

Numerical Integration of Differential Equations: The Damped Harmonic Oscillator

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Second Year Computational Physics Report

April 2020

Abstract

An analysis of four numerical integration methods has been conducted. It was done with a purpose of determining the method with the highest accuracy. The Verlet method was found to be most accurate for any time step used. The three other methods used were Euler's method, Improved Euler's method and Euler-Cromer method. The Verlet method and the Euler-Cromer method were found to be symplectic. The relationship between the size of a time step and the accuracy has been observed to be inversely proportional. Secondly the investigation into the effects of heavy damping onto oscillation has been conducted. This was done utilising the Verlet method. Lastly the investigation of an application of a constant and sinusoidal force onto an oscillator was conducted. This naturally led onto the investigation of resonance for a forced oscillator and the effect of the value of the damping coefficient on the resonance curve. It was found that heavy damping eliminates resonance.

1. Introduction

Simple harmonic motion (SHM) is a type of motion that occurs in every area of physics where a small oscillation around a potential well occurs. This means that this type of motion is present in nearly every field of physics due to the presence of energy potentials. The applications of SHM can range from the description of a quantum particle to the behaviour of a pendulum for small displacements [1].

Simple harmonic motion is specified as a motion of a particle of mass m that is subject to a force that depends on its displacement. Hence the characterising property of simple harmonic motion is that

$$a \propto -x. \quad \text{Equation 1}$$

For a being acceleration of a particle and x being the displacement from the point of equilibrium. It is a motion in one dimension along the line with the position x being the magnitude of the displacement from the equilibrium position. The acceleration of the particle a is obtained from the application of Newton's second law and the addition of vector forces that act on the particle. This system is going to be the baseline on which the investigation into numerical methods is going to be conducted. The method of an iterative integrator is going to be explored. The effects of its predictions on the energy of the system are going to be utilized to judge the accuracy of the methods.

2. Theory

If you consider a particle in a spring-mass system with a damping force being dependent on the velocity of the object it is easy to derive that the equation of motion [2] is given by

$$m\ddot{x} + b\dot{x} + kx = 0. \quad \text{Equation 2}$$

This equation of motion represents a damped simple harmonic oscillator with m being the mass of the particle, b being the damping coefficient and k being the spring constant. The single dot represents a derivative with respect to time. This equation has 3 analytic solutions depending on the values of the constants. Those are: heavily damped oscillator, critically damped oscillator and lightly damped oscillator. A fourth solution exists if the damping constant equals 0. That solution is just the solution of a simple harmonic oscillator given by

$$x(t) = A\sin(\omega_0 t) + B\cos(\omega_0 t) \quad \text{Equation 3}$$

with ω_0 being the natural angular frequency of the oscillation which equals to $\sqrt{\frac{k}{m}}$. This solution if multiplied by a damping term $e^{-\frac{b}{2m}t}$ gives the damped oscillator solution. However, the angular frequency is given by

$$\sqrt{-\frac{k}{m} + \frac{b^2}{4m^2}}. \quad \text{Equation 4}$$

The solutions of the damped simple harmonic oscillator depends on the value of $\alpha = -\frac{k}{m} + \frac{b^2}{4m^2}$. If $\alpha > 0$ the solution is heavily damped, but if $\alpha < 0$ the solution is lightly damped. However, if $\alpha = 0$ then the solution is critically damped and reaches the equilibrium point the fastest without any oscillation.

Due to the use of iteration it should be said that the acceleration of a particle at any point of the iteration is given by Equation 2 in the form

$$a_n = -\frac{b}{m}v_n - \frac{k}{m}x_n. \quad \text{Equation 5}$$

Furthermore, the energy of the particle at any given velocity and position is obtained by

$$E_n = \frac{1}{2}mv_n^2 + \frac{1}{2}kx_n^2. \quad \text{Equation 6}$$

The approximation described in section 3 is utilised during the investigation into a constant force being applied for a given time period. This approximation is given by

$$x_{n+1} = Ax_n + Bx_{n-1} + \frac{A' \sin(\omega' nh)}{m} h^2. \quad \text{Equation 7}$$

For ω' being the frequency of the force acting on the particle and A' being its amplitude. For the case of a constant force the sinusoidal force expression is replaced with a constant.

2.1 Euler's method

The first numerical method is Euler's method which is defined by

$$x_{n+1} = x_n + hv_n \text{ and} \quad \text{Equation 8}$$

$$v_{n+1} = v_n + ha_n. \quad \text{Equation 9}$$

The variable h (in seconds) is the time step chosen for the simulation. [3]

This method was found to not be a symplectic integrator. This can be confirmed by the application of Equation 7 and 8 with Equation 9. This obtains

$$E_{n+1} = E_n(1 + \frac{kh^2}{m}). \quad \text{Equation 10}$$

This shows that the energy of the oscillator increases with each step. This gives an error term to the energy of $O(h^2)$.

2.2 Improved Euler's method

The second numerical method is an improvement upon Euler's method. It is defined by

$$x_{n+1} = x_n + hv_n + \frac{1}{2}h^2a_n \text{ and} \quad \text{Equation 11}$$

$$v_{n+1} = v_n + ha_n. \quad \text{Equation 12}$$

This method also doesn't conserve energy and the energy error is given by

$$E_{n+1} = E_n + O(h). \quad \text{Equation 13}$$

With $O(h)$ being a polynomial with the dominant term being h (for h being small). This means that the energy will also be increasing as in the previous method. [4][5]

2.3 Euler-Cromer method

The third numerical method is a symplectic integrator which conserves average energy over the whole time period. It is defined by

$$x_{n+1} = x_n + hv_{n+1} \text{ and} \quad \text{Equation 14}$$

$$v_{n+1} = v_n - \frac{kh}{m}x_n. \quad \text{Equation 15}$$

Expanding the expression for energy we obtain

$$E_{n+1} = E_n - \frac{1}{2}h^2 \left(\frac{k^2x_n^2}{m} - kv_n^2 \right) - h^3 \left(\frac{k^2x_nv_n}{m} \right) + h^4 \left(\frac{k^3x_n^2}{2m^2} \right). \quad \text{Equation 16}$$

The second term averages out to be zero over the complete cycle what allows the energy to be conserved. [6]

2.4 Verlet's method

The last numerical integrator is also a symplectic integrator which provides an accurate projectile path with not much more computational cost than Euler's method [7]. It is defined by

$$x_{n+1} = Ax_n + Bx_{n-1}. \quad \text{Equation 17}$$

With $A = \frac{2(2m - kh^2)}{D}$ and $B = \frac{bh - 2m}{D}$. With $D = 2m + bh$. All the constants have the previously defined meaning. This method requires two initial positions. The value of x_1 is obtained utilising a Taylor expansion up to the second order.

$$\begin{aligned} x_1 &\cong x_0 + hv_0 + \frac{1}{2}a_0h^2 \\ &\cong x(h) + O(h^3) \end{aligned} \quad \text{Equation 18}$$

This allows for the computation of the positional terms. The velocity of the oscillator is computed utilising the Verlet-Störmer method along with the mean value theorem. The velocities are obtained by

$$v_n = \frac{x_{n+1} - x_{n-1}}{2h} + O(h^2). \quad \text{Equation 19}$$

This means that the error of the velocity utilising this method is given by the h^2 term in the expansion. The velocity is defined utilising an average. It means that the influence of an external force can be evaluated using Equation 7. It is later seen in section 4. [8]

The critical value of the damping factor is given by

$$b_{critical} = 2\sqrt{km} \quad \text{Equation 20}$$

3. Method

The values of the constants were $k = 0.93 \text{ Nm}^{-1}$ and $m = 5.44 \text{ kg}$. The initial position was chosen to be 1 m with the initial velocity of -1 ms^{-1} . No applied force was used for the first stage of the investigation. The methods have been compared utilising a maximum time of 100 seconds with the damping coefficient being equal to 0. This allows for the observation of the system at constant energy. The value of the time step was varied between 1 second to 0.001 seconds. This allowed for the observation of the relation between the magnitude of the time step and the accuracy of the simulation. The methods were compared visually along with the use of the amount of "fictitious energy" that each simulation generated. These are summarised in Figure 3. Those were then graphed to show the relation between the error and time. [9]

The program possesses an ability to write and load simulations. It does so by the means of a json file format and the imported json library. This was done to assure that the data is readable and can be utilised in different scripts. Furthermore, for the ease of data manipulation an object-oriented approach has been utilised. This was done to ensure maximum easiness in modifying the script later. This was done at the cost to memory usage that this script requires to run. This was addressed by resetting the variables after the results are obtained. This didn't cause any problems on the machine it was tested but it may cause problems on a weaker system.

The investigation into the behaviour of a constant push force has been done with the use of Equation 7. The position of the oscillator was then plotted against time. The resonance curve was obtained using the same function. Instead of fixing the time during which the sinusoidal force acts it was altered to be the entirety of the running time. The amplitude of the oscillation was then found and plotted against the frequency of the force. The natural frequency of the oscillator was found to be 0.41 rads^{-1} . This gave the time period of 15.2 s . The constant force was hence applied at

45.6 s, 57.0 s, 53.2 s. This corresponds to 3, 3.75, and 3.50 of the periods respectively. At these times a constant force of 2 N was applied. Meanwhile the sinusoidal force was investigated at the frequency of 0.065 Hz and 1 Hz with the magnitude of 2 N.

The effects of the alteration of the damping coefficient were tested by using the Verlet method to calculate the position, velocity and energy. Plotting them as functions of time was done with the goal of visualising the behaviour. The damping term values investigated were multiples of the critical damping coefficient obtained by Equation 20. The mass, spring constant and initial conditions were used as before.

4. Analysis

4.1 Comparison of the numerical methods

Figure 1 shows that the Verlet methods gives the best results. This can be further confirmed by the fact that the error in energy generated over the whole runtime was $< 1 J$. Meanwhile, the second lowest was the Euler-Cromer method which obtained $> 35 J$ for all the values tested. it is important to mention that the Improved Euler's method provided a small error in the range

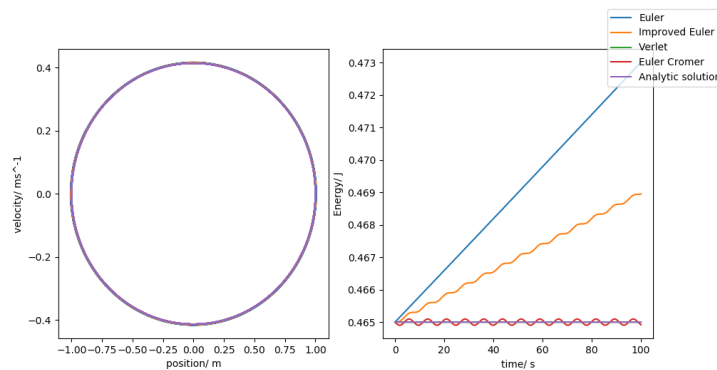


Figure 1 – The phase plot of the oscillation (left) and the energy as a function of time (right). This plot was obtained with the time-step being 0.001 s. The phase plots show the paths obtained from all methods being drawn over each other.

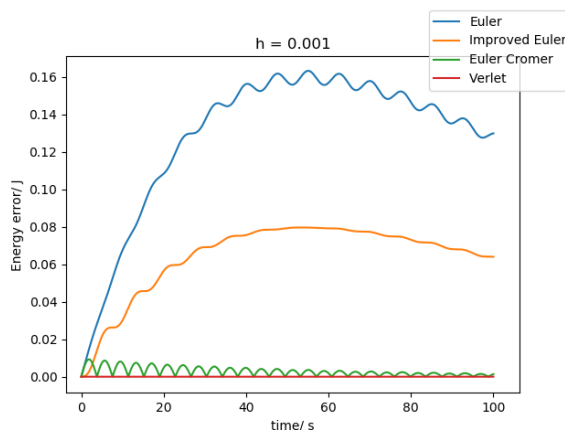


Figure 2 – Energy error curve. This plot was obtained by subtracting the analytical energy values from the model energy values with the damping coefficient $b = 0 \text{ kg s}^{-1}$.

of the damping coefficient passing the critical value. The inaccuracy in the Euler-Cromer method comes from the fact that its energy needs to be averaged over a complete cycle while the Verlet method doesn't. The testing of simulations running over complete cycles wasn't tested due to the desire to ensure fairness of the comparison method. The remaining two are not symplectic methods and it can be visually seen that they are not as accurate as the Verlet method in Figure 1 or 2. This means that the Verlet method was utilised for the rest of the investigation with modifications described in section 3.

Then the test was conducted for a damped oscillator. It consolidated that the Verlet method has the smallest error out of all the numerical integrators. As evidence by the error being barely visible on the graph. As visible in Figure 1 the error for the Verlet integrator is very close to zero for the entirety of the simulation while other integrators either increase or oscillate around the value. In Figure 2 the energy error curve is the least prominent for the Verlet method under damping as well. The time step value was chosen to be 0.001 s due to it being

the best balance between the time taken and the accuracy of the computation. The accuracy increases with the decrease in the step size. This was more apparent for both Euler's methods due to those two methods having the smallest degree of accuracy. It is important to note that all the methods perform very well when the value of the damping coefficient is large. Verlet's method produced errors which were not exponential in contrast with Euler's method. Verlet's method was also the one that presented the error not increasing exponentially with time. This is very important for long term simulations. Even despite only a second order Taylor expansion used in Equation 18 the accuracy was much higher than anticipated.

B ($kg s^{-1}$)	Euler (J)	Improved Euler (J)	Euler – Cromer (J)	Verlet (J)
0	42107.75	20448.09	609.92	0.98
0.1	13530.68	6591.90	279.68	0.48
1	225.95	94.98	41.51	0.08
2	47.52	7.78	34.82	0.07
3	23.80	11.49	34.84	0.08

Figure 3 – The table showing the error in the total energy for a given set of damping constants. The values are rounded to two decimal places.

4.2 Unforced Oscillations

A step size of 0.001 s and a maximum time of 100 s were used to model the unforced simple harmonic oscillator with damping terms half, double and equal to the critical damping coefficient. This was obtained from Equation 20. The value was calculated to be 4.5 kg s^{-1} . The energy plots show the expected behaviour. The critically damped oscillator returned to the equilibrium position in the shortest amount of time. The

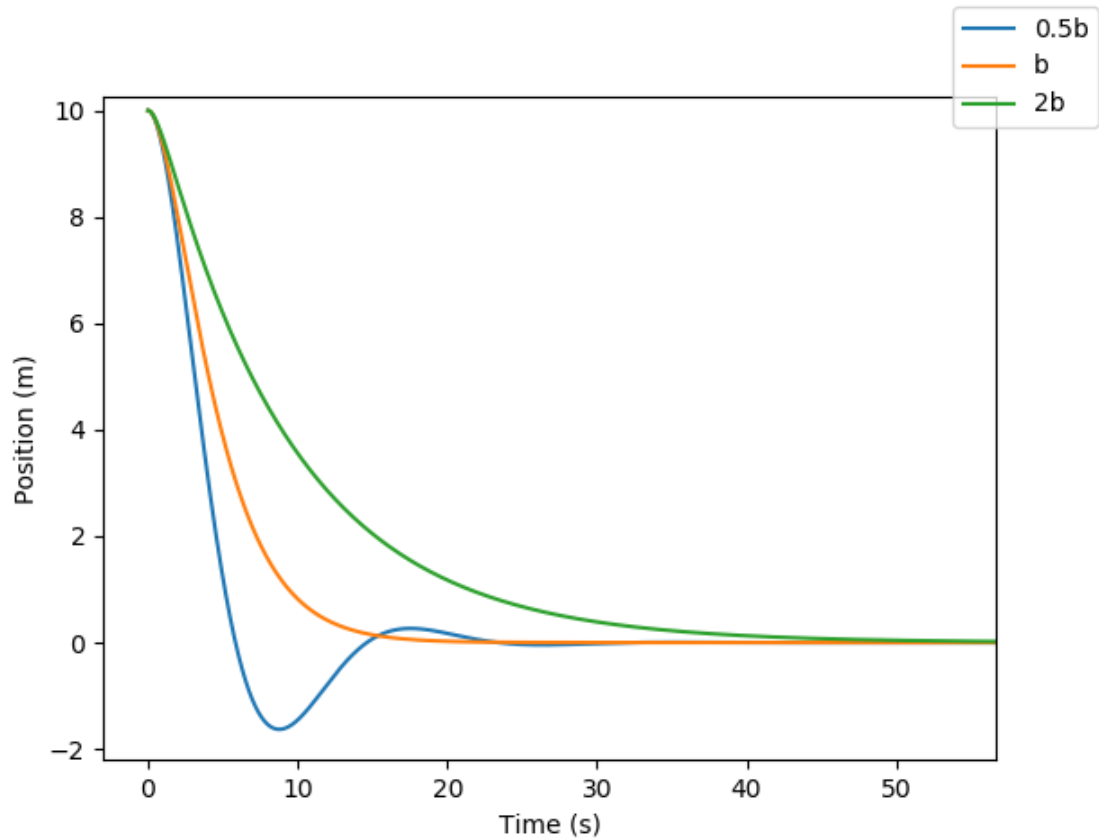


Figure 4 – The position against time plot for the critically damped oscillation with the value of the damping coefficient being half, equal to and double the critical value.

heavily damped oscillator took longer while the lightly damped oscillator moved past the equilibrium point twice. This can be observed in Figure 4. The energy was calculated using Equation 6. The energy decreased the fastest for the critically damped oscillator. This can be observed in Figure 5.

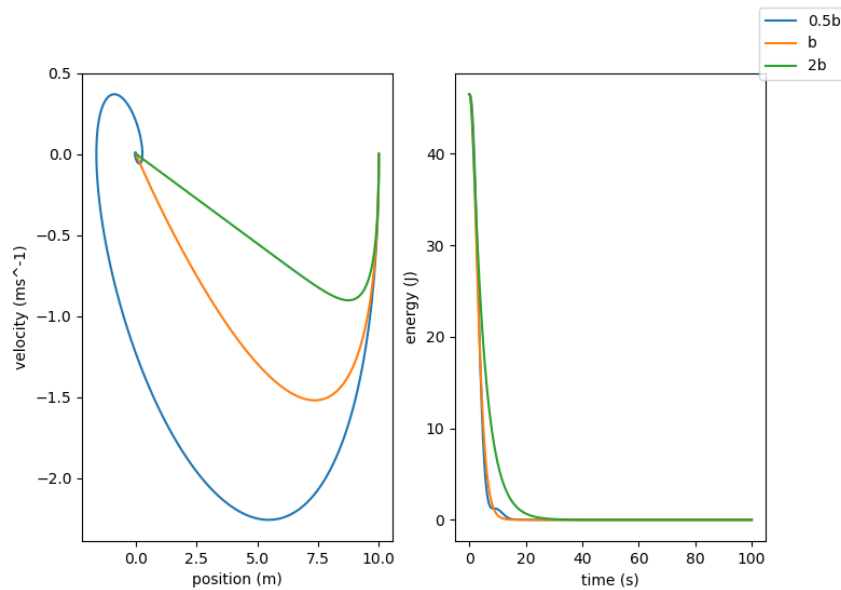


Figure 5 – The phase plot (left) and the energy plot (right) were obtained for the overdamped oscillator.

All the calculations were calculated utilising the Verlet method as it was chosen to be the most accurate as explained in section 4.1. The spiral behaviour presented in Figure 5 can be attributed to the damping forces present in the system arising from the damping coefficient.

4.3 Instantaneously Forced Oscillations

The effects of an application of a constant force and a sinusoidal force were investigated. In the case of a constant force the time at which it was applied determined the response. All the oscillations returned to a behaviour explained in section 4.1 after a short transition period. The amplitude changed after the force was applied. When the

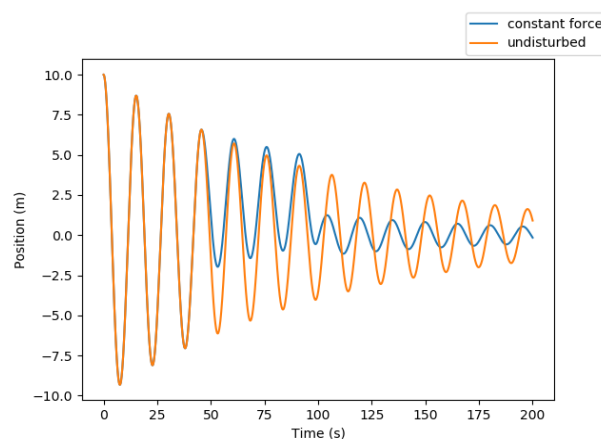


Figure 6 – The position against time plot shows the change in the oscillation after the constant force was being applied. This force was applied at the equilibrium point.

the velocity the amplitude increased. Furthermore, the opposite was true as well. The magnitude of the force was found to be not as dominant as expected. Towards the upper values of the force the oscillation still occurred however, reached a rest point

force was applied at the equilibrium point the amplitude changed the most as can be seen in Figure 6. For the half-cycle and quarter-cycle force applications the effect was weaker, but the amplitude still decreased. Hence the plots were not included.

The sinusoidal force caused an increase in the amplitude around the frequency of the force being equal to the natural frequency of the oscillation. This can be seen in Figure 7. This increase in amplitude occurred only around the natural frequency. If at the moment of the force being applied the direction of the force was in the same direction as

much faster. This force had the same effect as shifting the equilibrium point. This is analogous to a suspended spring system where gravity shifts the equilibrium point.

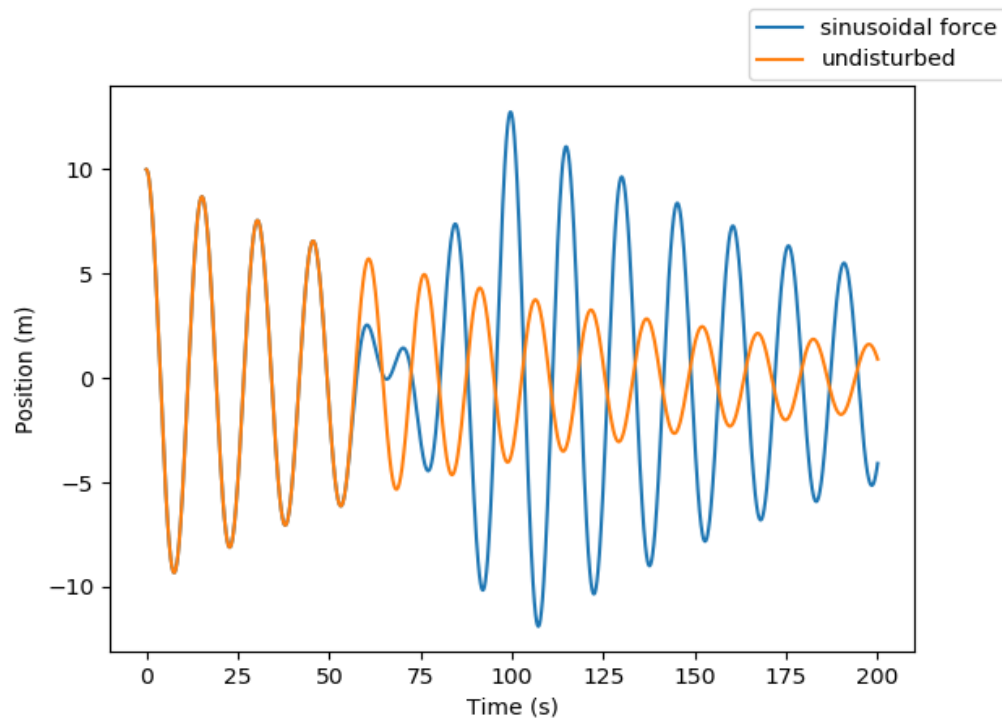


Figure 7 – The position against time plot for a forced oscillator. The force used is applied at the natural frequency of the oscillator hence its amplitude is increased. This is due to resonance.

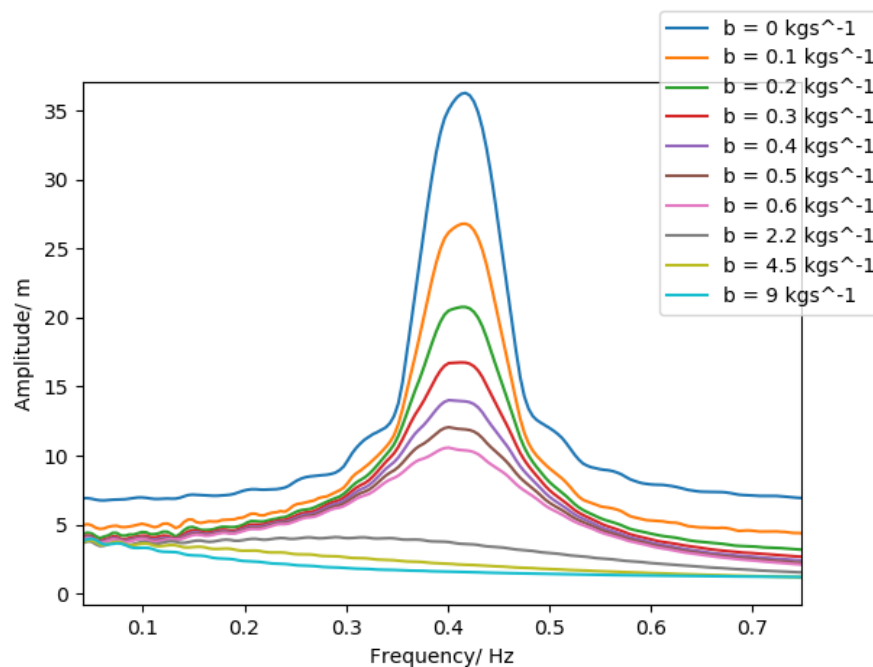


Figure 8 – The collection of resonance plots for different damping coefficients. The time step of 0.01 s was used with the total running time of 200 s.

Resonance was observed at the angular frequency calculated in section 3. It was characterised by a sharp increase in the amplitude of the oscillation at the natural

frequency. A resonance curve with different damping terms is shown in Figure 8. The value of the damping term increasing caused the lowering of the resonance peak. At the heavy/critical damping the damping coefficient was too high to allow for the resonance to occur. This meant that the motion dissipated too quickly for any oscillation to occur.

5. Error Analysis

The main sources of error in this simulation arose from the discretisation error, truncation error, initial conditions and the net growth in the error as the simulation progressed. The discretisation error arises from the fact that the time step at which values are evaluated isn't infinitesimal but of a fixed value. This was mitigated by the utilisation of a very small step size. The data obtained could have been developed at higher accuracies, but the time constraints didn't allow for this. The fact that the errors grow as the simulation progresses is since iterations utilise previous values to work out the path of the oscillator. This was mitigated by ensuring that the terms that were omitted, like in the case of a Taylor expansion, were powers of the time step allowing the convergence of the method to happen with minimal error present. The truncation error increased as the step interval decreased.

6. Conclusion

The models were utilised and compared as described in section 4.1. The Verlet method was found to be the most accurate for modelling forced and unforced oscillation. The investigation into the effects of the step size onto the accuracy has been conducted. It has shown the inverse proportionality between the step size and the accuracy. This was especially apparent for the less accurate Euler's method. It can be deduced that the smaller the size of the interval the better the prediction. It is advised to always use the smallest step size possible. However, it is apparent that the decrease of the time interval gives diminishing returns after passing the 10^{-4} s size.

The application of a constant force onto an oscillator produced a transition period after which the oscillation resumed with a shift in phase and amplitude. The sinusoidal force has been investigated with the use of a resonance plot. It was confirmed that the sharp amplitude increase occurs at the frequency close to the natural frequency of the oscillator. Lastly the Verlet method was utilised to observe the behaviour of the oscillator in the case of being heavily damped and critically damped (as described in section 4.2). It was also confirmed that resonance did not occur when the oscillator was heavily or critically damped.

7. References

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8. Appendix - Code

The history of changes can be found in the GitHub repository [10].

```
# -*- coding: utf-8 -*-
"""
Author: Tomasz Neska
Date: 06/03/2020
Description: Project 2 - utilising three different methods it evaluates the effects of time
step and accuracy on the
behaviour of the simple harmonic oscillator
"""

# initialisation
import string
from math import *
import numpy as np
import matplotlib.pyplot as plt
import random
import time
import cmath
import json
from scipy import optimize

plt.rcParams.update({'font.size': 14})
plt.style.use('default')
figure = plt.figure()
plt.rcParams.update({'errorbar.capsize': 2})

class SHO(object):
    def __init__(self, time_step, max_time, b=0.01, m=1.0, k=1.0, init_x=0.0, init_v=0.0,
fileNameSave="data.txt",
                fileNameLoad="data.txt"):
        """
        parameter name          type          description
        :param time_step:       float          the time step used in
calculations
        :param max_time:        float          the time the simulation lasts
        :param b:               float          damping coefficient
        :param m:               float          mass of the oscillator
        :param k:               float          spring constant
        :param init_x:          float          initial position
        :param init_v:          float          initial velocity
        :param fileNameSave:    string         filename used for saving data
        :param fileNameLoad:    string         filename used for loading data
        self.no_steps           float          number of iterations
        self.natural_angular_frequency float    the natural frequency of the
oscillation
        self.gamma              float          the damping constant
        self.quality_factor      float          quality factor
        self.analytic_series_pos array         the position of the analytical
solution
        self.analytic_series_vel array         the velocity of the analytical
solution
        self.analytic_energy     array         the array containing energy
data of the analytical solution
        self.__coefficients     array         coefficients used for the

```

<i>analytical solution</i>	<code>self.b_critical</code>	<i>float</i>	<i>the critical damping constant</i>
	<code>self.Euler_data</code>	<i>array</i>	<i>the data from Euler's method</i>
	<code>self.B_Euler_data</code>	<i>array</i>	<i>the data from improved Euler's</i>
<i>method</i>	<code>self.Verlet</code>	<i>array</i>	<i>the data from Verlet's method</i>
	<code>self.Euler_Cromer_data</code>	<i>array</i>	<i>the data from Euler-Cromer</i>
<i>method</i>	<code>self.analytical_data</code>	<i>array</i>	<i>variable for storing</i>
<i>analytical data</i>	<code>self.time</code>	<i>array</i>	<i>the time array used in all</i>
<i>simulations</i>	<code>self.disturbed_Verlet_data</code>	<i>array</i>	<i>the data of the Verlet method</i>
<i>with force applied</i>			

```

'''
self.fileNameSave = fileNameSave
self.fileNameLoad = fileNameLoad
self.b = b
self.m = m
self.k = k
self.h = time_step
self.init_v = init_v
self.init_x = init_x
self.no_steps = int(np rint(max_time / time_step))
self.natural_angular_frequency = np.sqrt(self.k / self.m)
if self.b != 0:
    self.gamma = self.b / self.m
    self.quality_factor = self.natural_angular_frequency / self.gamma
self.analytic_series_pos = []
self.analytic_series_vel = []
self.analytic_energy = []
self.__coefficients = []
self.solver()
self.analytic_solution()
self.data = []
self.b_critical = 2 * np.sqrt(self.k * self.m)
# data variables
self.Euler_data = []
self.B_Euler_data = []
self.Verlet_data = []
self.Euler_Cromer_data = []
self.analytical_data = []
self.time = np.array(range(0, self.no_steps, 1)) * self.h
self.disturbed_Verlet_data = []

def runSimulation(self):
'''
Runs the integrators as a single function
'''
self.Euler_integrator()
self.Better_Euler_integrator()
self.Verlet_integrator()
self.Euler_Cromer_integrator()
print("Simulation has been executed")

def getCoefficients(self):
# simple get function
return self.__coefficients

def Euler_integrator(self):
'''
parameter name          type          description
position_series          array          stores the position
temporarily
velocity_series          array          stores the velocity
temporarily
v_n                      float          stores the nth velocity term
x_n                      float          stores the nth position term
a_n                      float          stores the nth acceleration
term
'''
position_series = [self.init_x]
velocity_series = [self.init_v]
for counter in range(1, self.no_steps, 1):

```

```

v_n = velocity_series[len(velocity_series) - 1]
x_n = position_series[len(position_series) - 1]
a_n = (-self.b / self.m) * v_n + (-self.k / self.m) * x_n

position_series.append(x_n + self.h * v_n)
velocity_series.append(v_n + self.h * a_n)

self.Euler_data = [position_series, velocity_series,
self.energy_function(position_series, velocity_series)]

def Better_Euler_integrator(self):
    '''
        parameter name            type            description
        position_series            array            stores the position
    temporarily
        velocity_series            array            stores the velocity
    temporarily
        v_n                        float            stores the nth velocity term
        x_n                        float            stores the nth position term
        a_n                        float            stores the nth acceleration
    term
    '''
    position_series = [self.init_x]
    velocity_series = [self.init_v]
    for counter in range(1, self.no_steps, 1):
        v_n = velocity_series[len(velocity_series) - 1]
        x_n = position_series[len(position_series) - 1]
        a_0 = (-self.b / self.m) * v_n + (-self.k / self.m) * x_n

        position_series.append(x_n + self.h * v_n + 0.5 * self.h ** 2 * a_0)
        velocity_series.append(v_n + self.h * a_0)

    self.B_Euler_data = [position_series, velocity_series,
self.energy_function(position_series, velocity_series)]

def Euler_Cromer_integrator(self):
    '''
        parameter name            type            description
        position_series            array            stores the position
    temporarily
        velocity_series            array            stores the velocity
    temporarily
        v_n                        float            stores the nth velocity term
        x_n                        float            stores the nth position term
        a_n                        float            stores the nth acceleration
    term
        temp                        float            temporary variable - stores
    the v_n+1 term of the velocity
    '''
    position_series = [self.init_x]
    velocity_series = [self.init_v]
    for counter in range(1, self.no_steps, 1):
        v_n = velocity_series[len(velocity_series) - 1]
        x_n = position_series[len(position_series) - 1]

        a_0 = (-self.b / self.m) * v_n + (-self.k / self.m) * x_n

        temp = v_n + self.h * a_0 # v_n+1
        velocity_series.append(temp)
        position_series.append(x_n + self.h * temp)

    self.Euler_Cromer_data = [position_series, velocity_series,
self.energy_function(position_series, velocity_series)]

def Verlet_integrator(self):
    '''
        parameter name            type            description
        position_series            array            stores the position
    temporarily
        velocity_series            array            stores the velocity
    temporarily
        v_n                        float            stores the nth velocity term
        x_n                        float            stores the nth position term
        a_n                        float            stores the nth acceleration
    '''

```

```

term
    x_1                                float                stores the second position of
the oscillation
    D                                  float                temporary variable for ease of
calculation
    B                                  float                temporary variable for ease of
calculation
    A                                  float                temporary variable for ease of
calculation
    '''
    position_series = [self.init_x]
    velocity_series = [self.init_v]

    D = 2 * self.m + self.b * self.h
    B = ((self.b * self.h) - (2 * self.m)) / D
    A = 2 * (2 * self.m - (self.k * self.h ** 2)) / D

    a_0 = (-self.b / self.m) * self.init_v + (-self.k / self.m) * self.init_x
    x_1 = self.init_x + self.init_v * self.h + 0.5 * a_0 * self.h ** 2 # obtained using a
Taylor expansion of order 2
    position_series.append(x_1)

    for counter in range(1, self.no_steps, 1):
        position_series.append(A * position_series[counter] + B * position_series[counter
- 1])

    # calculating velocities using an approximation of O(h^2)
    # the velocity is estimated using the mean value theorem
    for counter in range(1, self.no_steps, 1):
        velocity_series.append(
            (position_series[counter + 1] - position_series[counter - 1]) / (2 * self.h))
# +O(h^2)
    position_series = position_series[:len(position_series) - 1]

    self.Verlet_data = [position_series, velocity_series,
self.energy_function(position_series, velocity_series)]

def energy_function(self, position, velocity):
    '''
    parameter name                type                description
    temp_pos                      numpy array          stores the position array
    temp_vel                      numpy array          stores the velocity array
    :return: the array containing energy values
    '''
    temp_pos = np.array(position)
    temp_vel = np.array(velocity)
    return 0.5 * self.m * temp_vel ** 2 + 0.5 * self.k * temp_pos ** 2

def convert_array(self, array):
    # operates on 1 dimensional arrays
    '''
    parameter name                type                description
    temp                          numpy array          the array holding the array
being converting
    :param array:
    :return: converted array
    '''
    temp = []
    for entry in array:
        temp.append(entry)

    return temp

def analytic_solution(self):
    # creates the analytic solution position series
    '''
    parameter name                type                description
    t_0                          float                the time that the simulation
is at
    '''
    t_0 = 0
    for counter in range(0, self.no_steps, 1):
        self.analytic_series_pos.append(self.ana_position(t_0))
        self.analytic_series_vel.append(self.ana_velocity(t_0))

```

```

        t_0 += self.h
        print("solution found")
        self.analytic_energy = self.energy_function(self.analytic_series_pos,
self.analytic_series_vel)

    def solver(self):
        """
        parameter name            type            description
        A                          float            function coefficient
        B                          float            function coefficient
        marker                      int              the marker indicating the type
of a solution
        p                          float            function coefficient
        q                          float            function coefficient
        K                          float            function coefficient
        """

        A = 0
        B = 0
        temp = (self.b ** 2 / (4 * self.m ** 2))
        if self.b == 0:
            marker = 1
            print("The analytic solution is a simple harmonic motion")
            omega = np.sqrt(self.k / self.m)
            A = self.init_v / omega
            B = self.init_x
            self.__coefficients = [omega, 0, marker, A, B]
        elif (self.k / self.m) > temp: # imaginary
            print("The solution is a lightly damped oscillation")
            marker = 3
            p = -1 * self.b / (2 * self.m)
            q = np.sqrt((self.k / self.m) - self.b ** 2 / (4 * self.m ** 2))
            # initial conditions
            A = (- self.init_x * p + self.init_v) / q
            B = self.init_x
            self.__coefficients = [p, q, marker, A, B]
        elif (self.k / self.m) == temp:
            print("The solution is a critically damped oscillation")
            marker = 2 # repeated real solutions
            K = -1 * self.b / 2 * self.m
            # initial conditions
            A = self.init_x
            self.__coefficients = [K, 0, marker, A, B]
        elif (self.k / self.m) < temp: # overdamped oscillation
            marker = 4
            print("The solution is an overdamped oscillation")
            p = -1 * self.b / 2 * self.m + np.sqrt(-(self.k / self.m) + self.b ** 2 / (4 *
self.m ** 2))
            q = -1 * self.b / 2 * self.m - np.sqrt(-(self.k / self.m) + self.b ** 2 / (4 *
self.m ** 2))

            # initial conditions
            A = (q * self.init_x - self.init_v) / (q - p)
            B = self.init_x - A
            self.__coefficients = [p, q, marker, A, B]

    def ana_position(self, t):
        """
        parameter name            type            description
        k_1                        float            function coefficient
        k_2                        float            function coefficient
        marker                      int              marker dictating the solution
type
        A                          float            function coefficient
        B                          float            function coefficient
        :return returns the position of an analytic solution at time t
        """
        k_1 = self.__coefficients[0]
        k_2 = self.__coefficients[1]
        marker = self.__coefficients[2]
        A = self.__coefficients[3]
        B = self.__coefficients[4]
        if marker == 1: # no damping solution
            return A * np.sin(self.natural_angular_frequency * t) + B *

```

```

np.cos(self.natural_angular_frequency * t)
    elif marker == 2: # regular damping (complex)
        return A * np.exp(k_1 * t)
    elif marker == 3: # repeated root
        return np.exp(k_1 * t) * (A * np.sin(k_2 * t) + B * np.cos(k_2 * t))
    elif marker == 4: # two real distinct solutions
        return A * np.exp(k_1 * t) + B * np.exp(k_2 * t)

def ana_velocity(self, t):
    """
    parameter name      type      description
    k_1                  float      function coefficient
    k_2                  float      function coefficient
    marker               int        marker dictating the solution
type
    A                    float      function coefficient
    B                    float      function coefficient
:return returns the velocity of an analytic solution at time t
    """
    k_1 = self.__coefficients[0]
    k_2 = self.__coefficients[1]
    marker = self.__coefficients[2]
    A = self.__coefficients[3]
    B = self.__coefficients[4]
    if marker == 1: # no damping solution
        return A * k_1 * np.cos(k_1 * t) - B * k_1 * np.sin(k_1 * t)
    elif marker == 2: # regular damping (complex)
        return A * k_1 * np.exp(k_1 * t)
    elif marker == 3: # repeated root
        return k_1 * np.exp(k_1 * t) * (A * np.sin(k_2 * t) + B * np.cos(k_2 * t)) +
np.exp(k_1 * t) * (
    A * k_2 * np.cos(k_2 * t) - B * k_2 * np.sin(k_2 * t))
    elif marker == 4: # two real distinct solutions
        return A * k_1 * np.exp(k_1 * t) + B * k_2 * np.exp(k_2 * t)

def plot_data(self): # plots all on separate graphs
    # analytical solution
    """
    parameter name      type      description
    axes_1               object     subplot object
    axes_2               object     subplot object
    axes_3               object     subplot object
    axes_4               object     subplot object
    axes_5               object     subplot object
    axes_6               object     subplot object
    axes_7               object     subplot object
    axes_8               object     subplot object
    axes_9               object     subplot object
    axes_10              object     subplot object
    figure                object     figure object
    figure2               object     figure object
    figure3               object     figure object
    figure4               object     figure object
    figure5               object     figure object
    """
    figure3 = plt.figure()
    axes_5 = figure3.add_subplot(121)
    axes_5.plot(self.analytic_series_pos, self.analytic_series_vel, label="Analytical")
    axes_5.set_xlabel("position/ m") # edit later if the functions don't exist
    axes_5.set_ylabel("velocity/ ms^-1") # as above
    # energy plotting
    axes_6 = figure3.add_subplot(122)
    axes_6.plot(self.time, self.analytic_energy)
    axes_6.set_xlabel("time/ s")
    axes_6.set_ylabel("energy/ J")
    figure3.legend()

    # Euler method
    # plotting
    figure = plt.figure()
    axes_1 = figure.add_subplot(121)
    axes_1.plot(self.Euler_data[0], self.Euler_data[1], label="Euler")
    axes_1.set_xlabel("position/ m") # edit later if the functions don't exist
    axes_1.set_ylabel("velocity/ ms^-1") # as above

```



```

# energy plotting

axes_2 = figure.add_subplot(122)
axes_2.plot(self.time, self.Euler_data[2])
axes_2.set_xlabel("time/ s")
axes_2.set_ylabel("energy/ J")
figure.legend()
# end plotting

# Better Euler method
# plotting
figure2 = plt.figure()
axes_3 = figure2.add_subplot(121)
axes_3.plot(self.B_Euler_data[0], self.B_Euler_data[1], label="Better Euler")
axes_3.set_xlabel("position/ m") # edit later if the functions don't exist
axes_3.set_ylabel("velocity/ ms^-1") # as above
# energy plotting
axes_4 = figure2.add_subplot(122)
axes_4.plot(self.time, self.B_Euler_data[2])
axes_4.set_xlabel("time/ s")
axes_4.set_ylabel("energy/ J")
figure2.legend()
# end plotting

# Verlet method
figure4 = plt.figure()
axes_7 = figure4.add_subplot(121)
axes_7.plot(self.Verlet_data[0], self.Verlet_data[1], label="Verlet")
axes_7.set_xlabel("position/ m")
axes_7.set_ylabel("velocity/ ms^-1")
# energy plotting
axes_8 = figure4.add_subplot(122)
axes_8.plot(self.time, self.Verlet_data[2])
axes_8.set_xlabel("time/ s")
axes_8.set_ylabel("energy/ J")
figure4.legend()

# Euler Cromer method
figure5 = plt.figure()
axes_9 = figure5.add_subplot(121)
axes_9.plot(self.Euler_Cromer_data[0], self.Euler_Cromer_data[1], label="Euler Cromer
Method")
axes_9.set_xlabel("position/ m")
axes_9.set_ylabel("velocity/ ms^-1")
# energy plotting
axes_10 = figure5.add_subplot(122)
axes_10.plot(self.time, self.Euler_Cromer_data[2])
axes_10.set_xlabel("time/ s")
axes_10.set_ylabel("energy/ J")
figure5.legend()

def plot_single(self):
    # plots the single
    """
    parameter name          type          description
    axes_1                  object          subplot object
    axes_2                  object          subplot object
    figure                  object          figure object
    """

    figure = plt.figure()
    axes_1 = figure.add_subplot(121)
    axes_1.set_xlabel("position/ m")
    axes_1.set_ylabel("velocity/ ms^-1")
    # energy_function
    energy = []
    axes_2 = figure.add_subplot(122)
    axes_2.set_xlabel("time/ s")
    axes_2.set_ylabel("Energy/ J")

    axes_1.plot(self.Euler_data[0], self.Euler_data[1], label="Euler")
    axes_2.plot(self.time, self.Euler_data[2])

    axes_1.plot(self.B_Euler_data[0], self.B_Euler_data[1], label="Improved Euler")
    axes_2.plot(self.time, self.B_Euler_data[2])

```

```

axes_1.plot(self.Verlet_data[0], self.Verlet_data[1], label="Verlet")
axes_2.plot(self.time, self.Verlet_data[2])

axes_1.plot(self.Euler_Cromer_data[0], self.Euler_Cromer_data[1], label="Euler
Cromer")
axes_2.plot(self.time, self.Euler_Cromer_data[2])

axes_1.plot(self.analytic_series_pos, self.analytic_series_vel, label="Analytic
solution")
axes_2.plot(self.time, self.analytic_energy)
figure.legend()

def save_data(self):
    """
    parameter name            type            description
    temp                      array            stores the analytic data
    data                      dictionary        stores the data to be saved
    """
    # all files are saved as json dictionaries in the format "name of method": [position,
    velocity, energy]
    # the header of the file headers named appropriately contains the
    temp = [self.analytic_series_pos, self.analytic_series_vel,
self.convert_array(self.analytic_energy)]

    data = {}
    data["Analytic"] = temp
    data["Euler"] = [self.Euler_data[0], self.Euler_data[1],
self.convert_array(self.Euler_data[2])]
    data["Better Euler"] = [self.B_Euler_data[0], self.B_Euler_data[1],
self.convert_array(self.B_Euler_data[2])]
    data["Verlet"] = [self.Verlet_data[0], self.Verlet_data[1],
self.convert_array(self.Verlet_data[2])]
    data["Euler Cromer"] = [self.Euler_Cromer_data[0], self.Euler_Cromer_data[1],
self.convert_array(self.Euler_Cromer_data[2])]
    data["coefficients"] = [self.h, self.no_steps, self.b, self.m, self.k, self.init_x,
self.init_v] # h, T, b, m, k, x, v

    with open("data.txt", 'w') as outfile:
        json.dump(data, outfile)
    outfile.close()

def load_data(self):
    """
    parameter name            type            description
    data                      dictionary        stores the data to be saved
    json_file                object            json file object
    """
    try:
        with open(self.fileNameLoad) as json_file:
            data = json.load(json_file)
            self.Euler_data = data["Euler"]
            self.B_Euler_data = data["Better Euler"]
            self.Verlet_data = data["Verlet"]
            self.analytic_series_pos = data["Analytic"][0]
            self.analytic_series_vel = data["Analytic"][1]
            self.analytic_energy = data["Analytic"][2]
            self.Euler_Cromer_data = data["Euler Cromer"]
            self.h, self.no_steps, self.b, self.m, self.k, self.init_x, self.init_v =
data["coefficients"]
            self.time = np.array(range(0, self.no_steps, 1)) * self.h

        json_file.close()
        return True
    except:
        print("The file was not found")
        return False

def find_accuracy(self):
    # finds the accuracy of the simulation by using the analytic energy as a baseline
    # this assigns a number of "fictitious energy" and also graphs the growth of the
    errors with time
    """
    parameter name            type            description

```

<code>axes_1</code>	<code>object</code>	<code>subplot object</code>
<code>fict_energy</code>	<code>numpy array</code>	<code>stores the error energy</code>
<code>baseline</code>	<code>numpy array</code>	<code>stores the analytic energy</code>
<code>temp</code>	<code>numpy array</code>	<code>stores the plotting value of</code>

the error energy

<code>figure</code>	<code>object</code>	<code>figure object</code>
---------------------	---------------------	----------------------------

```

figure = plt.figure()
axes_1 = figure.add_subplot(111)
axes_1.set_ylabel("Energy error/ J")
axes_1.set_xlabel("time/ s")
fict_energy = []
baseline = np.array(self.analytic_energy)

# Euler's method
temp = np.abs(np.array(self.Euler_data[2]) - baseline)
axes_1.plot(self.time, temp, label="Euler")
fict_energy.append(np.sum(temp))
# Better Euler
temp = np.abs(np.array(self.B_Euler_data[2]) - baseline)
axes_1.plot(self.time, temp, label="Improved Euler")
fict_energy.append(np.sum(temp))
# Cromer
temp = np.abs(np.array(self.Euler_Cromer_data[2]) - baseline)
axes_1.plot(self.time, temp, label="Euler Cromer")
fict_energy.append(np.sum(temp))
# Verlet
temp = np.abs(np.array(self.Verlet_data[2]) - baseline)
axes_1.plot(self.time, temp, label="Verlet")
fict_energy.append(np.sum(temp))
temp = 0
print("The energy errors for b = " + str(self.b))
print("Euler: " + str(fict_energy[0]) + " J")
print("Improved Euler: " + str(fict_energy[1]) + " J")
print("Euler Cromer: " + str(fict_energy[2]) + " J")
print("Verlet: " + str(fict_energy[3]) + " J")
figure.legend()
axes_1.set_title("h = " + str(self.h))

def const_dist_Verlet_integrator(self, force, min, max):
    """
    parameter name      type      description
    position_series      array      stores the position
    temporarily
    velocity_series      array      stores the velocity
    temporarily
    v_n                  float      stores the nth velocity term
    x_n                  float      stores the nth position term
    a_n                  float      stores the nth acceleration
    term
    x_1                  float      stores the second position of
    the oscillation
    D                    float      temporary variable for ease of
    calculation
    B                    float      temporary variable for ease of
    calculation
    A                    float      temporary variable for ease of
    calculation
    """
    position_series = [self.init_x]
    velocity_series = [self.init_v]

    D = 2 * self.m + self.b * self.h
    B = (self.b * self.h - 2 * self.m) / D
    A = 2 * (2 * self.m - self.k * self.h ** 2) / D

    a_0 = (-self.b / self.m) * self.init_v + (-self.k / self.m) * self.init_x
    x_1 = self.init_x + self.init_v * self.h + 0.5 * a_0 * self.h ** 2 # obtained using a
    Taylor expansion of order 2
    position_series.append(x_1)

    for counter in range(1, self.no_steps, 1):
        if (counter * self.h > min) and (counter * self.h < max):
            position_series.append(

```

```

        A * position_series[counter] + B * position_series[counter - 1] + (force /
self.m * self.h ** 2))
    else:
        position_series.append(A * position_series[counter] + B *
position_series[counter - 1])

    # calculating velocities using an approximation of  $O(h^2)$ 
    # the velocity is estimated using the mean value theorem
    # the velocity is independent of the equation of motion. It just utilises the
definition of velocity. If h is small
    # enough this approximation holds true
    for counter in range(1, self.no_steps, 1):
        velocity_series.append(
            (position_series[counter + 1] - position_series[counter - 1]) / (2 * self.h))
# +O(h^2)
position_series = position_series[:len(position_series) - 1]

self.disturbed_Verlet_data = [position_series, velocity_series,
self.energy_function(position_series, velocity_series)]

def funct_dist_Verlet_integrator(self, min, max, Amp, freq):
    """
    parameter name            type            description
    position_series            array            stores the position
temporarily
    velocity_series            array            stores the velocity
temporarily
    v_n                        float            stores the nth velocity term
    x_n                        float            stores the nth position term
    a_n                        float            stores the nth acceleration
term
    x_1                        float            stores the second position of
the oscillation
    D                          float            temporary variable for ease of
calculation
    B                          float            temporary variable for ease of
calculation
    A                          float            temporary variable for ease of
calculation
    """
    position_series = [self.init_x]
    velocity_series = [self.init_v]

    D = 2 * self.m + self.b * self.h
    B = (self.b * self.h - 2 * self.m) / D
    A = 2 * (2 * self.m - self.k * self.h ** 2) / D

    a_0 = (-self.b / self.m) * self.init_v + (-self.k / self.m) * self.init_x
    x_1 = self.init_x + self.init_v * self.h + 0.5 * a_0 * self.h ** 2 # obtained using a
Taylor expansion of order 2
    position_series.append(x_1)

    for counter in range(1, self.no_steps, 1):
        if (counter * self.h > min) and (counter * self.h < max):
            position_series.append(
                A * position_series[counter] + B * position_series[counter - 1] + (
                    Amp * np.sin(freq * counter * self.h) / self.m * self.h ** 2))
        else:
            position_series.append(A * position_series[counter] + B *
position_series[counter - 1])

    # calculating velocities using an approximation of  $O(h^2)$ 
    # the velocity is estimated using the mean value theorem
    # the velocity is independent of the equation of motion. It just utilises the
definition of velocity. If h is small
    # enough this approximation holds true
    for counter in range(1, self.no_steps, 1):
        velocity_series.append(
            (position_series[counter + 1] - position_series[counter - 1]) / (2 * self.h))
# +O(h^2)
position_series = position_series[:len(position_series) - 1]

self.disturbed_Verlet_data = [position_series, velocity_series,
self.energy_function(position_series, velocity_series)]

```

```

def push_testing(self, min, max, force, amp, freq):
    """
    parameter name          type          description
    constant_data            array         array of the data modified by
a constant force
    function_data            array         array of the data modified by
a sinusoidal force
    axes_1                   object        subplot object
    figure                    object        figure object
    """
    # constant force
    self.const_dist_Verlet_integrator(force, min, max)
    constant_data = self.disturbed_Verlet_data

    # sinusoidal force
    self.funct_dist_Verlet_integrator(min, max, amp, freq)
    function_data = self.disturbed_Verlet_data

    figure = plt.figure()
    axes_1 = figure.add_subplot(111)
    axes_1.set_xlabel("Time (s)")
    axes_1.set_ylabel("Position (m)")
    axes_1.plot(self.time, constant_data[0], label="constant force")
    axes_1.plot(self.time, function_data[0], label="sinusoidal force")
    axes_1.plot(self.time, self.Verlet_data[0], label="undisturbed")
    figure.legend()

def search(self, arr, x):
    """
    variable name          type          description
    i                      integer       counter
    arr                    list          array
    x                      float         the value being searched
    """
    # Linear search function
    for i in range(len(arr)):
        if arr[i] == x:
            return i
    return -1

def resonance_Plot(self):
    """
    parameter name          type          description
    temp                    float/array   temporary variable
    freq_array              array         array of angular frequencies
    amplitude               array         sum of the amplitudes for a
resonance plot
    figure                  object        figure object
    axes_1                  object        subplot object
    """
    temp = np.sqrt(self.k / self.m)
    freq_array = []
    for counter in range(0, 200, 1):
        freq_array.append(counter * temp * 0.01)

    amplitude = []
    for freq in freq_array:
        # get the data sets for a specific frequency
        self.funct_dist_Verlet_integrator(0, self.no_steps * self.h, self.init_x * 0.5,
freq)
        temp = self.disturbed_Verlet_data[0]
        # calculate amplitude
        temp = np.abs(np.array(temp))
        amplitude.append(np.mean(temp))
        # append to the arrays
    # plot the results
    figure = plt.figure()
    axes_1 = figure.add_subplot(111)
    axes_1.set_xlabel("Frequency/ Hz")
    axes_1.set_ylabel("Amplitude/ m")
    axes_1.plot(freq_array, amplitude)

```

```

def Critical(self):
    """
    parameter name            type            description
    temp                      float           stores the b value to avoid
change
    b                          float           damping coefficient
    data                      array           stores the data of all the
integrations
    figure4                   object          figure object
    figure                    object          figure object
    axes_1                   object          subplot object
    axes_7                   object          subplot object
    axes_8                   object          subplot object
    """
    temp = self.b
    b = [0.5 * self.b_critical, self.b_critical, 2 * self.b_critical]
    data = []
    for entry in b:
        self.b = entry
        self.Verlet_integrator()
        data.append(self.Verlet_data)

    # Verlet method
    # phase plots
    figure4 = plt.figure()
    axes_7 = figure4.add_subplot(121)
    axes_7.plot(data[0][0], data[0][1], label="0.5b")
    axes_7.plot(data[1][0], data[1][1], label="b")
    axes_7.plot(data[2][0], data[2][1], label="2b")
    axes_7.set_xlabel("position (m)")
    axes_7.set_ylabel("velocity (ms-1)")
    # energy plotting
    axes_8 = figure4.add_subplot(122)
    axes_8.plot(self.time, data[0][2])
    axes_8.plot(self.time, data[1][2])
    axes_8.plot(self.time, data[2][2])
    axes_8.set_xlabel("time (s)")
    axes_8.set_ylabel("energy (J)")
    # position plot
    figure = plt.figure()
    axes_1 = figure.add_subplot(111)
    axes_1.plot(self.time, data[0][0], label="0.5b")
    axes_1.plot(self.time, data[1][0], label="b")
    axes_1.plot(self.time, data[2][0], label="2b")
    axes_1.set_ylabel("Position (m)")
    axes_1.set_xlabel("Time (s)")
    figure4.legend()
    figure.legend()

    self.b = temp

def complete_Resonance(self):
    """
    parameter name            type            description
    temp 2                    float           stores the b value to avoid
change
    figure                   object          figure object
    axes_1                   object          subplot object
    freq_array               array           array of angular frequencies
    amplitude                array           sum of the amplitudes for a
resonance plot
    b_prime                  array           stores the damping
coefficients
    """
    figure = plt.figure()
    axes_1 = figure.add_subplot(111)
    axes_1.set_xlabel("Frequency/ Hz")
    axes_1.set_ylabel("Amplitude/ m")

    temp2 = self.b
    temp = np.sqrt(self.k / self.m)
    freq_array = []
    for counter in range(0, 220, 1):
        freq_array.append(counter * temp * 0.01)

```

```

b_prime = [0, 0.1, 0.2, 0.3, 0.4, 0.5, 0.6, 2.2, 4.5, 9]
amplitude = []
for b in b_prime:
    self.b = b
    for freq in freq_array:
        # get the data sets for a specific frequency
        self.funct_dist_Verlet_integrator(0, self.no_steps * self.h, self.init_x *
0.5, freq)

        temp = self.disturbed_Verlet_data[0]
        # calculate amplitude
        temp = np.abs(np.array(temp))
        amplitude.append(np.mean(temp))
        # append to the arrays

        axes_1.plot(freq_array, amplitude, label="b = " + str(b) + " kgs^-1")
        amplitude = []
    # plot the results

    figure.legend()
    self.b = temp2

def find_accuracy_complete(self):
    """
    parameter name          type          description
    temp                    float          stores the b variable to avoid
change
    """
    temp = self.b
    self.b = 0
    self.find_accuracy()
    self.b = 0.1
    self.find_accuracy()
    self.b = 1
    self.find_accuracy()
    self.b = 2
    self.find_accuracy()
    self.b = 3
    self.find_accuracy()
    self.b = temp

def getData(): # edit this function
    """
    parameter name          type          description
    m                       float         mass
    k                       float         spring constant
    b                       float         damping coefficient
    T                       float         total time
    h                       float         time step
    init_x                  float         initial position
    init_v                  float         initial velocity
    """
    m = float(input("Enter the value for the mass of the particle: "))
    k = float(input("Enter the value of the spring constant: "))
    b = float(input("Enter the value of the damping constant: "))
    T = float(input("Enter the time you want the simulation to run: "))
    h = float(input("Enter the time step in seconds: "))
    init_x = float(input("Enter the initial position: "))
    init_v = float(input("Enter the initial velocity: "))
    return m, k, b, T, h, init_x, init_v

def main():
    """
    parameter name          type          description
    option                  int          option chosen
    name                    string       file name
    os                      object       the class object
    check                   boolean      check variable
    """
    option = 0
    while option != "9":

        print("1. Run simulation")

```

```

print("2. Load old simulation")
option = input("Select an option:")
if option == "1" or option == "2":
    if option == "1":
        m, k, b, T, h, init_x, init_v = getData()
        os = SHO(h, T, b, m, k, init_x, init_v)
    elif option == "2":
        print("Enter the name of the file or type 'none' if you want to use default")
        name = input()
        if name == "none":
            os = SHO(0.01, 100)
            os.load_data()
        else:
            os = SHO(0.01, 100, fileNameLoad=name)
            check = os.load_data()
            if not check:
                print("Goodbye!")
                return 0

print("3. Run critical damping simulation")
print("4. Run simulation with the force applied")
print("5. Plot all of it")
print("6. Save the simulation")
print("7. Plot Resonance Curves")
print("8. Run default")
print("9. Leave")
option = input("Select an option:")
if option == "3":
    os.Critical()
elif option == "4":
    min = float(input("Minimum time: "))
    max = float(input("Maximum time: "))
    force = float(input("Force magnitude: "))
    Amp = float(input("Sinusoidal force amplitude: "))
    freq = float(input("Sinusoidal force frequency: "))
    os.push_testing(min, max, force, Amp, freq)
elif option == "5":
    os.plot_data()
    os.plot_single()
elif option == "6":
    os.save_data()
    print("File was saved as data.txt")
elif option == "7":
    os.complete_Resonance()
elif option == "8":
    print("Running default simulation")
    os.runSimulation()
    os.push_testing(45.6, 100, 2, 2, 0.062832) # at zero amplitude
    os.push_testing(57, 100, 2, 2, 6.2832) # at 3/4 of a cycle
    os.push_testing(53.2, 100, 2, 2, 0.41)
    os.find_accuracy_complete()
    os.Critical()
    os.complete_Resonance()
    os.plot_single()
    os.plot_data()
plt.show()

else:
    print("Goodbye!")

main()
plt.show()

```