

Diffusion, growth, and elasticity in batteries

A mathematical modelling perspective

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Lithium-ion Battery

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Background

Model

Buckling

Limiting
growth

Upscaling

Forerunning energy storage medium for portable electronic devices.

- Lithium is the lightest metal known.
- High energy capacity – run for long periods between chargings.

Biggest challenge: Use them for electric vehicles? Replace fossil fuels?

- ➡ Even higher energy capacity needed.
- ➡ Search for 'new' battery materials ...

Battery materials

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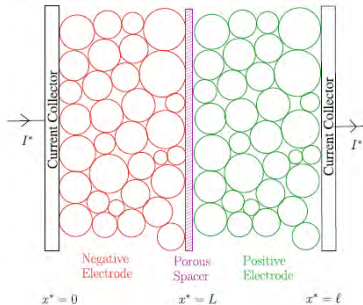
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Most common commercial battery materials:

Negative electrode (Anode): Graphite

Positive electrode (Cathode): Lithium Cobalt Oxide (LiCoO_2)

Electrolyte: Salt of lithium

Figure from G. Richardson, G. Denuault, C. Please, J. Eng. Math **72**, 41 (2012).

Silicon for anodes

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Why silicon?

Silicon can accommodate up to 4.4 atoms of Lithium in the fully lithiated state

➡ Theoretical specific capacity value of 4200 mAhg^{-1} (Graphite: 372 mAhg^{-1})

However ...

Silicon swells up during lithiation - 310% when fully lithiated

➡ Cyclic charging/discharging leads to mechanical failure (fracture)

A way out ...

Use nanostructured silicon electrode particles

Overview of the problem

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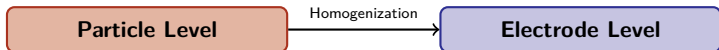
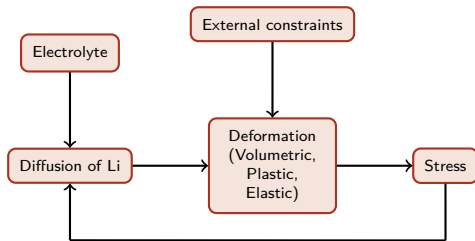
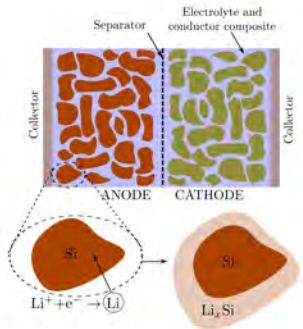
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Model Problem

Electrode particle level

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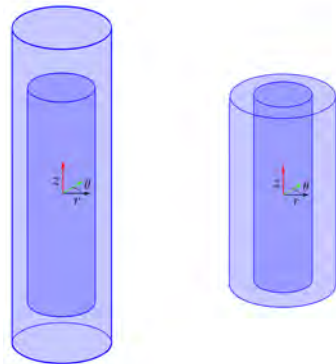


Figure: Left: Unconstrained case. Right: Constrained case – No net axial deformation, edges free to move laterally.

Description of Model

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Diffusion

Stress

Plastic stretch

Description of Model

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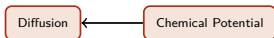
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Stress

Plastic stretch

$$\mathbf{j} = -Dc\nabla\mu$$

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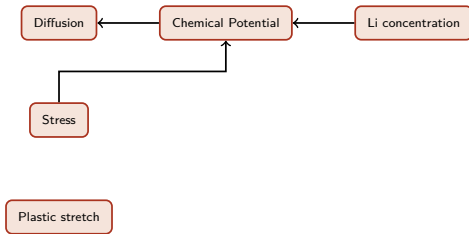
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$$\mathbf{j} = -Dc\nabla\mu$$

$$\mu = \mu_{\text{conc}} + \mu_{\text{stress}}$$

$$\mu_{\text{conc}} = \mu_{\text{conc}}^0 + \log(\gamma c)$$

$$\mu_{\text{stress}} = \frac{\partial W}{\partial c}$$

Description of Model

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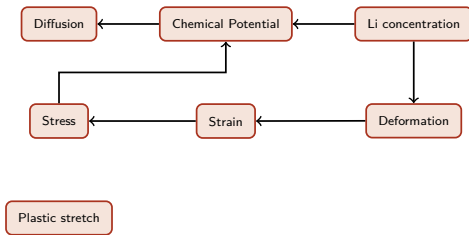
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$$\boldsymbol{\sigma}^0 = \frac{\partial W}{\partial \mathbf{F}} = f(c, \mathbf{E}^e)$$

$$\mathbf{F} = \mathbf{F}^p \mathbf{F}^e \mathbf{F}^c = \mathbf{I} + \nabla \mathbf{u}$$

$$\mathbf{E}^e = \frac{1}{2} \left(\mathbf{F}^{eT} \mathbf{F}^e - \mathbf{I} \right)$$

$$\mathbf{F}^c = (1 + 3\bar{\eta}c)^{1/3} \mathbf{I}$$

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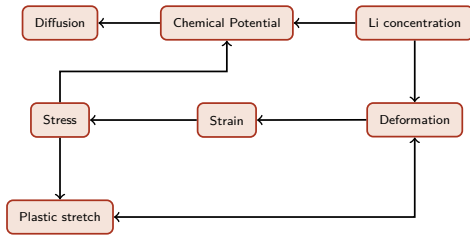
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$$\mathbf{F}^p = g(\boldsymbol{\sigma}^0)?$$

$$\det(\mathbf{F}^p) = 1$$

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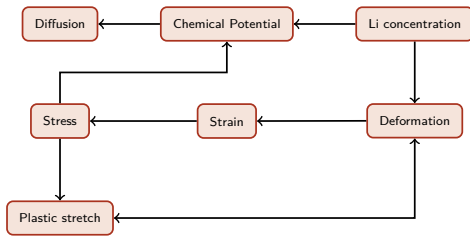
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Net axial force is zero.



Net axial displacement is zero.

$$\mathbf{j} = -Dc\nabla\mu$$

$$\mu = \mu_{\text{conc}} + \mu_{\text{stress}}$$

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Governing equations

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Diffusion:

$$\frac{\partial c}{\partial t} = -\mathbf{Div} \mathbf{j} \quad \rightarrow \quad \frac{\partial c}{\partial t} = -\frac{\partial j_r}{\partial r} - \frac{j_r}{r},$$
$$c(r, 0) = 0, \quad j_r(0, t) = 0, \quad j_r(1, t) = J_0(1 - c)$$

Mechanical equilibrium:

$$\mathbf{Div} \boldsymbol{\sigma}^0 = 0 \quad \rightarrow \quad \frac{\partial \sigma_r^0}{\partial r} + \frac{\sigma_r^0 - \sigma_\theta^0}{r} = 0,$$
$$u(0, t) = 0, \quad \sigma_r^0(1, t) = 0.$$

Plastic stretch evolution:

$$\frac{\partial \lambda_{r,\theta}}{\partial t} = \lambda_{r,\theta} \text{Pf} \left(\frac{\sigma_{\text{eff}}}{\sigma_f} - 1 \right)^m \frac{\tau_{r,\theta}}{||\boldsymbol{\tau}||} H \left(\frac{\sigma_{\text{eff}}}{\sigma_f} - 1 \right),$$
$$\lambda_{r,\theta}(r, 0) = 1.$$

Motivation to study buckling

An empirical evidence

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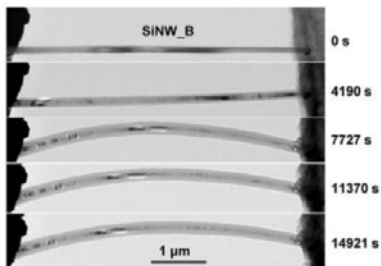


Figure: Buckling of a Si nanowire electrode particle as it undergoes lithiation.
(Reproduced from Liu *et al.*, ACS Nano, 7, 1495–1503, 2013.)

Buckling

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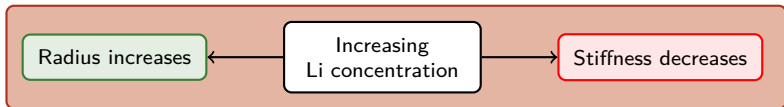
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General Considerations:

	LOW	HIGH
Length	✓	✗
Radius	✗	✓
Stiffness	✗	✓

For Si electrode particle:



↓
Modifications in the classical buckling criterion

Results

Buckling

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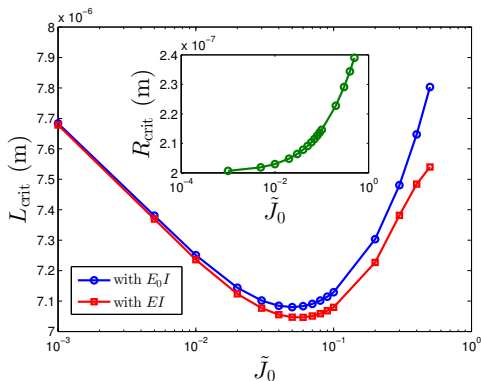
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- ➡ Competition between growing stress and growing radius
- ➡ Predominating influence of growing radius

Modified buckling criteria

$$F_{\text{crit}} = \frac{\pi^2 EI}{(KL)^2}$$

E : Modulus of elasticity

$$I = \frac{\pi R^4}{4} : 2^{\text{nd}} \text{ moment of area}$$

$$L_{\text{buck}} = \{L : F_{\text{crit}} \text{ exists}\}$$

$$L_{\text{crit}} = \min(L_{\text{buck}})$$

Modifications:

$$R = R_0 \left(1 + \frac{\partial u}{\partial r}\right)$$

$$E = E_0 (1 + \eta_E x_{\text{max}} c)$$

$E_0 I$: Only R changed

EI : Both E and R changed

Ref. : J. Chakraborty et al., [Int. J. Solids Struct.](#) **54**, 66–81 (2015).

An estimate for design

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Approximate the critical compressive force using the yield stress

➡ Determine critical length

$$\sigma_f \pi R_0^2 = \frac{\pi^2 E_0}{L^2} \left(\frac{\pi R_0^4}{4} \right)$$

$$\text{Or, } L = \frac{1}{2} \sqrt{\frac{E_0}{\sigma_f}} \pi R_0 \approx 8 \mu\text{m}$$

Size-effects

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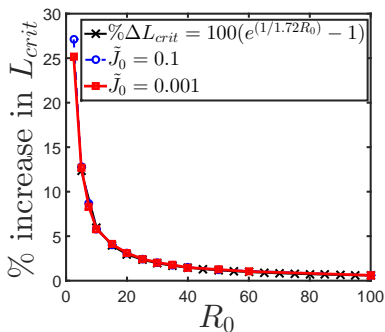
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Young's modulus influenced by lithium concentration

Young's modulus *can also* depend on particle size (radius < 30 nm)

$$E = E_0(1 + \chi_c)(1 + \chi_s)$$

χ_s determined from Bond-Order-Length-Strength (BOLS) theory



Sun et al., *J. Phys. Condens. Matt.* **14**, 7781 (2002).

Neogi and Chakraborty, *J. Appl. Phys.* **124**, 154302 (2018).

Limiting axial growth

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Is there any simple way to limit axial growth?



Use constraints

Two configurations

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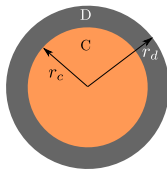
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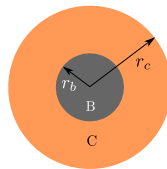
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Outer constraint



Inner constraint

■ Silicon

■ Constraining material

Phase diagrams

From numerical simulations

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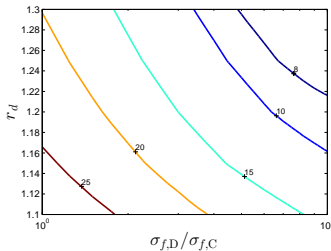
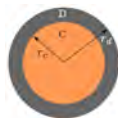
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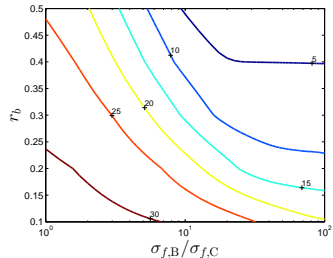
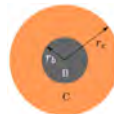
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Outer constraint



Inner constraint

Yield stress of Region D: $\sigma_{f,D}$

Yield stress of Region C: $\sigma_{f,C}$

Yield stress of Region B: $\sigma_{f,B}$

Contour values indicate percentage increase in length

Ref: J. Chakraborty et al., *J. Power Sources* **279**, 746–758 (2015).

Upscaling

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Electrode Particle Level

➡ Full set of equations



Simplified set of equations



Upscale using Homogenization



Battery level

An open question

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How can we incorporate large strains in the upscaling steps?

➡ Apply to silicon anodes

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Sourav Chattaraj
Jay Kishan



Prof. S. Jon Chapman
Prof. Alain Goriely
Prof. Colin P. Please



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THANK YOU