
PHYS-C0252 - Quantum Mechanics

Exercise set 3

Due date: May 15, 2024 by 23:59 on [MyCourses](#)

Return the exercises as a .pdf.

You can write by hand and take pictures, use digital note-taking or LaTeX etc.

1. Consider the temporal evolution of quantum states $\{|\psi\rangle\}$ determined by a Hamiltonian $\hat{H}(t)$. We define the evolution operator $\hat{U}(t, t_0)$ as

$$\hat{U}(t, t_0)|\psi(t_0)\rangle = |\psi(t)\rangle, \text{ for all } t \geq t_0 \in \mathbb{R} \text{ and } |\psi(t_0)\rangle \in \mathcal{H}$$

- (a) Prove the following identities:

- $\hat{U}(t_0, t_0) = \hat{I}$, for all $t_0 \in \mathbb{R}$.
- $\hat{U}^\dagger(t, t_0)\hat{U}(t, t_0) = \hat{I}$, for all $t > t_0 \in \mathbb{R}$. Hint: Calculate $\partial_t[\hat{U}^\dagger(t, t_0)\hat{U}(t, t_0)]$, use the Schrödinger equation equivalent for $\hat{U}(t, t_0)$, and use the result of i. as an initial condition. Note that we do not assume the Hamiltonian to be independent of time.
- $\hat{U}(t_2, t_0) = \hat{U}(t_2, t_1)\hat{U}(t_1, t_0)$, for all $t_0 < t_1 < t_2 \in \mathbb{R}$.

- (b) Construct the evolution operator $\hat{U}(t, t_0)$ for a time-dependent Hamiltonian $\hat{H}(t)$ using the following steps:

1. Using the Schrödinger equation for U , show that U is formally given by

$$U(t, t_0) = \hat{I} + \sum_{n=1}^{\infty} \left(\frac{-i}{\hbar}\right)^n \int_{t_0}^t dt_1 \int_{t_0}^{t_1} dt_2 \cdots \int_{t_0}^{t_{n-1}} dt_n \hat{H}(t_1)\hat{H}(t_2) \cdots \hat{H}(t_n).$$

2. Find $U(t, t_0)$ when \hat{H} is time-independent using the fact that

$$\int_{t_0}^t dt_1 \int_{t_0}^{t_1} dt_2 \cdots \int_{t_0}^{t_{n-1}} dt_n C = \frac{1}{n!} \int_{t_0}^t dt_1 \int_{t_0}^t dt_2 \cdots \int_{t_0}^t dt_n C$$

for a constant C .

3. Find an expression for $U(t, t_0)$ when \hat{H} is piecewise constant in time. Use the above result, divide the interval $[t_0, t]$ into n intervals $[t_0, t_1]$, $(t_1, t_2]$, \dots , $(t_{n-1}, t_n]$, where $t_k = t_0 + k\delta t$, $t_n = t$, and assume that \hat{H} is constant in each interval and that $\hat{H}(t_k)$ commute.
4. Take the limit $\delta t \rightarrow 0$ to convert the sum in the above result to an integral in order to obtain the final result for $U(t, t_0)$.

2. Let \vec{a} be any real-valued three-dimensional unit vector and $\theta \in \mathbb{R}$.

- (a) Prove that

$$e^{i\theta(\vec{a} \cdot \hat{\sigma})} = \hat{I} \cos(\theta) + i(\vec{a} \cdot \hat{\sigma}) \sin(\theta),$$

where $\vec{a} \cdot \hat{\sigma} = a_x \hat{\sigma}_x + a_y \hat{\sigma}_y + a_z \hat{\sigma}_z$, $\{a_k\}$ are the Cartesian components of \vec{a} , and $\{\hat{\sigma}_k\}$ are the Pauli matrices. Hint: Use the Taylor series definition of the operator exponential. Then, try to find a simplified expression for $(\theta \vec{a} \cdot \hat{\sigma})^k$, depending on whether k is even or odd.

- (b) Consider the operator $e^{-i\frac{\pi}{2}(\vec{a} \cdot \vec{\sigma})}$. Show that the case

$$\vec{a} = \vec{z} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

corresponds to a rotation by the angle π about the x axis of the Bloch sphere. Hint: operate on a general state $|\psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle$ and compare the Bloch vectors of the original and resultant states. Also note that you can always ignore the global phase, that is,

$$e^{i\phi_1} |0\rangle + e^{i\phi_2} |1\rangle = \underbrace{e^{i\phi_1}}_{\text{global phase}} (|0\rangle + e^{i(\phi_2 - \phi_1)} |1\rangle)$$

has the same Bloch sphere representation as

$$|0\rangle + e^{i(\phi_2 - \phi_1)} |1\rangle.$$

If you want, you can also show that $\vec{a} = \vec{x}$ and $\vec{a} = \vec{y}$ correspond to rotations about the x and y axis, respectively (but points are awarded only for the case $\vec{a} = \vec{z}$).

3. Consider a system described by the Hamiltonian $\hat{H} = \epsilon(-i|0\rangle\langle 1| + i|1\rangle\langle 0|)$, where $\{|0\rangle, |1\rangle\}$ form an orthonormal basis of the considered Hilbert space and ϵ is a real-valued constant with the dimension of energy. The eigenenergies of \hat{H} are $\pm\epsilon$ and the corresponding eigenstates are

$$|\epsilon\rangle = \frac{|0\rangle + i|1\rangle}{\sqrt{2}}, \quad |-\epsilon\rangle = \frac{|0\rangle - i|1\rangle}{\sqrt{2}}.$$

- (a) (2 points) What are the probabilities to measure ϵ and $-\epsilon$ if the system is in the state $|1\rangle$? What is the value of $\langle H \rangle$?
- (b) (4 points) Find the state $|\psi(t)\rangle$ at an arbitrary time t , when the system is initially in the state $|\psi(t=0)\rangle = |0\rangle$. What is the probability to find the system in the state $|0\rangle$ as a function of time? How does $\langle H \rangle$ change with time? Hint: you need to solve the Schrödinger equation.
- (4) Consider a peculiar gate known as square-root of NOT with the matrix representation

$$\sqrt{\text{NOT}} \triangleq \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}.$$

- (a) Is this gate unitary? Is it Hermitian? How does it act on the $|0\rangle$ and $|1\rangle$ basis states?
- (b) Show that the square-root of NOT gate is worthy of its name by verifying that $(\sqrt{\text{NOT}})^2 = \sigma_x$.

(c) Let us define the action of the following gates on a two-qubit system:

$$\text{NOT}^{(2)} |x0\rangle = |x1\rangle, \quad \text{NOT}^{(2)} |x1\rangle = |x0\rangle, \quad (1)$$

$$H^{(2)} |x0\rangle = \frac{1}{\sqrt{2}} (|x0\rangle + |x1\rangle), \quad H^{(2)} |x1\rangle = \frac{1}{\sqrt{2}} (|x0\rangle - |x1\rangle), \quad (2)$$

$$\text{CNOT}^{(2,1)} |x0\rangle = |x0\rangle, \quad \text{CNOT}^{(2,1)} |01\rangle = |11\rangle, \quad \text{CNOT}^{(2,1)} |11\rangle = |01\rangle, \quad (3)$$

where $x \in \{0, 1\}$, so for instance $\text{NOT}^{(2)} |00\rangle = |01\rangle$. The matrix representation of $\text{NOT}^{(2)}$ is

$$\text{NOT}^{(2)} \triangleq \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix},$$

which you can verify by operating on the basis states and checking that they match the definition (1). Find the matrix representations for $H^{(2)}$ and $\text{CNOT}^{(2,1)}$.

- (d) Either by using the action of the gates defined in equations (1-3) or by using the matrices derived in the previous problem, find a series of gates to construct the Bell state $\frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$ starting from $|00\rangle$.