# Fracture Mechanics Basic notions of solid mechanics

Luc St-Pierre

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# Reviewing notions of solid mechanics

There are a few important notions of solid mechanics to review before we tackle fracture mechanics. These include:

- ▶ Plane stress/strain
- ► Hooke's law
- ► Equilibrium equations
- Compatibility equations
- ► Airy stress functions

If this brief review is insufficient, refer to the book by Timoshenko and Goodier, *Theory of Elasticity*.

#### Plane stress

In continuum mechanics, 2D problems are usually treated as plane stress or plane strain.

Plane stress, which is used for thin plates, assumes that  $\sigma_{13} = \sigma_{23} = \sigma_{33} = 0$  so the stress tensor becomes:

$$\sigma_{ij} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & \sigma_{22} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & 0 \\ \sigma_{12} & \sigma_{22} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

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while the strains are:

$$\epsilon_{ij} = \begin{bmatrix} \epsilon_{11} & \epsilon_{12} & \epsilon_{13} \\ \epsilon_{12} & \epsilon_{22} & \epsilon_{23} \\ \epsilon_{13} & \epsilon_{23} & \epsilon_{33} \end{bmatrix} = \begin{bmatrix} \epsilon_{11} & \epsilon_{12} & 0 \\ \epsilon_{12} & \epsilon_{22} & 0 \\ 0 & 0 & \epsilon_{33} \end{bmatrix}$$

note that  $\epsilon_{33} \neq 0!$ 

#### Plane strain

Plane strain is used for **thick** structures and assumes that  $\epsilon_{13} = \epsilon_{23} = \epsilon_{33} = 0$ . Therefore, the strain tensor becomes:

$$\epsilon_{ij} = \begin{bmatrix} \epsilon_{11} & \epsilon_{12} & \epsilon_{13} \\ \epsilon_{12} & \epsilon_{22} & \epsilon_{23} \\ \epsilon_{13} & \epsilon_{23} & \epsilon_{33} \end{bmatrix} = \begin{bmatrix} \epsilon_{11} & \epsilon_{12} & 0 \\ \epsilon_{12} & \epsilon_{22} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

while the stress tensor is:

$$\sigma_{ij} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & \sigma_{22} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & 0 \\ \sigma_{12} & \sigma_{22} & 0 \\ 0 & 0 & \sigma_{33} \end{bmatrix}$$

where for a linear isotropic material  $\sigma_{33} = \nu(\sigma_{11} + \sigma_{22})$ , where  $\nu$  is the Poisson's ratio.

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#### Hooke's law

For an isotropic linear elastic material, the strain and stress components are related as:

$$\begin{bmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \epsilon_{zz} \\ \epsilon_{xy} \\ \epsilon_{xz} \\ \epsilon_{yz} \end{bmatrix} = \frac{1}{E} \begin{bmatrix} 1 & -\nu & -\nu & 0 & 0 & 0 \\ -\nu & 1 & -\nu & 0 & 0 & 0 \\ -\nu & -\nu & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2 + 2\nu & 0 & 0 \\ 0 & 0 & 0 & 0 & 2 + 2\nu & 0 \\ 0 & 0 & 0 & 0 & 0 & 2 + 2\nu \end{bmatrix} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \sigma_{xy} \\ \sigma_{xz} \\ \sigma_{yz} \end{bmatrix}$$

where E is the Young's modulus and  $\nu$  is the Poisson's ratio.

# Equilibrium equations

In 2D, the stress field should always respect equilibrium equations, and these are:

Cartesian coordinates:

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} = 0$$
$$\frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} = 0$$

Polar coordinates:

$$\frac{\partial}{\partial r}(r\sigma_{rr}) + \frac{\partial\sigma_{r\theta}}{\partial\theta} - \sigma_{\theta\theta} = 0$$
$$\frac{\partial\sigma_{\theta\theta}}{\partial\theta} + \frac{\partial}{\partial r}(r\sigma_{r\theta}) + \sigma_{r\theta} = 0$$

The above equations assume that body forces are negligible.

# Strains and compatibility equation

In 2D, and using cartesian coordinates, the strains are defined as:

$$\epsilon_{xx} = \frac{\partial u}{\partial x}$$
  $\epsilon_{yy} = \frac{\partial v}{\partial y}$   $\epsilon_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$ 

where u and v are the displacements in the x and y directions, respectively.

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where u and v are the displacements in the x and y directions, respectively. The three strain components are not independent since they are defined from only two values of displacement. They are related via the compatibility equation:

$$\frac{\partial^2 \epsilon_{xx}}{\partial y^2} + \frac{\partial^2 \epsilon_{yy}}{\partial x^2} = \frac{\partial^2 \epsilon_{xy}}{\partial x \partial y}$$

This equation can also be expressed in polar coordinates...but it is messy.

## Compatibility equation in terms of stress

For an isotropic linear elastic material loaded in plane stress, the strain and stress components are related as:

$$\epsilon_{xx} = \frac{1}{E}(\sigma_{xx} - \nu \sigma_{yy})$$

$$\epsilon_{yy} = \frac{1}{E}(\sigma_{yy} - \nu \sigma_{xx})$$

$$\epsilon_{xy} = \frac{1}{G}\sigma_{xy} = \frac{2(1+\nu)}{E}\sigma_{xy}$$

Subtituting these in the compatibility equation (and doing some algebra) returns:

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)(\sigma_{xx} + \sigma_{yy}) = 0$$

This compatibility equation in terms of stresses is also valid for plane strain.

Therefore, in 2D, a valid stress field as to respect both equilibrium and compatibility equations:

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} = 0$$
$$\frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} = 0$$
$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right) (\sigma_{xx} + \sigma_{yy}) = 0$$

Finding three unknown functions  $(\sigma_{xx}, \sigma_{yy}, \sigma_{xy})$  that respect these three equations is difficult; however, this can be made easier by introducing the Airy stress function.

The Airy stress function  $\phi(x,y)$  has no physical meaning, it is simply a mathematical trick. The function is related to the stress components as:

$$\sigma_{xx} = \frac{\partial^2 \phi}{\partial y^2}$$
  $\sigma_{yy} = \frac{\partial^2 \phi}{\partial x^2}$   $\sigma_{xy} = -\frac{\partial^2 \phi}{\partial x \partial y}$ 

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It can be shown that equilibirum and compatibility equations are both respected when:

$$\frac{\partial^4 \phi}{\partial x^4} + 2 \frac{\partial^4 \phi}{\partial x^2 \partial y^2} + \frac{\partial^4 \phi}{\partial y^4} = 0 \quad \Rightarrow \quad \nabla^4 \phi = 0$$

We have simplified the problem: there is now one function  $\phi$  and a single equation.

In polar coordinates, the Airy stress function  $\phi(r,\theta)$  is related to the stresses as:

$$\sigma_{rr} = \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} \quad \sigma_{\theta\theta} = \frac{\partial^2 \phi}{\partial r^2} \quad \sigma_{r\theta} = -\frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \phi}{\partial \theta} \right)$$

and the compatibility equation  $\nabla^4 \phi = 0$  becomes:

$$\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r}\frac{\partial}{\partial r} + \frac{1}{r^2}\frac{\partial^2}{\partial \theta^2}\right)\left(\frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r}\frac{\partial \phi}{\partial r} + \frac{1}{r^2}\frac{\partial^2 \phi}{\partial \theta^2}\right) = 0$$