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Stochastic Modeling and Simulation

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Preface

As a student of Scientific and Data Intensive Computing, I've created these notes while attending the **Stochastic Modeling and Simulation** course.

The course covers a wide range of topics, including Stochastic Nonlinear Models in different application fields (as physics, biomedicine, mathematics, ...), Stochastic Differential Equations, and Stochastic Simulation.

The topics covered in these notes include:

- Recap of Deterministic Models
- Stochastic Differential Equations and White Noise
- Fokker Planck Equation
- Noise-induced Transitions
- Colored noises
- Bounded Stochastic Processes
- Spatio-temporal Stochastic Processes
- Parameter Estimation from Data
- Stochasticity ...
- ...
- Continuous state space-discrete time Stochastic Processes
- Discrete Time Markov Chains
- Continuous Time Markov Chains
- Mean Field Approximation

While these notes were primarily created for my personal study, they may serve as a valuable resource for fellow students and professionals interested in Stochastic Modeling and Simulation.

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Introduction

A diverse array of fields, ranging from epidemiology and systems biology to finance and renewable energy, are fundamentally concerned with understanding systems that evolve over time. While simple, deterministic models can offer initial insights, many real-world phenomena are characterized by inherent randomness, uncertainty, and intricate interconnections. To capture this richness, we turn to the powerful framework of **Stochastic Modeling and Simulation**. Many systems we encounter are **complex**: disease spread modeled by the classic SIR (Susceptible, Infected, Recovered) framework, oscillating chemical reaction concentrations, or predator-prey ecosystem dynamics. In each case, the system's behavior emerges from interactions between countless individual components, that cannot be predicted from studying the parts in isolation.

Definition: *Complex System*

A **complex system** is a system composed of interconnected parts that, as a whole, exhibit one or more properties which are not obvious from the properties of individual parts alone.

A key feature of complex systems is **emergent behavior**, where macroscopic patterns arise from the local interactions of microscopic agents. This is closely related to the concepts of adaptation and self-organization. **Adaptation** refers to the process by which a system achieves a better fit with its environment. **Self-organization** is the spontaneous process through which a system's internal structure changes, often to facilitate this adaptation, without external control. A classic example is the formation of a flock of birds: each bird follows simple rules based on its neighbors, yet the flock as a whole displays coordinated, complex flight patterns that no single bird orchestrates.

Observation: *The Constructive Role of Noise*

When studying complex systems, we frequently discover that randomness, or *noise*, serves a purpose beyond merely interfering with deterministic signals. Surprisingly, the presence of noise combined with nonlinear dynamics can actually promote the formation of ordered structures and novel, persistent patterns.

1.1 The Modeling Framework

To build quantitative models of complex systems, we require a clear mathematical framework. This involves defining several key ingredients that form the foundation of any dynamical model:

- **Entities:** The fundamental components of the system. These can be modeled as *discrete* objects (e.g. individuals in a population) or as *continuous* quantities (e.g. the density of a population).
- **State:** A complete description of the system at a specific moment. The *state space* is the set of all possible states the system can occupy.
- **Time:** The independent variable against which the system's state evolves. Time can be treated as *discrete* (advancing in steps) or *continuous*.
- **Evolution Rules:** The laws or functions that dictate how the system's state changes from one moment to the next. These rules are often expressed using mathematical equations.

The Role of Differential Equations

For systems where the state and time are treated as continuous, the evolution rules are most naturally expressed using differential equations.

Definition: *Differential Equation*

A **differential equation** is a mathematical equation that relates one or more unknown functions to their derivatives. In modeling, it describes the instantaneous rate of change of a system's state:

$$\frac{dx}{dt} = f(x, t)$$

where x is the state variable and f is a function that describes how x changes over time.

In essence, differential equations are the language we use to describe a changing world. Perhaps the most famous and foundational example is Newton's Second Law of Motion, which relates the force on an object to its acceleration (the second derivative of its position).

$$F = ma \implies m \frac{d^2x}{dt^2} = F(x(t), \dot{x}(t))$$

This single equation forms the basis of classical mechanics and is a prime example of a **dynamical system**: a system whose state evolves over time according to a deterministic rule.

1.2 Bridging Models and Data in the Modern Era

In our era of big data and machine learning, an important question emerges: will data-driven methods make traditional mathematical modeling redundant? The answer is both more complex and more promising: the path forward involves **hybrid approaches** that harness the complementary strengths of both methodologies. Although real-world datasets are frequently noisy and incomplete, while models necessarily simplify reality, their synergistic integration can unlock insights that neither approach could achieve alone.

Hybrid strategies are rapidly becoming central to scientific discovery:

1. Physics-Informed Neural Networks (PINNs):

This is a cutting-edge deep learning framework for solving problems involving differential equations. A neural network is trained not only to fit observed data but also to obey the underlying physical laws of the system, encoded as differential equations. This helps the model generalize better from sparse or noisy data.

2. Machine Learning for Parameter Estimation:

While a mathematical model may capture the structure of a system (e.g., the SIR model), the specific parameters (infection rate, recovery rate) must be estimated from real-world data. Machine learning and statistical inference techniques are essential tools for this task, especially when dealing with high-dimensional and complex models.

3. Hybrid Modeling:

In some cases, parts of a system may be well-understood and easily described by equations, while other parts may be too complex. A hybrid model might use a traditional differential equation for the well-understood component and a neural network, trained on data, to represent the more complex, "black-box" component.

2 Stochastic Dynamics

In the previous chapter, we introduced the concept of modeling systems using differential equations. We begin our deeper exploration with **linear systems**. Many sophisticated systems, when analyzed locally around a point of equilibrium, can be effectively approximated by a linear model; this linearization is a cornerstone of dynamical systems theory.

2.1 Linear Systems: The Foundation of Dynamics

Linear Ordinary Differential Equations (ODEs) are characterized by equations where the dependent variable and its derivatives appear only to the first power. They are invaluable because they can often be solved analytically, providing clear insight into the system's behavior.

Consider a simple electrical **RL circuit** consisting of a resistor (R) and an inductor (L) in series. Kirchhoff's voltage law gives:

$$L \frac{di}{dt} + Ri = 0$$

This first-order linear homogeneous ODE has the solution:

$$i(t) = i_0 e^{-\frac{R}{L}t}$$

where $\tau = L/R$ is the *time constant* characterizing the exponential decay.

Adding a capacitor creates the **RLC circuit**, a second-order system. Using the state vector $\mathbf{y} = [i, Q]^T$ where Q is the charge, we obtain:

$$L \frac{di}{dt} + Ri + \frac{1}{C}Q = 0$$

This is a second-order differential equation. To analyze it as a system, we can define a state vector $\mathbf{y} = [i, Q]^T$. The system of first-order equations is:

$$\begin{cases} \frac{di}{dt} = -\frac{R}{L}i - \frac{1}{LC}Q \\ \frac{dQ}{dt} = i \end{cases} \Rightarrow A = \begin{bmatrix} -R/L & -1/(LC) \\ 1 & 0 \end{bmatrix}$$

The behavior of this system (e.g., oscillations, decay) is entirely determined by the eigenvalues of the matrix A .

Observation: The Ubiquity of the Linear Model

The same mathematical structure appears in completely different physical domains. Consider a damped harmonic oscillator, such as a mass on a spring moving through a viscous fluid. Newton's second law gives:

$$m\ddot{x} = -kx - \gamma\dot{x}$$

where k is the spring constant and γ is the damping coefficient. Rearranging gives:

$$m\ddot{x} + \gamma\dot{x} + kx = 0$$

This equation shares the same mathematical structure as the RLC circuit equation, illustrating how mathematical models can uncover similarities across entirely different physical systems.

2.2 The Need for Stochasticity: introducing noise

The deterministic models discussed so far are powerful idealizations, however, they assume that all parameters and forces are known and constant. In reality, systems are constantly subjected to small, unpredictable influences from their environment. The aggregate effect of these influences is termed ***noise***. Modeling this randomness is crucial for creating realistic descriptions of the world. A classic illustration of this is the phenomenon of **Brownian motion**. In 1827, botanist Robert Brown observed the erratic, random motion of pollen grains suspended in water. A simple deterministic model for a particle in a fluid, incorporating only a drag force ($m\dot{v} = -kv$), predicts that the particle's velocity should decay to zero almost instantly. This prediction is in contradiction with Brown's empirical observations of perpetual, jittery motion.

The solution to this paradox came from understanding what friction really is at the microscopic level. The drag force $-kv$ represents only the *average* effect of countless molecular collisions with the particle. Einstein and Langevin realized that each individual collision is random and unpredictable, creating fluctuating forces that keep the particle in perpetual motion.

In general, a particle moving through a fluid can experience three types of forces: the damping force we've already discussed, random collisions from molecules, and possibly some external deterministic force $F_d(x)$:

$$m\ddot{x} = -k\dot{x} + F_d(x) + F_s(t)$$

Let's consider the case where the particle is very light compared to the damping strength, meaning $m \ll k$. To make our notation cleaner, we can write $F_s(t) = k\xi(t)$ and $F_d(x) = kf(x)$:

$$m\ddot{x} = -k\dot{x} + kf(x) + k\xi(t)$$

When $m/k \ll 1$ (the so-called overdamped limit), the inertial term becomes negligible, and we arrive at:

$$\dot{x} = f(x) + \xi(t)$$

This is the famous **Langevin equation**. It represents a fundamental shift in how we think about physical systems—moving from purely deterministic descriptions to ones that embrace randomness. The equation captures the idea that what we observe as smooth, deterministic forces are often just the averaged effects of chaotic microscopic processes.

Example: Noise in a RLC circuit

In the previous example of the RLC circuit

$$Ri = -\frac{d}{dt}\Phi(B_{self}) = -\frac{d}{dt}(Li)$$

the derivation was incomplete!

Indeed, a circuit where a current i is flowing is embedded not only in the magnetic field B_{self} generated by $\Phi(Li)$ but also in other external random magnetic fields.

$$B_{total} = B_{self} + B_{external}$$

$$\Phi(B_{total}) = Li + \Phi(B_{external}) = Li + KB_{external}$$

with stochastic external disturbances reads as follows (by the current):

$$Ri = -L\frac{d}{dt}i - K\frac{d}{dt}B_{external}$$

which we may rewrite as

$$L\frac{d}{dt}i = -Ri + K\xi(t)$$

2.2.1 Modeling Impulsive Events: The Dirac Delta

Before we can properly define the stochastic force $F_s(t)$, we must first develop a mathematical tool to describe events that are instantaneous and intense. Consider the impact of a baseball bat on a ball. The force is immense but acts over a very short duration. We care about the net effect, the change in the ball's velocity, rather than the precise evolution of the force during the infinitesimal impact time. Let's model this using Newton's law, $mv = F(t)$. The total change in momentum is the integral of the force over the impact interval, say from $t = 0$ to $t = a$.

The velocity after the impact is:

$$v_{\text{after}} = \frac{1}{m} \int_0^a F(t) dt$$

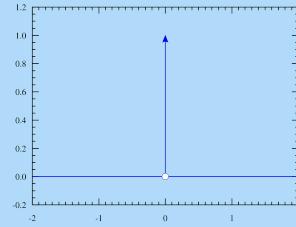
This integral, which represents the total impulse delivered, is finite. To model this, we introduce an idealized mathematical object known as the **Dirac delta function**, denoted $\delta(t)$. It is not a true function in the classical sense but a *distribution* defined by its effect under an integral.

Definition: The Dirac Delta Function

The Dirac delta function, $\delta(t)$, is a generalized function that describes an idealized impulse. Rather than being a function in the classical sense, $\delta(t)$ is a mathematical object with the following properties:

1. There exists a "very small" interval $J = (-\varepsilon/2, \varepsilon/2)$ outside which $\delta(t) \approx 0$.
2. $\delta(t) > 0$ within this interval.
3. The integral of $\delta(t)$ over J is equal to one:

$$\int_J \delta(t) dt = 1$$



The crucial consequence of these properties is the **sifting property**: for any function $f(t)$ that is continuous at $t = 0$, the Dirac delta "sifts out" its value at that point:

$$\int_{-\infty}^{\infty} f(t) \delta(t) dt = f(0)$$

To further illustrate the sifting property, consider a function $f(t)$ that is continuous and differentiable at $t = 0$, with $f(0) < \infty$ and $f'(0) < \infty$. Expanding $f(t)$ in a Taylor series around $t = 0$ gives:

$$f(t) = f(0) + f'(0)t + O(t^2)$$

Integrating $f(t)$ against the Dirac delta over the real line, we have:

$$\int_{-\infty}^{\infty} \delta(t) f(t) dt = \int_{-\infty}^{\infty} \delta(t) [f(0) + f'(0)t] dt$$

which can be split as:

$$= \int_{-\infty}^{\infty} \delta(t) f(0) dt + \int_{-\infty}^{\infty} \delta(t) f'(0)t dt$$

Now, recall the two fundamental properties of the Dirac delta: $\int_{-\infty}^{\infty} \delta(t) dt = 1$ and $\int_{-\infty}^{\infty} \delta(t)t dt = 0$ (since $t\delta(t)$ is an odd function). Applying these, we find:

$$\int_{-\infty}^{\infty} \delta(t) f(t) dt = f(0) \cdot 1 + f'(0) \cdot 0 = f(0)$$

This demonstrates that the Dirac delta "picks out" the value of f at $t = 0$, regardless of the behavior of f elsewhere.

Dirac Delta Classes of Functions

The Dirac delta function, being a mathematical idealization, can be approximated by various families of functions that become increasingly "spike-like" as a parameter grows large.

For $N \gg 1$, the family of rectangular pulse functions

$$\delta_N(t) = \begin{cases} N, & \text{if } t \in (-\frac{1}{2N}, \frac{1}{2N}) \\ 0, & \text{otherwise} \end{cases}$$

provides an example of Dirac delta approximation. As N increases, the function becomes taller and narrower while maintaining unit area, converging to the delta function in the distributional sense. Similarly, the exponential family

$$\delta_\gamma(t) = \frac{\gamma}{2} e^{-\gamma|t|}$$

for $\gamma \gg 1$ offers a smooth approximation that avoids the discontinuities of the rectangular pulse. This family has several advantages: it is infinitely differentiable, has exponentially decaying tails, and provides a more realistic model for physical phenomena.

Observation: Verification of the Sifting Property

Let's verify that the rectangular pulse function $\delta_N(t)$ satisfies the sifting property in the limit as $N \rightarrow \infty$. For any continuous function $f(t)$:

$$\int_{-\infty}^{\infty} f(t) \delta_N(t) dt = \int_{-1/(2N)}^{1/(2N)} f(t) \cdot N dt = N \int_{-1/(2N)}^{1/(2N)} f(t) dt$$

Since $f(t)$ is continuous at $t = 0$, we can use the mean value theorem.

There exists some $c \in [-1/(2N), 1/(2N)]$ such that:

$$N \int_{-1/(2N)}^{1/(2N)} f(t) dt = N \cdot f(c) \cdot \frac{1}{N} = f(c)$$

As $N \rightarrow \infty$, the interval shrinks to zero and $c \rightarrow 0$. By continuity of f , we have $f(c) \rightarrow f(0)$.

Therefore:

$$\lim_{N \rightarrow \infty} \int_{-\infty}^{\infty} f(t) \delta_N(t) dt = f(0)$$

This confirms that $\delta_N(t)$ approaches the Dirac delta function in the distributional sense.

2.2.2 Defining the Stochastic Force: White Noise

The stochastic force in the Langevin equation, commonly denoted as $\xi(t)$ (or sometimes $F_s(t)$, after normalization by mass and other constants), models the cumulative effect of a huge number of independent, microscopic collisions. This motivates the following statistical properties, which define what is known as **Gaussian white noise**.

- White Noise (Temporal Uncorrelation):** The values of $\xi(t)$ at different times are uncorrelated, reflecting the physical assumption that molecular collisions at different instants are independent. This is mathematically expressed as:

$$\langle \xi(t) \xi(q) \rangle = 0 \quad \text{for } t \neq q$$

More generally, the autocorrelation function is given by the Dirac delta:

$$\langle \xi(t) \xi(q) \rangle = \delta(t - q)$$

where $\delta(t - q)$ is the Dirac delta function, indicating that the noise is only correlated with itself at the same instant.

2. **Gaussian Distribution:** By the central limit theorem, the sum of many independent random impacts leads to a Gaussian distribution for $\xi(t)$ at any fixed time t . This ensures that all finite-dimensional distributions of the process are Gaussian, making the process analytically tractable.
3. **Zero Mean:** The random molecular collisions are, on average, isotropic and unbiased, so the mean of the stochastic force vanishes:

$$\langle \xi(t) \rangle = 0$$

where $\langle \cdot \rangle$ denotes the ensemble average.

4. **Infinite Instantaneous Variance:** The idealized white noise process is so irregular that its variance at any fixed time is formally infinite:

$$\langle \xi^2(t) \rangle \gg 1$$

This reflects the mathematical abstraction of white noise, which is not a function in the usual sense but a generalized function (distribution).

Warning: A Mathematical Abstraction

The concept of white noise, with its delta-correlated structure, is a powerful but physically unrealizable abstraction. It implies infinite variance ($\langle \xi(t)^2 \rangle \rightarrow \infty$) and infinite power. In any real system, correlations exist over some small but non-zero timescale. However, if this correlation time is much shorter than any other characteristic timescale of the system, modeling the noise as "white" is an excellent and mathematically convenient approximation.

Multiplicative Noise: The Case of the SIS Epidemic Model

Multiplicative noise arises when the amplitude of the noise depends on the state of the system itself. In the context of stochastic differential equations, this means the noise term is not simply additive, but is multiplied by a function of the state variable:

$$\frac{dx}{dt} = f(x, t) + g(x, t) \xi(t)$$

where $g(x, t)$ is not constant.

A classic example is found in epidemiology, specifically in the **SIS epidemic model**. In this model, individuals can become infected, recover, and then become susceptible again (as opposed to the SIR model, where recovered individuals are removed from the susceptible pool). The deterministic SIS model is given by:

$$\begin{aligned} \frac{dS}{dt} &= -\beta SI + \mu(N - S) \\ \frac{dI}{dt} &= \beta SI - \mu I \end{aligned}$$

where S and I are the numbers of susceptible and infected individuals, β is the infection rate, μ is the recovery rate, and N is the total population.

Incorporating Stochasticity

In reality, the contact rate β is not perfectly constant: it fluctuates due to random factors such as social behavior, environmental changes, or other sources of randomness. To model this, we

introduce a stochastic term:

$$\beta \rightarrow \beta + \sigma \xi(t)$$

where σ quantifies the strength of the noise and $\xi(t)$ is Gaussian white noise.

Substituting this into the SIS model, the equation for the number of infected individuals becomes:

$$\frac{dI}{dt} = [\beta SI - \mu I] + \sigma SI \xi(t)$$

Here, the noise term is ***multiplicative***: its amplitude depends on both S and I . This reflects the fact that random fluctuations in the infection rate have a larger effect when there are more susceptible and infected individuals interacting.

Observation: Multiplicative vs. Additive Noise

Additive noise affects the system independently of its state, while multiplicative noise depends on the current state. In epidemic models, multiplicative noise is more realistic, as the impact of random fluctuations in transmission is naturally proportional to the number of possible contacts.

2.2.3 The Simplest Langevin Equation (SLAE)

To build a deeper intuition for the behavior of stochastic systems, we analyze the simplest possible Langevin equation, where the deterministic drift term is zero ($f(x) = 0$). This corresponds to a free particle subject only to random kicks from its environment:

$$\dot{x} = \omega \xi(t)$$

Here, ω is a constant representing the intensity of the noise $\xi(t)$. While this equation is ill-defined in its derivative form, we can formally integrate it to find the particle's position $x(t)$, assuming an initial position $x(0)$:

$$x(t) = x(0) + \omega \int_0^t \xi(s) ds$$

From this expression, we can derive the fundamental statistical properties of the process $x(t)$.

1. Average Position

The average position, or the first moment of $x(t)$, is found by taking the ensemble average of the equation. Since the noise has zero mean, $\langle \xi(s) \rangle = 0$, the integral of the average noise vanishes:

$$\langle x(t) \rangle = \langle x(0) \rangle + \omega \int_0^t \langle \xi(s) \rangle ds = \langle x(0) \rangle$$

If the initial position is deterministic, $\langle x(0) \rangle = x_0$, the average position of the particle does not change over time, since the particle is equally likely to be pushed in any direction.

2. Autocorrelation and "Memory"

To understand how the position at one time relates to the position at another, we compute the autocorrelation function, $\langle x(t)x(q) \rangle$.

$$\langle x(t)x(q) \rangle = \left\langle \left(\omega \int_0^t \xi(s) ds \right) \left(\omega \int_0^q \xi(\theta) d\theta \right) \right\rangle = \omega^2 \int_0^t \int_0^q \langle \xi(s) \xi(\theta) \rangle ds d\theta$$

Using the white noise property $\langle \xi(s) \xi(\theta) \rangle = \delta(s - \theta)$, we get:

$$\langle x(t)x(q) \rangle = \omega^2 \int_0^t \left(\int_0^q \delta(s - \theta) d\theta \right) ds$$

The inner integral with respect to θ is 1 if s is within the interval $[0, q]$, and 0 otherwise. This simplifies the double integral. Assuming, without loss of generality, that $t \leq q$, the condition $s \in [0, q]$ is always met for the outer integral's range $s \in [0, t]$. The integral thus becomes:

$$\langle x(t)x(q) \rangle = \begin{cases} \omega^2(t-q) & \text{if } t > q \\ \omega^2(q-t) & \text{if } t < q \end{cases} = \omega^2|t-q|$$

We can combine these cases into a single elegant expression:

$$\langle x(t)x(q) \rangle = \omega^2 \min(t, q)$$

This result reveals that, unlike the driving noise $\xi(t)$, the position process $x(t)$ **does** have memory. Its position at time t is correlated with its position at all other times.

3. Mean Squared Displacement and Variance

Setting $q = t$ gives the second moment, or **mean squared displacement** from the origin (assuming $x(0) = 0$):

$$\langle x^2(t) \rangle = \omega^2 t$$

The **variance** of the process is then:

$$\text{Var}[x(t)] = \langle x^2(t) \rangle - \langle x(t) \rangle^2 = \omega^2 t$$

The variance grows linearly with time, a hallmark of diffusive processes. The particle, on average, wanders further and further from its starting point.

4. Mean Squared Increment and Non-Differentiability

Let's examine the behavior of the process over a small time increment. The mean squared change in position over an interval of length $|t - q|$ is:

$$\begin{aligned} \langle (x(t) - x(q))^2 \rangle &= \langle x^2(t) \rangle + \langle x^2(q) \rangle - 2\langle x(t)x(q) \rangle \\ &= \omega^2 t + \omega^2 q - 2\omega^2 \min(t, q) \\ &= \omega^2 |t - q| \end{aligned}$$

Now consider the incremental ratio, which approximates the derivative. Let $q = t + h$:

$$\left\langle \left(\frac{x(t+h) - x(t)}{h} \right)^2 \right\rangle = \frac{\langle (x(t+h) - x(t))^2 \rangle}{h^2} = \frac{\omega^2 h}{h^2} = \frac{\omega^2}{h}$$

In the limit as the interval h shrinks to zero, this quantity diverges:

$$\lim_{h \rightarrow 0^+} \left\langle \left(\frac{x(t+h) - x(t)}{h} \right)^2 \right\rangle = \lim_{h \rightarrow 0^+} \frac{\omega^2}{h} = +\infty$$

The mean squared value of the derivative is infinite. This is a profound result: it is the first formal evidence that the path $x(t)$, while continuous, is **nowhere differentiable**. This highly irregular, "jagged" nature is a fundamental property of processes driven by white noise.

We said that the measure of the memory of the white noise is represented by the following average:

$$\langle \xi(t)\xi(q) \rangle = \delta(qt)$$

and we showed some heuristic reasoning. However, for the sake of precision, the above formula is derived by a statistical concept that is very useful to verify the degree of similarity between two generic random variables: the **covariance**, which is defined as follows:

$$\text{COV}(x, y) = \langle (x - \langle x \rangle)(y - \langle y \rangle) \rangle = \langle xy \rangle - \langle x \rangle \langle y \rangle$$

For the sake of precision, in the cases where y is a temporal shift of x (as in our case), we say that we are computing the **autocovariance** of x .

2.2.4 The Wiener Process and its Properties

The mathematical pathologies of white noise $\xi(t)$ and the resulting non-differentiability of its integral motivated the development of a more rigorous framework. Instead of focusing on the ill-defined derivative form $\dot{x} = \xi(t)$, we work directly with the integrated process. This process, which formalizes the concept of Brownian motion, is named in honor of Norbert Wiener.

Definition: The Wiener Process

The **Wiener process** $W(t)$ is a continuous-time stochastic process that is the solution to the SDE:

$$\begin{cases} \frac{dW}{dt} = \xi(t) \\ W(0) = 0 \end{cases}$$

where $\xi(t)$ is Gaussian white noise with unit variance.

Note: While this differential form is intuitive, it should be understood that $\xi(t)$ is not a function in the classical sense, and the equation is interpreted in the sense of stochastic integration.

The Wiener process is a Gaussian process with the following key properties:

1. **Zero Mean:** The expected value of the process is zero:

$$\langle W(t) \rangle = 0$$

This property reflects the symmetric nature of the random fluctuations around the origin.

2. **Gaussian Distribution:** For any fixed time t , $W(t)$ is normally distributed with mean zero and variance t :

$$W(t) \sim \mathcal{N}(0, t)$$

This means the probability density function is given by:

$$p(W, t) = \frac{1}{\sqrt{2\pi t}} \exp\left(-\frac{W^2}{2t}\right)$$

3. **Autocorrelation:** The autocorrelation function is given by

$$\langle W(t)W(s) \rangle = \min(t, s)$$

This property captures the "memory" of the process.

4. **Independent Increments:** The increments of the process are independent. In particular, for any collection of non-overlapping time intervals $[t_1, s_1], [t_2, s_2], \dots, [t_n, s_n]$, the corresponding increments $W(s_1) - W(t_1), W(s_2) - W(t_2), \dots, W(s_n) - W(t_n)$ are independent random variables. The mean of each increment is:

$$\langle W(s) - W(t) \rangle = 0$$

and these increments are also Gaussian.

5. **Increment Variance:**

The variance of the increment over the interval $[t, s]$ is proportional to the time difference:

$$\langle (W(s) - W(t))^2 \rangle = |s - t|$$

This scaling with time difference is a fundamental property that distinguishes diffusive processes from other types of motion.

6. **Increment Distribution:** More precisely, for $s > t$, the increment is distributed as

$$W(s) - W(t) \sim \mathcal{N}(0, |s-t|)$$

This means the increment has a normal distribution with zero mean and variance equal to the time difference. In particular, the distribution of $W(t)$ itself (starting from $W(0) = 0$) is:

$$W(t) \sim \mathcal{N}(0, t) \quad \Rightarrow \quad p(W, t) = \frac{1}{\sqrt{2\pi t}} \exp\left(-\frac{W^2}{2t}\right)$$

These properties completely characterize the Wiener process. Any continuous-time stochastic process satisfying these conditions is a Wiener process, making it a fundamental building block in the theory of stochastic processes and stochastic differential equations.

Increment Analysis and Non-Differentiability

The properties of the Wiener process's increments lead to one of its most counter-intuitive and important features: its paths are continuous but almost surely nowhere differentiable. Let's analyze an infinitesimal increment $dW_t = W(t+dt) - W(t)$ over a small time interval $dt > 0$. From the properties above, this increment is a Gaussian random variable with mean 0 and variance dt :

$$dW_t \sim \mathcal{N}(0, dt)$$

This implies that the standard deviation of the increment is $\sigma = \sqrt{dt}$. Now, consider the finite difference quotient used to define a derivative:

$$\frac{dW_t}{dt} = \frac{W(t+dt) - W(t)}{dt}$$

This new random variable is also Gaussian, but its variance is $\text{Var}\left(\frac{dW_t}{dt}\right) = \frac{1}{dt^2} \text{Var}(dW_t) = \frac{dt}{dt^2} = \frac{1}{dt}$. As $dt \rightarrow 0$, this variance diverges to infinity. The notion of a derivative in the classical sense is therefore problematic.

We can investigate this further by considering the probability that the magnitude of the incremental ratio exceeds some large, arbitrary value M . Let h be a small, finite time step. We are interested in:

$$\Pr\left(\left|\frac{W(t+h) - W(t)}{h}\right| > M\right)$$

Since the increment $W(t+h) - W(t)$ follows a $\mathcal{N}(0, h)$ distribution, we can standardize it by dividing by its standard deviation, \sqrt{h} . Let $Z = \frac{W(t+h) - W(t)}{\sqrt{h}}$, where $Z \sim \mathcal{N}(0, 1)$. The probability becomes:

$$\Pr\left(\left|\frac{Z\sqrt{h}}{h}\right| > M\right) = \Pr(|Z| > M\sqrt{h})$$

As $h \rightarrow 0$, the threshold $M\sqrt{h}$ also goes to zero. The probability that a standard normal variable Z has a magnitude greater than an infinitesimally small number approaches 1.

$$\lim_{h \rightarrow 0} \Pr(|Z| > M\sqrt{h}) = \Pr(|Z| > 0) = 1$$

This means that for any large M , as we make the time interval smaller, it becomes a near certainty that the slope of the secant line will exceed M . This confirms that the path of a Wiener process is not differentiable at any point.

3

Stochastic Numerical Methods

The non-differentiability of the Wiener process path fundamentally breaks the foundation of classical calculus, rendering traditional analytical techniques inadequate for stochastic differential equations. This mathematical obstacle necessitates the development of an entirely new framework—stochastic calculus, with its own integration theory and differentiation rules.

3.1 The Differential Form of SDEs: Itô Equation

Let's reconsider Newton's second law. In its most common form, it is written as $F = ma$. However, its more fundamental statement relates force to the change in momentum $p = mv$, expressed in differential form as:

$$dp = Fdt$$

This form is more general. For instance, consider the motion of a rocket, whose mass $m(t)$ changes as it consumes fuel. In this case, $F = ma$ is incorrect. The correct formulation is:

$$d(m(t)v(t)) = Fdt$$

This differential way of writing physical laws is powerful and provides the foundation for correctly interpreting stochastic equations. A Langevin equation written as $\dot{x} = f(x, t) + g(x, t)\xi(t)$ is mathematically problematic. The rigorous approach is to express it in its differential form using the Wiener process increment dW_t , which represents the integral of the white noise $\xi(t)$:

$$dx = f(x, t)dt + g(x, t)dW_t$$

This is a **Stochastic Differential Equation (SDE)**. Its solution is understood in an integral sense:

$$x(t) = x_0 + \int_0^t f(x(\tau), \tau)d\tau + \int_0^t g(x(\tau), \tau)dW_\tau$$

The second integral is a stochastic integral, an object whose properties are fundamentally different from the standard Riemann integral. The infinitesimal increment dW_t is defined as $dW_t = W(t+dt) - W(t)$. As we have established, it is a Gaussian random variable with mean 0 and variance dt , so $dW_t \sim \mathcal{N}(0, dt)$. This can be expressed as:

$$dW_t = G(t)\sqrt{dt}$$

where $G(t)$ is a random variable drawn from the standard normal distribution, $\mathcal{N}(0, 1)$.

The stochastic integral $\int_0^t g(x(\tau), \tau)dW_\tau$ represents the cumulative effect of the random forcing over time. Unlike deterministic integrals, this integral cannot be evaluated using traditional calculus rules due to the irregular nature of the Wiener process paths. The integral must be understood in the sense of Itô or Stratonovich, with Itô integration being the more commonly used convention in stochastic differential equations.

Definition: Itô Equation

Given an SLAE, we can always rewrite it as an **Itô equation**, by defining $dW = G(t)\sqrt{dt}$:

$$dx = f(x)dt + g(x)dW$$

3.1.1 The Euler-Maruyama Method

The Itô equation can be solved numerically using the Maruyama algorithm for stochastic differential equations. This is essentially the stochastic version of Euler's algorithm. Given a time interval $[0, T]$ and $h = T/N$, so that $t = jh$ for $j = 0, \dots, N$. Suppose the Euler algorithm is:

$$dx = x(t_j + h) - x(t_j) = x(t_{j+1}) - x(t_j)$$

we can set $dt \approx h$ and write:

$$x(t_{j+1}) = x(t_j) + f(x(t_j))h$$

The Maruyama formula starts from this form by also considering the Gaussian effect, adding:

$$x(t_{j+1}) = x(t_j) + f(x(t_j))h + G_j\sqrt{h}$$

with $G_j \sim \mathcal{N}(0, 1)$. Obviously, starting from this algorithm, more precise variants have been successively created.

Definition: The Euler-Maruyama Method

For the stochastic differential equation $dX_t = a(X_t, t)dt + b(X_t, t)dW_t$ with initial condition $X(0) = x_0$ and uniform time step h , the Euler-Maruyama approximation is given by:

$$X_{j+1} = X_j + a(X_j, t_j)h + b(X_j, t_j)\sqrt{h}G_j$$

where $t_j = jh$ and $\{G_j\}_{j=0}^{N-1}$ is a sequence of independent standard normal random variables.

This numerical scheme provides a practical foundation for simulating stochastic processes, though more sophisticated methods have been developed to improve accuracy and stability for specific applications.

3.1.2 Itô's Lemma: The Stochastic Chain Rule

Having established the framework of Itô equations, we now turn our attention to the fundamental problem of change of variables in stochastic calculus. When dealing with deterministic differential equations, the chain rule provides a straightforward mechanism for transforming variables. However, in the stochastic setting, the irregular nature of Brownian motion necessitates a more sophisticated approach.

Consider an SDE in Itô form:

$$dx = a(x)dt + b(x)dW$$

Suppose we wish to perform a transformation from x to y defined by $y = \psi(x)$. Our objective is to derive an equation of similar form for the transformed variable y .

Following the classical approach, we expand dy in a Taylor series:

$$\begin{aligned} dy &= \psi'(x)dx + \frac{1}{2}\psi''(x)(dx)^2 + \dots \\ &= \Psi'(x)[a(x)dt + b(x)dW] + \frac{1}{2}\Psi''(x) \left[b^2(x)(dW)^2 + \underbrace{a^2(x)(dt)^2}_{O(dt^2)} + \underbrace{2a(x)b(x)dtdW}_{O(dt^{3/2})} \right] \end{aligned}$$

To derive a formula consistent with the Itô framework, we must carefully consider the order of magnitude of each term. Since $dW = O(\sqrt{dt})$, we can eliminate terms of order $(dt)^{3/2}$ and higher, including the mixed term $dtdW$ and $(dt)^2$. The remaining term $(dW)^2$ is of order dt , and by the fundamental property of Brownian motion, we have $(dW)^2 = dt$.

Substituting this, we get:

$$dy = \psi'(x)a(x)dt + \psi'(x)b(x)dW + \frac{b(x)^2}{2}\psi''(x)dt$$

Realigning the terms, we get:

$$dy = \left[\psi'(x)a(x) + \frac{b(x)^2}{2}\psi''(x) \right] dt + \psi'(x)b(x)dW$$

Definition: Itô's Lemma

Let X_t be an Itô process that satisfies the SDE $dX_t = a(X_t, t)dt + b(X_t, t)dW_t$. Let $\psi(x, t)$ be a twice-differentiable scalar function. Then the process $Y_t = \psi(X_t, t)$ is also an Itô process, and its differential dY_t is given by:

$$dY_t = \left[\psi'(x)a(x) + \psi''(x)\frac{b(x)^2}{2} \right] dt + \psi'(x)b(x)dW_t$$

This fundamental result encompasses the classical chain rule terms, namely $\psi'a$ and $\psi'b$, augmented by an additional stochastic correction term, $\psi''\frac{b^2}{2}$, known as the **Itô correction**. This correction term arises directly from the non-vanishing quadratic variation of the Wiener process and represents the fundamental distinction between stochastic and deterministic calculus.

3.2 The Stochastic Malthusian Model and its Paradox

The simplest model of population growth is the Malthusian model, which assumes an unlimited environment and a constant per capita growth rate, r . The dynamics are described by the ODE:

$$\dot{x} = rx$$

where $x(t)$ represents the population size. The solution is simple exponential growth or decay, $x(t) = x(0)e^{rt}$. A more realistic model acknowledges that the growth rate is not constant but fluctuates randomly due to environmental variations. We can model this by making the growth rate a stochastic process, $r \rightarrow r + \omega\xi(t)$. This leads to the **stochastic Malthusian model**, an SDE with multiplicative noise:

$$dx = (r + \omega\xi(t))x$$

where ω denotes the multiplicative noise amplitude. We can obtain the Ito formula:

$$dx = rxdt + \omega x dW$$

The solution of this stochastic equation can be obtained by applying Itô's formula. Setting $y = \ln x$, we find the drift and diffusion coefficients for the transformed process. Using Itô's rule:

$$dy = \left[\frac{1}{x}rx - \frac{1}{2x^2}\omega^2x^2 \right] dt + \frac{1}{x}\omega x dW \quad \xrightarrow{\text{simplifying}} \quad dy = \left[r - \frac{\omega^2}{2} \right] dt + \omega dW$$

This simplifies to a linear SDE which can be solved directly. Integrating from 0 to t :

$$y(t) = y_0 + \left(r - \frac{\omega^2}{2} \right) t + \omega W(t)$$

Supposing $W_0 = 0$, we can calculate the moments of $y(t)$:

$$\langle y(t) \rangle = \langle y_0 \rangle + \left(r - \frac{\omega^2}{2} \right) t$$

Therefore, if we have $\omega^2/2 > r$, then $y(t) \rightarrow -\infty$. We notice that, intuitively considering what $y(t)$ represents, if ω is large the population will tend to extinction regardless of r . In other words, a population highly subject to events, whether negative or positive, will tend to extinction.

We can then find the second moment of $y(t)$:

$$\langle y(t)^2 \rangle = \left\langle \left[y_0 + \left(r - \frac{\omega^2}{2} \right) t \right]^2 + 2\omega \left[y_0 + \left(r - \frac{\omega^2}{2} \right) t \right] W(t) + \omega^2 W(t)^2 \right\rangle$$

which, upon solving, becomes:

$$\dots = \left[y_0 + \left(r - \frac{\omega^2}{2} \right) t \right]^2 + \omega^2 t$$

Assuming for convenience that y_0 is deterministic, then we have that the variance of $y(t)$ is:

$$\text{Var}[y(t)] = \langle y(t)^2 \rangle - \langle y(t) \rangle^2 = \omega^2 t$$

Again, we can note that the variance tends to diverge over time. Returning now to $x(t)$, we have:

$$x(t) = e^{y(t)} = e^{y_0 + \left(r - \frac{\omega^2}{2} \right) t} e^{\omega W(t)} = x_0 e^{\left(r - \frac{\omega^2}{2} \right) t} e^{\omega W(t)}$$

Then its mean value will be:

$$\langle x(t) \rangle = x_0 e^{\left(r - \frac{\omega^2}{2} \right) t} \langle e^{\omega W(t)} \rangle$$

So, we must compute the mean of $\exp(\omega W(t))$. We know that $W(t)$ is distributed with the distribution:

$$W(t) \sim N(0, t) = \frac{1}{\sqrt{2\pi t}} e^{-\frac{W^2}{2t}}$$

so, to calculate the mean, we must compute:

$$\int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi t}} e^{\omega W} e^{-\frac{W^2}{2t}} dW$$

We can then combine the two exponentials and rewrite the resulting exponent as:

$$e^{\frac{-W^2 + 2t\omega W - t^2\omega^2 + t^2\omega^2}{2t}} = e^{-\frac{(W-t\omega)^2}{2t}} e^{\frac{\omega^2 t}{2}}$$

As a consequence, the integral rewritten this way becomes:

$$\dots = e^{\frac{\omega^2 t}{2}} \int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi t}} e^{-\frac{(W-t\omega)^2}{2t}} dW$$

This, however, is the integral of a translated Gaussian which is trivially equal to 1. Therefore:

$$\langle e^{\omega W(t)} \rangle = e^{\frac{\omega^2 t}{2}}$$

Thus,

$$\langle x(t) \rangle = \langle x_0 \rangle e^{\left(rt - \frac{\omega^2 t}{2} \right) + \frac{\omega^2 t}{2}}$$

In the Langevin hypothesis, it was the fact that the non-stochastic differential equation was nothing other than the differential equation of the mean, and here we find this result. Indeed, the two terms subtract, giving as a solution the solution of the non-stochastic case.

The Paradox

However, something strange emerges when we examine the long-term behavior. Returning to the definition of $y(t)$, we can rewrite it as:

$$y(t) = y_0 + \left(r - \frac{\omega^2}{2} \right) t + \omega\sqrt{t}G$$

As $t \rightarrow \infty$, the deterministic term $\left(r - \frac{\omega^2}{2} \right) t$ dominates, so:

$$\lim_{t \rightarrow +\infty} y(t) = \begin{cases} -\infty & \text{if } r < \frac{\omega^2}{2} \\ +\infty & \text{if } r > \frac{\omega^2}{2} \end{cases}$$

Since $x(t) = e^{y(t)}$, when $y(t) \rightarrow -\infty$, we have $x(t) \rightarrow 0$. This creates an apparent paradox: we showed that $\langle x(t) \rangle = x_0 e^{rt}$, which always grows exponentially regardless of ω , yet individual realizations $x(t)$ can tend to extinction when the noise is sufficiently large ($\omega^2/2 > r$).

The resolution lies in understanding the nature of the log-normal distribution. When $x(t) = e^{y(t)}$ where $y(t)$ is normally distributed, $x(t)$ follows a log-normal distribution, which is highly skewed. The mean and median can differ dramatically in such distributions.

Example: The Exponential Distribution

Consider the simple exponential distribution with density ae^{-aU} . Its mean is $\langle U \rangle = 1/a$, but its median satisfies:

$$\int_0^{\text{MED}} ae^{-aU} dU = \frac{1}{2}$$

which gives $\text{MED} = \log(2)/a < 1/a$. The median is smaller than the mean because the distribution has a long tail that pulls the mean upward.

Similarly, in our stochastic population model, $x(t)$ has a log-normal distribution. While the mean grows exponentially due to rare but extremely large population bursts, the typical behavior (represented by the median) can show decline when noise is large. The mean becomes dominated by infrequent but massive population explosions, making it a poor representative of typical outcomes.

Calculation of the Median

To quantify the typical behavior of the population, we need to calculate the median of $x(t)$. Since $y(t) \sim \mathcal{N}(\mu, \omega^2)$ where $\mu = y_0 + (r - \omega^2/2)t$ and the variance parameter is $\omega^2 t$, the random variable $x(t) = e^{y(t)}$ follows a log-normal distribution with probability density function:

$$f_X(x) = \frac{1}{\sqrt{2\pi}\omega\sqrt{t} \cdot x} \exp\left(-\frac{(\log x - \mu)^2}{2\omega^2 t}\right), \quad x > 0$$

The median m is defined as the value satisfying $P(X \leq m) = 1/2$, which gives us the integral equation:

$$\frac{1}{2} = \int_0^m \frac{1}{\sqrt{2\pi}\omega\sqrt{t} \cdot x} \exp\left(-\frac{(\log x - \mu)^2}{2\omega^2 t}\right) dx$$

To solve this integral, we employ the substitution $w = \log x$, which transforms $dx = e^w dw = x dw$. The limits of integration become $w \in (-\infty, \log m]$, and our integral becomes:

$$\frac{1}{2} = \int_{-\infty}^{\log m} \frac{1}{\sqrt{2\pi}\omega\sqrt{t}} \exp\left(-\frac{(w - \mu)^2}{2\omega^2 t}\right) dw$$

This is precisely the cumulative distribution function of a normal random variable $\mathcal{N}(\mu, \omega^2 t)$ evaluated at $\log m$. Since the median of any normal distribution equals its mean, we have:

$$\log m = \mu = y_0 + \left(r - \frac{\omega^2}{2} \right) t$$

Therefore, the median of $x(t)$ is:

$$\text{MED}[x(t)] = e^\mu = e^{y_0 + \left(r - \frac{\omega^2}{2} \right) t} = x_0 e^{\left(r - \frac{\omega^2}{2} \right) t}$$

This result elegantly resolves the paradox. While the mean $\langle x(t) \rangle = x_0 e^{rt}$ always grows exponentially, the median, which better represents typical population trajectories, follows the drift-corrected dynamics. The long-term behavior of the median is:

$$\lim_{t \rightarrow +\infty} \text{MED}[x(t)] = \begin{cases} 0 & \text{if } r < \frac{\omega^2}{2} \quad (\text{extinction regime}) \\ x_0 & \text{if } r = \frac{\omega^2}{2} \quad (\text{critical regime}) \\ +\infty & \text{if } r > \frac{\omega^2}{2} \quad (\text{growth regime}) \end{cases}$$

The threshold $r = \omega^2/2$ represents a critical noise level: below this threshold, typical populations grow indefinitely, while above it, typical populations face extinction despite the exponentially growing mean. This dichotomy between mean and median behavior is a hallmark of log-normal processes and highlights the importance of choosing appropriate summary statistics for highly skewed distributions.

3.3 The Perturbed Logistic Equation

We have discussed the fact that the Malthus model is not realistic because it assumes infinite resources. A much better model for population dynamics is the **Logistic model**, which incorporates a density-dependent growth rate. The deterministic logistic model can be written as follows:

$$\dot{x} = r(x)x$$

where $r(x)$ is a decreasing function of the population size x . The simplest choice for this function is linear:

$$r(x) = r_0 - \alpha x$$

Here, r_0 is the intrinsic growth rate at low densities, and α is a coefficient representing the strength of density-dependent regulation (e.g., competition for resources).

Introducing Stochasticity

In a realistic scenario, the intrinsic growth rate r_0 is not constant but fluctuates due to environmental variability. We can model its fast fluctuations as:

$$r_0 \rightarrow r_0 + \omega \xi(t)$$

where $\xi(t)$ is Gaussian white noise. This transforms the deterministic ODE into the stochastically perturbed logistic model:

$$dx = (r_0 - \alpha x)x dt + \omega x dW_t$$

To analyze this equation, we can again use the logarithmic transformation $y = \ln(x)$. Applying Itô's formula yields:

$$dy = \left(r_0 - \frac{\omega^2}{2} - \alpha e^y \right) dt + \omega dW$$

This transformed SDE can be formally integrated to give:

$$y(t) = y(0) + \left(r_0 - \frac{\omega^2}{2} \right) t + \omega W(t) - \alpha \int_0^t e^{y(s)} ds$$

We can now analyze the long-term behavior of the system. We know two key facts:

1. If $\omega^2 > 2r_0$, then the non-integral part will tend to $-\infty$ as $t \rightarrow \infty$.
2. The integral of the exponential is always positive since $x = e^y$ must be positive.

Slightly more precisely, the fact that $\int_0^t e^{y(s)} ds > 0$ gives us an upper bound on the process $y(t)$:

$$y(t) \leq y(0) + \left(r_0 - \frac{\omega^2}{2} \right) t + \omega W(t)$$

This leads us to a **sufficient condition** for population extinction. If the right-hand side of the inequality goes to $-\infty$, then $y(t)$ must also go to $-\infty$. This happens when the drift of the bounding process is negative. Thus, the condition for extinction is:

$$\frac{\omega^2}{2} > r_0 \implies \lim_{t \rightarrow \infty} x(t) = 0$$

This powerful result shows that if the environmental noise is sufficiently strong, the population will go extinct regardless of the density-dependent term. The noise effectively suppresses the intrinsic growth.

3.4 Foundations of Ito Calculus

3.4.1 Stochastic Equilibrium Points

The concept of an equilibrium point can be extended from deterministic to stochastic systems. For a deterministic system $\dot{x} = f(x)$, an equilibrium point x_e is defined by the condition $f(x_e) = 0$. For a stochastic differential equation (SDE) of the form

$$dx = f(x)dt + g(x)dW_t,$$

a **stochastic equilibrium point (SEP)**, denoted x_{se} , is a point where both the drift and diffusion terms vanish simultaneously:

$$f(x_{se}) = 0 \quad \text{and} \quad g(x_{se}) = 0.$$

This dual condition makes SEPs significantly rarer than their deterministic counterparts. As an example, consider the **stochastic Malthusian model**:

$$dx = r_0 x dt + \omega x dW_t.$$

In this case, $x_{se} = 0$ is a trivial SEP, as both $f(0) = r_0 \cdot 0$ and $g(0) = \omega \cdot 0$ are zero. The stability of this equilibrium is determined by the interplay between the growth rate r_0 and the noise intensity ω , as established previously:

- If $\omega^2/2 > r_0$, the median of the population size converges to zero, rendering the equilibrium point $x_{se} = 0$ **stochastically globally attractive**.
- If $\omega^2/2 < r_0$, the median grows exponentially, and the equilibrium point $x_{se} = 0$ acts as a **stochastic repulsor** (i.e., it is unstable).

This analysis extends to the **perturbed logistic model**, $dx = (r_0 x - \alpha x^2)dt + \omega x dW_t$, where $x_{se} = 0$ is also a SEP. The condition $\omega^2/2 > r_0$ remains sufficient for the global attractivity of the origin.

3.4.2 Derivation of Ito's Formula

We recall that during the analysis of the change of variable for Ito's formula, we substituted $(dW)^2 \rightarrow dt$ into the equation without formally proving its correctness. Let us now investigate this further, starting from the general SDE:

$$dx = a(x)dt + b(x)dW$$

We know that the Wiener increment can be expressed as $dW = \sqrt{dt}G(t)$, where $G(t) \sim \mathcal{N}(0, 1)$. Our goal was to apply a variable transformation $y = \psi(x)$ to this equation. We found that this transformation resulted in:

$$dy \equiv d\psi = dt \left[\psi'(x)a(x) + \psi''(x)\frac{b(x)^2}{2} \right] + \psi'(x)b(x)dW$$

This result was achieved through the "magical" substitution mentioned earlier. The objective was to obtain a new SDE for y that has the same form as the original one:

$$dy = q(y)dt + r(y)dW$$

This required having one term of order $O(dt)$ and another of order $O(\sqrt{dt})$, and to achieve this, we discarded all terms of higher order than dt . To verify this substitution, we must revisit a concept from earlier. Let us consider the increment of the Wiener process:

$$z = W(t+h) - W(t)$$

If we now consider the random variable $q = z^2$, we have $\langle q \rangle = h$, which is the variance of z . This was the rationale for the substitution we made in the derivation of dW , but in doing so, we were neglecting potentially important elements. Let us now evaluate the variance of q .

To formally establish this, we compute the fourth moment of a Gaussian random variable $z \sim \mathcal{N}(0, h)$, which represents the Wiener increment dW_t over a time step $h = dt$.

$$\langle z^4 \rangle = \int_{-\infty}^{+\infty} z^4 \frac{1}{\sqrt{2\pi h}} e^{-\frac{z^2}{2h}} dz.$$

Integration by parts, with $u = z^3$ and $dv = ze^{-z^2/2h}dz/\sqrt{2\pi h}$, yields:

$$\langle z^4 \rangle = \left[z^3 \left(-\frac{h}{\sqrt{2\pi h}} e^{-z^2/2h} \right) \right]_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} \left(-\frac{h}{\sqrt{2\pi h}} e^{-z^2/2h} \right) 3z^2 dz.$$

The boundary term vanishes, leaving:

$$\langle z^4 \rangle = 3h \int_{-\infty}^{+\infty} z^2 \frac{1}{\sqrt{2\pi h}} e^{-\frac{z^2}{2h}} dz = 3h \langle z^2 \rangle = 3h(h) = 3h^2.$$

With $h = dt$, this gives $\langle (dW_t)^4 \rangle = 3(dt)^2$. The variance of $(dW_t)^2$ is therefore:

$$\text{Var}[(dW_t)^2] = \langle (dW_t)^4 \rangle - (\langle (dW_t)^2 \rangle)^2 = 3(dt)^2 - (dt)^2 = 2(dt)^2.$$

Since the variance is of order $(dt)^2$, the fluctuations of $(dW_t)^2$ around its mean dt are of a higher order than dt itself. Consequently, in the limit $dt \rightarrow 0$, we can make the substitution $(dW_t)^2 = dt$. Substituting this result into the Taylor expansion for $d\Psi$ gives the final expression:

$$d\Psi = \left[\Psi'(x)a(x) + \frac{1}{2}\Psi''(x)b^2(x) \right] dt + \Psi'(x)b(x)dW,$$

which is the celebrated **Ito's formula**. The presence of the second-derivative term, arising from the non-zero quadratic variation of the Wiener process, is the fundamental feature that distinguishes stochastic from ordinary calculus.

3.4.3 Probability Density Function and its Evolution

Given a continuous-time, continuous-state stochastic process, $x(t)$, we often want to describe its behavior not by a single trajectory, but by the probability of finding the process in a certain state at a certain time. This is accomplished using the **probability density function (PDF)**, denoted $\rho(x, t)$. The PDF is defined such that the probability of the random variable $x(t)$ being in an infinitesimal interval $[\hat{x}, \hat{x} + d\hat{x}]$ is given by:

$$\text{Prob}(x(t) \in [\hat{x}, \hat{x} + d\hat{x}]) = \rho(\hat{x}, t)d\hat{x}$$

This implies that the probability of finding $x(t)$ in a finite interval $[a, b]$ is the integral of the PDF over that interval:

$$\text{Prob}(x(t) \in [a, b]) = \int_a^b \rho(x, t) dx$$

A central question in stochastic modeling is: if we know the initial distribution of the process, $\rho(x, 0)$, how does this distribution evolve for $t > 0$?

In the most general case, the law of evolution for $\rho(x, t)$ could depend on the entire history of the process, often denoted as $\Omega(x(\theta), 0 \leq \theta < t)$. This would mean that the future state depends on the full path taken to reach the present, leading to a complex law that could be described by an integro-differential equation.

However, for a large and very important class of processes, the situation is significantly simpler.

The Markov Property and Stochastic Differential Equations

Processes described by an Itô Stochastic Differential Equation (SDE) of the form

$$dx = f(x)dt + g(x)dW_t$$

have a special structure. The state of the system at an infinitesimal future time $t + dt$ is given by:

$$x(t + dt) = x(t) + f(x(t))dt + g(x(t))dW_t$$

Crucially, the statistical properties of $x(t + dt)$ depend only on the state $x(t)$ at the current time t , and not on the entire prior history. This is the hallmark of a **Markov process**.

Definition: Markov Process

A stochastic process $x(t)$ is called a **Markov process** if its future probability distribution, conditioned on its past and present values, depends only on the present value. In other words, the past and the future are conditionally independent given the present.

For a process modeled by an Itô SDE, the following key properties hold:

1. **The process is Markovian.** The future state depends only on the present, not the path taken to get there.
2. **The process has continuous paths.** The Itô equation implies that increments are infinitesimal. Since the drift term $f(x)dt$ is of order $O(dt)$ and the diffusion term $g(x)dW_t$ is of order $O(\sqrt{dt})$, the process $x(t)$ does not have finite jumps.

These two properties together, being Markovian and having continuous paths, have a profound implication: the law governing the evolution of the PDF, $\rho(x, t)$, must be local in both time and space. It cannot depend on spatially distant values or past temporal values. This means the evolution operator must be a local differential operator. Therefore, the evolution equation for $\rho(x, t)$ must be a **Partial Differential Equation (PDE)**.

3.4.4 The Fokker-Planck Equation

While Ito's Lemma allows us to find the SDE for a transformed variable, its most powerful application is in deriving a deterministic equation for the evolution of the probability density function (PDF), $\rho(x, t)$. This bridge from the stochastic world of individual paths to the deterministic world of distributions is the celebrated **Fokker-Planck equation**.

The derivation is a beautiful piece of mathematical physics that relies on a "weak" formulation. Instead of tracking $\rho(x, t)$ directly, we analyze how the expected value of an arbitrary, well-behaved "test function" $\psi(x)$ evolves over time.

Derivation from Ito's Lemma

Let $x(t)$ be a process governed by the SDE $dx = a(x)dt + b(x)dW_t$. Let $\psi(x)$ be an arbitrary, twice-differentiable function that vanishes (along with its derivatives) at the boundaries of the domain (e.g., at $\pm\infty$).

1. **Apply Ito's Lemma to $\psi(x)$**

From Ito's Lemma, we know the differential for $\psi(x(t))$ is:

$$d\psi(x) = \left[a(x)\psi'(x) + \frac{b(x)^2}{2}\psi''(x) \right] dt + b(x)\psi'(x)dW_t$$

2. Take the Expectation

Next, we take the ensemble average (expectation) of this equation.

$$\langle d\psi(x) \rangle = \left\langle \left[a(x)\psi'(x) + \frac{b(x)^2}{2}\psi''(x) \right] dt \right\rangle + \langle b(x)\psi'(x)dW_t \rangle$$

The expectation of the stochastic term vanishes, because dW_t has zero mean and is independent of the state $x(t)$ at the beginning of the infinitesimal step. This leaves us with an equation for the evolution of the mean of $\psi(x)$:

$$d\langle \psi(x) \rangle = \left\langle a(x)\psi'(x) + \frac{b(x)^2}{2}\psi''(x) \right\rangle dt$$

Dividing by dt , we get the time derivative:

$$\frac{d}{dt} \langle \psi(x) \rangle = \left\langle a(x)\psi'(x) + \frac{b(x)^2}{2}\psi''(x) \right\rangle$$

3. Introduce the PDF

The expectation of any function $h(x)$ can be written as an integral of that function against the PDF $\rho(x,t)$. Applying this to both sides of our equation:

$$\frac{d}{dt} \int_S \psi(x)\rho(x,t)dx = \int_S \left[a(x)\psi'(x) + \frac{b(x)^2}{2}\psi''(x) \right] \rho(x,t)dx$$

where S is the state space. Since $\psi(x)$ does not depend on time, we can bring the time derivative inside the integral on the left:

$$\int_S \psi(x) \frac{\partial \rho(x,t)}{\partial t} dx = \int_S \psi'(x)a(x)\rho(x,t)dx + \int_S \psi''(x) \frac{b(x)^2}{2} \rho(x,t)dx$$

4. Integration by Parts

The goal now is to remove the derivatives from the test function $\psi(x)$ and transfer them onto the other terms using integration by parts. The general formula is $\int u dv = [uv] - \int v du$.

For the first term on the right, let $u = a(x)\rho(x,t)$ and $dv = \psi'(x)dx$.

$$\int_S \psi'(x)[a(x)\rho(x,t)]dx = \boxed{\underbrace{[\psi(x)a(x)\rho(x,t)]}_{\partial S} - \int_S \psi(x) \frac{\partial}{\partial x} [a(x)\rho(x,t)]dx}$$

For the second term on the right, we must integrate by parts twice.

$$\int_S \psi''(x) \left[\frac{b(x)^2}{2} \rho(x,t) \right] dx = \boxed{\left[\psi'(x) \frac{b^2 \rho}{2} \right]_{\partial S} - \int_S \psi'(x) \frac{\partial}{\partial x} \left[\frac{b^2 \rho}{2} \right] dx}$$

Due to our assumption that ψ and its derivatives are zero at the boundary ∂S , all the boundary terms vanish. We apply integration by parts again to the remaining integral:

$$\boxed{- \int_S \psi'(x) \frac{\partial}{\partial x} \left[\frac{b^2 \rho}{2} \right] dx = - \left[\psi(x) \frac{\partial}{\partial x} \left[\frac{b^2 \rho}{2} \right] \right]_{\partial S} + \int_S \psi(x) \frac{\partial^2}{\partial x^2} \left[\frac{b(x)^2}{2} \rho(x,t) \right] dx}$$

The boundary term again vanishes, leaving only the final integral.

5. The Final Equation

Substituting these results back into our main equation, we have:

$$\int_S \psi(x) \frac{\partial \rho}{\partial t} dx = - \int_S \psi(x) \frac{\partial}{\partial x} [a\rho] dx + \int_S \psi(x) \frac{\partial^2}{\partial x^2} \left[\frac{b^2 \rho}{2} \right] dx$$

We can now group all terms under a single integral:

$$\int_S \psi(x) \left[\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} [a\rho] - \frac{\partial^2}{\partial x^2} \left[\frac{b^2 \rho}{2} \right] \right] dx = 0$$

This equation must hold for **any** valid choice of the test function $\psi(x)$. The only way an integral can be zero for every arbitrary test function is if the expression inside the square brackets is itself identically zero.

This gives us the celebrated Fokker-Planck equation.

Definition: **Fokker-Planck Equation**

Given a Stochastic Differential Equation of the type:

$$dx = a(x)dt + b(x)dW$$

its associated Probability Density Function, $\rho(x,t)$, will solve the Fokker-Planck equation:

$$\frac{\partial \rho}{\partial t} = - \frac{\partial}{\partial x} (a(x)\rho(x,t)) + \frac{\partial^2}{\partial x^2} \left(\frac{b(x)^2}{2} \rho(x,t) \right)$$

This must, of course, be supplemented with boundary conditions (e.g., Dirichlet) and an initial condition $\rho(x,0)$ to obtain a unique solution. Note that imposing a specific initial state $x(0)$ actually yields the conditional PDF, $\rho(x,t|x(0))$.

In summary, the probability density function $\rho(x,t)$ associated with the stochastic process governed by the SDE above satisfies the celebrated **Fokker-Planck equation** (also known as the forward Kolmogorov equation):

$$\frac{\partial}{\partial t} \rho(x,t) = - \frac{\partial}{\partial x} [a(x)\rho(x,t)] + \frac{\partial^2}{\partial x^2} \left(\frac{b^2(x)}{2} \rho(x,t) \right)$$

This partial differential equation describes the time evolution of the probability density $\rho(x,t)$ for the random variable x .

Since $\rho(x,t)$ is a probability density function, it must always satisfy the **normalization constraint**:

$$\int_{-\infty}^{\infty} \rho(x,t) dx = 1$$

This ensures that the total probability is conserved at all times.

To uniquely determine the solution, we must also specify the **initial distribution** of x :

$$\rho(x,0) = \theta(x)$$

where $\theta(x)$ is the given initial probability density function (for example, it could be a Dirac delta function if the initial state is known exactly, or a broader distribution if there is uncertainty).

Together, the Fokker-Planck equation, the normalization condition, and the initial condition fully characterize the time evolution of the probability density for the stochastic process.

③ Example: Particle in a Potential Well with Additive Noise

Let us describe the motion of a particle in a conservative potential field $U(x)$, subject to random thermal fluctuations. In the overdamped limit ($m \ll 1$), where inertial effects are negligible, the system's dynamics can be described by:

$$dx = f(x)dt + \omega dW_t$$

where $f(x) = -\frac{\partial U}{\partial x}$, and ω represents the constant intensity of the noise. Our goal is to find the **steady-state probability distribution**. The Fokker-Planck equation for this system is:

$$\frac{\partial \rho(x,t)}{\partial t} = -\frac{\partial}{\partial x}[f(x)\rho(x,t)] + \frac{\omega^2}{2}\frac{\partial^2 \rho(x,t)}{\partial x^2}$$

At steady state, the distribution no longer changes with time, so $\frac{\partial \rho}{\partial t} = 0$. We denote the stationary solution as $P(x)$. The equation simplifies to an ordinary differential equation:

$$\frac{d}{dx} \left[-f(x)P(x) + \frac{\omega^2}{2} \frac{dP(x)}{dx} \right] = 0$$

The term in the square brackets (the probability current $J(x)$) must be constant. For a system with reflecting or natural boundaries at infinity, this current must be zero everywhere:

$$-f(x)P(x) + \frac{\omega^2}{2} \frac{dP(x)}{dx} = 0$$

Substituting $f(x) = -U'(x)$, this becomes a separable first-order ODE:

$$\frac{\omega^2}{2} \frac{dP}{dx} = f(x)P = -U'(x)P \implies \frac{dP}{P} = -\frac{2}{\omega^2} U'(x) dx$$

Integrating both sides gives $\ln P = -\frac{2}{\omega^2} U(x) + \text{const.}$ Exponentiating yields the solution:

$$P(x) = C \exp \left(-\frac{2}{\omega^2} U(x) \right)$$

This is the famous **Boltzmann distribution**. C is determined by the condition $\int_{\mathbb{R}} P(x) dx = 1$:

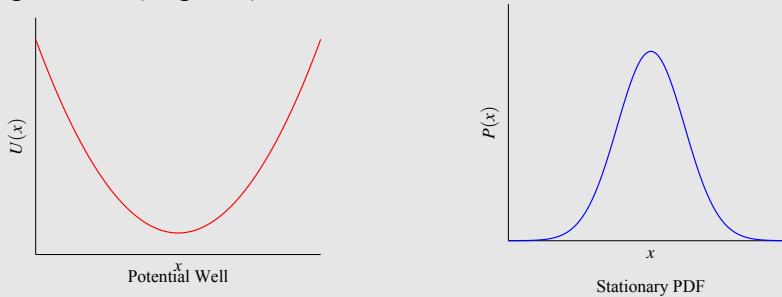
$$C = \frac{1}{\int_{-\infty}^{+\infty} \exp \left(-\frac{2}{\omega^2} U(x) \right) dx}$$

The particle is most likely to be found where $U(x)$ is lowest, though noise (ω) allows it to occasionally reach higher-energy states.

Let's consider the specific case of a harmonic potential well, $U(x) = \frac{1}{2}kx^2$. We have:

$$P(x) = C \exp \left(-\frac{k}{\omega^2} x^2 \right)$$

This is a Gaussian distribution centered at $x = 0$; the variance is $\sigma^2 = \omega^2 / (2k)$, showing that stronger noise (larger ω) leads to a wider distribution.



3.4.5 Special Cases of the Fokker-Planck Equation

Deterministic Systems with Random Initial Conditions

Suppose we have a deterministic system, $\dot{x} = a(x)$, but where the initial conditions are described by a probability distribution $\theta(x)$. Here, randomness comes only from the initial condition, not from noise in the dynamics: the diffusion coefficient is zero, $b(x) = 0$. The Fokker-Planck equation loses its second-order derivative term and becomes the **Liouville equation**:

$$\begin{cases} \frac{\partial P(x,t)}{\partial t} = -\frac{\partial}{\partial x}(a(x)P(x,t)) \\ P(x,0) = \theta(x) \end{cases}$$

This type of problem appears, for example, when studying the evolution of a population's distribution over time under deterministic laws, but starting from a known initial distribution.

Purely Diffusive Systems

The opposite case occurs when $a(x) = 0$, and the system is governed only by noise:

$$dx = b(x)dW$$

In population dynamics, this is useful in cases where the baseline growth rate is zero ($r = 0$) or constant. In this situation, the Fokker-Planck equation becomes a pure diffusion equation:

$$\partial_t P = \partial_x^2 \left(\frac{b(x)^2}{2} P(x,t) \right)$$

This is, in fact, a generalization of the process for the Wiener process, which I recall is $\dot{w} = \xi(t)$ (i.e., $a = 0, b = 1$). Indeed, using the Fokker-Planck equation, we can find the expression for the PDF of the Wiener process, which will be:

$$\partial_t P = \frac{1}{2} \partial_w^2 P$$

Or, analogously, for $\dot{x} = \omega\xi(t)$, which corresponds to the overdamped Brownian motion:

$$\partial_t P = \frac{\omega^2}{2} \partial_x^2 P$$

The Probability Current

The Fokker-Planck equation can also be viewed in a different way. Given the formula:

$$\partial_t P = \partial_x \left[-a(x)P(x,t) - \partial_x \left(\frac{b(x)^2}{2} P(x,t) \right) \right]$$

If we pull the outer ∂_x out and name the term inside the square brackets J , we get:

$$\partial_t P + \partial_x J = 0$$

which is the continuity (or flux) equation; in the multi-dimensional case, it is typically written as:

$$\partial_t \eta + \nabla \cdot J = 0$$

This makes intuitive sense: if we think of probability as a "substance" (like a population density η), the change in its distribution is nothing more than a redistribution, so the total amount does not change. The 1D equation found above is easily reduced from the multi-D form of the flux equation. If we take Brownian motion as an example, then:

$$J = -K \nabla P$$

For this reason, J in these cases is also known as the **current**.

Kac's Lemma

Kac's Lemma establishes a fundamental relationship between stationary probabilities and expected return times in ergodic stochastic processes. Consider a stationary Markov process with stationary distribution π and let A be a measurable subset of the state space. Define T_A as the first return time to A when starting from a point in A . Kac's Lemma states:

$$\mathbb{E}[T_A | X_0 \in A] = \frac{1}{\pi(A)}$$

where $\pi(A) = \int_A \pi(x)dx$ is the stationary measure of set A .

This result provides a direct connection between the equilibrium properties of the system (captured by π) and its dynamical properties (captured by return times). States or regions with low stationary probability have correspondingly long expected return times.

The lemma follows from the ergodic theorem: for an ergodic process, the long-run proportion of time spent in any set equals its stationary probability. Since visits to A occur approximately every $\mathbb{E}[T_A]$ time units, we have $\pi(A) \approx 1/\mathbb{E}[T_A]$.

This relationship has important applications in population dynamics, where rare population states correspond to long recovery times, and in statistical mechanics, where it connects equilibrium distributions to relaxation timescales near phase transitions.

3.5 The Ornstein-Uhlenbeck Process

We now consider one of the most important stochastic processes in science and engineering: the **Ornstein-Uhlenbeck (OU) process**. It appears in numerous contexts, from describing the velocity of a particle in a fluid (the original Langevin model) to modeling the voltage in a noisy electronic circuit. In essence, it is the canonical model for any noisy, linear system that tends to return to an equilibrium state.

The general form of the OU process is given by the SDE:

$$dz = -\gamma z dt + \omega dW$$

where $\gamma > 0$ is the *drift* or *relaxation rate*, which constantly pulls the process back towards its mean (in this case, zero), and ω is the noise intensity. This linear restoring force is what distinguishes the OU process from the free-wandering Wiener process.

Solving the Ornstein-Uhlenbeck SDE

To solve this linear SDE, we can use a technique analogous to the integrating factor method for ODEs. We define a new variable $Q(t) = e^{\gamma t} z(t)$. Our goal is to find an SDE for $Q(t)$. Using Ito's product rule (or Ito's lemma on $\psi(z, t) = e^{\gamma t} z$), we find:

$$\begin{aligned} dQ &= (d(e^{\gamma t}))z + e^{\gamma t}(dz) \\ &= (\gamma e^{\gamma t} dt)z + e^{\gamma t}(-\gamma z dt + \omega dW(\tau)) \\ &= \gamma e^{\gamma t} z dt - \gamma e^{\gamma t} z dt + \omega e^{\gamma t} dW(\tau) \\ &= \omega e^{\gamma t} dW(\tau) \end{aligned}$$

The drift terms have canceled perfectly, leaving a very simple SDE for $Q(t)$. We can now integrate it from 0 to t :

$$Q(t) - Q(0) = \omega \int_0^t e^{\gamma \tau} dW(\tau) \implies Q(t) = Q(0) + \omega \int_0^t e^{\gamma \tau} dW(\tau)$$

Substituting back $z(t) = e^{-\gamma t} Q(t)$ and $z(0) = Q(0)$, we arrive at the solution for the Ornstein-Uhlenbeck process:

$$z(t) = z(0)e^{-\gamma t} + \omega \int_0^t e^{\gamma(\tau-t)} dW(\tau)$$

The solution consists of two parts: a deterministic decay of the initial condition, $z(0)e^{-\gamma t}$, and a stochastic integral that represents the accumulated effect of the noise, with past noise contributions being exponentially "forgotten" over time due to the $e^{\gamma(\tau-t)}$ term.

Statistical Properties

From the solution, we can derive the key statistical properties of the OU process.

- **Mean**

Since the stochastic integral $\int e^{\gamma(\tau-t)} dW(\tau)$ has zero mean, the mean of the process is simply the deterministic part:

$$\boxed{\langle z(t) \rangle = \langle z(0) \rangle e^{-\gamma t}}$$

Assuming a deterministic starting point $z(0)$, the mean exponentially decays to zero, which is the equilibrium point of the system.

- **Variance**

The variance calculation is more involved but highly instructive. Starting from our solution:

$$z(t) = z(0)e^{-\gamma t} + \omega \int_0^t e^{\gamma(\tau-t)} dW(\tau)$$

we can calculate the variance by expanding $\text{Var}[z(t)] = \langle z(t)^2 \rangle - \langle z(t) \rangle^2$. Since $\langle z(t) \rangle = z(0)e^{-\gamma t}$ (assuming deterministic initial condition), we have:

$$\text{Var}[z(t)] = \left\langle \left(z(0)e^{-\gamma t} + \omega \int_0^t e^{\gamma(\tau-t)} dW(\tau) \right)^2 \right\rangle - z(0)^2 e^{-2\gamma t}$$

Expanding the square:

$$\text{Var}[z(t)] = \left\langle z(0)^2 e^{-2\gamma t} + 2z(0)e^{-\gamma t} \omega \int_0^t e^{\gamma(\tau-t)} dW(\tau) + \omega^2 \left(\int_0^t e^{\gamma(\tau-t)} dW(\tau) \right)^2 \right\rangle - z(0)^2 e^{-2\gamma t}$$

The cross term vanishes since $\langle dW(\tau) \rangle = 0$, leaving:

$$\text{Var}[z(t)] = \omega^2 \left\langle \left(\int_0^t e^{\gamma(\tau-t)} dW(\tau) \right)^2 \right\rangle$$

Applying the mean, we get:

$$\text{Var}[z(t)] = \omega^2 \left\langle \left(\int_0^t e^{\gamma(\tau-t)} dW(\tau) \right)^2 \right\rangle = \omega^2 \int_0^t \int_0^t e^{\gamma(\tau-t)} e^{\gamma(s-t)} \langle dW(\tau) dW(s) \rangle$$

Using the property that $\langle dW(\tau) dW(s) \rangle = \delta(\tau-s) d\tau ds$, the double integral becomes:

$$\text{Var}[z(t)] = \omega^2 \int_0^t \int_0^t e^{\gamma(\tau+s-2t)} \delta(s-\tau) ds d\tau = \omega^2 \int_0^t e^{2\gamma(\tau-t)} d\tau$$

Evaluating this integral:

$$\text{Var}[z(t)] = \omega^2 e^{-2\gamma t} \int_0^t e^{2\gamma\tau} d\tau = \omega^2 e^{-2\gamma t} \left[\frac{e^{2\gamma\tau}}{2\gamma} \right]_0^t = \frac{\omega^2}{2\gamma} (1 - e^{-2\gamma t})$$

The total variance becomes:

$$\text{Var}[z(t)] = \text{Var}[z(0)] e^{-2\gamma t} + \frac{\omega^2}{2\gamma} (1 - e^{-2\gamma t})$$

As $t \rightarrow \infty$, the first term vanishes, and the variance approaches a constant stationary value:

$$\lim_{t \rightarrow \infty} \text{Var}[z(t)] = \frac{\omega^2}{2\gamma}$$

This is a crucial difference from the Wiener process, whose variance grows linearly and without bound. The restoring force in the OU process balances the diffusive effect of the noise, leading to a finite stationary variance.

- **Autocovariance**

The autocovariance function for times t and q is defined as:

$$C(t, q) = \langle (z(t) - \langle z(t) \rangle)(z(q) - \langle z(q) \rangle) \rangle$$

For simplicity, let's assume the process starts at $z(0) = 0$. In this case, the mean is $\langle z(t) \rangle = 0$ for all t , so the autocovariance is simply the autocorrelation: $C(t, q) = \langle z(t)z(q) \rangle$.

Using the solution for $z(t)$ with $z(0) = 0$:

$$z(t) = \omega \int_0^t e^{\gamma(s-t)} dW(s)$$

The autocorrelation is:

$$\begin{aligned} \langle z(t)z(q) \rangle &= \left\langle \left(\omega \int_0^t e^{\gamma(s-t)} dW(s) \right) \left(\omega \int_0^q e^{\gamma(\tau-q)} dW(\tau) \right) \right\rangle \\ &= \omega^2 e^{-\gamma(t+q)} \left\langle \int_0^t e^{\gamma s} dW(s) \int_0^q e^{\gamma \tau} dW(\tau) \right\rangle \\ &= \omega^2 e^{-\gamma(t+q)} \int_0^t \int_0^q e^{\gamma(s+\tau)} \langle dW(s)dW(\tau) \rangle \end{aligned}$$

Using the property $\langle dW(s)dW(\tau) \rangle = \delta(s-\tau)dsd\tau$, the double integral collapses into a single integral. We must evaluate the integral over the domain $[0, t] \times [0, q]$.

$$\langle z(t)z(q) \rangle = \omega^2 e^{-\gamma(t+q)} \int_0^t \int_0^q e^{\gamma(\tau+s)} \delta(\tau-s) d\tau ds$$

For the evaluation of this integral, we need to consider two cases:

- **Case 1:** $t \leq q$

When $t \leq q$, the domain of integration is restricted by the smaller upper limit. The delta function $\delta(\tau-s)$ forces $\tau=s$, so we integrate over the overlap region $[0, t]$:

$$\langle z(t)z(q) \rangle = \omega^2 e^{-\gamma(t+q)} \int_0^t e^{2\gamma s} ds$$

Evaluating the integral:

$$\langle z(t)z(q) \rangle = \omega^2 e^{-\gamma(t+q)} \left[\frac{e^{2\gamma s}}{2\gamma} \right]_0^t = \frac{\omega^2}{2\gamma} e^{-\gamma(t+q)} (e^{2\gamma t} - 1) = \frac{\omega^2}{2\gamma} (e^{-\gamma(q-t)} - e^{-\gamma(q+t)})$$

- **Case 2:** $t > q$

When $t > q$, we integrate over $[0, q]$:

$$\langle z(t)z(q) \rangle = \omega^2 e^{-\gamma(t+q)} \int_0^q e^{2\gamma s} ds = \frac{\omega^2}{2\gamma} (e^{-\gamma(t-q)} - e^{-\gamma(t+q)})$$

Combining both cases, we can write the autocorrelation function as:

$$\boxed{\langle z(t)z(q) \rangle = \frac{\omega^2}{2\gamma} (e^{-\gamma|t-q|} - e^{-\gamma(t+q)})}$$

For a stationary process (when $t, q \rightarrow \infty$), the second term vanishes, and we obtain the stationary autocorrelation function:

$$R_z(\tau) = \lim_{t \rightarrow \infty} \langle z(t)z(t+\tau) \rangle = \frac{\omega^2}{2\gamma} e^{-\gamma|\tau|}$$

where $\tau = |t-q|$ is the time lag. This autocorrelation decays exponentially with a characteristic time $1/\gamma$. In the limit of rapid relaxation ($\gamma \rightarrow \infty$), the autocorrelation function becomes sharply peaked and tends toward a Dirac delta function.

We introduce now a scaling relation between the noise intensity and the relaxation rate, by setting:

$$\omega = c \gamma$$

with c a constant. Under this scaling, the SDE for z becomes

$$\dot{z} = -\gamma z + c \gamma \xi(t).$$

The autocorrelation function for z then reads

$$R_z(h) = \frac{c^2}{2} \gamma e^{-\gamma|h|}, \quad R_z(0) = \frac{c^2}{2} \gamma.$$

Its total area is given by

$$\int_{-\infty}^{+\infty} R_z(h) dh = \int_{-\infty}^{+\infty} \frac{c^2}{2} \gamma e^{-\gamma|h|} dh = \frac{c^2}{2} \gamma \cdot \frac{2}{\gamma} = c^2.$$

For small characteristic times ($\tau = 1/\gamma$ small), the autocorrelation function $R_z(h)$ approximates a white noise process:

$$\lim_{\gamma \rightarrow \infty} R_z(h; \gamma) = c^2 \delta(h).$$

This derivation shows how, by choosing the scaling $\omega = c \gamma$, the fluctuations in the linearized dynamics around a stable equilibrium effectively become white noise in the fast relaxation limit.

3.5.1 The OU Process as a Low-Pass Filter

One of the most important interpretations of the Ornstein-Uhlenbeck process is as a **low-pass filter**. It takes a "noisy" input signal (white noise) and produces a smoother, more physically realistic output by attenuating high-frequency fluctuations. To understand this, we can analyze the process in the frequency domain using the Fourier transform.

The Fourier Transform and Linear Systems

Let $f(t)$ be a well-behaved function. Its Fourier transform, $\hat{f}(\omega)$, is defined as:

$$\mathcal{F}[f(t)] = \hat{f}(\omega) = \int_{-\infty}^{+\infty} f(t) e^{-i\omega t} dt$$

A key property of the Fourier transform, which can be shown using integration by parts, is how it acts on derivatives:

$$\mathcal{F}[f'(t)] = i\omega \mathcal{F}[f(t)] = i\omega \hat{f}(\omega)$$

This property turns differentiation in the time domain into multiplication in the frequency domain, which is extremely useful for solving linear differential equations. For example, consider a general first-order linear system with input $y(t)$ and output $f(t)$:

$$f'(t) + \alpha f(t) = y(t)$$

Taking the Fourier transform of the entire equation gives:

$$i\omega \hat{f}(\omega) + \alpha \hat{f}(\omega) = \hat{y}(\omega) \implies \hat{f}(\omega) = \frac{\hat{y}(\omega)}{i\omega + \alpha}$$

The term $H(\omega) = 1/(i\omega + \alpha)$ is called the **transfer function** of the system. The power of the output signal is related to the power of the input signal by the squared magnitude of the transfer function:

$$|\hat{f}(\omega)|^2 = \frac{|\hat{y}(\omega)|^2}{\omega^2 + \alpha^2}$$

Power Spectrum of the OU Process

We can apply this framework to the OU process, which is described by the linear SDE:

$$\frac{dz}{dt} + \gamma z = \kappa \xi(t)$$

Taking the Fourier transform, we get:

$$(i\omega + \gamma)\hat{z}(\omega) = \kappa \mathcal{F}[\xi(t)] \implies \hat{z}(\omega) = \frac{\kappa}{i\omega + \gamma} \mathcal{F}[\xi(t)]$$

To analyze the filtering effect, we compare the **power spectral density** of the input signal, $\xi(t)$, with that of the output signal, $z(t)$. The power spectrum is the Fourier transform of the autocorrelation function.

For the input signal (white noise), the autocorrelation is $R_\xi(h) = \delta(h)$ (assuming unit variance for simplicity). Its power spectrum is therefore constant:

$$\mathcal{F}[R_\xi(h)] = \mathcal{F}[\delta(h)] = \int_{-\infty}^{+\infty} \delta(h) e^{-i\omega h} dh = 1$$

This flat spectrum is why it's called "white" noise: it contains equal power at all frequencies ω . For the output signal, $z(t)$, we previously found its stationary autocorrelation function to be:

$$R_z(h) = \frac{\kappa^2}{2\gamma} e^{-\gamma|h|}$$

The power spectrum of the output is the Fourier transform of this function:

$$\mathcal{F}[R_z(h)] = \frac{\kappa^2}{2\gamma} \int_{-\infty}^{+\infty} e^{-\gamma|h|} e^{-i\omega h} dh$$

We can split the integral over positive and negative values of h :

$$= \int_{-\infty}^0 e^{(\gamma-i\omega)h} dh + \int_0^{+\infty} e^{-(\gamma+i\omega)h} dh = \frac{1}{\gamma - i\omega} + \frac{1}{\gamma + i\omega} = \frac{2\gamma}{\gamma^2 + \omega^2}$$

Substituting this back, we get the output power spectrum:

$$\mathcal{F}[R_z(h)] = \frac{\kappa^2}{2\gamma} \left(\frac{2\gamma}{\gamma^2 + \omega^2} \right) = \frac{\kappa^2}{\gamma^2 + \omega^2}$$

This function, known as a **Lorentzian**, is peaked at $\omega = 0$ and decays as $1/\omega^2$ for high frequencies. The OU process acts as a low-pass filter: it suppresses high-frequency components of the input white noise while allowing low-frequency components to pass through, resulting in a smoother, more physically plausible random process.

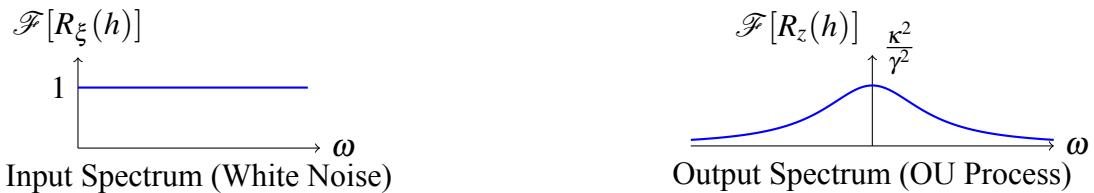


Figure 3.1: Power spectra: input white noise (flat, left) and output OU process (Lorentzian, right).

3.5.2 Fokker-Planck equation for Ornstein-Uhlenbeck

Given the Ito equation for the OU process:

$$dz = -\gamma z dt + \omega dW$$

we can find the Fokker-Planck equation. Recalling the general form of the equation for $P(z,t)$:

$$\partial_t P = -\partial_z [a(z)P(z,t)] + \frac{1}{2}\partial_z^2 [b(z)^2 P(z,t)]$$

In our case, with $a(z) = -\gamma z$ and $b(z) = \omega$, it becomes:

$$\partial_t P = \partial_z (\gamma z P) + \frac{\omega^2}{2} \partial_z^2 P$$

This equation is not simple to solve, even with the boundary conditions for a probability density:

$$P(\pm\infty, t) = 0 \quad \text{and} \quad \int_{\mathbb{R}} P(z, t) dz = 1$$

These are additional conditions compared to standard PDEs and are usually of great help.

The most tractable approach is to find the stationary solution $P_s(z)$, which is the long-time limit of $P(z,t)$:

$$P_s(z) = \lim_{t \rightarrow \infty} P(z, t)$$

In the stationary state, $\partial_t P = 0$, so the equation becomes an ODE:

$$0 = \frac{d}{dz} (\gamma z P_s) + \frac{\omega^2}{2} \frac{d^2 P_s}{dz^2}$$

which we can rearrange to:

$$\frac{d^2 P_s}{dz^2} = -\frac{2\gamma}{\omega^2} \frac{d}{dz} (z P_s)$$

Integrating once with respect to z , we obtain:

$$\frac{dP_s}{dz} = C - \frac{2\gamma}{\omega^2} z P_s(z)$$

The boundary conditions require that $P_s(z) \rightarrow 0$ and consequently $P'_s(z) \rightarrow 0$ as $z \rightarrow \pm\infty$. For this to hold, the integration constant C must be zero. This leaves us with a first-order linear ODE:

$$\frac{dP_s}{dz} = -\frac{2\gamma}{\omega^2} z P_s(z)$$

The solution is a Gaussian function, which is consistent with the properties of the Ornstein-Uhlenbeck process:

$$P_s(z) = A e^{-\frac{\gamma}{\omega^2} z^2}$$

This is nothing other than a Gaussian with $\mu = 0$ and $\sigma^2 = \omega^2/2\gamma$, with A a constant required for normalization, namely

$$A = \frac{1}{\sqrt{2\pi\sigma^2}}$$

If we now consider the unperturbed system $\dot{z} = -\gamma z$, we see that the solution tends to 0 for long times and, since it is deterministic, we obviously have

$$P_s^{\text{DET}}(z) = \delta(z)$$

which is consistent with the perturbed case. In fact, the only equilibrium point of the deterministic case is 0, which turns out to be the center of the Gaussian we found. Therefore, for long times, in the perturbed case, the solution will oscillate around 0.

3.6 The Ginzburg-Landau Model

We now move from the monostable Ornstein-Uhlenbeck process to a more complex and interesting class of systems: bistable systems. These are systems that possess two stable equilibrium states. The simplest and most iconic model for such a system is the **Ginzburg-Landau equation**.

The Deterministic Model

The deterministic Ginzburg-Landau model describes the dynamics of a particle in a "double-well" potential. The equation of motion is given by:

$$\dot{x} = x - x^3 = x(1-x)(1+x)$$

This system has three equilibrium points, found by setting $\dot{x} = 0$:

$$x_L = -1, \quad x_C = 0, \quad x_R = +1$$

A simple stability analysis reveals the nature of these points:

- $x_L = -1$ and $x_R = +1$ are **locally asymptotically stable**. If the system starts near one of these points, it will converge to it.
- $x_C = 0$ is an **unstable** equilibrium. If the system starts at this point, any infinitesimal perturbation will cause it to move away towards either x_L or x_R .

The force $F(x) = x - x^3$ can be derived from a potential $U(x)$ where $F = -dU/dx$:

$$U(x) = -\frac{x^2}{2} + \frac{x^4}{4}$$

This potential has a characteristic "double-well" shape, where the stable equilibria x_L and x_R correspond to the minima (the "wells") of the potential, and the unstable equilibrium $x_C = 0$ corresponds to the local maximum (the "barrier") between them.

If the initial condition $x(0)$ is drawn from a probability distribution $\theta(x)$, the long-term behavior of the system is deterministic. Any trajectory with $x(0) > 0$ will end up at $x_R = +1$, and any with $x(0) < 0$ will end up at $x_L = -1$. The stationary PDF will therefore consist of two Dirac delta peaks at the stable equilibria, with weights determined by the initial probability of being on either side of the unstable point:

$$\lim_{t \rightarrow \infty} \rho(x, t) = A_R \delta(x - 1) + A_L \delta(x + 1)$$

where $A_R = \Pr(x_0 > 0)$ and $A_L = \Pr(x_0 < 0)$.

Stochastic Perturbations of the Model

Now, let's add a stochastic white noise force to the system:

$$\dot{x} = x - x^3 + \omega \xi(t) \quad \text{or} \quad dx = (x - x^3)dt + \omega dW_t$$

This is a case of additive noise perturbing a conservative force. As we derived previously, the stationary probability distribution $P_s(x)$ for such a system is given by the Boltzmann distribution:

$$P_s(x) = C \exp\left(-\frac{2}{\omega^2} U(x)\right)$$

Substituting the Ginzburg-Landau potential $U(x) = -x^2/2 + x^4/4$, we get:

$$P_s(x) = C \exp\left(-\frac{2}{\omega^2}\left(-\frac{x^2}{2} + \frac{x^4}{4}\right)\right) = C \exp\left(\frac{x^2}{\omega^2} - \frac{x^4}{2\omega^2}\right)$$

where C is a normalization constant. The shape of this distribution is fundamentally controlled by the noise intensity ω .

The relationship between the potential $U(x)$ and the stationary PDF $P_s(x)$ is clear:

- The stable equilibria of the deterministic system (minima of $U(x)$) become the **modes** (peaks) of the stationary PDF.
- The unstable equilibrium (maximum of $U(x)$) becomes the **antimode** (trough) of the stationary PDF.

The noise intensity ω determines how closely the system is confined to these potential wells:

- **Weak Noise ($\omega \ll 1$)**: The distribution $P_s(x)$ is sharply bimodal, with two distinct peaks centered around $x = -1$ and $x = +1$. The probability of finding the particle near the unstable point $x = 0$ is extremely low. The system will exhibit small stochastic fluctuations around one of the two stable states for very long periods.
- **Strong Noise ($\omega \gg 1$)**: The noise provides enough energy for the particle to easily cross the potential barrier between the two wells. The two peaks in the PDF broaden and merge, and the distribution becomes unimodal, centered at $x = 0$. For very large x , the x^4 term still dominates, so the probability $P_s(x)$ must decay to zero.

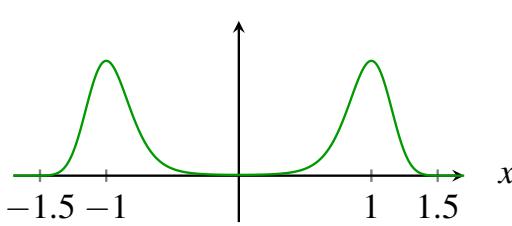


Figure 3.2: PDF (non normalized) for small ω

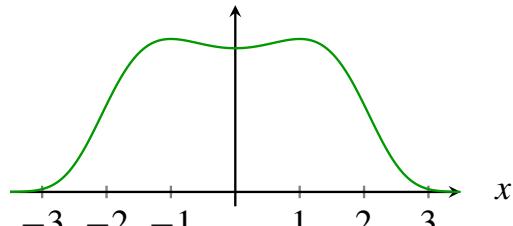


Figure 3.3: PDF (non normalized) for large ω

This "inheritance" of stability properties from the deterministic potential to the modes of the stationary PDF is a key feature of systems with additive noise. As we will see, this direct correspondence breaks down in the presence of multiplicative noise.

⌚ Observation: Stochastic Multistability

It is possible to say that after a very long period of time, also the first case we could see that the system moves from one peak to the other. This is due to the fact that the system is not stable and we have a non-zero probability of moving from one peak to the other.

3.7 Multiplicative Noise

While additive noise provides a good model for systems perturbed by external forces, many real-world systems feature fluctuations whose intensity depends on the state of the system itself. This leads to the concept of **multiplicative noise**, which can induce much richer and more complex behaviors than its additive counterpart.

General SDE and Stationary Fokker-Planck Equation

Consider a system governed by a parameter p . The simplest case is when it is possible to write the field in the form

$$\dot{x} = b(x) + p\hat{g}(x)$$

that is, with linear dependence on the parameter considered. The parameter p is usually influenced by the external world, so in some cases it is possible for it to become stochastic. We can then describe a stochastic p as

$$p \rightarrow p + \alpha\xi(t)$$

Substituting, we obtain

$$\dot{x} = b(x) + p\hat{g}(x) + \hat{g}(x)\alpha\xi(t)$$

We can then group the stochastic terms, obtaining

$$\dot{x} = f(x) + g(x)\xi(t)$$

The Fokker-Planck equation for the general case is:

$$\frac{\partial\rho}{\partial t} = -\frac{\partial}{\partial x}[f(x)\rho] + \frac{\partial^2}{\partial x^2}\left[\frac{g(x)^2}{2}\rho\right]$$

So far, we have only considered cases where $g(x)$ is constant. Here, however, we are dealing with the most general version of what we have seen so far. The boundary conditions remain the same as before, and as usual, we analyze the stationary case. We have:

$$\frac{d}{dx}\left[\frac{g(x)^2}{2}P_s(x)\right] = f(x)P_s(x)$$

This is a first-order ODE that can be solved for $P_s(x)$. Let's define an auxiliary function $Q(x) = \frac{g(x)^2}{2}P_s(x)$. The equation becomes:

$$\frac{dQ}{dx} = f(x)\frac{2}{g(x)^2}Q(x)$$

This is a standard linear ODE of the form $Q'(x) = a(x)Q(x)$, whose solution is $Q(x) = Ce^{A(x)}$, where $A(x) = \int a(s)ds$. Applying this, we find:

$$Q(x) = C \exp\left(\int \frac{2f(s)}{g(s)^2}ds\right)$$

Substituting back $P_s(x) = \frac{2}{g(x)^2}Q(x)$, we arrive at the general solution for the stationary PDF:

$$P_s(x) = \frac{C'}{g(x)^2} \exp\left(\int^x \frac{2f(z)}{g(z)^2}dz\right)$$

where C' is a normalization constant. Let's remark that, being a probability density, the integral of $P_s(x)$ over the entire state space must be equal to one.

We can distinguish two cases:

- If the integral exists (is finite), then everything is fine and we can calculate C to normalize $P_s(x)$.
- If instead the integral is infinite, we can only find a formal solution that has no physical meaning: in this case it is necessary to study the steady state directly starting from the SDE.

3.7.1 Noise-Induced Transitions

In the case of additive noise ($g(x) = \text{const}$), the extrema of the stationary PDF $P_s(x)$ coincide with the equilibria of the deterministic system (where $f(x) = 0$). With multiplicative noise, this simple correspondence breaks down.

Recall the following identity, which will be useful for manipulating the stationary distribution:

$$\frac{1}{g^2(x)} = e^{-\ln g^2(x)} = e^{-2 \ln g(x)}$$

Using this, we can rewrite the stationary probability density function in a more transparent exponential form:

$$P_s(x) = 2C \exp \left\{ -2 \log(g(x)) + \int_{\alpha}^x \frac{2f(s)}{g^2(s)} ds \right\}$$

This highlights how both the noise amplitude $g(x)$ and the drift $f(x)$ contribute to shaping the stationary distribution.

To investigate the location of the extrema of $P_s(x)$, we differentiate with respect to x . Applying the chain rule to the exponent, we obtain:

$$P'_s(x) = 2C \exp \left\{ -2 \frac{g'(x)}{g(x)} + \frac{2f(x)}{g^2(x)} \right\}$$

This expression will allow us to determine where the stationary distribution attains its maxima and minima, as these correspond to the points where $P'_s(x) = 0$.

To find the extrema of $P_s(x)$, we must find where its derivative is zero, $P'_s(x) = 0$. After some algebraic manipulation of the general solution, this condition simplifies to:

$$f(x) \geq g(x)g'(x)$$

This time, unfortunately, we no longer have the close connection we had before; everything now depends on α and on the structure of $g(x)$. In fact, I can rewrite the relation as

$$f(x) \geq \alpha^2 g'(x)g(x)$$

This remarkable result shows that the modes of the stationary distribution are no longer located at the deterministic equilibria ($f(x) = 0$). Instead, their locations are shifted by a term that depends on the noise function $g(x)$ and its derivative.

This phenomenon, discovered by Horsthemke and Lefever in the 1970s, is called a **Noise-Induced Transition (NIT)**. It means that by simply varying a parameter that controls the noise intensity, new extrema can appear in the PDF that were completely absent in the unperturbed deterministic system.

Before the discovery of NITs, it was commonly believed that noise could only act in two ways:

- In unimodal systems, it creates a "cloud" of fluctuations around the single deterministic equilibrium.
 - In multimodal systems, it can cause "jumps" between the existing deterministic equilibria.
- NITs showed that noise can play a much more constructive role, inducing oscillations and creating entirely new states that are not present in the unperturbed deterministic system.

③ Example: Noise-Induced Bifurcation in Gene Expression

Consider a simple model for protein production, where a protein X acts as its own transcription factor. The production rate is $Q(x)$ and the degradation rate is dx .

$$\dot{x} = Q(x) - dx$$

A common form for $Q(x)$ that models self-regulation is:

$$Q(x) = R + K \frac{x^2}{K_d + x^2}$$

where R is a basal production rate and the second term is a Hill function describing feedback. The deterministic system is:

$$\dot{x} = \underbrace{R + K \frac{x^2}{K_d + x^2}}_{f(x)} - dx$$

For small degradation rates d , this system has a single, globally stable equilibrium point. Now, suppose the degradation rate d is subject to stochastic fluctuations: $d \rightarrow d + \alpha \xi(t)$. The SDE becomes:

$$dx = \left(R + K \frac{x^2}{K_d + x^2} - dx \right) dt - \alpha x dW_t$$

Here, the noise is multiplicative, with $g(x) = -\alpha x$, so $g'(x) = -\alpha$. The condition for the extrema of the stationary PDF, $f(x) \geq g(x)g'(x)$, is:

$$R + K \frac{x^2}{K_d + x^2} - dx \geq (-\alpha x)(-\alpha) = \alpha^2 x$$

Rearranging gives:

$$R + K \frac{x^2}{K_d + x^2} \geq (d + \alpha^2)x$$

Depending on the value of the noise parameter α , this algebraic equation can have one or three solutions. This means that by changing the noise intensity, the stationary PDF can transition from being unimodal (with one peak) to bimodal (with two peaks), even though the underlying deterministic system only had one equilibrium point. This is a classic example of a noise-induced transition: the noise itself creates new, stable macroscopic states for the system.

3.8 Perturbation of the Logistic Growth Model

Let's revisit the logistic growth model, which serves as another prime example of how multiplicative noise can fundamentally alter a system's behavior. We start with the deterministic logistic equation:

$$\dot{z} = r_1 z - r_2 z^2$$

Through a linear change of variables $x = r_2 z$, we can adimensionalize the equation to a simpler form:

$$\dot{x} = rx - x^2$$

Now, we introduce stochasticity by perturbing the intrinsic growth rate, $r \rightarrow r + \omega \xi(t)$. This leads to the corresponding Itô equation with multiplicative noise:

$$dx = (rx - x^2)dt + \omega x dW_t$$

This defines our drift and diffusion coefficients as $f(x) = rx - x^2$ and $g(x) = \omega x$.

The stationary Fokker-Planck equation for this system is:

$$\frac{d}{dx} \left[\frac{g(x)^2}{2} P_s(x) \right] = f(x) P_s(x)$$

Substituting our specific $f(x)$ and $g(x)$ gives:

$$\frac{d}{dx} \left[\frac{\omega^2 x^2}{2} P_s(x) \right] = (rx - x^2) P_s(x)$$

To solve this differential equation, we can use the substitution method. Let $Q = x^2 P_s$, then:

$$Q' = \frac{2r - x}{\omega^2 x} Q$$

This can be rewritten as:

$$\frac{dQ}{Q} = \frac{2r - x}{\omega^2 x} dx = \left[\frac{2r}{\omega^2 x} - \frac{1}{\omega^2} \right] dx$$

Integrating both sides:

$$\ln Q = \frac{2r}{\omega^2} \ln x - \frac{x}{\omega^2} + C$$

Taking the exponential:

$$Q(x) = A \exp \left[\frac{2r}{\omega^2} \ln x - \frac{x}{\omega^2} \right] = Ax^{\frac{2r}{\omega^2}} e^{-\frac{x}{\omega^2}}$$

Since $Q = x^2 P_s$, we obtain:

$$P_s(x) = \frac{Q(x)}{x^2} = Ax^{\frac{2r}{\omega^2} - 2} e^{-\frac{x}{\omega^2}}$$

where A is a normalization constant.

A crucial step is to check whether this candidate solution is a valid, normalizable probability distribution. The integral $\int_0^\infty P_s(x) dx$ must be finite.

$$\int_0^\infty x^{\frac{2r}{\omega^2} - 2} e^{-\frac{x}{\omega^2}} dx$$

This is a form of the Gamma function integral, $\int_0^\infty t^{k-1} e^{-t} dt$. For this integral to converge at the lower bound ($x \rightarrow 0$), the exponent of x must be greater than -1.

$$\frac{2r}{\omega^2} - 2 > -1 \implies \frac{2r}{\omega^2} > 1 \implies r > \frac{\omega^2}{2}$$

This is valid only if $r > \omega^2/2$ as we have already seen in the past. If this does not hold, we cannot integrate but what does this actually mean? To understand this, let's return to the SDE:

$$dx = x(r-x)dt + \omega x dW$$

Applying now the transformation $y = \psi(x) = \log x$, then $\psi'(x) = x^{-1}$ and $\psi''(x) = -x^{-2}$. Therefore:

$$dy = \left[\frac{1}{x}x(r-x) - \frac{1}{x^2} \frac{\omega^2 x^2}{2} \right] dt + \omega dW$$

Therefore, substituting in y :

$$dy = \left[r - \frac{\omega^2}{2} - e^y \right] dt + \omega dW$$

Integrating then:

$$y(t) = y_0 + \left(r - \frac{\omega^2}{2} \right) t - \int_0^t e^{y(s)} ds + \omega [W(t) - W(0)]$$

Notice that the integral term, $-\int_0^t e^{y(s)} ds$, is always negative (since $e^{y(s)} > 0$ for all s).

- **Case 1: $r < \omega^2/2$**

If $r < \omega^2/2$, the drift term $\left(r - \frac{\omega^2}{2}\right)t$ is negative and dominates for large t , while the stochastic term becomes negligible in comparison. As a result, $y(t) \rightarrow -\infty$ as $t \rightarrow \infty$, which means $x(t) = e^{y(t)} \rightarrow 0^+$. In this regime, the probability distribution $P(x,t)$ collapses to a delta function at zero:

$$\lim_{t \rightarrow \infty} P(x,t) = \delta(x)$$

This shows that, in this case, the Fokker-Planck equation predicts a stationary solution, but in reality, all probability accumulates at extinction ($x = 0$). Thus, the Fokker-Planck approach alone can be misleading for this parameter regime.

- **Case 2: $r > \omega^2/2$**

On the other hand, when $r > \omega^2/2$, the stationary solution is normalizable and given by:

$$P_s(x) = Ax^{\frac{2r}{\omega^2}-2} e^{-\frac{2}{\omega^2}x}$$

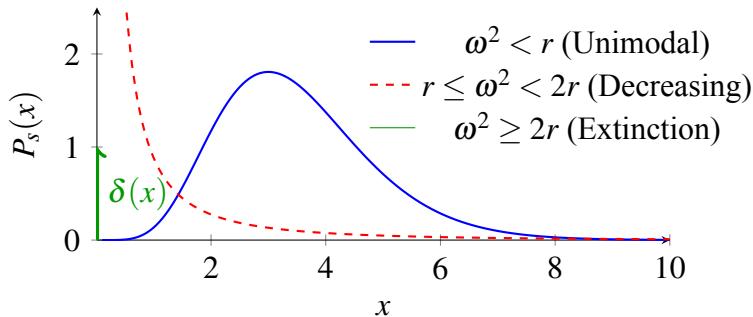


Figure 3.4: Noise-Induced Transition in the logistic model.

Let us also consider two special cases:

- If $\omega^2 = 2r$ (i.e., $2r/\omega^2 = 1$), the stationary distribution simplifies to:

$$P_s(x) = Ax^{-1}e^{-\frac{x}{r}}$$

- If there is no noise, $\omega^2 = 0$, the stationary distribution becomes:

$$P_s(x) = Ae^{-\frac{x}{r}}$$

This change has a significant effect on the shape of the stationary distribution. Specifically, when the parameters satisfy

$$\frac{2r}{\omega^2} < 2 \quad \text{or equivalently} \quad r < \frac{\omega^2}{2}$$

the distribution $P_s(x)$ is no longer unimodal (with a peak at some $x > 0$), but instead becomes a monotonically decreasing function of x . In this regime, the probability density is highest near $x = 0$ and decreases for larger x , so the system is most likely to be found close to $x = 0$.

Furthermore, it can be shown that the stationary distribution $P_s(x)$ is normalizable (i.e., its total probability integrates to 1) only if $\omega^2 < 2r$. This is consistent with the more general result discussed in the previous section.

3.9 From Discrete Jumps to Continuous Diffusion

In the previous sections, we modeled stochastic systems using Stochastic Differential Equations (SDEs), which describe the infinitesimal evolution of processes with continuous paths, like the Ornstein-Uhlenbeck process. However, many systems in physics, chemistry, and biology are more naturally described by discrete "jumps" between states. Examples include a molecule switching between chemical species, an ion channel opening or closing, or a population changing by single-birth or death events.

These phenomena are governed by **jump processes**, where the system remains in a given state for a random period before instantaneously jumping to a new state. The evolution of the probability distribution for such a process is not described by a Fokker-Planck equation but by a more general framework: the **Master Equation**.

3.9.1 The Master Equation

Let's consider a process where the state variable x can take any value in a continuous state space. The core of a jump process is the **transition rate** (or jump kernel) $\Omega(y,x)$, which defines the probability per unit time of a jump occurring *from* state y *to* state x .

The probability of a system in state y jumping to any state within a small region around x in an infinitesimal time interval dt is given by:

$$\text{Prob(jump } y \rightarrow x \text{ in } dt) = \Omega(y,x)dt$$

Conversely, the probability that the system, currently in state x , does **not** jump to any other state during dt is one minus the total probability of jumping *out* of x :

$$\text{Prob(no jump from } x \text{ in } dt) = 1 - \left(\int \Omega(x,s)ds \right) dt$$

The Master Equation describes the time evolution of the probability density $P(x,t)$ by balancing the probability flows into and out of state x . The rate of change of $P(x,t)$ is the sum of all probability

flowing *in* from other states s (the gain term) minus the total probability flowing *out* to all other states s (the loss term). This balance gives the integro-differential Master Equation:

$$\frac{\partial P(x,t)}{\partial t} = \underbrace{\int P(s,t)\Omega(s,x)ds}_{\text{Gain: Jumps from } s \text{ to } x} - \underbrace{P(x,t) \int \Omega(x,s)ds}_{\text{Loss: Jumps from } x \text{ to } s}$$

3.9.2 Deriving the Diffusion Equation

A profound connection exists between discrete jump processes and continuous diffusion processes. Under the right conditions, a process consisting of many small, frequent jumps becomes indistinguishable from a continuous process described by a Fokker-Planck equation. We can demonstrate this by considering two simple jump models.

Example: The Uniform Jump Process

Consider a particle that performs random jumps where the jump destination is uniformly distributed within a small interval of length 2ϵ around its current position. If the particle is at s , it can jump to any $x \in [s - \epsilon, s + \epsilon]$. The jump kernel can be written as:

$$\Omega(s,x) = C \cdot \mathbb{I}_{[s-\epsilon, s+\epsilon]}(x)$$

where C is a constant representing the overall jump frequency and \mathbb{I} is the indicator function, therefore \mathbb{I} will be 1 in the set and 0 outside. Thus I have the Master Equation:

$$\partial_t P = -2\epsilon C P(x,y) + \int_{x-\epsilon}^{x+\epsilon} P(y,t) dy$$

Now define $y = z + x$, then

$$\int_{x-\epsilon}^{x+\epsilon} P(y,t) dy = \int_{-\epsilon}^{\epsilon} P(x+z,t) dz$$

Since ϵ is small, I can Taylor expand

$$P(x+z,t) = P(x,t) + z\partial_x P(x,t) + \frac{z^2}{2}\partial_x^2 P(x,t) + \dots$$

Therefore

$$\begin{aligned} C \int_{-\epsilon}^{\epsilon} P(x+z,t) dz &= C \left[2\epsilon + 0 \cdot \partial_x + \frac{1}{2} \frac{1}{3} 2\epsilon^3 \partial_x^2 \right] P(x,t) \\ &= C \left[2\epsilon + \frac{1}{3} \epsilon^3 \partial_x^2 \right] P(x,t) \end{aligned}$$

Putting everything together, we obtain

$$\partial_t P(x,t) = \frac{C\epsilon^3}{3} \partial_x^2 P(x,t)$$

To have $\frac{C\epsilon^3}{3} = O(1)$, we need $C \propto 1/\epsilon^3$, which is actually a huge jump rate.

Note that this is the formula for the PDF associated with a Wiener process, so a random walk process as described is described by a Wiener process.

③ Example: The Fixed-Step Random Walk

Now consider a different model where jumps are of a fixed size $\pm \varepsilon$. The transition rate is described by Dirac delta functions:

$$\Omega(y, x) = C \cdot \delta(|x - y| - \varepsilon)$$

Therefore,

$$\partial_t P = \int (P(y, t)\Omega(y, x) - P(x, t)\Omega(x, y)) dy$$

As before,

$$\partial_t P = C [P(x - \varepsilon, t) + P(x + \varepsilon, t) - 2P(x, t)]$$

and thus

$$\partial_t P = C (P(x + \varepsilon, t) - P(x, t) + P(x - \varepsilon, t) - P(x, t))$$

We can now divide above and below by ε^2 and take the incremental ratio as $\varepsilon \rightarrow 0$, obtaining

$$\partial_t P = \varepsilon^2 C \frac{P(x + \varepsilon, t) - 2P(x, t) + P(x - \varepsilon, t)}{\varepsilon^2} = C\varepsilon^2 \partial_x^2 P(x, t)$$

Even for this discrete case, we obtain again the equation for the PDF of a Wiener process. However, there is something subtle: if C is finite, the process does not diffuse in the limit $\varepsilon \rightarrow 0$. To have a finite rate of change of probability, the diffusion coefficient must be defined as

$$\lim_{\varepsilon \rightarrow 0} C\varepsilon^2 = O(1)$$

So we must have $C \propto 1/\varepsilon^2$.

The Diffusion Limit and the Wiener Process

In both examples, we found that a process of discrete jumps can be described by a diffusion equation in the limit of small jump sizes ($\varepsilon \rightarrow 0$). However, for the diffusion coefficient D to remain finite and non-zero in this limit, the jump rate C must scale appropriately:

- For uniform jumps: $D = C\varepsilon^3/3 \implies C \propto 1/\varepsilon^3$.
- For fixed-step jumps: $D = C\varepsilon^2 \implies C \propto 1/\varepsilon^2$.

In both cases, the jump rate must diverge as the jump size shrinks. This is the **diffusion limit**: a macroscopic diffusion process emerges from an infinite number of infinitesimally small, infinitely frequent jumps.

The resulting equation,

$$\frac{\partial P}{\partial t} = D \frac{\partial^2 P}{\partial x^2}$$

is precisely the Fokker-Planck equation for a scaled **Wiener process** described by the SDE $dx = \sqrt{2D}dW_t$. This provides a powerful microscopic justification for the SDE framework: continuous diffusion models can be viewed as the large-scale limit of underlying discrete jump phenomena.

3.10 The Stratonovich Integral

We started with the Itô formula, which arises from discretizing an SDE of the form $dx = a(x)dt + b(x)dW_t$ as:

$$x(t+dt) - x(t) = a(x(t))dt + b(x(t))dW_t$$

Here, the diffusion term $b(x)$ is evaluated at the beginning of the time interval, t . This choice is mathematically convenient because it ensures that the integrand is independent of the future increment dW_t , which makes the resulting stochastic integral a martingale.

However, this is not the only possible choice. An alternative, proposed by Ruslan Stratonovich, evaluates the diffusion term at the midpoint of the time interval, $t + dt/2$. This gives rise to the **Stratonovich SDE**, denoted with a \circ :

$$dx = a(x)dt + b(x) \circ dW_t$$

Its discretized form is:

$$x(t+dt) - x(t) = a(x(t))dt + b\left(x\left(t + \frac{dt}{2}\right)\right)dW_t$$

This subtle difference in the evaluation point has profound consequences, leading to a different set of calculus rules. While the Itô integral is a martingale, the Stratonovich integral is not. However, the Stratonovich chain rule follows the rules of ordinary calculus, which can make it more intuitive for physicists and engineers modeling systems derived from physical principles.

Connecting the Itô and Stratonovich SDEs

We can find a direct relationship between the two formalisms by expanding the midpoint term in the Stratonovich definition. First, we approximate the state at the midpoint:

$$x\left(t + \frac{dt}{2}\right) \approx x(t) + a(x(t))\frac{dt}{2} + b(x(t))\left(W\left(t + \frac{dt}{2}\right) - W(t)\right)$$

Next, we expand the function b around $x(t)$:

$$b\left(x\left(t + \frac{dt}{2}\right)\right) \approx b(x(t)) + b'(x(t))\left(x\left(t + \frac{dt}{2}\right) - x(t)\right)$$

Substituting the first expression into the second gives:

$$b\left(x\left(t + \frac{dt}{2}\right)\right) \approx b(x(t)) + b'(x(t))\left[a(x(t))\frac{dt}{2} + b(x(t))d\hat{W}\right]$$

where $d\hat{W} = W(t + \frac{dt}{2}) - W(t)$. Now we substitute this back into the Stratonovich SDE definition:

$$dx = a(x)dt + \left(b(x) + b'(x)\left[a(x)\frac{dt}{2} + b(x)d\hat{W}\right]\right)dW_t$$

Expanding this, we get three terms involving dW_t :

$$dx = a(x)dt + b(x)dW_t + a(x)b'(x)\frac{dt}{2}dW_t + b(x)b'(x)d\hat{W}dW_t$$

The term $dt dW_t$ is of order $O(dt^{3/2})$ and can be neglected. The final term requires us to evaluate the expectation $\langle d\hat{W}dW_t \rangle$.

$$\langle d\hat{W}dW_t \rangle = \left\langle \left(W(t + \frac{dt}{2}) - W(t) \right) (W(t + dt) - W(t)) \right\rangle = \min \left(t + \frac{dt}{2}, t + dt \right) - t = \frac{dt}{2}$$

where the minimum comes from the autocorrelation of the Wiener process $\langle W(t)W(s) \rangle = \min(t, s)$.

Therefore, the term $b(x)b'(x)d\hat{W}dW_t$ contributes a drift of $\frac{1}{2}b(x)b'(x)dt$. Combining all terms, the Stratonovich SDE is equivalent to the following Itô SDE:

$$dx = \left(a(x) + \frac{1}{2}b(x)b'(x) \right) dt + b(x)dW_t$$

Definition: Stratonovich-to-Itô Conversion

A Stratonovich SDE

$$dx = a(x)dt + b(x) \circ dW_t$$

is equivalent to an Itô SDE with a modified drift term:

$$dx = \left(a(x) + \frac{1}{2}b(x)b'(x) \right) dt + b(x)dW_t$$

This conversion formula allows us to switch between the two calculi, leveraging the strengths of each. The extra drift term $\frac{1}{2}b(x)b'(x)$ is often called the "noise-induced drift" or "spurious drift."

3.10.1 Fokker-Planck Equation for Stratonovich SDEs

Given the conversion formula, we can easily find the Fokker-Planck equation for a Stratonovich SDE. We simply take the Fokker-Planck equation for an Itô SDE and replace the drift $a(x)$ with the effective drift $a(x) + \frac{1}{2}b(x)b'(x)$.

$$\frac{\partial \rho}{\partial t} = -\frac{\partial}{\partial x} \left[\left(a(x) + \frac{1}{2}b(x)b'(x) \right) \rho \right] + \frac{1}{2} \frac{\partial^2}{\partial x^2} [b(x)^2 \rho]$$

This equation can be rewritten in a more compact and elegant form. Noting that $b(x)b'(x) = \frac{1}{2} \frac{d}{dx}(b(x)^2)$, we can combine the terms:

$$\frac{\partial \rho}{\partial t} = -\frac{\partial}{\partial x} [a(x)\rho] + \frac{1}{2} \left(-\frac{\partial}{\partial x} \left[\frac{d(b^2)}{dx} \rho \right] + \frac{\partial^2}{\partial x^2} [b^2 \rho] \right)$$

Let us use the product rule on the second derivative:

$$\frac{\partial^2}{\partial x^2} [b^2 \rho] = \frac{\partial}{\partial x} \left(\frac{\partial(b^2)}{\partial x} \rho + b^2 \frac{\partial \rho}{\partial x} \right) = \frac{\partial^2(b^2)}{\partial x^2} \rho + 2 \frac{\partial(b^2)}{\partial x} \frac{\partial \rho}{\partial x} + b^2 \frac{\partial^2 \rho}{\partial x^2}$$

Now, consider the term

$$-\frac{\partial}{\partial x} \left[\frac{\partial(b^2)}{\partial x} \rho \right] + \frac{\partial^2}{\partial x^2} [b^2 \rho]$$

Substituting the expanded form, we get:

$$\frac{\partial \rho}{\partial t} = - \left(\frac{\partial^2(b^2)}{\partial x^2} \rho + \frac{\partial(b^2)}{\partial x} \frac{\partial \rho}{\partial x} \right) + \left(\frac{\partial^2(b^2)}{\partial x^2} \rho + 2 \frac{\partial(b^2)}{\partial x} \frac{\partial \rho}{\partial x} + b^2 \frac{\partial^2 \rho}{\partial x^2} \right)$$

Cancelling the common terms, we are left with:

$$\frac{\partial \rho}{\partial t} = \frac{\partial(b^2)}{\partial x} \frac{\partial \rho}{\partial x} + b^2 \frac{\partial^2 \rho}{\partial x^2} = \frac{\partial}{\partial x} \left[b \frac{\partial}{\partial x} (b \rho) \right]$$

Therefore, the equation becomes:

$$\frac{\partial \rho}{\partial t} = - \frac{\partial}{\partial x} [a(x) \rho] + \frac{1}{2} \frac{\partial}{\partial x} \left[b(x) \frac{\partial}{\partial x} (b(x) \rho) \right]$$

This is the common form of the Fokker-Planck equation found in many physics textbooks. It highlights that in the Stratonovich interpretation, the diffusion term has a simpler structure.

The Stationary Distribution for Stratonovich SDEs

The Stratonovich form turns out to be quite convenient when dealing with the stationary distribution P_s , since

$$0 = - \frac{d}{dx} [a(x) P_s] + \frac{1}{2} \frac{d}{dx} \left[b(x) \frac{d}{dx} (b(x) P_s) \right]$$

By simplifying and rearranging, we easily obtain

$$\frac{b(x)}{2} \frac{d}{dx} [b(x) P_s] = a(x) P_s$$

Now, define $\mathcal{H} = b(x) P_s$ so that

$$\frac{d\mathcal{H}}{dx} = \frac{2a(x)}{b(x)^2} \mathcal{H} \implies \mathcal{H} = C \exp \left[\int_x \frac{2a(z)}{b(z)^2} dz \right]$$

Therefore,

$$P_s(x) = \frac{C}{b(x)} \exp \left[\int_x \frac{2a(z)}{b(z)^2} dz \right] = C \exp \left[-\log(b(x)) + \int_x \frac{2a(z)}{b(z)^2} dz \right]$$

Once again, we want a formula that allows us to calculate the stationary points, so we can verify that it is sufficient to consider the stationary points of the exponent:

$$\frac{2a(x)}{b(x)^2} - \frac{b'(x)}{b(x)} \geq 0 \implies a(x) \geq \frac{1}{2} b(x) b'(x)$$

The only difference compared to the Itô formulation is the factor $1/2$. The difference is therefore only quantitative, not qualitative. We have seen that the transitions due to noise can emerge in a new way that depends on the parameterization, even if the positions of the maxima and minima and their values change by different factors.

Applications of SSDEs

Let us revisit the Malthusian growth model, which is described by

$$\dot{x} = rx$$

where r is the growth rate.

Suppose now that the growth rate is subject to random fluctuations, so that $r \rightarrow r + \alpha\xi(t)$, where $\xi(t)$ is white noise.

Now, let us assume that the noise is such that it is described in the Stratonovich sense. The equation becomes

$$dx = rxdt + \alpha x \circ dW$$

where \circ denotes the Stratonovich integral. The corresponding Ito version of this formula is

$$dx = \left(rx + \frac{\alpha^2}{2}x \right) dt + \alpha x dW$$

If we now apply the logarithmic transformation $y = \ln x$, we obtain

$$dy = \left(r + \frac{\alpha^2}{2} - \frac{\alpha^2}{2} \right) dt + \alpha dW = rdt + \alpha dW$$

However, if we instead use the Ito interpretation from the start, the equation for x is

$$dx = rxdt + \alpha x dW$$

which, under the transformation $y = \ln x$, gives

$$dy = \left(r - \frac{\alpha^2}{2} \right) dt + \alpha dW$$

This leads to a crucial difference between the two interpretations.

- **Itô case:** If $\frac{\alpha^2}{2} > r$, then the drift term in the equation for $y = \ln x$ becomes negative, so $y(t) \rightarrow -\infty$ as $t \rightarrow \infty$. This means $x(t) \rightarrow 0$, and the population will eventually go extinct even if $r > 0$.
- **Stratonovich case:** The noise-induced drift term cancels out, so extinction only occurs if $r < 0$. In other words, the population persists as long as the deterministic growth rate r is positive, regardless of the noise strength.

This difference is significant: in the Itô interpretation, strong enough noise can drive the population extinct even when the average growth rate is positive, while in the Stratonovich interpretation, only a negative growth rate leads to extinction.

Which interpretation should be used? The answer depends on the nature of the noise and the modeling context. The Itô calculus is appropriate when the noise is idealized as perfectly white and uncorrelated in time. The Stratonovich calculus is more suitable when the noise has a finite correlation time or comes from a smooth, physical process. In real-world applications, the correct choice depends on the details of the system being modeled.

In summary, the choice between Itô and Stratonovich calculus is not just a technicality: it can fundamentally change the predicted behavior of stochastic models, especially for questions like extinction thresholds and long-term outcomes.

⚠ Warning: *Raw Lecture Notes*

The following chapters contain unstructured notes taken during lectures and may require further organization and refinement.

4

Lecture: 04/04/2025

...

$$\frac{\partial P}{\partial t}(x,t) = \int_{S \in s} P(s,t) \Omega(s,x) ds - P(x,t) \int_{a \in s} \Omega(x,a) da$$

...

$$\frac{\partial P}{\partial t}(x,t) = \underbrace{rP(x+1,t)}_{backward} + \underbrace{rP(x-1,t)}_{forward} - 2rP(x,t)$$

...

More in general we can write

$$\begin{aligned} \Omega(s,a) &= \sum_{j \in \mathbb{Z}} K_{s,a} \delta(s-a-j) \\ \Rightarrow \quad \frac{\partial P}{\partial t}(x,t) &= \sum_{j \in \mathbb{Z}} \int P(s,t) K_{s,x} \delta(s-x-j) ds - P(x,t) \sum_{j \in \mathbb{Z}} \int K_{x,a} \delta(x-a-j) da \\ \frac{\partial P(x,t)}{\partial t} &= \left(\sum_{j \in \mathbb{Z}} K_{x,x+j} K_{x+j,x} P(x+j,t) \right) - P(x,t) \left(\sum_{j \in \mathbb{Z}} K_{x,x-j} \right) \end{aligned}$$

...

$$\begin{cases} S' = -\beta \frac{I}{N} S \\ I' = \beta \frac{I}{N} S - \gamma I \end{cases}$$

This system represents the dynamics of a population of individuals that can be in one of two states: susceptible (S) or infected (I). The parameter β represents the rate at which susceptible individuals become infected, while γ represents the rate at which infected individuals recover.

We have ($R = N - S - I$) and:

$$I(t) \geq 0; \quad S(t) \geq 0; \quad R(t) \geq 0$$

$$(t, t+dt) X(t) = (S(t), I(t)) \in \mathbb{R}^2$$

The "removal" of individuals accounts for the recovery of infected individuals, its probability is given by:

$$\Pr \left[(S(t+dt), I(t+dt)) = (S(t), I(t)-1) \mid (S(t), I(t)) \right] = \gamma dt$$

...

$$\underbrace{\begin{pmatrix} S \\ I \end{pmatrix}}_t \rightarrow \underbrace{\begin{pmatrix} S \\ I-1 \end{pmatrix}}_{t+dt} = \begin{pmatrix} S \\ I \end{pmatrix} + \begin{pmatrix} 0 \\ -1 \end{pmatrix}$$

The contagion process instead is given by:

$$\underbrace{\begin{pmatrix} S \\ I \end{pmatrix}}_t \rightarrow \underbrace{\begin{pmatrix} S \\ I \end{pmatrix} + \begin{pmatrix} -1 \\ 1 \end{pmatrix}}_{t+dt}$$

...

We can think of σ and α as the states of the system at time t and $t + dt$, respectively:

$$\sigma = \begin{pmatrix} S_\sigma \\ I_\sigma \end{pmatrix}, \quad \alpha = \begin{pmatrix} S_\alpha \\ I_\alpha \end{pmatrix}$$

...

$$\Omega(\sigma, \alpha) = \gamma I_\sigma \delta \left(\sigma - \alpha - \begin{pmatrix} 0 \\ -1 \end{pmatrix} \right) + \beta \frac{I_\sigma}{N} S_\sigma \Omega \left(\sigma - \alpha - \begin{pmatrix} -1 \\ 1 \end{pmatrix} \right)$$

...

$$\frac{\partial P(S, I, t)}{\partial t} = \beta \frac{(I-1)}{N} (S+1) P(S+1, I-1, t) + \gamma (I+1) P(S, I+1, t) - \left(\gamma I + \beta \frac{I}{N} S \right) P(S, I, t)$$

Your starting point is an integer position; whatever jump you take, $x + dt$ must be another integer position. In fact, the probability distribution is non-zero only for integer positions.

4.1 Discrete time, Discrete state space

State s is discrete

Time $t \subseteq \mathbb{N} \cup \{0\}$

$$P\{x(t+1) = \alpha | x(t) = \sigma\} = \theta_{\sigma\alpha} \in [0, 1]$$

$$P\{x(t) = \omega\} = P_\omega(t)$$

The probability of being in a state α at time $t + 1$ is given by the probability of being in state σ at time t multiplied by the transition probability from σ to α :

$$P_\alpha(t+1) = \sum_{\sigma \in S} \Pr(\sigma) \theta_{\sigma\alpha}$$

$$P_\alpha(t+1) = P_\alpha(t) \theta_{\alpha\alpha} + \sum_{\sigma \in S \setminus \{\alpha\}} P_\sigma(t) \theta_{\sigma\alpha}$$

$$\theta_{\alpha,\alpha} = 1 - \sum_{\beta \in S \setminus \{\alpha\}} \theta_{\alpha,\beta}$$

$$P_\alpha(t+1) = P_\alpha(t) \left[1 - \sum_{\beta \in S \setminus \{\alpha\}} \theta_{\alpha,\beta} \right] + \sum_{\sigma \in S \setminus \{\alpha\}} P_\sigma(t) \theta_{\sigma,\alpha}$$

$$\begin{aligned}
P_\alpha(t+1) &= P_\alpha(t) + \sum_{\sigma} P_\sigma(t) \theta_{\sigma,\alpha} - P_\alpha(t) \sum_{\beta \in S \setminus \{\alpha\}} \theta_{\alpha,\beta} \\
P_\alpha(t+1) - P_\alpha(t) &= \sum_{\sigma} P_\sigma(t) \theta_{\sigma,\alpha} - P_\alpha(t) \sum_{\beta \in S \setminus \{\alpha\}} \theta_{\alpha,\beta} \\
\frac{P_\alpha(t+1) - P_\alpha(t)}{U} &= \sum_{\sigma} P_\sigma(t) \frac{\theta_{\sigma,\alpha}}{U} - P_\alpha(t) \sum_{\beta} \frac{\theta_{\alpha,\beta}}{U} \\
&\dots
\end{aligned}$$

4.2 SIS model

A SIS model is a simple model of disease spread in a population. In this model, individuals can be in one of two states: susceptible (S) or infected (I). The dynamics of the system are governed by two parameters: the infection rate β and the recovery rate γ . The model assumes that individuals can move between these two states, with susceptible individuals becoming infected at a rate proportional to the number of infected individuals they come into contact with, and infected individuals recovering at a constant rate.

Moreover, μ is the natural death rate of the population, which is assumed to be the same as the birth rate. This means that the population size remains constant over time, and the total number of individuals in the population is given by $N = S + I$.

$$\begin{cases} S' = \mu - \mu S - \beta IS \\ I' = \beta IS - (\mu + \gamma)I \end{cases}, \quad \begin{cases} \mu \rightarrow \mu + \omega_\mu \xi_\mu \\ \beta \rightarrow \beta + \omega_\beta \xi_\beta \end{cases}, \quad \xi(t) = \begin{pmatrix} \xi_\mu \\ \xi_\beta \end{pmatrix}.$$

This is actually a stochastic model of disease spread, where the parameters μ, β , are subject to random fluctuations.

...

[a lot of stuff missing]

...

Case whit indipendent noise and ...

$$\begin{aligned}
x_j &= x(jh) \\
x_{j+1} &= x_j + \alpha(x_j)h + \beta(x_j) \begin{bmatrix} G_j \\ G_{j+1} \end{bmatrix} \sqrt{h}
\end{aligned}$$

...

4.3 ma che cazzo ne so

Suppose we are in a region of R^N and all the points follows the following law:

$$\begin{aligned}
\dot{x}_i &= f(x_i) \\
n(x, 0) &= \tilde{O}(x) \\
\inf n(x, t) dx &= N \\
\int n(x, 0) dx &= \int \tilde{O}(x) dx = N \\
c &= [a, b]
\end{aligned}$$

$$N_c(t) = \int_a^b n(x,t)dx$$

$$P[x(t) \in [a,b]] \simeq \frac{N_c(t)}{n}$$

...

4.4 ma che cazzo ne so pt.2

Let's consider an interval $[t, t + dt]$ and some particles that follows the law:

$$\dot{x} = f(x) \equiv v(x)$$

Then, the number of particles that enter and exit the interval $[x, x + dx]$ at time t is given by:

$$\begin{cases} \text{Enter : } n(x,t)v(x)dt \\ \text{Exit : } n(x+dt,t)v(x+dt)dt \end{cases}$$

...

[missing a lot of stuff]

...

the product $n \cdot v$ is called **current** density and is denoted by $J(x,t)$:

$$\begin{cases} \frac{\partial n}{\partial t} + \frac{\partial}{\partial x} J(x,t) = 0 \\ J(x,t) = n(x,t)v(x) \end{cases}, \quad \frac{\partial n}{\partial t} + \text{div}J(x,t) = 0$$

...

... probability current ...

5

Lecture 11/04/2025

$$\dot{x} = a(x) + b(x)\eta_h(t)$$

$$\langle \eta_h(t) \rangle = 0$$

\mathbb{R}_h is a function with a peak at $|z| < h$ and $\mathbb{R}_h(z) = 0$ for $|z| > h$.

If the limit $\lim_{h \rightarrow 0^+} \mathbb{R}_h(z) = \delta(\tau)$, then:

$$\dot{x}_h = a(x_h) + b(x_h)\eta_h(t) \rightarrow \dot{x} = a(x) + b(x)\circ\xi(t)$$

Let's consider a population and two opinions:

$$x + y = 1$$

People can change their opinion with a rate R and the ratio of changing from x to y is θ , and from y to x is k .

$$\begin{cases} \dot{x} = +\theta xy - kyx - \varepsilon x + \varepsilon y \\ \dot{y} = -\theta xy + kyx + \varepsilon x - \varepsilon y \end{cases}$$

where ε is the rate of changing opinion. (?)

Substituting $y = 1 - x$ we can rewrite the first equation as:

$$\dot{x} = x(1-x)(\theta - k) + \varepsilon(1-x) - \varepsilon x$$

$$\frac{dx}{dt} = \lambda x(1-x) + 1 - 2x, \quad \lambda \rightarrow \lambda + \alpha \xi(t)$$

$$dx = \underbrace{\{\lambda x(1-x) + 1 - 2x\}}_{a(x)} dt + \underbrace{\alpha x(1-x)}_{b(x)} \circ dW$$

...

$$dx = \{1 - 2x\} dt + \alpha x(1-x) \circ dW$$

$$a(x) = \frac{1}{2}b'(x)b(x), \quad b(x) = \alpha(x - x^2), \quad b'(x) = \alpha(1 - 2x)$$

so

$$1 - 2x = \frac{\alpha^2}{2}x(1-x)(1-2x)$$

which has two equilibrium points:

$$x_1 = \frac{1}{2}, \quad x_2 : 1 = \frac{\alpha^2}{2}x(1-x) \quad (?)$$

...

$$m\ddot{x} = -\gamma_T \dot{x} + F_T(x) + \omega_T \xi(t)$$

$$\begin{cases} \dot{x} = v \\ \dot{v} = -\frac{\gamma_T}{m}v + \frac{F_T(x)}{m} + \frac{\omega_T}{m}\xi(t) \end{cases}$$

to simplify the notation we can write:

$$\gamma = \frac{\gamma_T}{m}, \quad F = \frac{F_T(x)}{m}, \quad \omega = \frac{\omega_T}{m}, \quad U = \underbrace{\int F_T(x)dx}_{?}, \quad U' = \frac{dU}{dx} = F_T(x)$$

we get:

$$\begin{aligned} & \begin{cases} \dot{x} = v \\ \dot{v} = -\gamma v + F(x) + \omega \xi(t) \end{cases} \\ & \frac{\partial p}{\partial t} = -v \frac{\partial p}{\partial x} - \frac{\partial}{\partial v} [(F(x) - \gamma v)p] + \frac{\partial^2 \omega^2}{\partial v^2} \frac{v^2}{2} \\ & \frac{\partial p}{\partial t} = -v \frac{\partial p}{\partial x} - F(x) \frac{\partial p}{\partial v} + \gamma \frac{\partial}{\partial v} (vp) + \frac{\partial^2 \omega^2}{\partial v^2} \frac{v^2}{2} \\ & 0 = -v \frac{\partial p}{\partial x} + U'(x) \frac{\partial p}{\partial v} + \gamma p + \gamma v \frac{\partial p}{\partial v} + \frac{\partial^2 \omega^2}{\partial v^2} \frac{v^2}{2} \\ & p(x, v) = A(x)B(v) \\ & -vA'(x)B(v) + U'(x)A(x)B'(v) + \gamma A(x)B(v) + \gamma v A(x)B'(v) + \frac{\omega^2}{2} B''(v) = 0 \\ & \underbrace{-v \frac{A'(x)}{A(x)} + U'(x) \frac{B'(v)}{B(v)}}_{=0} + \underbrace{1 \gamma + \gamma v \frac{B'(v)}{B(v)} + \frac{\omega^2}{2} \frac{B''(v)}{B(v)}}_{2^{nd} \text{ term}} = 0 \end{aligned}$$

We have to options:

1. set the second term to zero
2. boh

Let's consider the first option and let's define some "test" variables B_T and B'_T . We have:

$$\frac{B'_T(v)}{B_T(v)} = -\eta v \Rightarrow B'_T(v) = -\eta v B_T(v) \Rightarrow B(v) = C e^{-\eta v^2/2}$$

we have

$$B'(v) = -\eta v B(v), \quad B''(v) = -\eta v B'(v) = -\eta v (-\eta v B(v))$$

MISSING: boh

$$P_s = \frac{1}{z} e^{-\frac{\gamma}{\omega^2} v^2 - \frac{2\gamma}{\omega^2} U(x)} = \frac{1}{z} e^{-\frac{2\gamma}{\omega^2} \left[\frac{v^2}{2} + U(x) \right]}$$

Applying back the transformation we have:

$$p(x, v) = \frac{1}{z} e^{-\frac{2\gamma_T}{\omega_T^2} \left[\frac{mv^2}{2} + U_T(x) \right]}$$

so:

$$\iint p_s(x, v) dx dv = 1, \quad \frac{1}{z} \iint e^{-\frac{2\gamma_T}{\omega_T^2} E_T(x, v)} dx dv = 1$$

TODO: check if this is correct

MISSING: fishes example ?

$$dx = f(x)dt - \underbrace{(cxdt + \omega x dW)}_{\text{#fishes killed in } (t, t+dt)}$$

We want the number of fishes to be positive.

MISSING: end of the lecture

6

Lecture: 05/05/2025

...

... if there is no linearity, ...

$$\frac{\partial P}{\partial t} = -\frac{\partial}{\partial x} \left\{ (\theta \int zP(z,t)dz + (1-\theta)x - x^3)P \right\} + \frac{\omega^2}{2} \frac{\partial^2 P}{\partial x^2} \quad N \gg 1$$

$$M(t) = \int_{\mathbb{R}} xP(x,t)dx$$

$$P_s(x, M_s) = C(M_s) \exp \{ \theta M_s x + \dots \}$$

MISSING: end of the formula above

This solution is not actually so "usable"

$$M_s = \int_{\mathbb{R}} xP_s(x; M_s)dx \Rightarrow M_s = \Psi(M_s)$$

$$M_s = \Psi(M_s) \rightarrow \text{"unique solution"}$$

There are more interesting cases, for instance when $\Psi(M_s)$ has more than one solution:

In this case, our system has more than one steady states. It means that we loose the unicity of the solution (so there is no more global attractiveness)

E.g.

$$\dot{X}_i = f(x_i, < x >) + g(x_i)\xi_i \quad N \gg 1$$

$$\dot{x} = f(x, M(t)) + g(x)\xi(t)$$

$$M(t) = \int zP(z,t)dt$$

The Fokker-Plank equation will be:

$$\frac{\partial P}{\partial t} = -\frac{\partial}{\partial x}[f(x, M(t))] + \frac{1}{2} \frac{\partial^2}{\partial x^2}[g^2(x)P]$$

The steady state solutions will be the solutions of the following equation:

$$\begin{cases} 0 = -\frac{d}{dx}[t(x, M_s)P] + \frac{d^2}{dx^2} \left[\frac{g^2(x)}{2} P \right] \\ M_s = \int zP_s(z; M_s)dx \end{cases}$$

$M_s = \Psi(M_s)$

TODO: add linking sentence

$$P(x, M_s, \theta) = C(M, \theta) \exp \left[\frac{2}{\omega} \left(\theta M_s x + (1 - \theta) \frac{x^2}{2} - \frac{x^4}{4} \right) \right]$$

$$\boxed{M_s = 0}$$

... (?)

$$0 < \theta < \theta_c$$

... (?)

So we have two solutions:

$$\begin{aligned} M_s &= a \\ M_s &= -a \end{aligned}$$

➊ Example:

$$M_s = \Psi(M_s; \theta)$$

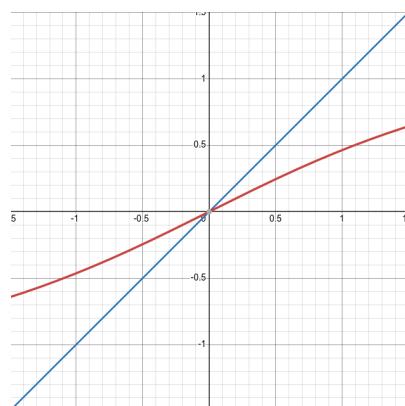
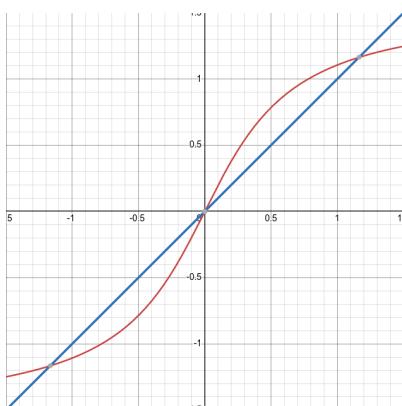
$$\begin{cases} y = M_s \\ y = \Psi(M_s; \theta) \end{cases} \Rightarrow \begin{array}{l} P_s(x; M_1) \\ P_s(x; M_2) \\ P_s(x; M_3) \end{array}$$

so for $\theta = \theta_1$ we have **multistability**, while for $\theta = \theta_2$ we have **monostability**.

Theorem 1.

$$\left| \frac{d\Psi}{dM_s} \Big|_{M_s=M_c} \right| < 1 \Rightarrow P_s(x; M_1, \theta^1) \text{ is locally stable}$$

$$\left| \frac{d\Psi}{dM_s} \Big|_{M_s=M_c} \right| > 1 \Rightarrow P_s(x; M_2, \theta^2) \text{ is locally unstable}$$



...

$$\dot{x} = (ax + x^3 - x^5) - D(x - M(t)) + \alpha(1 + x^2) \odot \xi(t)$$

$$M = \Psi(M; D, \alpha)$$

we have that for small D and α we have a unique solution, while for large D and α we have 5 different solutions.

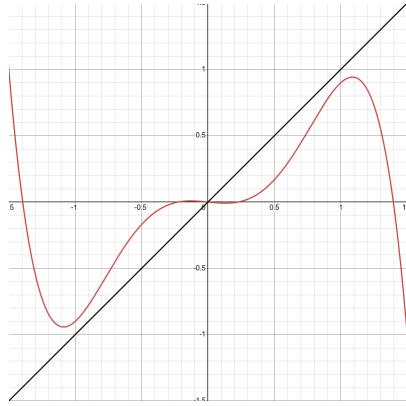


Figure 6.1: 1 solution

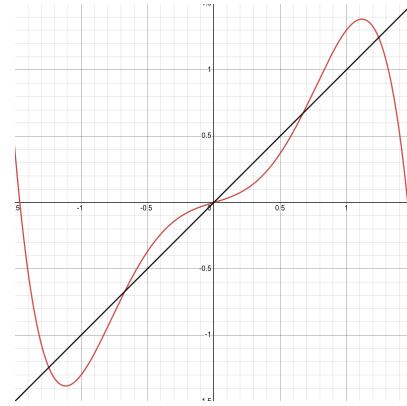


Figure 6.2: 5 solutions

So 0 is always a solution, and from a certain value of D we have 5 solutions, 3 of which are stable and 2 are unstable.

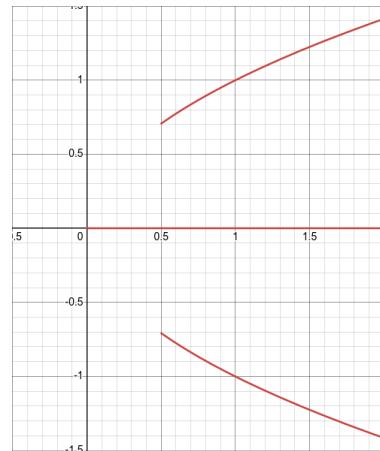


Figure 6.3: Stable solutions of the system

...

$$\dot{x} = F(x_i, \langle x \rangle) + g(x_i, \langle x \rangle) \xi_i(t)$$

An example is the movement of a guitar string that vibrates.

$$m_i \ddot{x}_i = -\gamma \dot{x}_i$$

If the deviation is big, we have a function in the complex space, but if the deviation is small, we can use the linear approximation.

$$m_j \ddot{z}_j = -\gamma \dot{z}_j - k(z_j - z_{j-1}) - k(z_j - z_{j+1}) = -\gamma \dot{z}_j - k(z_{j-1} - 2z_j + z_{j+1})$$

we have now a discretization of the position of the string $z(t, x)$:

$$m_j \frac{d^2 z}{dt^2}(t, x_j) = -\gamma \frac{dz}{dt}(t, x_j) + k[z(t, x_j + D) - 2z(t, x_j) + z(t, x_j - D)]$$

$$\mu \frac{\partial^2 z}{\partial t^2} = -\gamma \frac{\partial z}{\partial t} + c \frac{\partial^2 z}{\partial x^2} + \hat{\omega} \xi(x, t)$$

We can see the stochastic term as the wind that moves the string.

7

Lecture: 09/05/2025

7.1 Spatiotemporal noisy model

$$\frac{\partial \phi}{\partial t} = f(\phi) + g(\phi)\xi_m(r,t) + D\mathcal{L}[\phi] + h(\phi)F(t) + \xi_a(r,t)$$

- $f(\phi)$: deterministic part
- $g(\phi)$: multiplicative noise
- $D\mathcal{L}[\phi]$: linear part
- $h(\phi)$: additive noise

$\mathcal{L}[\phi]$ is a Laplacian or a integral operator

Examples:

- $\mathcal{L}[\phi] = \nabla^2 \phi$
- $\mathcal{L}[\phi] = -a_0 \nabla^2 \phi - \nabla^4 \phi$
- $\mathcal{L}[\phi] = -(\nabla^2 + k_0^2)^2 \phi = -(K_0^2 + 2k_0 \nabla^2 + \nabla^4) \phi$

Observation:

If we apply the fourier transform of:

$$\mathcal{L}[\phi] = -(\nabla^2 + k_0^2)^2 \phi = -(K_0^2 + 2k_0 \nabla^2 + \nabla^4) \phi$$

we get:

$$...F(\phi)$$

$$\mathcal{L}[\phi(r)] = \int \phi(r') \omega(r - r') dr'$$

How do we simulate this equation? We can simply obtain the domain and discretize it.
Lattice-based Approximation:

...

Field coupling approximation:

$$l(\phi_i, \phi_j) = w_i \phi_i + \sum_{j \in nn(i)} w_j \phi_j$$

For example:

$$\mathcal{L}[\phi] = \nabla^2 \phi \approx l(\phi_i, \phi_j) = \frac{1}{\Delta^2} \sum_{j \in nn(i)} (\phi_j - \phi_i)$$

If we have a stochastic process which is discrete in time and space, we have:

$$\langle \xi(r, t) \xi(r', t') \rangle = sC \left(\frac{|r - r'|}{d}, \frac{|t - t'|}{\tau_c} \right)$$

As in the purely temporal noise, τ_c is a measure of the temporal memory of the noise, d is the spatial memory of the noise.

The spatiotemporal brother of the ornstein-uhlenbeck noise is "Ojalvo e al" noise.

$$\frac{\partial \phi}{\partial t} = a\phi + D\nabla^2 \phi + \xi_{gn}$$

⌚ Observation: Ojalvo e al and the Ornstein-Uhlenbeck process

If we set $D = 0$ we have a series of Ornstein-Uhlenbeck processes at each point of the domain.

$$\frac{\partial \phi}{\partial t} = a\phi + \xi_{gn}$$

Noise induced pattenrs §

$$\frac{\partial \phi}{\partial t} = f(\phi) + g(\phi) \xi_m(r, t) + D\mathcal{L}[\phi] + \dots$$

...

Perturbed Swift-Hohenberg model:

$$\frac{\partial \phi}{\partial t} = a\phi + D\mathcal{L}[\phi] + \xi_{gn} \dots$$

$$\frac{\partial \phi}{\partial t} = f(\phi) + D\mathcal{L}[\phi] = a\phi - D(\nabla^2 + k_0^2)^2 \phi$$

Transitory pattern that disappear.

Additive noise generate pattenrs

$$\frac{\partial \phi}{\partial t} = a\phi + D\mathcal{L}[\phi] + \xi_{gn}, \quad \mathcal{L}[\phi] = -(\nabla^2 + k_0^2)^2 \phi$$

Permanent patterning: details change in time.

we can distinguish two cases:

- $a < 0$
- $a > 0$

with multiplicative noise it can induce bimodality in the pdf of ϕ :

$$\frac{\partial \phi}{\partial t} = a\phi - phi^3 + \phi \xi_{gn} + D\mathcal{L}[\phi]$$

A bad model of glaciations

$$dx = [x(a - x^2) + A \cos \Omega t] dt$$

where

- x : is the (normalized) Earth's temperature
- $A \cos(\Omega t)$: small periodic variations of the solar irradiation.

We have that if A is small, $x(t)$ fluctuates around $+\sqrt{a}$.

The model fails.

Including stochastic noise:

$$dx = [x(a - x^2) + A \cos \Omega t] dt + \varepsilon dW$$

where ε is the noise intensity.

This time we have a white noise, according to:

- ε is small: the noise is negligible
- ε is large: the noise is dominant

they finally managed to model the glaciations.

Spatial Stochastic Resonance

$$\frac{\partial \phi}{\partial t} = a\phi - \phi^3 + D \frac{\partial^2 \phi}{\partial x^2} + F(t) + \varepsilon \xi_{gn}$$

TODO: check the formula

...

8

Discrete Time Markov Chains

We consider the case where the time is a subset of the integers ($t \subset \mathbb{Z}$ or $t \subset \mathbb{N}_0$)

The state space is a finite set $S = \{s_1, s_2, \dots, s_N\}$

$$P\{x(t+1)|x(0), x(1), \dots, x(t)\} = P\{x(t+1)|x(t)\}$$

The probability that the process in $t+1$ is σ is the sum of the probability that the process in t is δ and the probability that the process in $t+1$ is σ given that the process in t is δ .

$$P\{x(t+1) = \sigma\} = \sum_{\delta \in S} P\{x(t+1) = \sigma | x(t) = \delta\} P\{x(t) = \delta\} = \sum_{\delta \in S} P\{x(t) = \delta\} \theta_{\delta\sigma}$$

with $\theta_{\delta\sigma} \in [0, 1]$.

So we have:

$$P_\sigma(t+1) = \sum_{\delta \in S} P_\delta(t) \theta_{\delta\sigma} \Rightarrow P(t+1) = P(t)\Theta(t)$$

Where $P(t) = [P_1(t), P_2(t), \dots, P_N(t)]$ is the probability vector at time t and $\Theta(t)$ is the transition matrix at time t .

We have:

$$P(1) = P(0)\Theta, \quad P(2) = P(1)\Theta = P(0)\Theta^2, \quad P(3) = P(2)\Theta = P(0)\Theta^3, \quad \dots$$

We can write:

$$P(t) = P(0)\Theta^t = P(0) \prod_{q=0}^{t-1} \Theta(q)$$

We have two properties:

$$\sum_{\sigma \in S} \theta_{\delta\sigma} = 1, \quad \sum_{\delta \in S} P_\delta(t) = 1$$

So we have:

$$P_\sigma(t+1) = \sum_{\delta \in S} P_\delta(t) \theta_{\delta\sigma} \Rightarrow \sum_{\sigma \in S} P_\sigma(t+1) = \sum_{\sigma} \sum_{\delta} P_\delta(t) \theta_{\delta\sigma} = \sum_{\delta} P_\delta(t) \sum_{\sigma} \theta_{\delta\sigma} = 1$$

So

MISSING: something

Let's consider again the discrete time Markov chain.

$$P_\sigma(t+1) = P(t)\Theta \Rightarrow P(t) = P(0)\Theta^t$$

...

$$P^\infty = P^\infty \Theta$$

Let's study the eigenvalues of Θ :

$$\Theta v = \lambda v$$

$$\sum_{\sigma} \theta_{\delta\sigma} = 1$$

(... non ho capito perchè ma l'autovettore di Θ è formato da tutti 1 ...)

If the multiplicity of the eigenvalue is more than 1, we have multiple solutions to our system.



Sometimes this equation is written as:

$$P_a(t+1) = \sum_{s \in S} W_{as} P_s(t)$$

where W_{as} ...

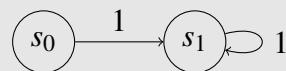
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The simplest markov chain is the one with only one state, with a transition with probability 1 from the state to itself.

$$P(0) = [1], \quad P(t) = [1]$$

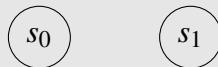
Another simple markov chain is the one with two states, with a transition with probability 1 from the state s_0 to the state s_1 and with probability 1 from the state s_1 to the state s_1 itself.



We have the situation below:

$$P(0) = [1, 0], \quad P(t) = [0, 1]$$

A more complex example is the following:



TODO: add the transitions in the figure

$$P_0(t+1) = \theta_{00} P_0(t) + P_1(t) \theta_{10}$$

$$P_1(t) = 1 - P_0(t) \Rightarrow P_0(t+1) = \theta_{00}P_0(t) + \theta_{10}(1 - P_0(t))$$

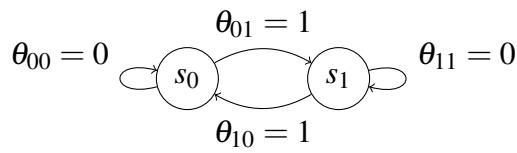
$$P_0(t+1) = (\theta_{00} - \theta_{10})P_0(t) + \theta_{10}$$

If we have a case like:

$$x(t+1) = ax + b$$

then we have a term b that don't allow us to use the resolution formula we are used to ($x(t) = x(0)a^t$), but ...

...



$$P(0) = [1, 0] = x(0) = 0$$

$$P(1) = [0, 1] = x(1) = 1$$

$$P(2) = [1, 0] = x(2) = 0$$

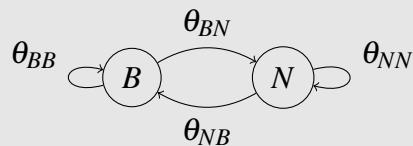
⋮

Which is periodic.

...

② Example: *Bora example*

Let's consider the case of Bora in Trieste. We have two states: B and N .



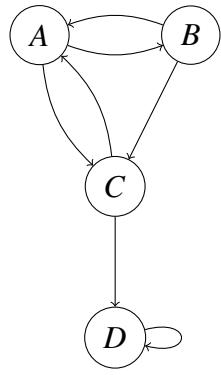
We have the following transition probabilities:

$$\begin{aligned} \theta_{BB} &\approx \frac{NT^{BB}}{N_B} & \theta_{BN} &\approx \frac{NT^{BN}}{N_B} \\ \theta_{NB} &\approx \frac{NT^{NB}}{N_N} & \theta_{NN} &\approx \frac{NT^{NN}}{N_N} \end{aligned}$$

This model is actually too artificial, because the transition probabilities are not independent, and also depends on other parameters, like the temperature.

The most natural representation of Markov Chains are oriented graphs.

Let's consider a 4 states model:



...boh...

Boh

9

Lecture: 12/05/2025

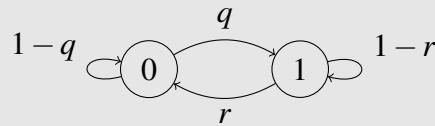
$$P(t+1) = P(t)\theta = \dots$$

$$P(t) = P(0)\theta^t = P(0)\dots$$

???????

Example: 2 States Markov Chain

Let's consider the following example (similar to the previous one):



We have two states (0 and 1), and a probability q of going from 0 to 1, and a probability r of going from 1 to 0. Let $P_0(t)$ be the probability that the system is in state 0 at time t , and $P_1(t)$ the probability of being in state 1 at time t . The evolution of these probabilities is given by:

$$\begin{cases} P_0(t+1) = (1-q)P_0(t) + rP_1(t) \\ P_1(t+1) = qP_0(t) + (1-r)P_1(t) \end{cases}$$

Since there are only two states, the probabilities must sum to 1:

$$P_1(t) + P_0(t) = 1$$

So, $P_1(t) = 1 - P_0(t)$.

Stationary (Equilibrium) Distribution: In the long run, the probabilities reach a steady state (stationary distribution), where $P_0(t+1) = P_0(t) = P_0^e$ and $P_1(t+1) = P_1(t) = P_1^e$. Setting the evolution equations to equilibrium, we get:

$$\begin{cases} P_0^e = (1-q)P_0^e + rP_1^e \\ P_1^e = qP_0^e + (1-r)P_1^e \end{cases}$$

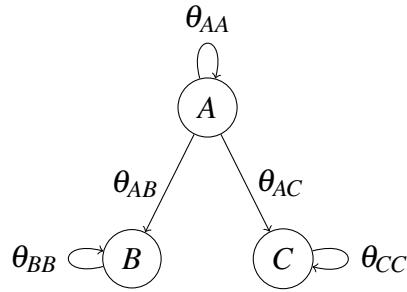
Using $P_1^e = 1 - P_0^e$, substitute into the first equation, we obtain:

$$P_0^e = \frac{r}{q+r}$$

Similarly:

$$P_1^e = \frac{q}{q+r} = 1 - P_0^e$$

Let's consider a Markov chain with three states: A , B , and C . The transitions are as follows:



Let $P_A(t)$, $P_B(t)$, $P_C(t)$ be the probabilities of being in states A , B , C at time t . The initial state is:

$$P(0) = (P_A(0), P_B(0), P_C(0))$$

Evolution Equations:

$$P_A(t+1) = \theta_{AA}P_A(t) + \theta_{AB}P_B(t) + \theta_{AC}P_C(t)$$

...

$$P_B(t+1) = P_B(t) + \theta_{AB}P_A(t) = P_B(t) + \theta_{AB}P_A(0)\theta^t$$

$$P_B(1) = P_B(0) + \theta_{AB}P_A(0)$$

$$P_B(2) = P_B(1) + \theta_{AB}P_A(1) = P_B(0) + \theta_{AB}P_A(0) + \theta_{AB}P_A(1)$$

$$P_B(t) = P_B(0) + \theta_{AB}P_A(0) [1 + \theta_{AA} + \theta_{AA}^2 + \dots + \theta_{AA}^{t-1}]$$

...

$$P_B^{eq} = P_B(0) + P_A(0) \frac{\theta_{AB}}{1 - \theta_{AA}}$$

$$P_C^{eq} = P_C(0) + P_A(0) \frac{\theta_{AC}}{1 - \theta_{AA}}$$

$$\det(\theta - \lambda z) = \begin{vmatrix} \theta_{AA} - \lambda & \theta_{AB} & \theta_{AC} \\ 0 & 1 - \lambda & 0 \\ 0 & 0 & 1 - \lambda \end{vmatrix} = 0$$

...

$$P^e = H \text{Diag}[1, 0, \dots, 0] H^{-1} P(0)$$

and we have:

$$P^e = W P^e$$

Properties of W :

1. At least $\lambda_1 = 1$
2. $|\lambda_j| \leq 1$ for all j

3. $\sum_{j=1}^n V_j = 0$
 where V_j is the eigenvector of W associated with λ_j .

...

$$WV_j = \lambda_j V_j$$

$$\sum_{C=1}^n W_{RC}(V_j)_C = \lambda_j (V_j)_R$$

$$\sum_{R=1}^n \sum_{C=1}^n W_{RC}(V_j)_C = \lambda_j \left(\sum_{R=1}^n (V_j)_R \right)$$

$$\underbrace{\sum_{C=1}^n (V_j)_C}_S = \lambda_j \left(\underbrace{\sum_{R=1}^n (V_j)_R}_S \right) \Rightarrow (1 - \lambda_j)S = 0$$

...

$$P^{eq} = H \text{ Diag}[1, 0, \dots] P(0)$$

$$P(0) = C_2 V_1 + \sum_{j=2}^N C_j V_j$$

...

$$P(0) = U_1 + \sum_{j=2}^N C_j V_j$$

Where U_1 is the normalized vector

...

$$P(0) = C_1 V_1 + C_2 V_2 + \sum_{j=3}^N C_j V_j$$

$$1 = C_1 \|V_1\|_1 + C_2 \|V_2\|_1 \Rightarrow C_2 = \frac{1 - C_1 \|V_1\|_1}{\|V_2\|_1}$$

$$P(0) = C_1 U_1 + (1 - C_1) U_2 + \sum_{j=3}^N C_j U_j$$

$$P(t) = W^t P(0) = C_1 U_1 + (1 - C_1) U_2 + \sum_{j=3}^N C_j \lambda_j^t V_j$$

So C_1 depends on the initial state $P(0)$:

$$C_1 = P(P(0))$$

...

$$P(t) = C_1 P(0) U_1 + (1 - C_1(P(0))) U_2 + \sum_j C_j \lambda_j^t V_j$$

$$P^e = C_1 U_1 + (1 - C_1) U_2$$

If we have a problem and we want to study the time of remaining in a state A , we can simply consider all the transitions from A to other states as a single transition from A to a new state B (which represents the rest of the world), and it is given by the sum of all the initial transitions. So we have now only 2 transitions: the loop from A to itself, and the new AB transition.

We have:

$$x(0) = A, \quad P(0) = (1, 0)$$

$$P_A(t) = \theta_{AA}^t$$

...

$$P_A(T) = \theta_{AB} = \theta_{AB} \theta_{AA}^t = (1 - \theta_{AA}) \theta_{AA}$$

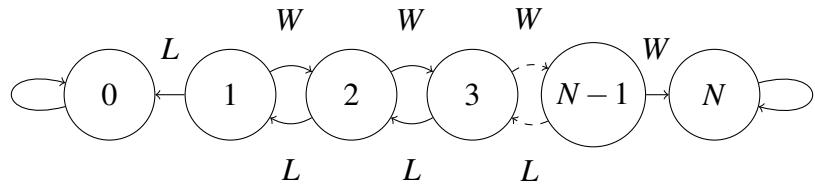
So the average value of T is:

$$\begin{aligned} \langle T \rangle &= \sum T \theta_{AB} \theta_{AA}^t = \theta_{AB} \sum_{t=0}^{\infty} t \theta_{AA}^t \\ &= \theta_{AB} [0 + \theta_{AA}^1 + 2\theta_{AA}^2 + 3\theta_{AA}^3] = \theta_{AB} \theta_{AA} [1 + 2\theta_{AA} + 3\theta_{AA}^2 + \dots] \\ &= \theta_{AB} \theta_{AA} \frac{d}{d\theta_{AA}} [-1 + 1 + \theta_{AA} + \theta_{AA}^2 + \dots] \\ &= \theta_{AB} \theta_{AA} \frac{d}{d\theta_{AA}} \left[\frac{1}{1 - \theta_{AA}} \right] \end{aligned}$$

10

Lecture: 16/05/2025

Consider a game, you have a probability p of winning, and a probability $1 - p$ of losing. You can bet 1 euro each time, and if you win you gain 1 euro, if you lose you lose 1 euro, if you run out of money you stop playing.



We can find that $L < W$.

$$\dots P(t) = W^t P(0)$$

$$\dots (\theta^t)_{AB} > 0$$

0 and N are **absorbing** states.

We have that a state σ is **transient** if (for $t \gg 1$) $P_\sigma(t) = 0$.

Another important property is the **ergodicity**: All the states are visited during the process lifetime, and there is no periodicity.

We can approximate the probability distribution of the transitions as follows:

$$P_A^{eq} = \frac{\#(x(t) = A)}{T}, \quad P_B^{eq} = \frac{\#(x(t) = B)}{T}$$

Which is *the number of times the process is in state A at time t divided by the total number of transitions T* (Same for B).

The **return time** is the time it takes for a process to return to a given state. For a state σ , we can define the return time T_σ as:

$$T_\sigma = \min\{t > 0 : x(t) = \sigma | x(0) = \sigma\}$$

The **mean return time** $\langle T_\sigma \rangle$ is the average time it takes for the process to return to state σ after leaving it. For an ergodic Markov chain, the mean return time is related to the equilibrium probability by:

$$\langle T_\sigma \rangle = \frac{1}{P_\sigma^\infty}$$

This makes intuitive sense: if a state has a high equilibrium probability, it will be visited frequently, leading to a short mean return time. Conversely, states with low equilibrium probabilities will have longer mean return times.

All the problems continuous in time but that involves a discrete state space can be solved using a Continuous Time Markov Chain (CTMC), which are a subset of all the stochastic processes with discrete state space and continuous time.

$$\Pr \{x(t) = \sigma \mid x(\theta), \quad \theta \in [0, t]\}$$

If we consider an infinitesimal time interval dt , we can write:

$$\Pr \{x(t + dt) = \alpha \mid x(t) = \sigma\}$$

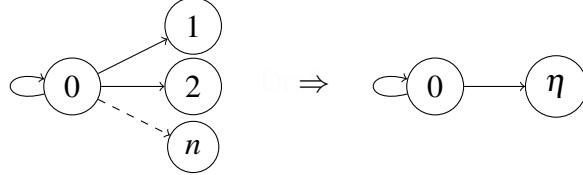
So the probability of eving more than one transition in the interval is:

$$\begin{aligned} \Pr \{\geq 2 \text{ events } (t, t + dt)\} &= 0 \\ \Pr \{1 \text{ event } (t, t + dt)\} &= O(dt) \\ \Pr \{0 \text{ events } (t, t + dt)\} &= 1 - O(dt) \simeq 1 \end{aligned}$$

We can define the **transition rate** as:

$$\Pr \{x(t + dt) = \alpha \mid x(t) = \sigma\} = W_{\alpha\sigma} dt$$

$W_{\alpha\sigma}$ is always positive and can be $>$ or $<$ than 1.



To compute the probability of remaining in the state 0, we can collapse all the other states into a single state η .

$$x(t) = \begin{bmatrix} S(t) \\ I(t) \end{bmatrix}$$

With: $(S, I) \rightarrow (S - 1, I + 1)$,

$$S(t + dt) = S(t) - 1, I(t + dt) = I(t) + 1$$

$$x(t + dt) = x(t) + (-1, 1)$$

...

$$\begin{cases} \Pr(Cont) = \Pr \left\{ x(t + dt) = x(t) + \binom{-1}{1} \right\} = \beta \frac{I(t)}{N} S(t) dt \\ \Pr(Rec) = \Pr \left\{ x(t + dt) = x(t) + \binom{0}{-1} \right\} = \gamma I(t) dt \end{cases} \Rightarrow \begin{cases} W_{Cont} = \beta \frac{I(t)}{N} S(t) \\ W_{Rec} = \gamma I(t) \end{cases}$$

...

$$W_{\eta\sigma} = \sum_{\alpha \in \mathcal{S} \setminus \{\sigma\}} W_{\alpha\sigma}$$



We have that:

$$\begin{cases} \Pr \{x(t + dt) = B | x(t) = A\} = W_{AB}dt \\ \Pr \{x(t + dt) = B | x(t) = B\} = 1 \end{cases}$$

$$P_A(t + dt) = (1 - W_{AB}dt)P_A(t) \Rightarrow P_A(t + dt) = P_A(t) - WP_A(t)dt$$

$$P'_A(t) = -WP_A(t)$$

$P_A(t) = P_A(0)e^{-Wt}$

Let's define T the time of remaining in the state A :

...

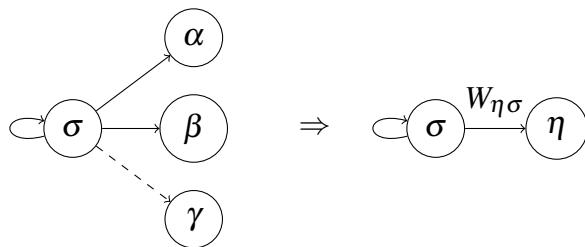
$$\mathcal{P}_I(T) = We^{-WT}$$

Since this is an exponential distribution, we can compute the mean return time as:

$$\langle T \rangle = \frac{1}{W}$$

...

—



If we want to compute the probability of remaining in the state σ for a time T , we can write:

$$P\{x(T + dt) \in \mathcal{S} \setminus \{\sigma\} | x(T) = \sigma\}$$

Also in this case we can collapse the states into a single state η .

$$W_{\eta\sigma} = \sum_{\alpha \neq \sigma} W_{\alpha\sigma} = W_{sum}$$

So we can say that:

$$\begin{cases} t_\eta & x(t_\eta) = \sigma \\ T \sim \exp(\lambda = W_{sum}) \end{cases} \Rightarrow t_{next} = t_\eta + T$$

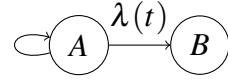
$$\Pr\{x(t_\eta + T + dt) = \alpha | x(t+T) = \sigma\} = W_{\alpha\sigma} \cdot 1 \cdot dt,$$

$$\Pr\{x(t_\eta + T + dt) = \beta | x(t+T) = \sigma\} = W_{\beta\sigma} \cdot 1 \cdot dt,$$

...

...

Final



$$P_A(t + dt) = (1 - \lambda(t)dt)P_A(t)$$

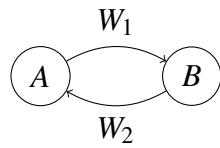
$$P'_A(t) = -\lambda(t)P_A(t) = P_A(t) = P_A(0)e^{-L(t)}, \quad L(t) = \int_0^t \lambda(s)ds$$

$$\mathcal{P}(T) = \lambda(T)e^{-L(t)}$$

$$T \sim \mathcal{P}(T) = \lambda(T)e^{-L(t)}$$

$$P(\sigma \rightarrow \alpha) = \frac{\lambda_\sigma(T)}{L_{sum}(t)}$$

—



$$\begin{cases} \dot{P}_1(t) = W_2 P_B(t) - W_1 + P_A(t) \\ \dot{P}_2(t) = W_1 P_A(t) - W_2 P_B(t) \end{cases} \Rightarrow \begin{cases} \dot{P}_A + \dot{P}_B = 0 \\ P_A + P_B = 1 \end{cases}$$

$$\begin{cases} \dot{P}_A = \sum(W_{Ay}P_y - W_{yA}P_A) \\ \sum_{A \in \mathcal{S}} P_A(t) = 1 \end{cases}$$

And we have that:

$$\dot{P}_A = -W_1 P_A + W_2 (1 - P_A) = W_2 - (W_1 + W_2)P_A$$

So, for an equilibrium we have that:

$$0 = W_2 - (W_1 + W_2)P_A^{eq}$$

$$P_A^{eq} = \frac{W_2}{W_1 + W_2}, \quad P_B^{eq} = \frac{W_1}{W_1 + W_2}$$

We can write our transition matrix:

$$W = \begin{bmatrix} -W_1 & W_2 \\ W_1 & -W_2 \end{bmatrix}$$

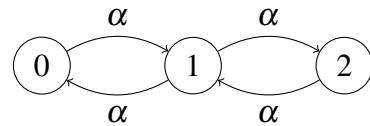
Which is singular, so one of the eigenvalue is 0.

...

$$\begin{cases} \dot{P}_\sigma = \sum (W_{\sigma y} P_y - W_{y\sigma} P_\sigma) \\ \dot{P} = A P \\ \sum_\sigma \dot{P}_\sigma(t) = 0 \end{cases}$$

$$0 = A p^{eq}$$

...



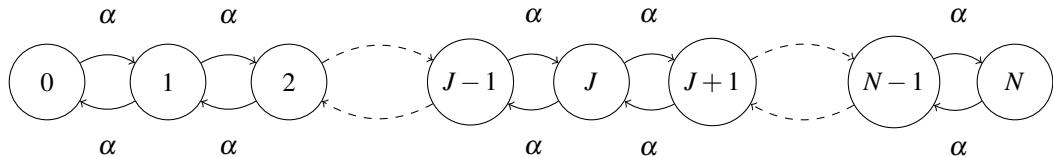
The steady state is given by:

$$\begin{aligned} \dot{P}_0 &= -\alpha P_0 + \alpha P_1 \\ \dot{P}_1 &= \alpha P_0 - 2\alpha P_1 + \alpha P_2 \\ \dot{P}_2 &= \alpha P_1 - \alpha P_2 \end{aligned}$$

...

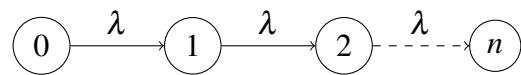
$$3\phi = 1 \Rightarrow \boxed{\phi = \frac{1}{3}}$$

we can now generalize this to a system with N states:



...

Let's consider an unidirectional chain, with a single transition rate λ :



$$\begin{aligned} \dot{P}_0 &= -\lambda P_0 \\ \dot{P}_1 &= \lambda P_0 - \lambda P_1 \\ &\dots \\ \dot{P}_\sigma &= \lambda P_\sigma - \lambda P_{\sigma-1} \end{aligned}$$

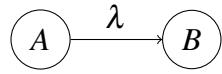
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The probability distribution $P_n(t) = \Pr\{N(t) = n\}$ has a symmetry property:

$$P_n(t) = \frac{(\lambda t)^n}{n!} e^{-\lambda t}$$

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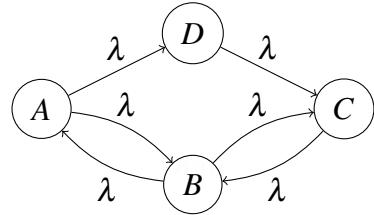
Lecture 19/05/2025



We have two properties:

$$\sum_{n=0}^{\infty} P_n(t) = 1$$

$$\lim_{t \rightarrow \infty} P_n(t) = 0 \quad \forall n \geq 0$$



To solve a general case, we need to solve the following system of equations:

$$\dot{P} = A(t)P$$

we have:

$$A_{ij} = \lambda(t)\alpha_{ij}$$

$$d\theta = \lambda(t)dt$$

$$\frac{d\theta}{dt} = \lambda(t)$$

$$\theta = \int_0^t \lambda(s)ds = L(t)$$

$$\frac{dP}{dt} = \lambda(t)\alpha P$$

So:

$$P(t) = e^{\lambda(t)\alpha} P(0)$$

—
Let's consider now a case where we have a periodic

$$A(t) = \begin{cases} A_1 & 0 < \text{mod}(t, T) < Q \\ A_2 & Q < \text{mod}(t, T) < T \end{cases}$$

$$t = 0 \Rightarrow P(0) = P_0$$

$$0 < t < Q \Rightarrow \dot{P} = A_1 P \Rightarrow P(t) = e^{A_1 t} P_0$$

$$Q < t < T \Rightarrow \dot{P} = A_2 P \Rightarrow P(t) = e^{A_2(t-Q)} P(Q)$$

$$T < t < T + Q \Rightarrow \dot{P} = A_1 P \Rightarrow P(t) = \underbrace{e^{A_2(T-Q)} e^{A_1 Q}}_{B(T,Q)} P(0)$$

so

$$P(T) = B(T, Q) P(0)$$

$$P(2T) = B^2(T, Q) P(0)$$

$$P(nT) = B^n(T, Q) P(0)$$

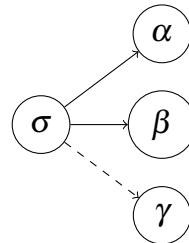
—
CTMC model:

$$f = \{\sigma_1, \dots, \sigma_n\}$$

$$\dot{P} = A(t) P$$

$$P(0) = \begin{bmatrix} 0 \\ \vdots \\ 1 \\ \vdots \\ 0 \end{bmatrix} = e_y$$

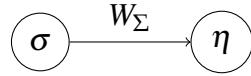
—
Gillespie algorithm:



$$X(0) = \sigma_0$$

We want to find when where will be the jump and which will be the next state.

Also in this case we can collapse the states in a single one; the transition rate is given by the sum of the transition rates of the original states.



So we can write:

$$\mathcal{P}(T) = W_\Sigma e^{-W_\Sigma T}$$

And we have:

$$t_{n+1} = t_n + T_n, \quad \text{where} \quad T_n \sim \mathcal{P}(T)$$

$$P(\sigma_i) = \Pr(Event(\sigma \rightarrow \sigma_i) \text{ara } t_n + T_n | oneE \sim t_n + T_n)$$

TODO: Che cazzo ha scritto il prof? ↑↑↑↑↑

$$\Pr\{\sigma \rightarrow \sigma_i \ (T, T + dt)\} = (W_x dt) e^{-W_\Sigma T} = \boxed{P(\sigma_i)(W_\Sigma e^{W_\Sigma T} dt)}$$

MISSING: qualcosa

...

1.

$$T_n \sim W_\Sigma e^{-W_\Sigma T} \Rightarrow t_{n+1} = t_n + T_n$$

2.

$$P(\sigma_i) = \frac{W_{\sigma_i}}{W_\Sigma} \quad t_n \rightarrow n$$

A first application of the Gillespie algorithm is the following:

Example: Gillespie Algorithm: Contagion and Recovery (SIR Model)

Consider a system with two possible events:

- **Contagion:** $(S, I) \rightarrow (S - 1, I + 1)$, with rate $W_{con} = \beta \frac{I}{N} S$
- **Recovery:** $(S, I) \rightarrow (S, I - 1)$, with rate $W_{rec} = \gamma I$

The Gillespie algorithm proceeds as follows:

1. **Initialize:** $t_0 = 0, x_0 = (S_0, I_0)$

2. **Compute the rates:**

- $W_{con} = \beta \frac{I_n}{N} S_n$
- $W_{rec} = \gamma I_n$
- $W_\Sigma = W_{con} + W_{rec}$

3. **Draw two random numbers** $U_1, U_2 \sim \mathcal{U}[0, 1]$

4. **Determine the time to the next event:**

The waiting time T_n is exponentially distributed:

$$T_n = \frac{-\ln U_1}{W_\Sigma}$$

Update the time: $t_{n+1} = t_n + T_n$

5. Determine which event occurs:

Compute the probabilities:

$$\Pr(\text{contagion}) = \frac{W_{con}}{W_\Sigma}, \quad \Pr(\text{recovery}) = \frac{W_{rec}}{W_\Sigma}$$

If $U_2 < \Pr(\text{contagion})$, a contagion event occurs; otherwise, a recovery event occurs.

6. Update the state:

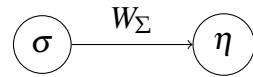
- If contagion: $(S_{n+1}, I_{n+1}) = (S_n - 1, I_n + 1)$
- If recovery: $(S_{n+1}, I_{n+1}) = (S_n, I_n - 1)$

7. Repeat from step 2 until $I_n = 0$ or another stopping criterion is met.

Note: The formula for T_n comes from inverting the CDF of the exponential distribution:

$$CDF(T) = 1 - e^{-W_\Sigma T} \implies T = \frac{-\ln(1 - U_1)}{W_\Sigma}$$

Since U_1 is uniformly distributed, so is $1 - U_1$, and it is common to write $T = \frac{-\ln U_1}{W_\Sigma}$.



$$X(t_n) = \sigma$$

$$\mathcal{P}(T_n) = W_\Sigma(t_n + T_n) e^{-\int_{t_n}^{t_n+T_n} W_\Sigma(z) dz}$$

$$CDF(T_n) = 1 - e^{-\int_{t_n}^{t_n+T_n} W_\Sigma(z) dz}$$

...
calling $\Psi(T) = \int_{t_n}^{t_n+T_n} W_\Sigma(z) dz$

$$1 - e^{-\Psi(T_n)} = U_n \quad \Rightarrow \quad e^{-\Psi(T_n)} = 1 - U_n \quad \Rightarrow \quad \Psi(T_n) = -\ln(1 - U_n)$$

...

$$\beta(t) = \beta_n(1 + \delta \cos(\omega t))$$

$$x(t_n) = (S_n, I_n)$$

$$\int_{t_n}^{t_n+T_n} W_\Sigma(z) dz = -\ln(1 - U_n)$$

$$\int_{t_n}^{t_n+T_n} \left[\gamma I_n + \beta_n \frac{I_n}{N} S_n (1 + \delta \cos(\omega z)) \right] dz = -\ln(1 - U_n)$$

$$\left[\gamma I_n + \beta_n \frac{I_n}{N} S_n \right] T_n + \beta_n \frac{I_n}{N} S_n \frac{1}{\omega} [\sin(t_n + T_n) - \sin(t_n)] = -\ln(1 - U_n)$$

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Lecture 26/05/2025

MISSING: everything

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Lecture 30/05/2025

13.1 Random Walks

Let's consider a particle that moves on a line.

$$\rho(x,t) = \Pr[\text{Particle is at location } (x,x+dx) \text{ at time } t]$$

we have:

$$\frac{\partial \rho}{\partial t} = \int [\rho(y,t)\Omega(y,x) - \rho(x,t)\Omega(x,y)] dy$$

Let's consider now the probability that the particle just jumped in the location $(x,x+dx)$ at time t :

$$p(x,t) = \Pr[\text{Particle just jumped in the location } (x,x+dx) \text{ at time } t]$$

We define a jump weight as:

$$\eta(y)$$

so:

$$p(x,t) = \int_0^t \int_{-\infty}^{+\infty} dy \eta(y) p(x-y, t-T) \omega(t) dt + \delta(x)\delta(t)$$

Definition: Fourier and Laplace Transforms

- **Fourier transform**

Let's recall the Fourier transform:

$$\mathcal{F}_x[f(x)] = \int_{-\infty}^{+\infty} f(x) e^{-ikx} dx$$

we have:

$$\mathcal{F}_x[f(x) \star g(x)] = \hat{f}(k)\hat{g}(k)$$

- **Laplace transform**

Let's recall the Laplace transform:

$$\mathcal{L}_x[f(x)] = \int_0^{\infty} f(x) e^{-st} dt$$

we have:

$$\mathcal{L}_x[f(t) \star g(t)] = \tilde{f}(s)\tilde{g}(s)$$

$$p(x,t) = [p(x,t) \star_s \eta(x)] \star_t \omega(t)$$

Let's call $\bar{p}(k, s) = \hat{p}(x, t)$, then we can write:

$$\bar{p}(k, s) = \hat{\eta}(k)\tilde{\omega}(s)\bar{p}(k, s) + 1 \quad \Rightarrow \quad \bar{p}(k, s)\{1 - \hat{\eta}(k)\tilde{\omega}(s)\} = 1$$

$$\bar{p}(k, s) = \frac{1}{1 - \hat{\eta}(k)\tilde{\omega}(s)}$$

$$\Psi(T) = \int_T^{+\infty} \omega(q) dq = 1 - \int_0^T \omega(q) dq$$

and we can rewrite $\rho(x, t)$ as:

$$\rho(x, t) = \int_0^t p(x, t-T) \Psi(T) dT$$

The Laplace transform of $\Psi(T)$ is:

$$\begin{aligned} \mathcal{L}[\Psi(T)] &= \mathcal{L}[1 - \int_0^T \omega(q) dq] = \frac{1}{s} - \frac{1}{s} \tilde{\omega}(s) \\ \bar{\rho}(k, s) &= \left(\frac{1 - \tilde{\omega}(s)}{s} \right) \left(\frac{1}{1 - \hat{\eta}(k)\tilde{\omega}(s)} \right) \end{aligned}$$

Cauchy distribution (????)

$$\begin{aligned} \omega(T) &\sim \frac{1}{T^{1+\alpha}} \quad T \gg 1 \\ \omega(T) &= \frac{A}{1+B T^2} \quad T \gg 1 \quad \Rightarrow \quad \omega(T) \sim \frac{A}{B} \frac{1}{T^2} \end{aligned}$$

Then its integral is infinite (cause compares a logarithm).

$$\frac{\partial \rho}{\partial t} = D \frac{\partial^2 \rho}{\partial x^2} \quad \dot{x} = c\xi(t)$$

$$\langle x^2(t) \rangle \equiv t^\alpha$$

for $0 < \alpha < 1$ we have subdiffusion, for $1 < \alpha < 2$ we have superdiffusion.