Quantum Information Spring 2023 Problem Set 2

1. Mixed States

Consider the 2×2 matrix

$$\rho = \begin{pmatrix} 3/4 & \sqrt{2}e^{-i\phi}/4\\ \sqrt{2}e^{i\phi}/4 & 1/4 \end{pmatrix} . \tag{1}$$

(a) Is the matrix a density matrix? Solution.

We have to check that the matrix is Hermitian, has $\operatorname{tr} \rho = 1$, and is positive semi-definite. The matrix is clearly Hermitian $\rho^{\dagger} = \rho$. Its trace is $\operatorname{tr} \rho = 3/4 + 1/4 = 1$. To check the last condition, we solve the characteristic equation for the eigenvalues

$$0 = \det(\rho - \lambda \mathbb{I}) = \left(\frac{3}{4} - \lambda\right) \left(\frac{1}{4} - \lambda\right) - \frac{1}{8} \tag{2}$$

$$=\lambda^2 - \lambda + \frac{1}{16} \ . \tag{3}$$

The roots of this equation are $\lambda = (2 \pm \sqrt{3})/4$. Both of the eigenvalues are positive, so the matrix ρ is positive (semi-)definite and therefore a valid density matrix.

(b) If so do we have a pure state or a mixed state? Solution.

Since

$$\rho^2 = \begin{pmatrix} 11/16 & e^{-i\phi}/(2\sqrt{2}) \\ e^{i\phi}/(2\sqrt{2}) & 3/16 \end{pmatrix} \neq \rho , \qquad (4)$$

we have a mixed state.

2. Entangling gates

Consider the following quantum circuit

$$q_0: |0\rangle$$
 H $q_1: |0\rangle$ Z

(a) Find the state vector $|\psi\rangle$ produced by this circuit. Solution.

The first Hadamard transforms

$$|0\rangle \otimes |0\rangle \mapsto \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes |0\rangle = \frac{1}{\sqrt{2}}|0\rangle \otimes |0\rangle + \frac{1}{\sqrt{2}}|1\rangle \otimes |0\rangle .$$
 (5)

After the controlled-NOT gate the state is

$$\frac{1}{\sqrt{2}}|0\rangle \otimes |0\rangle + \frac{1}{\sqrt{2}}|1\rangle \otimes |0\rangle \mapsto \frac{1}{\sqrt{2}}|0\rangle \otimes |0\rangle + \frac{1}{\sqrt{2}}|1\rangle \otimes |1\rangle . \tag{6}$$

After the final gate we have

$$\frac{1}{\sqrt{2}}|0\rangle \otimes |0\rangle + \frac{1}{\sqrt{2}}|1\rangle \otimes |1\rangle \mapsto \frac{1}{\sqrt{2}}|0\rangle \otimes |0\rangle - \frac{1}{\sqrt{2}}|1\rangle \otimes |1\rangle . \tag{7}$$

So, the state vector produced by this circuit is

$$\frac{1}{\sqrt{2}}|0\rangle \otimes |0\rangle - \frac{1}{\sqrt{2}}|1\rangle \otimes |1\rangle \quad \text{or} \quad \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\0\\0\\-1 \end{pmatrix} . \tag{8}$$

(b) Compute the corresponding density operator. *Solution*.

The density operator/matrix is obtained as the product $\rho = |\psi\rangle\langle\psi|$

$$\rho = \frac{1}{2} \begin{pmatrix} 1\\0\\0\\-1 \end{pmatrix} \begin{pmatrix} 1&0&0&-1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1&0&0&-1\\0&0&0&0\\0&0&0&0\\-1&0&0&1 \end{pmatrix} . \tag{9}$$

3. Entanglement Entropy

(a) Consider the Hilbert space $\mathcal{H}_A \otimes \mathcal{H}_B$, where $\mathcal{H}_A = \mathcal{H}_B = \mathbb{C}^2$ and the state

$$|\psi\rangle = \frac{1}{2}(|00\rangle - |01\rangle - |10\rangle + |11\rangle) = \frac{1}{2} \begin{pmatrix} 1\\ -1\\ -1\\ 1 \end{pmatrix} .$$
 (10)

Calculate the entanglement entropies

$$-\operatorname{tr}(\rho_A \log \rho_A)$$
 and $-\operatorname{tr}(\rho_B \log \rho_B)$, (11)

where the reduced density matrices are

$$\rho_A = \operatorname{tr}_{\mathcal{H}_B} |\psi\rangle\langle\psi| \quad \text{and} \quad \rho_B = \operatorname{tr}_{\mathcal{H}_A} |\psi\rangle\langle\psi| .$$
(12)

Solution.

The density matrix $\rho = |\psi\rangle\langle\psi|$ is

Therefore, the reduced density matrix ρ_A is

$$\rho_A = (\mathbb{I}_A \otimes \langle 0|) |\psi\rangle\langle\psi| (\mathbb{I}_A \otimes |0\rangle) + (\mathbb{I}_A \otimes \langle 1|) |\psi\rangle\langle\psi| (\mathbb{I}_A \otimes |1\rangle)$$
(14)

$$= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \end{pmatrix} |\psi\rangle\langle\psi| \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$
 (15)

$$+\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 1 \end{pmatrix} |\psi\rangle\langle\psi| \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$
 (16)

$$= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} |\psi\rangle\langle\psi| \begin{pmatrix} 1 & 0 \\ 0 & 0 \\ 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} |\psi\rangle\langle\psi| \begin{pmatrix} 0 & 0 \\ 1 & 0 \\ 0 & 0 \\ 0 & 1 \end{pmatrix}$$
(17)

$$=\frac{1}{2}\begin{pmatrix}1 & -1\\ -1 & 1\end{pmatrix}. \tag{18}$$

In order to find the entanglement entropy of ρ_A , we'll solve for its eigenvalues λ from the characteristic equation

$$\det (\rho_A - \lambda \mathbb{I}_A) = \det \begin{pmatrix} 1/2 - \lambda & -1/2 \\ -1/2 & 1/2 - \lambda \end{pmatrix} = \left(\frac{1}{2} - \lambda\right)^2 - \frac{1}{4} = 0 , \qquad (19)$$

which has the solutions $\lambda_1 = 0$ and $\lambda_2 = 1$. Therefore the entanglement entropy is

$$S(\rho_A) = -\operatorname{tr}(\rho_A \log \rho_A) = -\sum_i \lambda_i \log \lambda_i = -0 \log 0 - 1 \log 1 = 0.$$
 (20)

We could do the same computation to find first ρ_B and then its entropy, but since $\rho = |\psi\rangle\langle\psi|$ is a pure state we must have $S(\rho_B) = S(\rho_A) = 0$.

(b) Is this state separable? Why or why not? Solution.

This state is separable and therefore can be written $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$ for some $|\psi_A\rangle \in \mathcal{H}_A$ and $|\psi_B\rangle \in \mathcal{H}_B$. This is because the reduced states ρ_A and ρ_B are pure because $S(\rho_A) = S(\rho_B) = 0$. Explicitly, this can be shown by noting that $|\psi\rangle = |-\rangle \otimes |-\rangle$.

4. Trace Distance

Consider the density operators

$$\rho = \frac{3}{4} |0\rangle\langle 0| + \frac{1}{4} |1\rangle\langle 1| \quad \text{and} \quad \sigma = \frac{2}{3} |+\rangle\langle +| + \frac{1}{3} |-\rangle\langle -| , \qquad (21)$$

where $|\pm\rangle = (|0\rangle \pm |1\rangle)/\sqrt{2}$. Calculate the trace distance $D(\rho, \sigma)$. Solution.

Let's write the density operators as matrices

$$\rho = \begin{pmatrix} 3/4 & 0\\ 0 & 1/4 \end{pmatrix}
\tag{22}$$

$$\sigma = \frac{2}{3} \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} + \frac{1}{3} \begin{pmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$
(23)

$$= \frac{2}{3} \begin{pmatrix} 1/2 & 1/2 \\ 1/2 & 1/2 \end{pmatrix} + \frac{1}{3} \begin{pmatrix} 1/2 & -1/2 \\ -1/2 & 1/2 \end{pmatrix}$$
 (24)

$$= \begin{pmatrix} 1/2 & 1/6 \\ 1/6 & 1/2 \end{pmatrix} . \tag{25}$$

Since $D(\rho, \sigma) = \frac{1}{2} \operatorname{tr} |\rho - \sigma|$, where $|A| \equiv \sqrt{A^{\dagger} A}$, we first compute $|\rho - \sigma|$

$$\rho - \sigma = \begin{pmatrix} 1/4 & -1/6 \\ -1/6 & -1/4 \end{pmatrix} . \tag{26}$$

Then

$$(\rho - \sigma)^{\dagger}(\rho - \sigma) = \begin{pmatrix} 13/144 & 0\\ 0 & 13/144 \end{pmatrix} . \tag{27}$$

Since this is diagonal, the square root is easy to compute

$$\sqrt{(\rho - \sigma)^{\dagger}(\rho - \sigma)} = \begin{pmatrix} \sqrt{13/12} & 0\\ 0 & \sqrt{13/12} \end{pmatrix} = |\rho - \sigma|.$$
 (28)

Therefore, the trace distance is

$$D(\rho, \sigma) = \frac{1}{2} \operatorname{tr} |\rho - \sigma| = \frac{1}{2} \frac{\sqrt{13}}{6} = \frac{\sqrt{13}}{12} . \tag{29}$$

5. Fidelity

Consider the density operators

$$\rho = \frac{3}{4} |0\rangle\langle 0| + \frac{1}{4} |1\rangle\langle 1| \quad \text{and} \quad \sigma = \frac{2}{3} |0\rangle\langle 0| + \frac{1}{3} |1\rangle\langle 1| . \tag{30}$$

(a) Compute the fidelity $F(\rho, \sigma) = \operatorname{tr} \sqrt{\sqrt{\rho}\sigma\sqrt{\rho}}$. Solution.

We start by writing the density operators as matrices

$$\rho = \begin{pmatrix} 3/4 & 0 \\ 0 & 1/4 \end{pmatrix} \quad \text{and} \quad \sigma = \begin{pmatrix} 2/3 & 0 \\ 0 & 1/3 \end{pmatrix} . \tag{31}$$

Both matrices are diagonal, so the matrix square roots are again simple:

$$\sqrt{\rho} = \begin{pmatrix} \sqrt{3}/2 & 0\\ 0 & 1/2 \end{pmatrix} \tag{32}$$

$$\sqrt{\rho}\sigma\sqrt{\rho} = \begin{pmatrix} 1/2 & 0\\ 0 & 1/12 \end{pmatrix} \tag{33}$$

$$\sqrt{\sqrt{\rho}\sigma\sqrt{\rho}} = \begin{pmatrix} 1/\sqrt{2} & 0\\ 0 & 1/(2\sqrt{3}) \end{pmatrix} . \tag{34}$$

Therefore, the fidelity is

$$F(\rho, \sigma) = \frac{1}{\sqrt{2}} + \frac{1}{2\sqrt{3}}$$
 (35)

(b) A bit flip channel flips the qubit from $|0\rangle$ to $|1\rangle$ (and vice versa) with probability 1-p. It can be described with the quantum channel

$$\rho \mapsto \mathcal{E}(\rho) = E_0 \rho E_0^{\dagger} + E_1 \rho E_1^{\dagger} , \qquad (36)$$

where

$$E_0 = \sqrt{p} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} , \quad E_1 = \sqrt{1-p} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} .$$
 (37)

The bit flip channel has the following effect on our states ρ and σ :

$$\rho \mapsto \mathcal{E}(\rho) = \begin{pmatrix} (1+2p)/4 & 0\\ 0 & (3-2p)/4 \end{pmatrix}$$
 (38)

$$\sigma \mapsto \mathcal{E}(\sigma) = \begin{pmatrix} (1+p)/3 & 0\\ 0 & (2-p)/3 \end{pmatrix} . \tag{39}$$

Calculate the fidelity $F(\mathcal{E}(\rho), \mathcal{E}(\sigma))$. Is the new fidelity smaller or larger than the original fidelity $F(\rho, \sigma)$? Can you explain why?

Solution. The states $\mathcal{E}(\rho)$ and $\mathcal{E}(\sigma)$ after the bit flip error has occurred are still diagonal, so computing the square roots is still trivial. We compute like before

$$\sqrt{\rho} = \begin{pmatrix} \sqrt{1+2p/2} & 0\\ 0 & \sqrt{3-2p/2} \end{pmatrix} \tag{40}$$

$$\sqrt{\rho}\sigma\sqrt{\rho} = \begin{pmatrix} (1+p)(1+2p)/12 & 0\\ 0 & (3-2p)(2-p)/12 \end{pmatrix}$$
 (41)

$$\sqrt{\sqrt{\rho}\sigma\sqrt{\rho}} = \begin{pmatrix} \sqrt{(1+p)(1+2p)}/(2\sqrt{3}) & 0\\ 0 & \sqrt{(3-2p)(2-p)}/(2\sqrt{3}) \end{pmatrix} . (42)$$

This gives us a fidelity

$$F(\mathcal{E}(\rho), \mathcal{E}(\sigma)) = \frac{1}{2\sqrt{3}} \left(\sqrt{(1+p)(1+2p)} + \sqrt{(3-2p)(2-p)} \right). \tag{43}$$

One can check that $F(\mathcal{E}(\rho), \mathcal{E}(\sigma)) \geq F(\rho, \sigma)$. We can see from the attached plot that if no error happens p = 1, then the fidelity is unchanged as it should. Also if the error happens with certainty p = 0, the fidelities are equal again. Otherwise the presence of the bit flip error makes the states more similar to each other.

