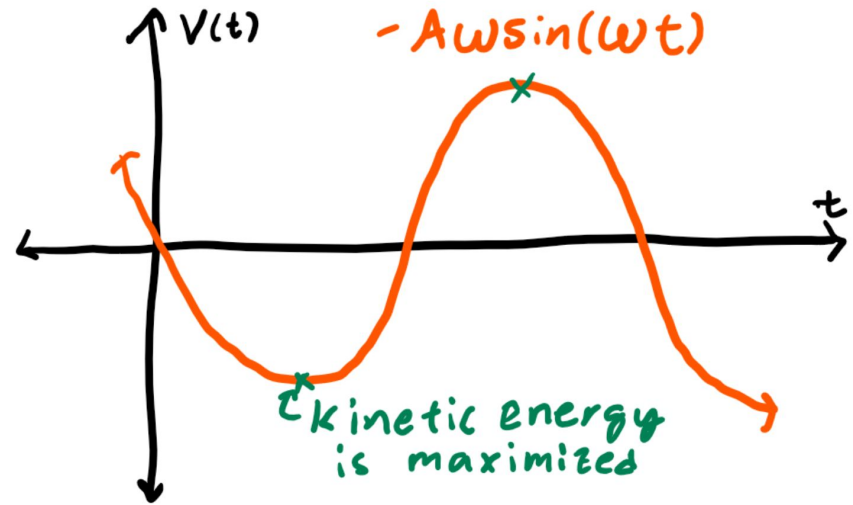


Velocity

We know the derivative of position gives velocity, so if we take the derivative of the position function: $x(t) = A\cos(\omega t)$, we get $v(t) = -A\omega\sin(\omega t)$.

Notice that the maximum velocity is $A\omega$. In circular motion, $A = r$, so maximum velocity is $r\omega = v$.

Maximum kinetic energy is thus $\frac{1}{2}mv^2$, or $\frac{1}{2}mA^2\omega^2$.



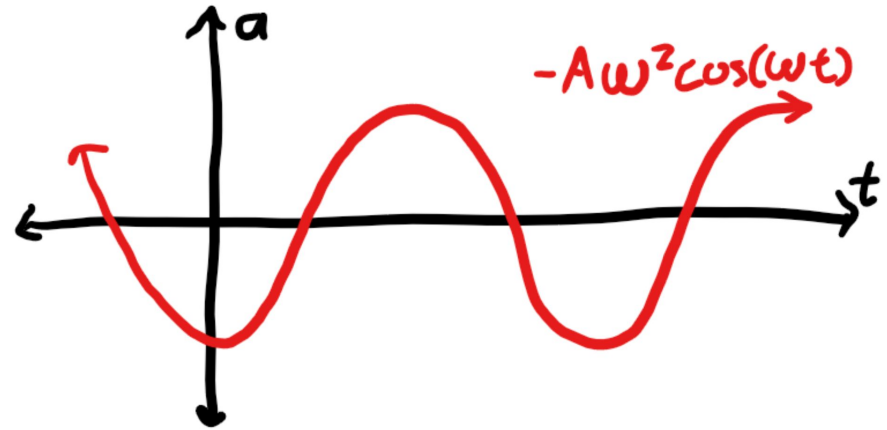
Acceleration

We take the derivative of the position function to get $a(t) = -A\omega^2\cos(\omega t)$.

By Hooke's law, $F = -kx = ma$.

Let's plug in the equations we have for $a(t)$ and $x(t)$: $-kA\cos(\omega t) = m[-A\omega^2\cos(\omega t)]$
rearranging into $k=m\omega^2$.

Thus, $\omega = \text{sqrt}(k/m)$.

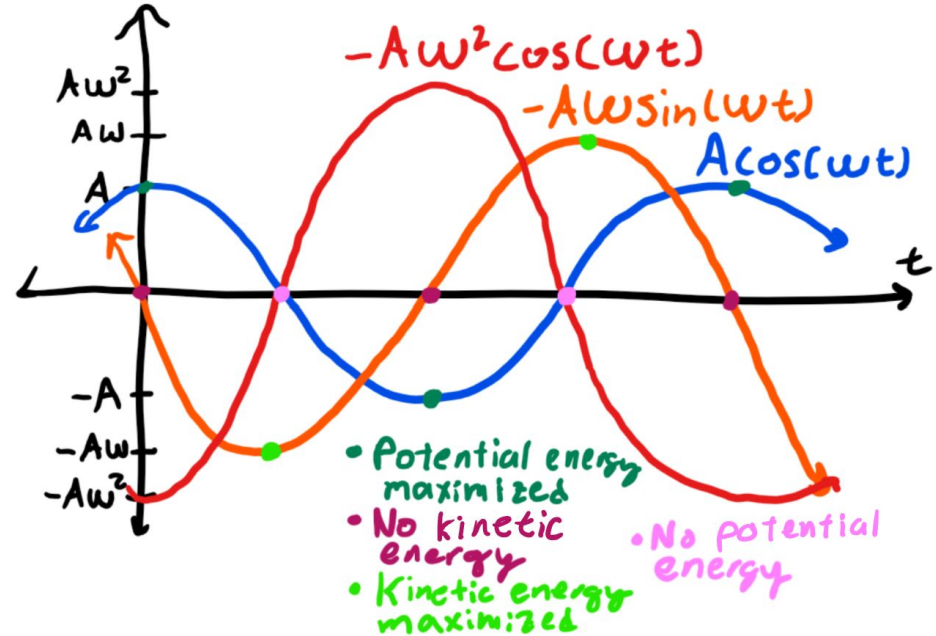


More graphs

Note how the acceleration is maximized when position is furthest from equilibrium, but the velocity is zero there.

Instead, velocity is maximized when position and acceleration are both zero.

Thus, potential energy is maxed out when the spring is fully stretched (kinetic energy is zero) but kinetic energy is maxed when the spring is non-stretched out (potential energy is zero).

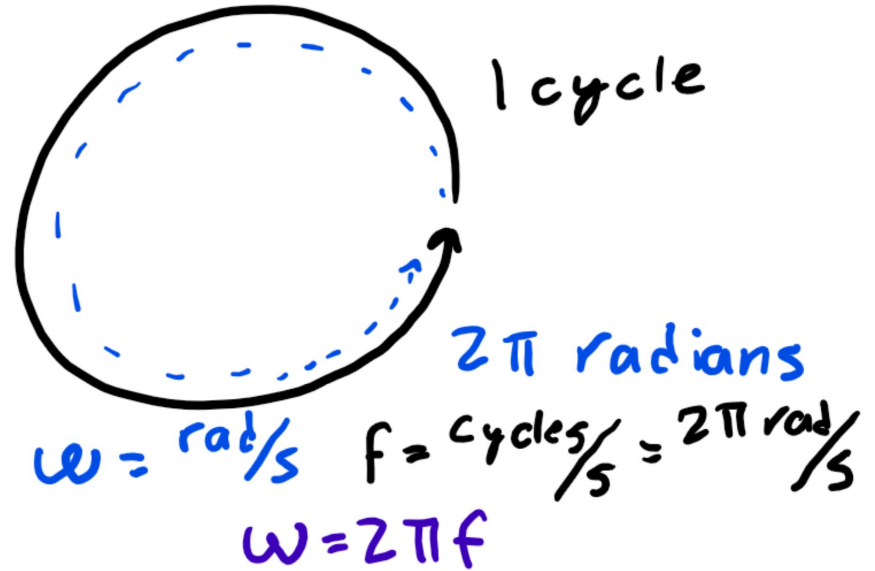


More definitions

Angular velocity, ω , is how many radians an object travels per second, but what if we wanted to know how many cycles it traveled in a second?

We call the amount of cycles per second frequency, represented f , and it also has units $1/s$, or Hertz (Hz).

Since there are 2π radians in a complete cycle, $f = \omega/(2\pi)$.




Even more definitions

Period, symbolized T , is how long it takes to complete a cycle, measured in seconds.

Since frequency is cycles per second and period is seconds per cycle, the period is the inverse of frequency, so that means $T = 1/f$.

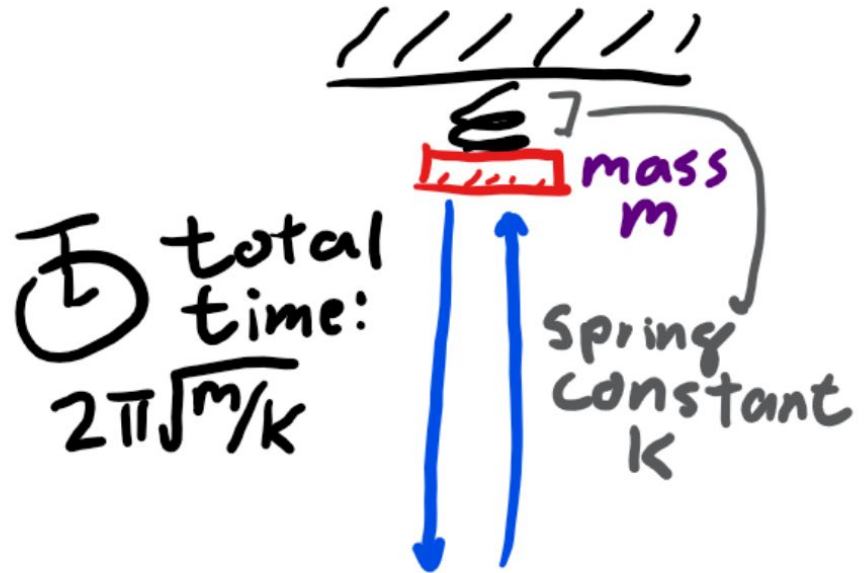
Thus, $T = 2\pi/\omega$.


$$T = \text{period} = \frac{\text{seconds}}{\text{cycle}}$$
$$f = \frac{\text{cycles}}{\text{second}} = 1/T$$

Frequency and period of a spring

Since $\omega = \sqrt{k/m}$, and $f = \omega/2\pi$, that means $f = \sqrt{k/m}/(2\pi)$.

Period is the inverse, so the period of a spring oscillation is $T = 2\pi\sqrt{m/k}$.



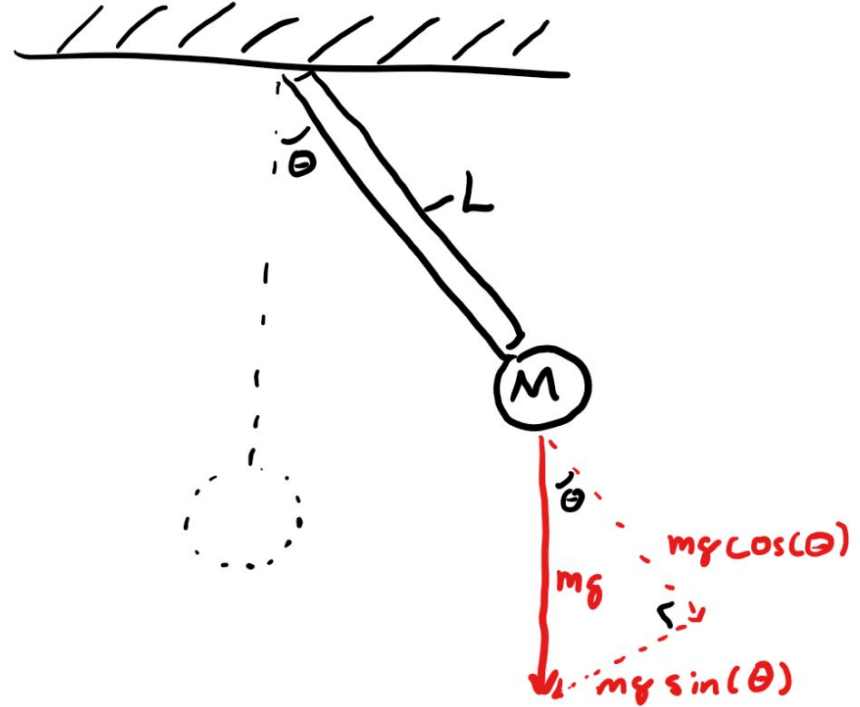
Simple pendulum

Forces on a pendulum

A simple pendulum is a mass m held a distance L away from the pivot.

When the pendulum is θ away from the rest position, $-mg$ acts on it, but $-mg\cos(\theta)$ pulls on the string.

We assume the tension of the string counteracts that part of the force, so the only force acting on the pendulum is $-mg\sin(\theta)$.

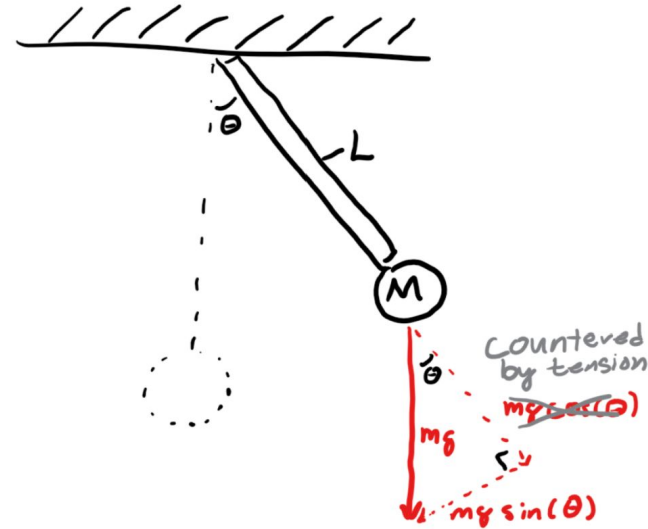


Torque on a pendulum

Since this force acts perpendicularly to the string, it is a torque, and we know that $T = I\alpha$.

I for a point mass is mL^2 and we know that $rF = T = -Lmg\sin(\theta)$, and plugging this in gives $-Lmg\sin(\theta) = mL^2\alpha$.

Rearranging gives $\alpha = -g\sin(\theta)/L$. Since α is also $d^2\theta/(dt)^2$, we get a differential equation $d^2\theta/(dt)^2 = -g\sin(\theta)/L$.



$$\begin{aligned} \tau &= I\alpha = ML^2\alpha \\ \tau &= -Lmg\sin(\theta) = ML^2\alpha \\ \alpha &= -g\sin(\theta)/L \end{aligned}$$

Small angle approximation

If you plug in a few values for $\sin(\theta)$ for very small values of θ , you'll notice that $\sin(\theta) \approx \theta$.

You'll also notice that for $\cos(\theta)$, small values will be approximately $1 - \theta^2/2$, and for $\tan(\theta)$, small values are also approximately θ .

What is considered “small” is around 15 degrees, or $\pi/12$ radians.

θ	$\sin(\theta)$	$\tan(\theta)$	$1 - \theta^2/2$	$\cos(\theta)$
0	0	0	1	1
0.01	0.009999833334	0.01000033335	0.99995	0.9999500004
0.02	0.01999866669	0.02000266709	0.9998	0.9998000067
0.03	0.0299955002	0.03000900324	0.99955	0.9995500337
0.04	0.03998933419	0.040021347	0.9992	0.9992001067
0.05	0.04997916927	0.05004170838	0.99875	0.9987502604
0.06	0.05996400648	0.06007210383	0.9982	0.9982005399
0.07	0.06994284734	0.07011455787	0.99755	0.9975510003
0.08	0.07991469397	0.08017110471	0.9968	0.9968017063
0.09	0.0898785492	0.09024378991	0.99595	0.995952733
0.1	0.09983341665	0.1003346721	0.995	0.9950041653

Differential equation

Using the small angle approximation,
we get $d^2\theta/(dt)^2 = -(g/L)(\theta)$.

So we need to find some function
whose second derivative is the negative
of itself times g/L .

Wait, the $\theta(t) = \cos(\omega t)$ thing's second
derivative was $-\omega^2 \cos(\omega t)$, so if $\omega^2 = g/L$,
this would work out.

Thus, $\omega = \sqrt{g/L}$, so $T = 2\pi\sqrt{L/g}$.

$$\frac{d}{dt^2} \underline{\cos(\omega t)} = \frac{d}{dt} [-\omega \sin(\omega t)]$$
$$= -\omega^2 \underline{\cos(\omega t)}$$

$$\text{If } \omega = \sqrt{g/L}$$

$$-\omega^2 \underline{\cos(\omega t)} = -g/L \underline{\cos(\omega t)}$$

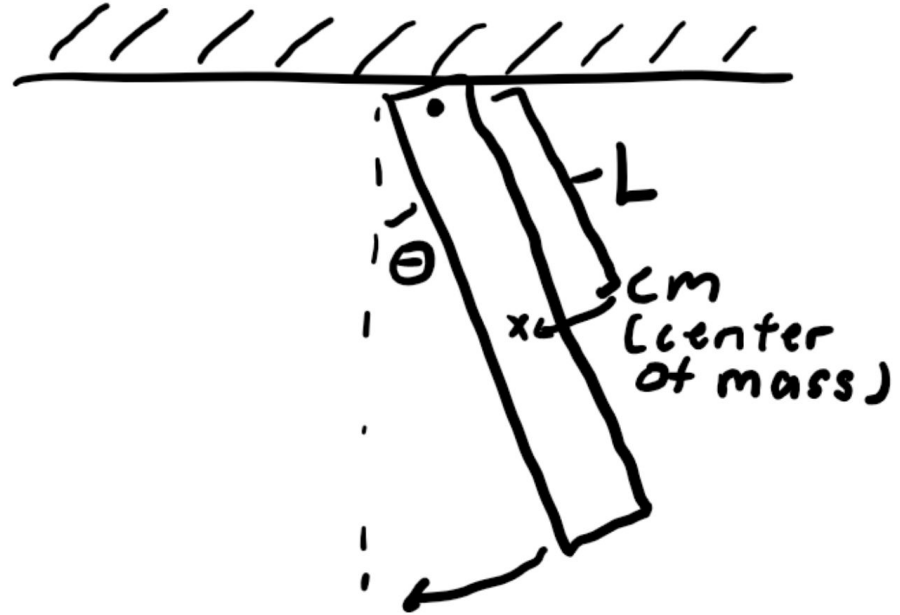
Physical pendulum

Terminology (aka definitions)

A physical pendulum is just any item being rotated from some pivot point.

We can imagine the center of mass as the point of pivoting, so L will not be the length of the pendulum but the distance the center of mass is from the pivot.

I will represent the moment of inertia at the center of mass.



Equation

We can write the same $T = -Lmg\sin(\theta)$ equation as before where $T = I\alpha$.

Thus, $\alpha = d^2\theta/(dt)^2 = -Lmg\sin(\theta)/I$.

The solution is again $\theta(t) = \cos(\omega t)$ but this time $\omega = \sqrt{mgL/I}$, so we can find the period to be $T = 2\pi\sqrt{I/(mgL)}$.



Kepler's laws (non-simple/harmonic motion)



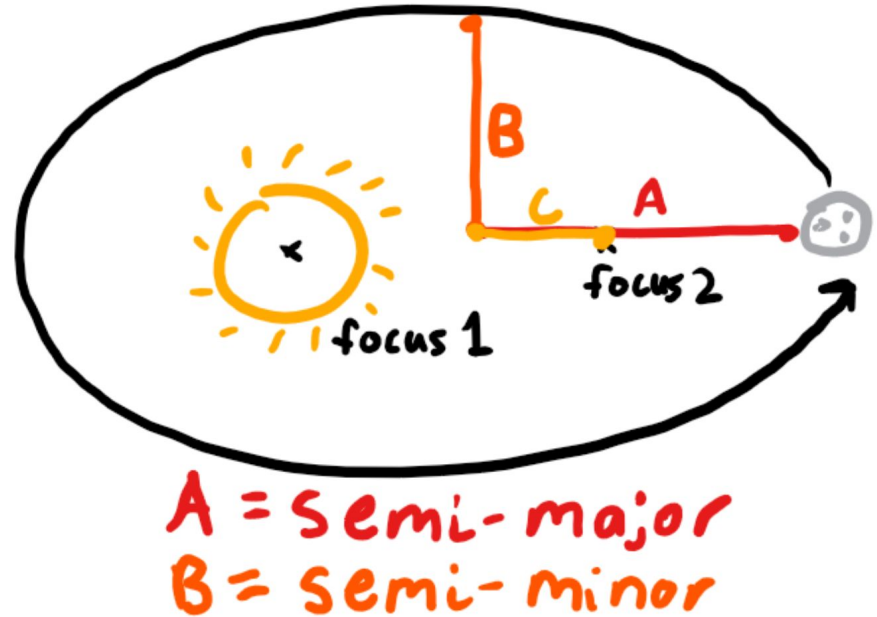
First law

A planet orbits its sun in an ellipse with the sun on one of the foci of the ellipse.

We can say there are two axes to the ellipse: the longer one, the semi-major axes that pass through the foci, and the semi-minor one, perpendicular to the foci.

We say the length of the semi-major is A , semi-minor is B , and the distance a focus is from the center is C .

A focus is defined such that $A^2 - B^2 = C^2$.



Second law

Let's draw a line from the planet to the sun and let it sweep out a pizza shape in a certain amount of time t . Let's say the area of that piece is K .

Like simple harmonic motion, the planet moves faster the closer it is from the sun, but faster the further away it is (conceptually, this is because gravitational force further away).

In fact, this works out so well that for any pizza made in a time t , the area swept out will always be K .



Third law

$T = n\sqrt{A^3}$ for some constant n . In other words, T^2 is proportional to A^3 .

Let's say a planet orbiting with a semi-major axis of A is struck so its new semi-major axis is $4A$.

If the old period was $T = n\sqrt{A^3}$, the new period is $n\sqrt{(4A)^3} = 8n\sqrt{A^3}$.

Thus, the new period is 8 times as long as the old one since n is a constant.