

• 独立气体的热力学性质

重新考虑 $k=0$ 项:

$$\ln Q = -\frac{V}{(2\pi)^3} \int d^3k \cdot \ln(1 - ze^{-\beta \epsilon_k}) = \underbrace{\ln(1-z)}_{k=0 \text{ 项, } O(\ln N), \text{ 可去!}}$$

$$= \frac{V}{\lambda^3} g_{5/2}(z).$$

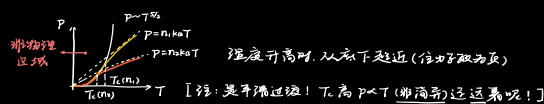
$$PV = k_B T \ln Q \Rightarrow P = \frac{k_B T}{\lambda^3} g_{5/2}(z).$$

当 $T < T_c: z \rightarrow 1$.

$$P = \frac{k_B T}{\lambda^3} \zeta\left(\frac{5}{2}\right) \simeq 1.341 \cdot \frac{k_B T}{\lambda^3}. \quad \text{与 } n \text{ 无关! (注: } T_c \text{ 依赖于 } \rho \text{ 或 } \mu)$$

$$P \sim T^{5/2} = T^{5/2}.$$

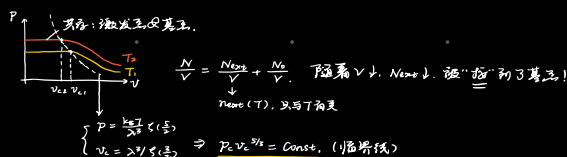
解释: 基态粒子对压强没有贡献, 只有激发态粒子(只与 T 相关)有贡献!



$P \sim V^{-5/3}$ 更甚.

$$n \lambda^3 = \zeta\left(\frac{5}{2}\right).$$

固定 T , 减小 V ($n \uparrow$) \Rightarrow 临界温度 $n_c(T)$, $v_c(T)$.



相变潜热

$$\frac{dP}{dT} = \frac{L}{T(v_2 - v_1)}, \quad \begin{array}{l} 1: \text{激发态} \\ 2: \text{基态 } v_2 = 0 \end{array}$$

$$= -\frac{L}{T v_c}$$

$$P = \frac{k_B T}{\lambda^3} \zeta\left(\frac{5}{2}\right) \Rightarrow \frac{dP}{dT} = \frac{5}{2} \frac{P}{T}$$

$$L = -\frac{5}{2} P v_c = -\frac{5}{2} k_B T \cdot \frac{\zeta(5/2)}{\zeta(3/2)} < 0.$$

相变潜热为负: 激 \rightarrow 凝相变潜热, 合理!

内能与比热

$$E = \frac{3}{2} k_B T \frac{V}{\lambda^3} g_{3/2}(z).$$

$$T \leq T_c \text{ 时: } E = \frac{3}{2} k_B T \cdot \frac{V}{\lambda^3} \zeta\left(\frac{3}{2}\right)$$

$$C_V = \frac{dE}{dT} = \frac{5}{2} \frac{E}{T} = \frac{15}{4} k_B \cdot \frac{V}{\lambda^3} \zeta\left(\frac{3}{2}\right).$$

$T > T_c$ 时:

$$C_V = \frac{15}{4} k_B \cdot \frac{V}{\lambda^3} g_{3/2}(z) + \frac{3}{2} k_B T \cdot \frac{V}{\lambda^3} g'_{3/2}(z) \cdot \frac{dz}{dT}$$

$$g_m(z) = \sum_{l=1}^{\infty} \frac{z^l}{l^m}, \quad g'_m(z) = \sum_{l=1}^{\infty} \frac{z^{l-1}}{l^{m-1}} = \frac{1}{z} g_{m-1}(z).$$

$$\frac{dz}{dT} = ? \quad \frac{dN}{dT} = \frac{d}{dT} \left(\frac{V}{\lambda^3} g_{3/2}(z) \right) = 0 \Rightarrow \frac{dz}{dT} = 0.$$

