



Aalto University

# **MS-C1350 Partial differential equations**

## **Chapter 5.1-5.3**

### **Heat equation – physical interpretation, fundamental solution and nonhomogeneous problem**

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## 5.1 Heat equation

- Now we study the heat equation

$$u_t - \Delta u = 0$$

- and the nonhomogeneous heat equation

$$u_t - \Delta u = f$$

with appropriate initial and boundary conditions.

- The Laplace is taken with respect to the spatial variable  $x$

$$\Delta u = \sum_{j=1}^n \frac{\partial^2 u}{\partial x_j^2}.$$

- Let  $\Omega \subset \mathbb{R}^n$  and  $T > 0$ . The problem is to find a function  $u = u(x, t)$  such that it is a solution to the heat equation in  $\Omega \times (0, T)$ .

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## 5.1 Heat equation

- ▶ Physically, a solution  $u = u(x, t)$  of the heat equation represents the temperature of the body  $\Omega$  at the point  $x$  and time  $t$ .
- ▶ Observe that any solution  $v = v(x)$  of the Laplace equation induces a time independent solution  $u = u(x, t) = v(x)$  of the heat equation
- ▶ This suggests that for every claim about solutions to the Laplace equation there should be a corresponding claim for the solutions of the heat equation.
- ▶ However, the dependence in time leads to new phenomena and challenges which are not visible in the stationary case.

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## 5.1 Initial and boundary conditions

- ▶ The appropriate **initial condition** is

$$u(x, 0) = g(x), \quad x \in \Omega.$$

This describes the initial temperature distribution at the time  $t = 0$ .

- ▶ In addition, we may have a Dirichlet type boundary condition

$$u(x, t) = g(x, t), \quad x \in \partial\Omega, \quad t > 0,$$

which describes the temperature on the boundary,

- ▶ or a Neumann type boundary condition

$$\frac{\partial u}{\partial \nu}(x, t) = \nabla_x u(x, t) \cdot \nu(x) = h(x, t), \quad x \in \partial\Omega, \quad t > 0,$$

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## 5.2 Fundamental solution

- ▶ We already derived (with Fourier techniques) a formula for the solution of the Cauchy problem

$$\begin{cases} u_t - \Delta u = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ u = g & \text{on } \mathbb{R}^n \times \{t = 0\}, \end{cases}$$

where  $g \in C_0^\infty(\mathbb{R}^n)$ .

- ▶ We obtained

$$u(x, t) = \int_{\mathbb{R}^n} \Phi(x - y, t) g(y) dy = \frac{1}{(4\pi t)^{\frac{n}{2}}} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4t}} g(y) dy.$$



$$\Phi(x, t) = \begin{cases} \frac{1}{(4\pi t)^{\frac{n}{2}}} e^{-\frac{|x|^2}{4t}}, & x \in \mathbb{R}^n, \quad t > 0, \\ 0, & x \in \mathbb{R}^n, \quad t \leq 0, \end{cases}$$

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## 5.2 Fundamental solution

- ▶  $\Phi$  is a solution to the heat equation in the upper half-space.
- ▶  $\Phi$  is unbounded in any neighbourhood of  $(0, 0)$ .
- ▶ We proved earlier that the solution attains the initial values  $g$  in the sense

$$\lim_{t \rightarrow 0} u(x, t) = g(x) \quad \text{for every } x \in \mathbb{R}^n$$

(“good kernels”)

- ▶ This means that we have existence of a solution to the Cauchy problem.
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- Formally the fundamental solution solves the following PDE:

$$\begin{cases} \frac{\partial \Phi}{\partial t} - \Delta \Phi = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ \Phi = \delta_0 & \text{in } \mathbb{R}^n \times \{t = 0\}, \end{cases}$$

where  $\delta_0$  is the Dirac measure in  $\mathbb{R}^n$  giving unit mass to the point 0.

## 5.2 Remarks about heat equation

- ▶ The heat equation has infinite speed of propagation.
- ▶ For example, if initial values  $g \in C(\mathbb{R}^n)$  are non-negative  $g \geq 0$  and there is a point  $y \in \mathbb{R}^n$  such that  $g(y) > 0$ , then  $u > 0$  everywhere in the upper half-space.
- ▶ From the formula

$$u(x, t) = \int_{\mathbb{R}^n} \Phi(x - y, t) g(y) dy = \frac{1}{(4\pi t)^{\frac{n}{2}}} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4t}} g(y) dy.$$

we get the decay estimate

$$\sup_{x \in \mathbb{R}^n} |u(x, t)| \leq \frac{1}{(4\pi t)^{\frac{n}{2}}} \int_{\mathbb{R}^n} |g(y)| dy, \quad t > 0.$$

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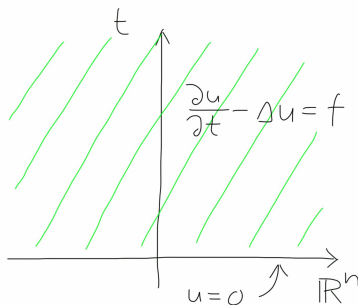
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## 5.3 Nonhomogeneous problem

- Consider the nonhomogeneous Cauchy problem

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$$u(x, t) = \int_{\mathbb{R}^n} \Phi(x - y, t) g(y) dy = \frac{1}{(4\pi t)^{\frac{n}{2}}} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4t}} g(y) dy.$$

- we obtain that for every fixed  $s$ , with  $0 < s < t$ , the function

$$u(x, t; s) = \int_{\mathbb{R}^n} \Phi(x - y, t - s) f(y, s) dy$$

solves the translated initial value problem

$$\begin{cases} \frac{\partial u}{\partial t}(x, t; s) - \Delta u(x, t; s) = 0, & x \in \mathbb{R}^n, \quad t > s, \\ u(x, s; s) = f(x, s), & x \in \mathbb{R}^n. \end{cases}$$

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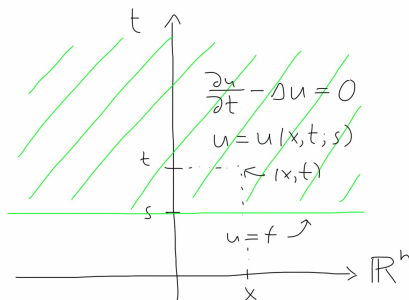
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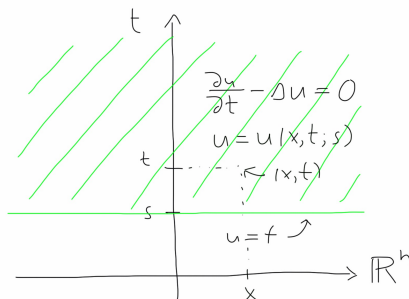


- Duhamel's principle suggests that we can construct a solution to the nonhomogeneous problem by integrating solutions  $u(x, t; s)$  over  $s \in (0, t)$  and have

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## 5.3 Duhamel's principle

- ▶ Duhamel's principle is a process of expressing the solution of a nonhomogeneous problem as an integral of the solutions to the homogeneous problem in the way that the source term is interpreted as the initial condition.
- ▶ It does not depend on the specific structure of the equation and it applies to other linear ODEs and PDEs as well.
- ▶ We can verify it by differentiating the function  $u$  defined by an integral.

$$u(x, t) = \int_0^t u(x, t; s) ds, \quad x \in \mathbb{R}^n, \quad t > 0.$$

- ▶ Recall that (Lemma 5.3)

$$\frac{\partial u}{\partial t}(x, t) = u(x, t; t) + \int_0^t \frac{\partial u}{\partial t}(x, t; s) ds.$$

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## 5.3 Duhamel's principle

### Theorem

*The solution of the nonhomogeneous Cauchy problem is*

$$\begin{aligned} u(x, t) &= \int_0^t \int_{\mathbb{R}^n} \Phi(x - y, t - s) f(y, s) dy ds \\ &= \int_0^t \frac{1}{(4\pi(t - s))^{\frac{n}{2}}} \left( \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, s) dy \right) ds. \end{aligned}$$

► Formal proof:

$$\begin{aligned} \frac{\partial u}{\partial t}(x, t) &= u(x, t; t) + \int_0^t \frac{\partial u}{\partial t}(x, t; s) ds \\ &= f(x, t) + \int_0^t \Delta u(x, t; s) ds \\ &= f(x, t) + \Delta \left( \int_0^t u(x, t; s) ds \right) \\ &= f(x, t) + \Delta u(x, t). \end{aligned}$$

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## 5.3 Nonhomogeneous problem with general initial data

- To solve the general problem

$$\begin{cases} u_t - \Delta u = f & \text{in } \mathbb{R}^n \times (0, \infty), \\ u = g & \text{in } \mathbb{R}^n \times \{t = 0\}, \end{cases}$$

we can use the same approach as for the Laplace equation and write  $u = u_1 + u_2$  with

$$\begin{cases} \frac{\partial u_1}{\partial t} - \Delta u_1 = f & \text{in } \mathbb{R}^n \times (0, \infty), \\ u_1 = 0 & \text{in } \mathbb{R}^n \times \{t = 0\}, \end{cases}$$

and

$$\begin{cases} \frac{\partial u_2}{\partial t} - \Delta u_2 = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ u_2 = g & \text{in } \mathbb{R}^n \times \{t = 0\}. \end{cases}$$