

· 例子: 谐振子

$$E = \sum_{i=1}^N \frac{p_i^2}{2m}$$

实际上的能级是离散的, 但能级间隔很小时, 可以认为是连续.

对应的条件:  $k_B T \gg \Delta E$

计算配分函数:

$$\begin{aligned} Z &= \int \frac{\prod_{i=1}^N dp_i d^3 q_i}{h^{3N} N!} \cdot e^{-\beta \sum_{i=1}^N \frac{p_i^2}{2m}} \\ &= \frac{V^N}{N! h^{3N}} \cdot \left[ \int dp \cdot e^{-\beta \frac{p^2}{2m}} \right]^{3N} \\ &= \frac{V^N}{N! h^{3N}} \cdot \left( \frac{2\pi m}{\beta} \right)^{\frac{3N}{2}} = \frac{V^N}{N! \lambda_T^{3N}}, \quad \lambda_T := \frac{h}{\sqrt{2\pi m k_B T}}, \text{ 热波长} \end{aligned}$$

这里实际上我们并不需要计算  $\sum \Omega(E) \cdot e^{-\beta E}$ ! 只需要用 [2.9.3] 把 E 表示出来, 这里  $\Omega(E)$  主要用于澄清概念.

$$\begin{aligned} \ln Z &= N \cdot \ln \left( \frac{V}{h^3} \left( \frac{2\pi m}{\beta} \right)^{\frac{3}{2}} \right) - N \ln N + N \\ &= N \cdot \ln \left( \frac{V}{N} \left( \frac{2\pi m}{\beta h^3} \right)^{\frac{3}{2}} \right) + N. \end{aligned}$$

热力学量:

$$\langle E \rangle = -\frac{\partial}{\partial \beta} \ln Z = -N \cdot \left( -\frac{\frac{3}{2}}{\beta} \right) = \frac{3}{2} N k_B T.$$

$$P = \frac{1}{\beta} \frac{\partial}{\partial V} \ln Z = \frac{N}{\beta} \cdot \frac{1}{V} = \frac{N k_B T}{V}, \quad \Rightarrow pV = N k_B T$$

$$\begin{aligned} S &= k_B (\ln Z - \beta \frac{\partial}{\partial \beta} \ln Z) \\ &= N k_B \left[ \ln \left( \frac{V}{N} \left( \frac{2\pi m}{\beta h^3} \right)^{\frac{3}{2}} \right) + \frac{5}{2} \right]. \end{aligned}$$

· 例子: 二能级系统

$$\text{伊辛磁体} \quad \sigma_i = \pm 1, \quad \mathcal{H} = J \sum_{i=1}^N \sigma_i$$

$$\text{能量} \quad \varepsilon_i = -\mu_B \sigma_i$$

$$\text{总能量} \quad E = \sum_{i=1}^N \varepsilon_i = -\mu_B \sum_{i=1}^N \sigma_i$$

$$\begin{aligned} Z &= \sum_{\{\varepsilon_i\}} e^{-\beta \sum_{i=1}^N \varepsilon_i} \quad (\text{对所有可能状态求和}) \\ &= \sum_{\{\sigma_i\}} e^{\beta \mu_B \sum_{i=1}^N \sigma_i} \\ &= \sum_{\{\sigma_i\}} \prod_{i=1}^N e^{\beta \mu_B \sigma_i} \quad (\text{这一步是把对 } 2^N \text{ 项求和拆成 } N \text{ 个 } 2 \text{ 项的积}) \\ &= \prod_{i=1}^N \left( \sum_{\sigma_i = \pm 1} e^{\beta \mu_B \sigma_i} \right) = [2 \cosh(\beta \mu_B)]^N. \end{aligned}$$

$$\ln Z = N \cdot \ln(2 \cosh(\beta \mu_B))$$

热力学量:

$$\begin{aligned} E &= -\frac{\partial}{\partial \beta} \ln Z = -N \mu_B \cdot \tanh(\beta \mu_B) \\ &= -N \mu_B \cdot \frac{e^{\beta \mu_B} - e^{-\beta \mu_B}}{e^{\beta \mu_B} + e^{-\beta \mu_B}} \\ &= N \cdot \frac{(-\mu_B) \cdot e^{-(-\beta \mu_B)} + (\mu_B) \cdot e^{-\beta \mu_B}}{e^{-(-\beta \mu_B)} + e^{-(-\beta \mu_B)}} \\ &= N \cdot \langle \varepsilon \rangle. \end{aligned}$$

式中:  $\langle \varepsilon \rangle$  为单个自旋的平均能量.

$$\langle \varepsilon \rangle = P_+ \varepsilon_+ + P_- \varepsilon_-$$

$$P_{\pm} = \frac{e^{\mp \beta \mu_B}}{Z}, \quad Z := e^{\beta \mu_B} + e^{-\beta \mu_B}, \quad \text{单个自旋配分函数.}$$

· 注意: 这两个例子, 都是  $N$  个互相无作用的自旋!

$$\text{这种情况下: } E = \sum_{i=1}^N \varepsilon_i.$$

$$\begin{aligned} Z &= \sum_{\{\varepsilon_i\}} e^{-\beta \sum_{i=1}^N \varepsilon_i} = \prod_{i=1}^N \left( \sum_{\varepsilon_i} e^{-\beta \varepsilon_i} \right) \\ &= Z^N. \end{aligned}$$

$\Rightarrow$  体系总的配分函数, 等于单自旋配分函数之积!

一对可全同粒子:  $Z = \frac{1}{N!} Z^N$ . 例: 理想气体,  $Z = \frac{V}{\lambda_T^3}$ .

能量与平均能量 (Sloppy Language)

$\langle E \rangle$ ,  $E^*$ , 有多接近?

看  $E$  的涨落:  $\langle E^2 \rangle - \langle E \rangle^2$ . (记作  $\langle E^2 \rangle_c$ )

$$\langle E \rangle = -\frac{\partial}{\partial \beta} \ln Z.$$

$$\frac{\partial^2}{\partial \beta^2} \ln Z = \frac{\partial}{\partial \beta} \left( \frac{1}{Z} \frac{\partial Z}{\partial \beta} \right)$$

$$= \frac{1}{Z} \frac{\partial^2 Z}{\partial \beta^2} - \frac{1}{Z^2} \left( \frac{\partial Z}{\partial \beta} \right)^2$$

$$= \langle E^2 \rangle - \langle E \rangle^2$$

$$\Rightarrow \langle E^2 \rangle_c = \frac{\partial^2}{\partial \beta^2} \ln Z$$

$$= -\frac{\partial \langle E \rangle}{\partial \beta} \quad \left( \frac{\partial}{\partial \beta} = \frac{\partial}{\partial T} \frac{\partial T}{\partial \beta} = -k_B T^2 \frac{\partial}{\partial T} \right)$$

$$= k_B T^2 \cdot \frac{\partial \langle E \rangle}{\partial T} = k_B T^2 C_V \propto N.$$

$$\text{相对涨落: } \frac{\sqrt{\langle E^2 \rangle_c}}{\langle E \rangle} \sim \frac{\sqrt{N}}{N} = \frac{1}{\sqrt{N}} \rightarrow 0, N \rightarrow \infty.$$

在概率统计当中,  $Z$  为“特征函数”, 是“矩”的生成函数,  $\ln Z$  生成的是“中心矩”.

继续求:

$$\langle E^3 \rangle_c = (-1)^3 \left( \frac{\partial}{\partial \beta} \right)^3 \ln Z \propto N. \quad [Z = \sum_n e^{-\beta E_n} \text{ 和概率统计中矩的生成函数 (特征函数) 有相同的形式}]$$

中心极限定理:

对于  $N$  个独立同分布的随机变量  $x_i$ , ( $x_i$  的均值为 0, 方差为 1)

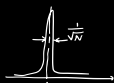
$$X = \sum_{i=1}^N x_i, \quad \langle X \rangle = N \langle x \rangle.$$

$$\langle X^2 \rangle_c = N \langle x^2 \rangle_c.$$

$$\langle X^n \rangle_c = N \langle x^n \rangle_c \propto N.$$

定义  $y = \frac{1}{N} X$ , 则当  $N \rightarrow \infty$  时,

$$p(y) \simeq \frac{1}{\sqrt{2\pi \frac{\langle x^2 \rangle_c}{N}}} \cdot e^{-\frac{(y - \langle x \rangle)^2}{2 \frac{\langle x^2 \rangle_c}{N}}}.$$



[高维性的证明用到了“特征函数”, 核心是说: 把  $e^{itx}$  展开, 领头非零项是  $-\frac{1}{2} \frac{t^2}{N!} x^2$ .  $N \rightarrow \infty$  时又和这一致了]

[和之前讨论的“勒让德分”很相似:]

$$\langle X^n \rangle_c = \langle (X - \langle X \rangle)^n \rangle$$

$$= \left\langle \sum_{k=0}^n \binom{n}{k} X^k (-\langle X \rangle)^{n-k} \right\rangle$$

$$= \sum_{k=0}^n \binom{n}{k} \langle X^k \rangle (-1)^{n-k} \langle X \rangle^{n-k}$$

设  $X$  的均值为 0  $\Rightarrow$  中值为  $\langle X \rangle = 0$

$$\langle X^n \rangle \propto N?$$

$$\langle X^n \rangle \propto N$$

$$f(x), F(t) = \int_{-\infty}^{+\infty} f(x) e^{itx} dx$$

$$F^{(n)}(t) = \int_{-\infty}^{+\infty} (-ix)^n f(x) e^{-itx} dx$$

$$\frac{F^{(n)}(0)}{n!} = (-i)^n \langle X^n \rangle$$

$$p(E_n) = \frac{e^{-\beta E_n}}{Z}$$

$$\langle E \rangle = \frac{1}{Z} \sum E_n e^{-\beta E_n}$$

$$= \frac{1}{Z} \cdot \frac{\partial}{\partial \beta} Z$$

$$\langle E^2 \rangle = \sum (E_n - \langle E \rangle)^2 e^{-\beta E_n}$$

$$\langle E \rangle = \int \mathcal{N}(E) dE \cdot e^{-\beta E} \cdot \frac{1}{Z}$$

$$F(\beta) = \int \frac{\mathcal{N}(E)}{Z} e^{-\beta E} dE$$

子完全变量地，能级相解

E.

$$E. \varphi_E(t) = \int f(E) e^{iEt} dt.$$

$$y = \sum x_i$$

$$\langle y \rangle = N \langle x \rangle$$

$$\langle y^2 \rangle = 0$$

$$\langle y^2 \rangle_c = \left\langle \sum_i (x_i - \langle x \rangle)^2 \right\rangle$$

$$= N \langle x^2 \rangle + N \langle x \rangle^2 - 2N \langle x \rangle^2 \propto N$$

$$= N (\langle x^2 \rangle - \langle x \rangle^2) \propto N$$

$$\left\langle \sum_i (x_i - \langle x \rangle)^n \right\rangle$$

$$= N \cdot \left( \sum_{k=0}^n (-1)^{n-k} \langle x^k \rangle \langle x \rangle^{n-k} \right)$$

$$\varphi_y(t) = \int_{-\infty}^{+\infty} f(y) e^{-ity} dy$$

卷积  $\rightarrow$  卷积,  $\mu=0, \sigma \approx 1$

$$\varphi_N(t) = \int_{-\infty}^{+\infty} f(x) e^{-itx} dx$$

$$x \rightarrow x/N$$

$$f(x) dx = f\left(\frac{x}{N}\right) d\left(\frac{x}{N}\right) \quad f(y) = f(x_1) + f(x_2) + \dots + f(x - \sum x_i)$$

$$\Rightarrow f\left(\frac{x}{N}\right) = N \cdot f(x)$$

$$\varphi_{\frac{x}{N}}(t) = \int_{-\infty}^{+\infty} f\left(\frac{x}{N}\right) e^{-it \frac{x}{N}} d\frac{x}{N} \Rightarrow \varphi_y(t) = (\varphi_N(t))^N$$

$$\int \left(1 - it \frac{x}{N} - \frac{1}{2} t^2 \frac{x^2}{N^2} + O\left(\frac{1}{N^3}\right)\right) f(x) dx$$

$$= 1 - \frac{1}{2} t^2 \frac{1}{N^2} \sigma^2 = 1 - \frac{t^2}{2N^2} + O\left(\frac{1}{N^3}\right)$$

$$\left(1 - \frac{t^2}{2N^2}\right)^N \rightarrow \left(e^{-\frac{t^2}{2N}}\right)^N \Rightarrow \sigma = \frac{1}{N}, \sigma^2 = \frac{1}{N}$$

$$\Rightarrow y \sim e^{-\frac{t^2}{2} N}$$

$$y \sim e^{-N \frac{t^2}{2}} e^{-\dots}$$

$$- \frac{x^2}{2\sigma^2} - itx$$

$$= \int e^{-\frac{1}{2}(x+it)^2} dx e^{-\frac{t^2}{2}}$$

$$= \frac{1}{\sqrt{2\pi}} (x+it\sigma)^2 - \frac{t^2\sigma^2}{2}$$