

### III 经典气体

• 理想气体:  $H = \sum_{i=1}^N \frac{\vec{p}_i^2}{2m}$

$$Z(T, V, N) = \frac{1}{N! h^{3N}} \int \prod_{i=1}^N d^3 p_i d^3 q_i e^{-\beta \sum_{i=1}^N \frac{\vec{p}_i^2}{2m}} = \frac{1}{N!} \left( \frac{V}{\lambda^3} \right)^N, \quad \lambda = \frac{h}{\sqrt{2\pi m k_B T}}$$

$$\langle E \rangle = -\frac{\partial}{\partial \beta} \ln Z = \frac{3}{2} N k_B T, \quad \langle E \rangle = \sum_n P_n \langle E_n \rangle$$

#### • 能均分定理

该子系统的哈密顿量可以写为

$$H = \sum_{a=1}^M A_a q_a^2 + B_a p_a^2 \quad [\text{注意这里 } p, q \text{ 是广义量}]$$

问:  $\langle A_a q_a^2 \rangle = ?$

$$\begin{aligned} \langle A_a q_a^2 \rangle &= \frac{\int \prod_b d p_b d q_b A_a q_a^2 e^{-\beta \sum (A_b q_b^2 + B_b p_b^2)}}{\int \prod_b d p_b d q_b e^{-\beta \sum (A_b q_b^2 + B_b p_b^2)}} \\ &= \frac{\int d q_a A_a q_a^2 e^{-\beta A_a q_a^2}}{\int d q_a e^{-\beta A_a q_a^2}} \\ &= -\frac{\partial}{\partial \beta} \ln \left[ \int d q_a e^{-\beta A_a q_a^2} \right] \\ &= -\frac{\partial}{\partial \beta} \ln \sqrt{\frac{\pi}{\beta A_a}} = \frac{1}{2} \frac{1}{\beta} = \frac{1}{2} k_B T \end{aligned}$$

同理,

$$\langle B_a p_a^2 \rangle = \frac{1}{2} k_B T$$

能均分定理: 哈密顿量中的独立平方项, 平均值为  $\frac{1}{2} k_B T$ .

例: 理想气体  $E = 3N \cdot \frac{1}{2} k_B T = \frac{3}{2} N k_B T$  (单原子)

#### • 双原子分子气体

$O_2$  振动, 转动, ...

$$H = H_{com} + H_{vib} + H_{rot}$$

质心 振动 转动

$$Z = \underline{Z}_{com} \cdot Z_{vib} \cdot Z_{rot}$$

↳ 仅前级级含  $V \Rightarrow pV = N k_B T$  级成立 (推导过程相同)

每个粒子的平均能量?

$$\text{质心: } E_{com} = \frac{3}{2} k_B T \quad C_{com} = \frac{3}{2} k_B$$

$$\text{振动: } H_{vib} = \sum_{i=1}^N \frac{p_i^2}{2\mu} + \frac{1}{2} k \omega^2 q_i^2$$

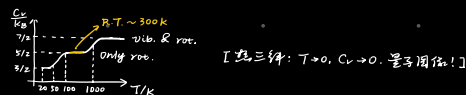
$$E_{vib} = 2 \cdot \frac{1}{2} k_B T = k_B T, \quad C_{vib} = k_B$$

$$\text{转动: } H_{rot} = \sum_{i=1}^N \frac{L_i^2}{2I}, \quad \text{两个方向 (不可忽略)}$$

$$E_{rot} = 2 \cdot \frac{1}{2} k_B T = k_B T, \quad C_{rot} = k_B$$

热容 单原子  $C_V = \frac{3}{2} k_B$  双原子  $C_V = \frac{7}{2} k_B$

实验 ( $H_2$ )



自由度随着降温依次“冻结”!

经典无法解释  $\Rightarrow$  量子效应! [量子效应的“冻结”要求是相对能量尺度而言的!]

↳ 解释: 类似“光电效应” 能量有“阈值”! “热激发”有门槛!

接下来详细分析

#### • 振动

$$E_n = (n + \frac{1}{2}) \hbar \omega, \quad n = 0, 1, 2, \dots$$

库特子配分函数:

$$Z_{vib} = \sum_{n=0}^{\infty} e^{-\beta E_n} = e^{-\frac{1}{2} \beta \hbar \omega} \frac{1}{1 - e^{-\beta \hbar \omega}}$$

$$\ln Z_{vib} = -\frac{1}{2} \beta \hbar \omega - \ln(1 - e^{-\beta \hbar \omega})$$

$$E_{vib} = -\frac{\partial}{\partial \beta} \ln Z_{vib} = \frac{1}{2} \hbar \omega + \frac{\hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}}$$

$$C_{vib} = \frac{dE_{vib}}{dT} = -\frac{1}{k_B T^2} \cdot \frac{d}{d\beta} \left( \frac{\hbar \omega}{e^{\beta \hbar \omega} - 1} \right)$$

$$= \frac{(\hbar \omega)^2}{k_B T^2} \cdot \frac{e^{\beta \hbar \omega}}{(e^{\beta \hbar \omega} - 1)^2}$$

$$= k_B \cdot \frac{(\beta \hbar \omega)^2 e^{\beta \hbar \omega}}{(e^{\beta \hbar \omega} - 1)^2}$$

决定性质的是  $\beta \hbar \omega$ !

- 高温极限:  $\beta \hbar \omega \ll 1$ .

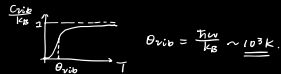
$$C_{vib} \simeq k_B \cdot \frac{(\beta \hbar \omega)^2}{(\beta \hbar \omega)^2} \cdot 2 = k_B. \text{ 回归经典.}$$

$$\text{定义 } \frac{\hbar \omega}{k_B T} = \frac{\Theta_{vib}}{T}. \quad \Theta_{vib}: \text{振动特征温度.}$$

- 低温极限:  $\beta \hbar \omega \gg 1$ .

$$C_{vib} \simeq k_B \cdot \frac{(\beta \hbar \omega)^2 e^{\beta \hbar \omega}}{e^{2\beta \hbar \omega}} = k_B \cdot (\beta \hbar \omega)^2 \cdot \underline{e^{-\beta \hbar \omega}}. \text{ 指数衰减!}$$

它的图像如下:



• 转动

$$L^2 \text{ 量子化. } L^2 = l(l+1)\hbar^2, \quad l=0,1,2,\dots$$

$$\text{简并度 } g_l = 2l+1.$$

配分函数:

$$Z_{rot} = \sum_{l=0}^{\infty} (2l+1) \cdot e^{-\frac{\beta}{2I} \hbar^2 l(l+1)} \quad (\text{子球})$$

- 高温极限:  $\frac{\beta \hbar^2}{2I} \ll 1$ .

$$\text{定义: } \Theta_{rot} = \frac{\hbar^2}{2Ik_B}, \quad \text{转动特征温度.}$$

$$\text{令 } x = l(l+1), \quad dx = 2l+1.$$

$$Z_{rot} \simeq \int_0^{+\infty} dx \cdot e^{-\frac{\Theta_{rot}}{T} x} = \frac{T}{\Theta_{rot}}.$$

$$E_{rot} \simeq -\frac{\partial}{\partial \beta} \ln Z_{rot} = k_B T, \quad C_{rot} \simeq k_B.$$

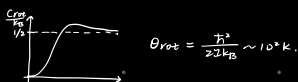
- 低温极限:  $\frac{\beta \hbar^2}{2I} \gg 1$ .

$$Z_{rot} \simeq 1 + 3 \cdot e^{-\frac{\Theta_{rot}}{2T}}$$

$$E_{rot} \simeq -\frac{\partial}{\partial \beta} \ln Z_{rot} = \frac{\hbar^2}{2I} \cdot \frac{3e^{-\beta \hbar^2/2I}}{1 + 3e^{-\beta \hbar^2/2I}}$$

$$C_{rot} \simeq \frac{dE}{dT} = 3k_B \cdot \left( \frac{2\Theta_{rot}}{T} \right)^2 e^{-2\Theta_{rot}/T}$$

它的图像如下:



• 一些说明

- 化学键本质上束缚电子, 原子核相互作用, 为什么考虑电子的能量?

$$\text{电子能量量级: } E_e \sim \text{eV}, \quad \underline{\text{eV} \sim 10^4 \text{ K.}} \text{ 远高于室温!}$$



$$E_e \simeq \frac{\hbar^2}{2ma^2}, \quad E_{rot} \simeq \frac{\hbar^2}{2I} \sim \frac{\hbar^2}{2Ma^2}$$

$$\frac{E_{rot}}{E_e} \simeq \frac{m}{M} \sim 10^{-4}, \quad \Rightarrow E_{rot} \sim 1 \text{ K!}$$

$$E_{vib} \simeq \frac{1}{2} \hbar \omega^2 \delta^2, \quad \frac{E_{vib}}{E_e} \simeq \frac{\hbar \omega}{\frac{\hbar^2}{2ma^2}} \sim m\omega \sim \sqrt{\frac{m}{M}} \sim 10^{-2}$$

$$\delta = a. \text{ 化学键键长, } \underline{E_{vib} \simeq E_e} \quad \Rightarrow \quad \omega \simeq \frac{1}{\sqrt{mM}}$$