Principles of Robot Autonomy I

State space dynamics – definitions and modeling





Agenda

- State space dynamics
 - Definitions
 - Modeling (kinematic and dynamic models)
 - Special case: LTI systems and linearization
- Readings
 - Chapter 1, sections 1.1 1.3 in D. Gammelli, J. Lorenzetti, K. Luo, G. Zardini, M. Pavone. *Principles of Robot Autonomy*. 2026.

State space models

- We can control a robot through the inputs to the system (e.g., motor torques, rotor thrusts, etc.)
- The *state* of a robot is a collection of variables (e.g., position, velocity) that change over time in response to the inputs
- A state space model

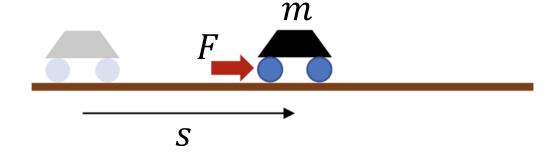
$$\dot{x}(t) = f(x(t), u(t))$$

is a mathematical description of how the state x evolves over time (i.e., \dot{x} or $dx/_{dt}$) in response to the inputs u

Example: double-integrator

- Suppose we can control the force pushing on a cart
- Newton's second law tells us that

$$F = m\ddot{s}$$



• Let x = (s, v) with $v = \dot{s}$, and u = F/m. Then we can write

$$\dot{x} = \begin{pmatrix} v \\ u \end{pmatrix} = \underbrace{\begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} x + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u}_{f(x,u)}$$

Kinematic models

- Kinematic models are mathematical models that describe the motion of a system without consideration of forces
- Kinematic models typically result from geometric constraints on the motion of a system, before considering any forces
- For example, the "unicycle" with generalized coordinates $q=(x,y,\theta)$ should not slip sideways, i.e.,

Holonomic and nonholonomic constraints

• More broadly, constraints on degrees of freedom come in various forms:

$$h(q) = 0$$

$$\underbrace{h(q) = 0}_{g(q, \dot{q}) = 0} \qquad \underbrace{g(q, \dot{q}) = 0}_{g(q, \dot{q}) = 0}$$

$$\underline{G(q)\dot{q}=0}$$

holonomic nonholonomic semi-holonomic / Pfaffian

Pfaffian constraints are a special, yet common case of nonholonomic constraints

• If G(q) has k rows (constraints) and d columns (DOFs), then

Kinematic model of the constrained system

$$\dot{q} = \sum_{j=1}^{d-k} u_j b_j(q) = [b_1(q) \quad b_2(q) \quad \cdots \quad b_{d-k}(q)] u = B(q) u_j b_j(q)$$

 $\dot{q} = \sum_{j=1}^{d-k} u_j b_j(q) = [b_1(q) \quad b_2(q) \quad \cdots \quad b_{d-k}(q)] u = B(q) u$ where $\{b_j(q)\}_{j=1}^{d-k}$ is a basis for admissible velocities, i.e., the null space of G(q).

Back to unicycle example

• The "unicycle" with DOFs $q=(x,y,\theta)$ should not slip sideways, i.e.,

$$\frac{\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} \cdot \begin{pmatrix} \sin \theta \\ -\cos \theta \end{pmatrix}}{\left[\sin \theta - \cos \theta \quad 0\right]} \dot{q} = 0$$

$$\dot{q} = \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} u_1 + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} u_2 = \begin{bmatrix} \cos \theta & 0 \\ \sin \theta & 0 \\ 0 & 1 \end{bmatrix} u$$

• Physically, $u_1=v$ is the forward velocity of the wheel, and $u_2=\omega$ is its rotational steering velocity

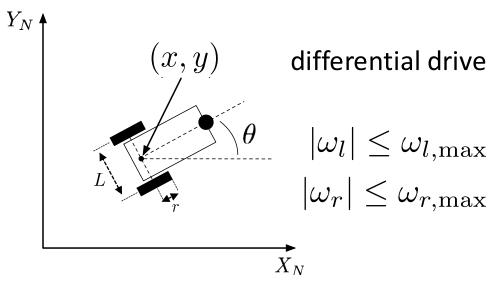
Unicycle and differential drive models

$$\begin{pmatrix} \dot{x} \\ \dot{y} \\ \dot{\theta} \end{pmatrix} = \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} v + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \omega$$

$$(x, y) \qquad \text{unicycle}$$

$$|v| \le v_{\text{max}}$$

$$\begin{pmatrix} \dot{x} \\ \dot{y} \\ \dot{\theta} \end{pmatrix} = \begin{pmatrix} \frac{r}{2}(\omega_l + \omega_r)\cos\theta \\ \frac{r}{2}(\omega_l + \omega_r)\sin\theta \\ \frac{r}{L}(\omega_r - \omega_l) \end{pmatrix}$$



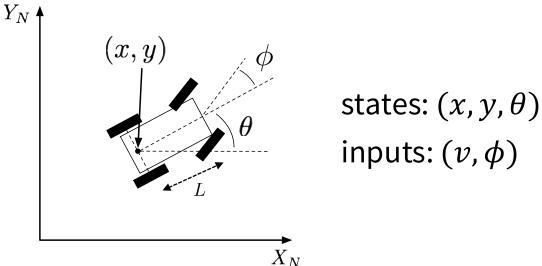
We can alternate between these kinematic models via the one-to-one input mappings:

$$v = \frac{r}{2}(\omega_r + \omega_l)$$
 $\omega = \frac{r}{L}(\omega_r - \omega_l)$

Simple car model

$$\begin{pmatrix} \dot{x} \\ \dot{y} \\ \dot{\theta} \end{pmatrix} = \begin{pmatrix} v \cos \theta \\ v \sin \theta \\ \frac{v}{L} \tan \phi \end{pmatrix}$$

$$|v| \le v_{\text{max}}, \ |\phi| \le \phi_{\text{max}} < \frac{\pi}{2}$$
 $v \in \{-v_{\text{max}}, v_{\text{max}}\}, \ |\phi| \le \phi_{\text{max}} < \frac{\pi}{2}$
 $v = v_{\text{max}}, \ |\phi| \le \phi_{\text{max}} < \frac{\pi}{2}$



- Simple car model
- Reeds-Shepp car
- Dubins car

References:

- J.-P. Laumond. Robot motion planning and control. 1998.
- S. LaValle. *Planning algorithms*. 2006.

From kinematic to dynamic models

- A kinematic state space model should be interpreted only as a subsystem of a more general dynamical model
- Improvements to the previous kinematic models can be made by placing integrators in front of action variables
- For example, for the unicycle model, one can set the speed as the integration of an action a representing acceleration, that is

$$\dot{x} = v\cos\theta, \quad \dot{y} = v\sin\theta, \quad \dot{\theta} = \omega, \quad \dot{v} = a$$

states: (x, y, θ, v) inputs: (ω, a)

Linear time-invariant models

- In general, $\dot{x} = f(x, u)$ is nonlinear, which can make it difficult to analyze
- Linear time-invariant (LTI) models take the form

$$\dot{x} = Ax + Bu$$

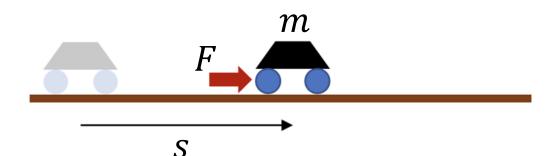
with constant matrices A and B

- For $\dot{x} = \alpha x$ with $x(0) = x_0$, the solution is $x(t) = x_0 e^{\alpha t}$. If $\alpha < 0$, the system is stable, i.e., x(t) converges to zero over time
- For $\dot{x} = Ax$ with $x(0) = x_0$, the solution is $x(t) = x_0 e^{At}$, where e^{At} is the matrix exponential
- Analogously to the scalar case, if $\operatorname{Real}(\lambda) < 0$ for each eigenvalue λ of A, then the system is stable

Example: PD control for a double-integrator

• Let x = (s, v) with $v = \dot{s}$, and u = F/m. Then

$$\dot{x} = \begin{pmatrix} v \\ u \end{pmatrix} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} x + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u$$



• Choose $u = -\kappa_p s - \kappa_d v$. Then

$$\dot{x} = \begin{bmatrix} 0 & 1 \\ -\kappa_p & -\kappa_d \end{bmatrix} x$$

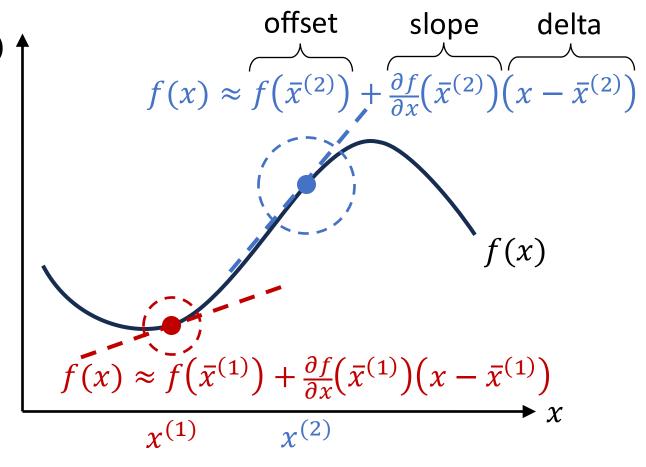
with eigenvalues $\lambda = -\frac{\kappa_d}{2} \pm \frac{1}{2} \sqrt{\kappa_d^2 - 4\kappa_p}$. If $\kappa_p > 0$ and $\kappa_d > 0$, then

 $\operatorname{Real}(\lambda) < 0$ for each eigenvalue, so the cart converges to a stand-still at s=0

• This is nice, can we use linear control tools if the system is non-linear?

Linearization

- Linearization approximates a f(x) of nonlinear function f near \bar{x} by a line, i.e., linear function
- The "slope" of the line is the derivative of f at \bar{x} . The change in $f(\bar{x})$ near \bar{x} is the slope multiplied by the distance from \bar{x}
- The quality of the approximation can vary with the linearization point \bar{x} and distance from \bar{x}



Linearization of non-linear state-space models

• For the nonlinear system $\dot{x} = f(x, u)$, the linearization around (\bar{x}, \bar{u}) is

$$\dot{x} \approx f(\bar{x}, \bar{u}) + \underbrace{\frac{\partial f}{\partial x}(\bar{x}, \bar{u})}_{}(x - \bar{x}) + \underbrace{\frac{\partial f}{\partial u}(\bar{x}, \bar{u})}_{}(u - \bar{u})$$

Since x and u can be vectors, we generalize derivatives to Jacobian matrices

• If (\bar{x}, \bar{u}) is an *equilibrium*, i.e., $f(\bar{x}, \bar{u}) = 0$, we can consider an LTI approximation of the system near (\bar{x}, \bar{u}) , with state $\Delta x = x - \bar{x}$ and input $\Delta u = u - \bar{u}$:

$$\dot{\Delta x} = A\Delta x + B\Delta u$$

• When (x, u) is near (\bar{x}, \bar{u}) , we can use tools from linear systems analysis and control on nonlinear systems -- more on this later with LQR control!

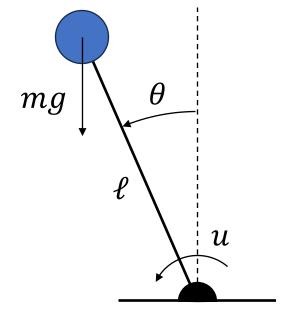
Example: Inverted pendulum

• The dynamics are described by $m\ell^2\ddot{\theta} = mg\ell\sin\theta + u$. In state space form with $x = (\theta, \dot{\theta})$, they are

$$\dot{x} = f(x, u) = \begin{pmatrix} x_2 \\ \frac{g}{\ell} \sin x_1 + \frac{1}{m\ell^2} u \end{pmatrix}$$

• Since (x, u) = 0 is an equilibrium, the linearization here is

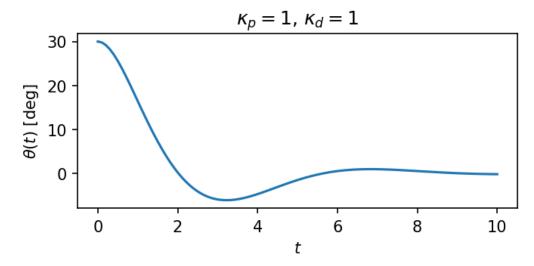
$$\dot{x} \approx \begin{pmatrix} \dot{\theta} \\ \frac{g}{\ell}\theta + u \end{pmatrix} = \begin{bmatrix} 0 & 1 \\ \frac{g}{\ell} & 0 \end{bmatrix} x + \begin{bmatrix} 0 \\ \frac{1}{m\ell^2} \end{bmatrix} u$$



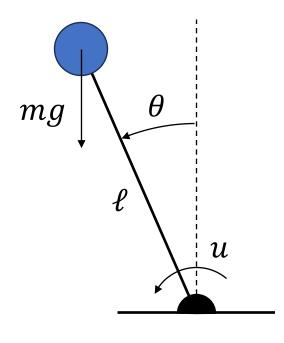
• This is close to a double-integrator! We could try $\frac{1}{m\ell^2}u = -\left(\frac{g}{\ell} + \kappa_p\right)\theta - \kappa_d\dot{\theta}$ to stabilize the pendulum near the upright equilibrium

Example: Inverted pendulum

• We try $\frac{1}{m\ell^2}u=-\left(\frac{g}{\ell}+\kappa_p\right)\theta-\kappa_d\dot{\theta}$ to stabilize the pendulum near the upright equilibrium:



 We will later discuss how we actually simulate this system on a computer



Next time

