

Entanglement purification of any stabilizer state

S. Glancy* and E. Knill†

*Mathematical and Computing Science Division, Information Technology Laboratory, National Institute of Standards and Technology,
Boulder, Colorado 80301, USA*

H. M. Vasconcelos‡

Department of Physics, University of Notre Dame, Notre Dame, Indiana 46556, USA

(Received 19 June 2006; published 19 September 2006)

We present a method for multipartite entanglement purification of any stabilizer state shared by several parties. In our protocol each party measures the stabilizer operators of a quantum error-correcting code on his or her qubits. The parties exchange their measurement results, detect or correct errors, and decode the desired purified state. We give sufficient conditions on the stabilizer codes that may be used in this procedure and find that Steane's seven-qubit code is the smallest error-correcting code sufficient to purify any stabilizer state. An error-detecting code that encodes two qubits in six can also be used to purify any stabilizer state. We further specify which classes of stabilizer codes can purify which classes of stabilizer states.

DOI: 10.1103/PhysRevA.74.032319

PACS number(s): 03.67.Mn, 03.67.Pp

I. INTRODUCTION

In this paper we describe a method for entanglement purification that is able to purify any stabilizer or graph state. The goal of a purification protocol is to increase the purity of a mixed state whose qubits are divided among several parties able to communicate only through classical channels. Specifically, suppose the state $|\psi\rangle$ is a pure entangled state of n qubits and each qubit of $|\psi\rangle$ is sent to a different party (Alice, Bob, Charlie, ...). In this initial stage, the qubits are transmitted through noisy quantum channels, so there is some probability that each may be affected by Pauli σ_x , σ_y , or σ_z errors. After the transmission, $|\psi\rangle$ becomes the mixed state $\hat{\rho}$. If m copies of $|\psi\rangle$ are prepared and transmitted, so that each party holds m qubits, it may be possible to obtain one or more copies of $|\psi\rangle$ with less noise (more purity). For this purpose, the parties may use local quantum operations and measurements on their own qubits and classical communication.

The entanglement purification (or “entanglement distillation”) problem was first studied by Bennett and co-authors in 1995 [1]. They describe a method for two separate parties to purify a Bell pair by use of local operations on copies of noisy Bell pairs and classical communication between the two parties. Since then many other researchers have studied this problem, introducing new protocols for the purification of Bell pairs and providing methods for purifying other classes of entangled states. For examples of this research see Refs. [2–11].

In the remainder of this section we give a brief review of stabilizers and introduce our method of using operator arrays to characterize many copies of entangled states shared by

several parties. In Sec. II we discuss the use of quantum error-detecting and -correcting codes for purification. In Sec. III we give classes of stabilizer codes that can be used to purify specific classes of states including the class of all stabilizer states. In Sec. IV we discuss these results and make some concluding remarks.

A. Stabilizer review

When using the stabilizer formalism, instead of writing out vectors in a Hilbert space, we specify a quantum state $|\psi\rangle$ using (i) a set of operators that have $|\psi\rangle$ as an eigenstate and (ii) the eigenvalues of $|\psi\rangle$ under each such operator (see Chap. 10 of Ref. [12]). This set of operators is the stabilizer of $|\psi\rangle$, and we say “ $|\psi\rangle$ is stabilized by” this set. Although the stabilizer often refers to a set of operators whose eigenvalue is one, in this paper we allow the stabilizer to include operators with other eigenvalues. Several states may share the same stabilizer, but have different eigenvalues for the members of the stabilizer. We are interested in states determined by stabilizers consisting of the set of stabilizing Pauli products—tensor products of Pauli matrices including the identity matrix. Such states are called “stabilizer states.” Stabilizer states are equivalent to “graph states” [13], where the latter are specified by means of graphs rather than stabilizers. The stabilizer of stabilizer states necessarily consists of commuting Pauli products. It forms a projective group (closed under multiplication up to a phase), so it is sufficient to specify its generators. The number of independent generators must equal the number of qubits. Because Pauli matrices have eigenvalues ± 1 , each generator must also have eigenvalue ± 1 .

For example, the Bell state $|B_{00}\rangle = |00\rangle + |11\rangle$ (with normalization omitted) has stabilizer generators XX and ZZ and eigenvalues $+1$ and $+1$, respectively. Here we use the abbreviations X , Y , and Z for the Pauli σ_x , σ_y , and σ_z matrices.

*Electronic address: sglancy@boulder.nist.gov

†Electronic address: knill@boulder.nist.gov

‡Electronic address: Hilma.M.DeVasconcelos.2@nd.edu

Expressions such as XY refer to the Pauli product where σ_x and σ_y act on the first and second qubits, respectively. We can specify $|B_{00}\rangle$ by means of its stabilizer and eigenvalues as follows:

$$|B_{00}\rangle = \begin{bmatrix} X & X \\ Z & Z \end{bmatrix}, \begin{bmatrix} +1 \\ +1 \end{bmatrix}. \quad (1)$$

Note that these are not numeric matrices in the square brackets. We are simply listing each generator and its corresponding eigenvalue on each row. The stabilizers of the three other Bell states are also generated by XX and ZZ , but they have different eigenvalues. The three qubit Greenberger-Horne-Zeilinger (GHZ) [14] state is

$$|000\rangle + |111\rangle = \begin{bmatrix} X & X & X \\ Z & Z & I \\ Z & I & Z \end{bmatrix}, \begin{bmatrix} +1 \\ +1 \\ +1 \end{bmatrix}. \quad (2)$$

The Bell states and the GHZ states are both in the class of CSS states (named after Calderbank, Shore, and Steane). A stabilizer state is a CSS state if it can be transformed with unitary single qubit operations into a form in which each of its stabilizer generators can be written using only X 's and I 's or Z 's and I 's (the “CSS form”). The class of CSS states is equivalent to the two-colorable graph states [15]. Methods for purifying any CSS state are already known [5,9].

Consider the state

$$|\Delta\rangle = \begin{bmatrix} X & Z & Z \\ Z & X & Z \\ Z & Z & X \end{bmatrix}, \begin{bmatrix} +1 \\ +1 \\ +1 \end{bmatrix}. \quad (3)$$

This state cannot be written in the CSS form, and it is impossible to transform $|\Delta\rangle$ into a CSS state by using single qubit operations. Readers familiar with the techniques of graph states may recognize that this state has a triangle graph.

We can also use the stabilizer formalism to describe quantum error-detecting and -correcting stabilizer codes. (Since all codes considered here are stabilizer codes, from now on we omit the modifier “stabilizer.”) To define a code, the stabilizer generators are used to specify a (more than one-dimensional) subspace into which one may encode quantum information. In this case, the subspace consists of the states with identical eigenvalues for each generator. If we want to encode m logical qubits using n physical qubits, the code is specified by $n-m$ independent generators of the stabilizer. For example the four qubit error-detecting code C_4 encodes two logical qubits using four physical qubits [16] and is described by the stabilizer with generators

$$C_4 \Leftrightarrow \begin{bmatrix} X & X & X & X \\ Z & Z & Z & Z \end{bmatrix}, \begin{bmatrix} +1 \\ +1 \end{bmatrix}. \quad (4)$$

The Hilbert space of the four physical qubits contains 16 dimensions, but the set of states with +1 eigenvalues of the

above operators is a four-dimensional subspace. Four dimensions are sufficient to contain two logical qubits, which we specify by giving their logical Pauli operators:

$$X_L^{(1)} = XXII, \quad (5a)$$

$$Z_L^{(1)} = ZIZI, \quad (5b)$$

$$X_L^{(2)} = IXIX, \quad (5c)$$

$$Z_L^{(2)} = IIZZ. \quad (5d)$$

The encoded logical operators have all the algebraic relationships that we expect from the X and Z operators on two qubits. Because the state of this four qubit system is always confined to the +1 eigenstate of the stabilizers, there are many equivalent representations of the logical qubit operators, which we can obtain by multiplying the logical operators by members of the stabilizer group. For example we can also write

$$X_L^{(1)} \simeq XXXX \times XXII \simeq IIXX, \quad (6)$$

because each of these alternatives has the same effect on the encoded logical qubits.

C_4 is an error-detecting code. An X error on any single physical qubit results in a state for which the eigenvalue of the generator $ZZZZ$ is changed to -1 . A Z error will change the eigenvalue of $XXXX$ to -1 , and a Y error changes both eigenvalues. To detect errors we simply measure each of the code’s generators. We call each such measurement a “parity check,” and the operator being measured is the “parity check operator.” We assume that the measurements are projective so that the state after a measurement is a projection of the initial state onto one of the two eigenspaces of the parity check operator. From each parity check we obtain an eigenvalue, which must be ± 1 . The “syndrome” is the vector of eigenvalues for the generators. The eigenvalue for any parity check operator can be inferred from the syndrome. Because for C_4 , the syndrome does not tell us which qubit received the error, we are unable to correct one-qubit errors.

Codes such as C_4 with the property that there is a choice of stabilizer generators that can be written with only X 's and I 's or Z 's and I 's are called “CSS codes.”

The smallest error-correcting code that protects a single logical qubit requires five physical qubits [2,17]. It has the generators

$$C_5 \Leftrightarrow \begin{bmatrix} X & Z & Z & X & I \\ I & X & Z & Z & X \\ X & I & X & Z & Z \\ Z & X & I & X & Z \end{bmatrix}, \begin{bmatrix} +1 \\ +1 \\ +1 \\ +1 \end{bmatrix}. \quad (7)$$

This code has the logical operators

$$X_L = XXXXX, \quad (8a)$$

$$Z_L = ZZZZZ, \quad (8b)$$

and is not a CSS code.

Another code that we use is Steane's seven qubit code \mathcal{C}_7 , which encodes one logical qubit in seven physical qubits [18]. It has the generators

$$\mathcal{C}_7 \Leftrightarrow \begin{bmatrix} X & I & X & I & X & I & X \\ I & X & X & I & I & X & X \\ I & I & I & X & X & X & X \\ Z & I & Z & I & Z & I & Z \\ I & Z & Z & I & I & Z & Z \\ I & I & I & Z & Z & Z & Z \end{bmatrix}, \quad (9)$$

and the logical qubit operators

$$X_L = XXXXXX, \quad (10a)$$

$$Z_L = ZZZZZZ. \quad (10b)$$

Any single qubit error changes the syndrome. The syndrome gives sufficient information to identify which of the qubits received the error, so it can then be corrected. This code is a CSS code, and it has the additional property that it is Hadamard invariant. In particular, if we apply the Hadamard H operator to every qubit ("transversal Hadamard"), then the code is unchanged because H performs the transformation

$$H: X \rightarrow Z, \quad (11a)$$

$$Z \rightarrow X. \quad (11b)$$

Consequently the stabilizer is transformed into itself with no change in eigenvalues for the generators. Furthermore, the transversal Hadamard exchanges logical X and Z operators, from which we infer that its effect on the code is a logical Hadamard operator. We call CSS codes whose set of stabilizers is Hadamard invariant and for which there are logical operators in CSS form with respect to which the transversal Hadamard acts as a logical Hadamard on each encoded qubit, "CSS- H codes." Note that \mathcal{C}_4 is not a CSS- H code. The seven-qubit code is the smallest nontrivial CSS- H error-correcting code. However, there is a six-qubit CSS- H error-detecting code encoding two logical qubits, which was described in Ref. [19]. It has generators

$$\mathcal{C}_6 \Leftrightarrow \begin{bmatrix} X & I & X & X & I & X \\ I & X & X & I & X & X \\ Z & I & Z & Z & I & Z \\ I & Z & Z & I & Z & Z \end{bmatrix}, \quad (12)$$

and the logical qubit operators

$$X_L^{(1)} = XXXIII, \quad (13a)$$

$$Z_L^{(1)} = ZZZIII, \quad (13b)$$

$$X_L^{(2)} = IIIXXX, \quad (13c)$$

$$Z_L^{(2)} = IIIZZZ. \quad (13d)$$

We have chosen different logical operators from those used in Ref. [19], so that the code is explicitly H invariant. \mathcal{C}_6 detects errors, and if we know whether the error is in the first or second group of three qubits, we can also correct it. Otherwise we know what type of error has affected the logical qubits, but not which logical qubit was affected. \mathcal{C}_6 is the smallest CSS- H error-detecting code.

B. Operator arrays

In entanglement purification we have n parties and m copies of a large entangled state. The copies are prepared and distributed so that each party holds one qubit of each copy, for a total of m qubits per party. The stabilizer group of the full state of this $m \times n$ qubit system has $m \times n$ generators, each of which is composed of $m \times n$ Pauli matrices. For example, if two copies of the Bell pair $|B_{00}\rangle$ are shared by Alice and Bob, the generators for this four-qubit system are

$$\left[\begin{array}{cccc} A1 & A2 & B1 & B2 \\ \hline X & I & X & I \\ Z & I & Z & I \\ I & X & I & X \\ I & Z & I & Z \end{array} \right], \quad (14)$$

where we use the top line in the chart to label each qubit as belonging to Alice or Bob and copy one or two of the shared state. Notice that qubits $A1$ and $B1$ are entangled with one another but not with $A2$ and $B2$.

We would like a method for representing the stabilizer generators that emphasizes the structure of these states. Instead of writing each generator as a single row in a table, we write each generator as an array whose columns belong to a particular party and whose rows represent a particular copy of the shared entangled state. The set of generators is a list of such arrays. Using these operator arrays, we represent the two copies of the Bell pair in Eq. (14) with

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline Z & Z \\ I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ X & X \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ Z & Z \end{array} \right]. \quad (15)$$

Here we can easily see that entanglement exists only within rows; row one is never entangled with row two. We call the generators of the state that has been copied (in this case XX and ZZ) the “master generators.” The members of Eq. (15) are “single-copy generators.” We can see that each single-copy generator has identities in all rows except one, which contains a master generator. The objects listed in Eq. (15) generate the “multicopy stabilizer.” A useful subset of the multicopy stabilizer is the set of “parallel stabilizer elements,” which are products of single-copy generators having the same master generator. These parallel stabilizer elements have rows equal to the identity or only one of the master generators.

II. ERROR-CORRECTING CODES FOR PURIFICATION

In this section we describe a general method for understanding entanglement purification protocols using quantum error-correcting codes. These ideas were described in Refs. [4,6]. The essence of the scheme is that each party measures the parity check operators of an error-correcting code on his or her qubits. By comparing the results of these measurements the parties learn about the errors that have corrupted their states. They correct the errors and then decode logical qubits into the desired entangled state.

To begin a purification protocol, the parties apply random stabilizer elements to the noisy states they hope to purify. Each party just applies an agreed-upon single qubit Pauli matrix to his or her qubit of a particular state. These Pauli matrices make up a random stabilizer element that the parties determine by classical communication before the start of the protocol. After doing this, they can treat any noisy state as a probabilistic mixture of states that have the same stabilizer but different eigenvalues for the generators. This fact was proven for Bell states in Ref. [1] and was extended to any stabilizer state by Aschauer, Dür, and Briegel in Ref. [5]. We include the proof here for pedagogical completeness.

Let $|\psi_0\rangle$ be the desired state for which all of the generators eigenvalues are +1. Any pure noisy state can be written as

$$|\psi\rangle = \sum_i \alpha_i P_i |\psi_0\rangle, \quad (16)$$

where i extends over all possible syndromes, P_i is a tensor product of Pauli matrices that moves the state from the all +1 syndrome to syndrome i , and α_i is an amplitude. If there is little noise, the α_i are small for P_i not equal to identity. The density matrix for this state is

$$|\psi\rangle\langle\psi| = \sum_{i,j} \alpha_i \alpha_j^* P_i |\psi_0\rangle\langle\psi_0| P_j^\dagger. \quad (17)$$

An arbitrary noisy mixed state is just a probabilistic mixture of noisy pure states, which we can write as

$$\rho = \sum_k p_k \sum_{i,j} \alpha_{ik} \alpha_{jk}^* P_i |\psi_0\rangle\langle\psi_0| P_j^\dagger, \quad (18)$$

where $\sum_k p_k = 1$. Alice and her friends now apply a random element of the stabilizer group to this state. After “forgetting” which of the N stabilizer elements they applied, the state becomes

$$\rho_Q = \frac{1}{N} \sum_{Q \in \text{stab.}} Q \rho Q^\dagger, \quad (19)$$

$$= \frac{1}{N} \sum_{ijk} p_k \alpha_{ik} \alpha_{jk}^* \sum_{Q \in \text{stab.}} Q P_i |\psi_0\rangle\langle\psi_0| P_j^\dagger Q^\dagger. \quad (20)$$

Because all of the Q ’s and P ’s are made of tensor products of Pauli matrices, they must either commute or anticommute with one another. Let $\langle Q, P \rangle = 0$ if P and Q commute and $\langle Q, P \rangle = 1$ if they anticommute. After commuting the P ’s and Q ’s we can use the fact that Q is in the stabilizer of $|\psi_0\rangle$ to obtain

$$\rho_Q = \frac{1}{N} \sum_{ijk} p_k \alpha_{ik} \alpha_{jk}^* \sum_{Q \in \text{stab.}} (-1)^{\langle Q, P_i P_j \rangle} P_i |\psi_0\rangle\langle\psi_0| P_j^\dagger. \quad (21)$$

Let us examine the sum $\sum_{Q \in \text{stab.}} (-1)^{\langle Q, P_i P_j \rangle}$ over all Q in the stabilizer of $|\psi_0\rangle$ in the case where $i \neq j$. The elements of the stabilizer that commute with $P_i P_j$ form a subgroup, \mathcal{Q}_{ij} . There must be some element of the stabilizer that anticommutes with $P_i P_j$ (otherwise $P_i P_j$ would itself be in the stabilizer, in which case $i=j$), let us call this element q . We can now generate every element of the stabilizer that anticommutes with $P_i P_j$ by multiplying every element of \mathcal{Q}_{ij} by q . (The elements that anticommute are a coset of \mathcal{Q}_{ij} .) Therefore the number of elements of the stabilizer that commute with $P_i P_j$ is equal to the number of elements of the stabilizer that anticommute with $P_i P_j$. Consequently the terms of Eq. (21) in the sum over all Q for which $i \neq j$ must all cancel one another leaving us with

$$\rho_Q = \frac{1}{N} \sum_{ik} p_k |\alpha_{ik}|^2 P_i |\psi_0\rangle\langle\psi_0| P_i^\dagger. \quad (22)$$

This is just a mixture, each of whose terms are eigenstates of the stabilizer operators with different eigenvalues. Therefore it would be sufficient to measure each of the generators to extract an error syndrome, correct errors and, if all the errors are corrected, obtain a pure state. However, because each party holds only a single qubit, they cannot measure the generators directly.

In the picture of the operator arrays, each party can measure any operator that has nonidentity entries only along his or her column, but the master generators describing the state are oriented along rows. We can use techniques of error-correcting codes to overcome this problem. Each party will use an error-detecting or -correcting code that encodes one or more logical qubits on the m physical qubits. We will assume

that every party will use the same code. They each measure the generators of that code and then share the measurement results. Depending on the construction of the code and the original state they are trying to purify, they expect a particular pattern of correlations between their measurement results. Errors in the transmission of the m entangled states should appear as aberrations in the syndrome patterns, so they can

be detected or corrected. The parties then have an encoded copy of a state, which they can decode.

Let us consider an example. Suppose Alice and Bob wish to purify Bell pairs $|B_{00}\rangle$, and they share four copies of noisy Bell pairs. We first examine the case in which no errors are present. The single-copy generators of the full eight qubit state are

$$\begin{aligned} & \left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & I & I \\ 3 & I & I \\ 4 & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & Z & Z \\ I & I & I \\ I & I & I \\ I & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & I \\ X & X & X \\ I & I & I \\ I & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & I \\ Z & Z & Z \\ I & I & I \\ I & I & I \end{array} \right], \\ & \left[\begin{array}{c|cc} & A & B \\ \hline 1 & I & I \\ 2 & I & I \\ 3 & X & X \\ 4 & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & I \\ I & I & I \\ Z & Z & Z \\ I & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & I \\ I & I & I \\ I & I & I \\ X & X & X \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & I \\ I & I & I \\ I & I & I \\ Z & Z & Z \end{array} \right]. \end{aligned} \quad (23)$$

Alice and Bob can use the \mathcal{C}_4 code to purify, so they each measure the two generators of \mathcal{C}_4 on his and her own qubits. The stabilizer arrays describing these single-party parity checks are

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & I \\ 2 & X & I \\ 3 & X & I \\ 4 & X & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & Z & I \\ Z & I & I \\ Z & I & I \\ Z & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & X \\ I & X & X \\ I & X & X \\ I & X & X \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & I & Z \\ I & Z & Z \\ I & Z & Z \\ I & Z & Z \end{array} \right]. \quad (24)$$

The result of each single-party parity check will be ± 1 with probability $\frac{1}{2}$. However,

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & X & X \\ 3 & X & X \\ 4 & X & X \end{array} \right] \text{ and } \left[\begin{array}{c|cc} & A & B \\ \hline & Z & Z \\ Z & Z & Z \\ Z & Z & Z \\ Z & Z & Z \end{array} \right] \quad (25)$$

are in the multicopy stabilizer of the states Alice and Bob are purifying. The operators given in Eq. (25) are parallel stabilizer elements made by repeating the master generators on multiple rows. They are also examples of “parallel parity checks” in which the same generator of the purifying code is repeated on multiple columns. In the absence of error, the eigenvalues of these operators are $+1$ (because they are in the

stabilizer), so the eigenvalues Alice and Bob obtain for their $XXXX$ measurements must match and similarly for their $ZZZZ$ measurements. Let us assume that after these measurements, Alice and Bob transform their states by applying known Pauli products so that they all have $+1$ eigenvalues for the operators in Eq. (24). Note that this is not strictly necessary as long as Alice and Bob keep track of their Pauli frame, where the Pauli frame is defined by a Pauli product that would restore the system so that all stabilizer generators have eigenvalue $+1$. If Alice and Bob know the Pauli frame they can simply adjust future manipulations of their states to compensate for the changes in eigenvalues without ever applying Pauli product compensations.

We can now find the generators of the new stabilizer that Alice and Bob share after their measurements. The new stabilizer must include all of the measurement operators and all elements of the old stabilizer that commute with the measurements. Eight generators are required and they include all of Eq. (24) and

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & X & X \\ 3 & I & I \\ 4 & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline & Z & Z \\ I & I \\ Z & Z \\ I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ X & X \\ I & I \\ X & X \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ Z & Z \\ Z & Z \end{array} \right]. \quad (26)$$

These are parallel stabilizer elements, so they are in the original multicopy stabilizer. Alice and Bob now each possess two logical qubits encoded in \mathcal{C}_4 . We can see the state of these logical qubits using the logical encoded operators given in Eq. (5). We rewrite the operators in Eq. (26) using the logical qubit operators as

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X_L & X_L \\ 2 & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline Z_L & Z_L \\ I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ X_L & X_L \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ Z_L & Z_L \end{array} \right]. \quad (27)$$

These are the master generators for two Bell pairs shared between Alice and Bob—exactly the state they wanted to purify.

Let us now examine the behavior of this scheme in the presence of an error. An error changes the eigenvalue of one of the single copy generators in Eq. (23) to -1 . Alice and Bob detect this change when they use their single-party parity checks to compile the multiparty parity checks. For example, if Alice's first qubit suffers from a Z error, the eigenvalue of the first single-copy generator in Eq. (23)'s list is -1 . Also, the eigenvalue of the first operator in Eq. (25)'s list is -1 . Therefore the product of the eigenvalues Alice and Bob obtain from their $XXXX$ single-party parity checks must be -1 . They detect this error when they compare measurement results and obtain their multiparty parity check. A Z error to any of the eight qubits will cause this same error syndrome, so Alice and Bob cannot correct it.

This analysis allows us to formulate sufficient conditions on the success of purification schemes of this form: (i) The multiparty parity checks that the parties perform must be sensitive to any change in the eigenvalues of the generators of the states they wish to purify. (ii) The stabilizer of the desired encoded state must be in the original multicopy stabilizer of the qubits to be purified.

III. PURIFYING STABILIZER STATES WITH STABILIZER CODES

In this section we discuss which classes of states may be purified using specified classes of error-correcting codes. The

simplest class of states are the CSS- H states, which can be transformed using local operations into states whose stabilizer generators can be written in CSS form and are H invariant as a set. More complex states are CSS, but not H invariant. The most general class of states we consider includes all stabilizer states. We similarly classify codes as being CSS- H , CSS, or any stabilizer code.

A. Purifying CSS- H states

All CSS- H states must contain an even number of qubits because there are equal numbers of Z -type and X -type generators. The only two, four, and six qubit CSS- H states are collections of Bell pairs, but more complicated states can be formed from eight or more qubits. In the following we describe how any CSS- H state can be purified by use of any error-detecting stabilizer code. We do this by showing that the multicopy stabilizer contains enough versatility to allow any error-detecting stabilizer code to meet the conditions (i) and (ii) stated at the end of the previous section. Matsumoto has already demonstrated that any stabilizer code can be used to purify Bell pairs and maximally entangled bipartite states of qudits [4].

Let us examine the example of Alice and Bob purifying a Bell pair using \mathcal{C}_5 as their purifying code. The Bell pair has master generators XX and ZZ , and the multi-copy stabilizer is generated by the single-copy generators

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & I & I \\ 3 & I & I \\ 4 & I & I \\ 5 & I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline Z & Z \\ I & I \\ I & I \\ I & I \\ I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ X & X \\ I & I \\ I & I \\ I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ Z & Z \\ I & I \\ I & I \\ I & I \end{array} \right], \left[\begin{array}{c|cc} & A & B \\ \hline I & I \\ I & I \\ X & X \\ I & I \\ I & I \end{array} \right],$$

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & I & I \\ 2 & I & I \\ 3 & Z & Z \\ 4 & I & I \\ 5 & I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ I & I \\ I & I \\ X & X \\ I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ I & I \\ I & I \\ Z & Z \\ I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ I & I \\ I & I \\ I & I \\ X & X \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ I & I \\ I & I \\ I & I \\ Z & Z \end{array} \right]. \quad (28)$$

Alice and Bob obtain single-party parity checks by measuring

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & I \\ 2 & Z & I \\ 3 & Z & I \\ 4 & X & I \\ 5 & I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ X & I \\ I & I \\ Z & I \\ X & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline X & I \\ I & I \\ X & I \\ Z & I \\ Z & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline Z & I \\ X & I \\ X & I \\ X & I \\ Z & I \end{array} \right],$$

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & I & X \\ 2 & I & Z \\ 3 & I & Z \\ 4 & I & X \\ 5 & I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ I & X \\ I & I \\ I & Z \\ I & X \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & X \\ I & I \\ I & X \\ I & X \\ I & Z \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & Z \\ I & X \\ I & I \\ I & X \\ I & Z \end{array} \right]. \quad (29)$$

Alice and Bob must now construct multiparty parity checks that are in the multiparty stabilizer and are sensitive to any error that would change the eigenvalue of any of the single-copy generators. They can do this by multiplying their single-party parity checks in parallel, so that they obtain the eigenvalues of

$$\left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & Z & Z \\ 3 & Z & Z \\ 4 & X & X \\ 5 & I & I \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline I & I \\ X & X \\ I & I \\ Z & Z \\ X & X \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline X & X \\ I & I \\ X & X \\ I & I \\ Z & Z \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline Z & Z \\ X & X \\ X & X \\ X & X \\ Z & Z \end{array} \right], \left[\begin{array}{cc} A & B \\ \hline Z & Z \\ Z & Z \\ Z & Z \\ Z & Z \\ Z & Z \end{array} \right]. \quad (30)$$

Each of these operators is in the multicopy stabilizer because each has rows equal to the master generators.

If no errors have occurred, Alice and Bob find that all of the multiparty parity checks give eigenvalues of +1. Suppose for example that Alice's first qubit has suffered from a Z error. Then the eigenvalue of the first member of Eq. (28) will be -1. When Alice and Bob examine their multiparty parity checks, they find that the first and third members of Eq. (30) have eigenvalue -1. In practice it is not even necessary to identify the particular qubit that suffered the error.

They need only to know how to correctly return the encoded space to the subsystem with all +1 eigenvalues for the generators of the stabilizers. They can therefore choose to apply Z to Alice's or Bob's first qubit to correct this error. Any single qubit error will give a different syndrome pattern. However, multiqubit errors (such as an X error to Alice's second qubit and an X error to Bob's fifth qubit) can result in the same syndromes as single qubit errors and would cause Alice and Bob to incorrectly "correct" the errors. However, any combination of errors restricted to a single row of the array can be corrected. Depending on their error model and the particular syndrome result they obtain, Alice and Bob may decide to simply discard their states and try again.

The X and Z Pauli operators for Alice's and Bob's encoded qubits are

$$X_L^{(A)} = \left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & I \\ 2 & X & I \\ 3 & X & I \\ 4 & X & I \\ 5 & X & I \end{array} \right], \quad Z_L^{(A)} = \left[\begin{array}{cc} A & B \\ \hline Z & I \\ Z & I \\ Z & I \\ Z & I \\ Z & I \end{array} \right],$$

$$X_L^{(B)} = \left[\begin{array}{c|cc} & A & B \\ \hline 1 & I & X \\ 2 & I & X \\ 3 & I & X \\ 4 & I & X \\ 5 & I & X \end{array} \right], \quad Z_L^{(B)} = \left[\begin{array}{cc} A & B \\ \hline I & Z \\ I & Z \\ I & Z \\ I & Z \\ I & Z \end{array} \right]. \quad (31)$$

Alice and Bob would like to have purified the state whose encoded generators are the master generators, i.e.,

$$X_L^{(A)} X_L^{(B)} = \left[\begin{array}{c|cc} & A & B \\ \hline 1 & X & X \\ 2 & X & X \\ 3 & X & X \\ 4 & X & X \\ 5 & X & X \end{array} \right], \quad \text{and } Z_L^{(A)} Z_L^{(B)} = \left[\begin{array}{cc} A & B \\ \hline Z & Z \\ Z & Z \\ Z & Z \\ Z & Z \\ Z & Z \end{array} \right]. \quad (32)$$

They will have this state after projecting the multicopy state into the encoded subspace by measuring the code's generators, provided that the encoded master generators were in the

original multicopy stabilizer. This is surely the case because the rows of the encoded master generators' arrays contain only the master generators.

Any other CSS- H state can be purified in a similar manner. We need only specify a method for constructing the multiparty parity checks from the single-party parity checks. Each multiparty parity check should include columns that contain the identity or only one of the generators of the purifying code. Which columns have identity and which contain the generator of the purifying code are determined so that the rows of the multiparty parity check array match a particular master generator (containing X 's), its Hadamard pair (containing Z 's), or the product of a master generator and its Hadamard pair (containing Y 's). For each Hadamard pair of master generators we construct a number of multiparty parity checks equal to the number of generators for the purifying code (four for \mathcal{C}_5). These generators give a syndrome that tells us which copy has received an X , Y , or Z error affecting that Hadamard pair of master generators. Similar syndromes are obtained for each pair of master generators. We now know which copy of the original state received an error and how that error affected each of that copy's pairs of generators. We can use this information to determine a Pauli product to apply to this copy to restore the correct syndrome and thus fix the error. For each Hadamard pair of master generators, we can correct an error that affects only one Hadamard pair of single-copy generators. Multiple errors may be corrected provided that they each affect single-copy generators associated with different Hadamard pairs of master generators. If qubits in multiple rows receive errors affecting the same Hadamard pair of master generators, this procedure may be confused, so Alice and her friends may want to use a more powerful error-correcting code.

The encoded master generator must also be in the multicopy stabilizer. We can see that this is always true because each master generator contains only X 's or Z 's (and I 's), and each master generator has a Hadamard pair. An encoded generator's array contains columns corresponding to the encoded X_L (or Z_L) operators or the identity, and each row therefore contains that particular master generator, its Hadamard pair, or the identity. Every array whose rows are master generators or the identity are in the multicopy stabilizer, thus the encoded master generator is in the multicopy stabilizer. Therefore any stabilizer code can be used to purify a CSS- H state.

B. Purifying CSS states

Methods for purifying any multiparty CSS state have been obtained by Aschauer, Dür, and Briegel in Ref. [5] and by Hostens, Dehaene, and De Moor in Ref. [9]. Our goal here is to show that any CSS state can be purified by use of any CSS error-detecting or -correcting code.

Let us use for an example the three qubit GHZ state with master generators ZZI , XXX , and IZZ . This is the simplest CSS state that is not H invariant and has more than one qubit. Its graph is a three-node path. It is instructive to consider why the non-CSS code \mathcal{C}_5 is unable to purify this state. Suppose that five copies of this state are distributed to Alice,

Bob, and Charlie. The multicopy stabilizer is generated by

$$\begin{array}{c} \left[\begin{array}{c|ccc} & A & B & C \\ \hline 1 & Z & Z & I \\ 2 & I & I & I \\ 3 & I & I & I \\ 4 & I & I & I \\ 5 & I & I & I \end{array} \right], \left[\begin{array}{c|ccc} & A & B & C \\ \hline X & X & X & \\ I & I & I & \end{array} \right], \left[\begin{array}{c|ccc} & A & B & C \\ \hline I & Z & Z & \\ I & I & I & \end{array} \right], \\ \vdots \quad \vdots \quad \vdots \\ \left[\begin{array}{c|ccc} & A & B & C \\ \hline 1 & I & I & I \\ 2 & I & I & I \\ 3 & I & I & I \\ 4 & I & I & I \\ 5 & Z & Z & I \end{array} \right], \left[\begin{array}{c|ccc} & A & B & C \\ \hline I & I & I & \\ X & X & X & \end{array} \right], \left[\begin{array}{c|ccc} & A & B & C \\ \hline I & I & I & \\ I & Z & Z & \end{array} \right]. \end{array} \quad (33)$$

Using \mathcal{C}_5 as their purifying code, Alice, Bob, and Charlie measure the single-party parity checks

$$\begin{array}{c} \left[\begin{array}{c|ccc} & A & B & C \\ \hline 1 & X & I & I \\ 2 & Z & I & I \\ 3 & Z & I & I \\ 4 & X & I & I \\ 5 & I & I & I \end{array} \right], \left[\begin{array}{c|ccc} & A & B & C \\ \hline I & I & I & \\ X & I & I & \\ Z & I & I & \\ Z & I & I & \\ X & I & I & \end{array} \right], \left[\begin{array}{c|ccc} & A & B & C \\ \hline X & I & I & \\ I & I & I & \\ X & I & I & \\ Z & I & I & \\ Z & I & I & \end{array} \right], \\ \left[\begin{array}{c|ccc} & A & B & C \\ \hline 1 & Z & I & I \\ 2 & X & I & I \\ 3 & I & I & I \\ 4 & X & I & I \\ 5 & Z & I & I \end{array} \right], \dots, \left[\begin{array}{c|ccc} & A & B & C \\ \hline I & I & Z & \\ I & I & X & \\ I & I & I & \\ I & I & X & \\ I & I & Z & \end{array} \right]. \end{array} \quad (34)$$

Using these arrays, Alice, Bob, and Charlie must now construct multiparty parity checks that are in the multicopy stabilizer and sensitive to the eigenvalues of all of the generators of the multicopy stabilizer. How can they check the eigenvalue of the operator ZZI acting on the first copy? They might attempt to combine Alice's and Bob's measurement of $ZXIXZ$ to produce the multiparty parity check

$$\left[\begin{array}{c|ccc} & A & B & C \\ \hline 1 & Z & Z & I \\ 2 & X & X & I \\ 3 & I & I & I \\ 4 & X & X & I \\ 5 & Z & Z & I \end{array} \right]. \quad (35)$$

This is unfortunately not in the multicopy stabilizer because the second row, XXI , is neither one of the master generators

nor a product of some of them. The result of this measurement will then be ± 1 with probability $\frac{1}{2}$, regardless of any errors. In fact it is not possible for Alice, Bob, and Charlie to construct a set of multiparty parity checks using this code, which is sufficient for detecting errors, so their attempt at purification fails. However, in the absence of errors, in this particular case the encoded state is still a GHZ state.

Suppose instead that Alice, Bob, and Charlie use a CSS code such as \mathcal{C}_4 . Then their single-party parity checks are columns containing only X 's or Z 's. They can combine parallel single-party parity checks (by which they have measured the same generator of the purifying code) to form multiparty parity checks whose rows are repetitions of the same master generator or the identity. These are parallel elements of the multicopy stabilizer, and by construction of the error-detecting (or -correcting) code they can detect (or correct) some set of errors on the initial states. For each master generator containing X 's we have multiparty parity checks for each generator of the purifying code that contains X 's. This gives a particular syndrome pattern for diagnosing errors of that master generator on each copy. If the purifying code is one-error-correcting, we can tell which copy has received a Z (or Y) error affecting that master generator. Each master generator containing Z 's is also matched with a particular syndrome pattern by measuring the code's Z -containing generators. This can tell us which copy has received an X (or Y) affecting that master generator. When using a one-error-correcting code, this scheme can correct a single error affecting one single-copy generator associated with each of the master generators. Multiple errors can be corrected provided that they each affect single-copy generators associated with different master generators. If the error is correctable, the parties can determine a Pauli product to correct it.

The encoded logical operators also contain only X 's or Z 's and they can similarly be combined in parallel to match the master generators, so they are also in the multicopy stabilizer. This provides the method to purify any CSS state with any CSS code.

C. Purifying any stabilizer state

One can purify any stabilizer state in a manner similar to that just described for CSS- H states and CSS states; we just need to find an appropriate class of error-detecting or -correcting codes. For an independent approach to purifying any stabilizer state, see Ref. [11].

Let us use the three qubit state with master generators XZZ , ZXZ , and ZZX as an example. The graph of this state is a triangle, and it is impossible to transform it into a CSS state by use of single qubit operations. The multicopy stabilizer for this state includes all arrays with rows given by the master stabilizers. Alice, Bob, and Charlie must ensure that their purifying code can produce multiparty parity checks and encoded master generators that are in the multicopy stabilizer. They cannot use \mathcal{C}_5 because there is no method of combining the single-party parity checks to produce a nontrivial element of the multicopy stabilizer. They also cannot use \mathcal{C}_4 because the encoded logical Pauli operators form arrays whose rows are not equal to master generators and are therefore not in the multicopy stabilizer.

Suppose Alice, Bob, and Charlie use \mathcal{C}_7 as their purifying code. The multicopy stabilizer has the generators

$$\begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ \begin{array}{c} 1 \\ 2 \\ \vdots \\ 7 \end{array} & \begin{array}{ccc} X & Z & Z \\ I & I & I \\ \vdots & \vdots & \vdots \\ I & I & I \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} Z & X & Z \\ I & I & I \\ \vdots & \vdots & \vdots \\ I & I & I \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} Z & Z & X \\ I & I & I \\ \vdots & \vdots & \vdots \\ I & I & I \end{array} \end{bmatrix}, \\ \vdots \quad \vdots \quad \vdots \\ \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ \begin{array}{c} 1 \\ \vdots \\ 6 \\ 7 \end{array} & \begin{array}{ccc} I & I & I \\ \vdots & \vdots & \vdots \\ I & I & I \\ X & Z & Z \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} I & I & I \\ \vdots & \vdots & \vdots \\ I & I & I \\ Z & X & Z \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} I & I & I \\ \vdots & \vdots & \vdots \\ I & I & I \\ Z & Z & X \end{array} \end{bmatrix}. \end{aligned} \quad (36)$$

Alice measures the single-party parity checks

$$\begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ \begin{array}{c} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \end{array} & \begin{array}{ccc} X & I & I \\ I & I & I \\ X & I & I \\ I & I & I \\ X & I & I \\ I & I & I \\ X & I & I \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} I & I & I \\ X & I & I \\ X & I & I \\ I & I & I \\ I & I & I \\ X & I & I \\ X & I & I \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} I & I & I \\ I & I & I \\ I & I & I \\ X & I & I \end{array} \end{bmatrix}, \\ \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ \begin{array}{c} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \end{array} & \begin{array}{ccc} Z & I & I \\ I & I & I \\ Z & I & I \\ I & I & I \\ I & I & I \\ Z & I & I \\ Z & I & I \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} I & I & I \\ Z & I & I \\ Z & I & I \\ I & I & I \\ I & I & I \\ Z & I & I \\ Z & I & I \end{array} \end{bmatrix}, \begin{bmatrix} & \begin{array}{ccc} A & B & C \end{array} \\ & \hline & \begin{array}{ccc} I & I & I \\ I & I & I \\ I & I & I \\ Z & I & I \end{array} \end{bmatrix}. \end{aligned} \quad (37)$$

Bob and Charlie obtain similar parity checks, except that the nonidentity operators are shifted to Bob's and Charlie's columns. Each single-party parity check has a Hadamard pair. Alice, Bob, and Charlie can obtain multiparty parity checks by combining in parallel the same measurements or their Y or Z variants to match a single master generator repeated on several rows. For example, they can check the parity of

$$\left[\begin{array}{c|ccc} & A & B & C \\ \hline 1 & X & Z & Z \\ 2 & I & I & I \\ 3 & X & Z & Z \\ 4 & I & I & I \\ 5 & X & Z & Z \\ 6 & I & I & I \\ 7 & X & Z & Z \end{array} \right] \quad (38)$$

using Alice's measurement of $XIXIXIX$, Bob's measurement of $ZIZIZIZ$, and Charlie's measurement of $ZIZIZIZ$. They can use this method to obtain three multiparty parity checks for each master generator because \mathcal{C}_7 has three Hadamard pairs of generators. If they were trying to purify a state with Y 's in the master generators, they could make multiparty parity checks with columns of Y 's by use of products of Hadamard pairs. This is sufficient for Alice, Bob, and Charlie to locate and correct any single qubit error.

Now Alice, Bob, and Charlie each have a single qubit encoded in the \mathcal{C}_7 subspace, and this qubit has the encoded logical operators $X_L^{(A,B, \text{ or } C)} = XXXXXXX$ and $Z_L^{(A,B, \text{ or } C)} = ZZZZZZZ$ oriented in columns. X_L and Z_L are a Hadamard pair. The encoded master generators have arrays that contain X_L 's and Z_L 's in columns, and they match a master generator in each row. Therefore the encoded master generators are in the multicopy stabilizer, and Alice, Bob, and Charlie can use \mathcal{C}_7 to purify this state.

Alice and her friends with alphabetized names can use any error-detecting or -correcting CSS-H code to purify any stabilizer state of many qubits. The smallest example is the error-detecting code \mathcal{C}_6 , which purifies two copies of the desired state from six. The multiparty parity checks are built from Hadamard variants of the code's generators to match each of the master generators of the state they wish to purify. For each master generator they construct a number of multiparty parity checks equal to half of the number of generators of the purifying code. If the code can correct errors, these parity checks give a syndrome that is sufficient to identify which copy suffered an error affecting that master generator. They repeat the procedure diagnosing a syndrome for each master generator, so that any single qubit error is at least detected. Any error combination that affects only one single-copy generator in each set of single-copy generators associated with a single master generator can be corrected by a one-error-correcting code.

If their code is an effective error-detecting or -correcting code, they can detect or correct (insofar as the code is able to correct) errors on their qubits. The encoded master generators are formed from Hadamard variants of encoded logic operators, so they are present in the multicopy stabilizer.

IV. CONCLUSIONS AND DISCUSSION

We introduced a method for understanding entanglement purification with stabilizer codes using operator arrays. We

explained how one can purify (i) any CSS Hadamard invariant state using any error-detecting stabilizer code, (ii) any CSS state using any error-detecting CSS code, and (iii) any stabilizer state using any error-detecting CSS Hadamard invariant code. The smallest code that can purify any stabilizer state is \mathcal{C}_6 , which is an error-detecting code encoding two logical qubits in six. The smallest error-correcting code that can purify any stabilizer state is Steane's code \mathcal{C}_7 , which encodes one logical qubit in seven. These state/code combinations are sufficient for purification, because they ensure that the parties purifying the state can construct multiparty parity checks that are in the stabilizer of the copies of the state they hope to purify, and the generators describing the desired encoded state are also in the original stabilizer. We expect that the results we described here can be extended to purify states of d -dimensional quantum systems using stabilizer codes such as those described in Ref. [20] and the appropriate generalizations of CSS and CSS-H codes.

These results raise many questions. For example, it is clear that there is great freedom in choosing codes to purify states. While we have given some sufficient conditions for choosing codes, we have not studied how to match codes and states to give high efficiency of purified state production. We expect each state/code combination to have its own conditions on the required fidelity of the input states. Strategies for maximizing thresholds and efficiencies are likely to be as rich as those used in other fault-tolerant quantum information processing tasks.

We anticipate that the entanglement purification methods we have described using stabilizer codes can be translated into the languages of permutation or hashing protocols and the graph state methods of Ref. [11]. This might be accomplished using the lexicon given by Hostens, Dehaene, and De Moor in Ref. [6]. Such a translation may deepen our understanding of all of these methods.

The scheme we outlined here may also have uses that extend beyond entanglement purification. Some interesting effects can occur even when we choose state/code combinations that fail. For example, if we try to purify two copies of the non-CSS triangle state with stabilizer generators XZZ , ZXZ , and ZZX using the code \mathcal{C}_4 , we can detect any errors to the multicopy stabilizer because the parity checks are of CSS-H form. However, the encoded logical Pauli operators are not. The state obtained after the "purification" is not two copies of the triangle state, but instead an encoded entangled six qubit CSS state whose graph is shaped like a hexagon. With clever choices for states and codes we can use this procedure for a type of remote state preparation that includes error-detection and -correction capabilities.

ACKNOWLEDGMENTS

We thank John Jost and Ryan Epstein for helpful comments on this manuscript. This paper is a contribution to the National Institute of Standards and Technology and not subject to U.S. copyright.

- [1] C. H. Bennett, G. Brassard, S. Popescu, B. Schumacher, J. A. Smolin, and W. K. Wootters, Phys. Rev. Lett. **76**, 722 (1996).
- [2] C. H. Bennett, D. P. DiVincenzo, J. A. Smolin, and W. K. Wootters, Phys. Rev. A **54**, 3824 (1996).
- [3] E. N. Maneva and J. A. Smolin, in *Quantum Computation and Information*, edited by S. J. Lomonaco, Jr. and H. E. Brandt (American Mathematical Society, Providence, 2000).
- [4] R. Matsumoto, J. Phys. A **36**, 8113 (2003).
- [5] H. Aschauer, W. Dür, and H.-J. Briegel, Phys. Rev. A **71**, 012319 (2005).
- [6] E. Hostens, J. Dehaene, and B. DeMoor, e-print quant-ph/0406017.
- [7] S. Watanabe, R. Matsumoto, and T. Uyematsu, J. Phys. A **39**, 4237 (2006).
- [8] A. Miyake and H. J. Briegel, Phys. Rev. Lett. **95**, 220501 (2005).
- [9] E. Hostens, J. Dehaene, and B. DeMoor, Phys. Rev. A **73**, 042316 (2006).
- [10] Y. W. Cheong, S.-W. Lee, J. Lee, and H.-W. Lee, e-print quant-ph/0512173.
- [11] C. Kruszynska, A. Miyake, H. J. Briegel, W. Dür, e-print quant-ph/0606090.
- [12] M. A. Nielsen and I. L. Chuang, *Quantum Computation And Quantum Information* (Cambridge University Press, Cambridge, 2000).
- [13] D. Schlingemann and R. F. Werner, Phys. Rev. A **65**, 012308 (2001).
- [14] D. M. Greenberger, M. Horne, and A. Zeilinger, in *Bell's Theorem, Quantum Theory and Conceptions of the Universe*, edited by M. Kafatos (Springer, Berlin, 1989).
- [15] K. Chen and H.-K. Lo, e-print quant-ph/0404133.
- [16] D. Gottesman, Ph.D. thesis, Caltech, 1997, e-print quant-ph/9705052.
- [17] R. Laflamme, C. Miquel, J. P. Paz, and W. H. Zurek, Phys. Rev. Lett. **77**, 198 (1996).
- [18] A. Steane, Proc. R. Soc. London, Ser. A **452**, 2551 (1996).
- [19] E. Knill, Nature (London) **434**, 39 (2005).
- [20] D. Gottesman, Chaos, Solitons Fractals **10**, 1749 (1998).